A TALE OF TWO PORTALS: LIGHT, NEW PHYSICS AT FUTURE E+E- COLLIDERS

Felix Yu
JGU Mainz

with Jia Liu, Xiao-Ping Wang, JHEP **1706** (2017) 077 [1704.00730]

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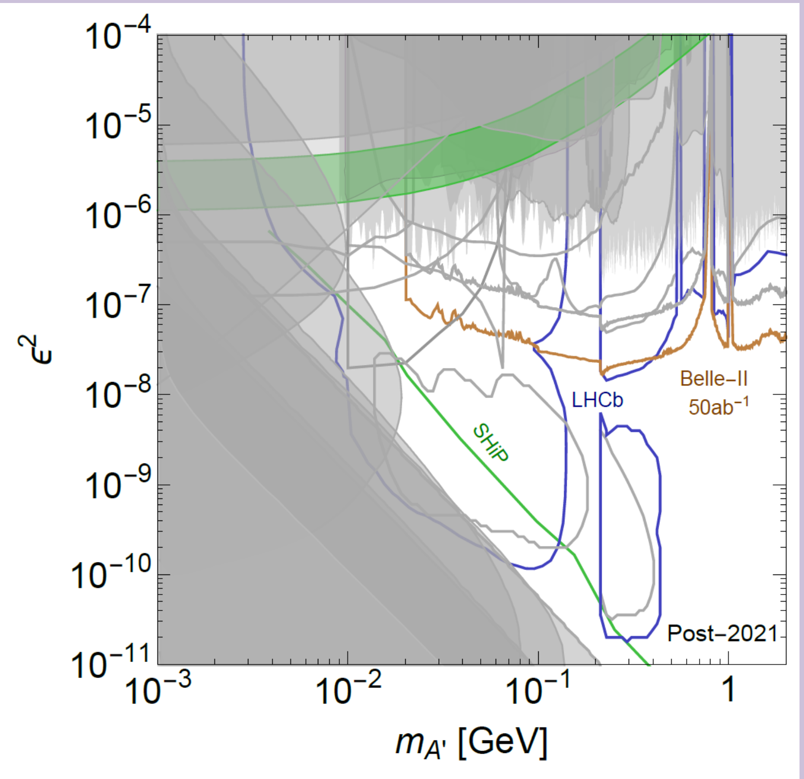
Introduction and Motivation

- After Higgs discovery, particle physics has entered a distinct, exploratory phase
  - Outstanding problems, including dark matter and naturalness, are more acute
  - Mass scales of new physics unknown
- In conjunction, couplings of new physics particles are also unknown
    - TeV-scale strongly-coupled particles with prompt, cascade decays are strongly constrained
    - Weak-scale, weakly-coupled particles less constrained
    - Very weakly-coupled particles are very weakly constrained

Portal couplings

- Given direct probes at a given energy scale, sensitivity to UV scales follows NDA
  - Renormalizable, “portal” couplings are few (e.g. scalar Higgs portal, neutrino portal, vector portal, axion portal)

Alexander, et al. [1608.08632]



- Nevertheless, Higgs factory energies probe new portal couplings at mass scales untested by beam-dump experiments or LHC

Future e+e- machines can produce new particles ***directly*** *here*

Outline

- Theory review: Double Dark Portal
  - Simultaneous kinetic mixing and scalar Higgs portal
- Phenomenology: dark matter probes
  - Direct detection and indirect detection probes
- Phenomenology: collider signatures
  - Unique capabilities of e^+e^- machine for probing dark vector, dark scalar production
- Conclusions

Double Dark Portal model

Kinetic mixing of *K* with hypercharge gauge boson *B*

$$\mathcal{L} \supset -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} - \frac{1}{4}W^{i}_{\mu\nu}W^{i\,\mu\nu} - \frac{1}{4}K_{\mu\nu}K^{\mu\nu} + \frac{\epsilon}{2\cos\theta_W}B_{\mu\nu}K^{\mu\nu}$$
$$+ |D_\mu H|^2 + |D_\mu \Phi|^2 + \mu_H^2|H|^2 - \lambda_H|H|^4 + \mu_D^2|\Phi|^2 - \lambda_D|\Phi|^4 - \lambda_{HP}|H|^2|\Phi|^2$$
$$+ \bar{\chi}(i\not{D} - m_\chi)\chi$$

$U(1)_D$ charges $\Phi \sim +1 \,, \; \chi \sim +1$

Scalar Higgs portal between dark Higgs Φ and SM *H*

- Two marginal operators: simultaneous vector portal and scalar portal couplings
  - Constraints driven by searches, not known from first principles (possible in UV completions)

Double Dark Portal model

- Steps for solving the neutral vector Lagrangian (pedagogical)
  - Diagonalize gauge boson mass matrix
    - Usual $t_W = g'/g$ rotation corresponds to

$$\mathcal{L} \supset \frac{-1}{4} \begin{pmatrix} Z^{\mu\nu}_{\text{SM}} & A^{\mu\nu}_{\text{SM}} & K^{\mu\nu} \end{pmatrix} \begin{pmatrix} 1 & 0 & \epsilon t_W \\ 0 & 1 & -\epsilon \\ \epsilon t_W & -\epsilon & 1 \end{pmatrix} \begin{pmatrix} Z_{\mu\nu,\text{ SM}} \\ A_{\mu\nu,\text{ SM}} \\ K_{\mu\nu} \end{pmatrix}$$

$$+ \frac{1}{2} \begin{pmatrix} Z^{\mu}_{\text{SM}} & A^{\mu}_{\text{SM}} & K^{\mu} \end{pmatrix} \begin{pmatrix} m^2_{Z,\text{ SM}} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & m^2_K \end{pmatrix} \begin{pmatrix} Z_{\mu,\text{ SM}} \\ A_{\mu,\text{ SM}} \\ K_{\mu} \end{pmatrix}$$

  - Require |ε|< c_W for positive kinetic mixing determinant
  - Field strengths are Abelian kinetic terms, non-Abelian interactions inherited from transformations

Double Dark Portal model

- Steps for solving the neutral vector Lagrangian (pedagogical)
    - Remove kinetic mixing and canonically normalize

$$U_1 = \begin{pmatrix} 1 & 0 & 0 \\ -\epsilon^2 t_W & 1 & \epsilon \\ -\epsilon t_W & 0 & 1 \end{pmatrix} \qquad U_2 = \begin{pmatrix} \sqrt{\dfrac{1-\epsilon^2}{1-\epsilon^2 c_W^{-2}}} & 0 & 0 \\ 0 & 1 & 0 \\ \dfrac{-\epsilon^3 t_W}{\sqrt{(1-\epsilon^2)(1-\epsilon^2 c_W^{-2})}} & 0 & \dfrac{1}{\sqrt{1-\epsilon^2}} \end{pmatrix}$$

$$\mathcal{L} \supset \frac{-1}{4} \begin{pmatrix} Z^{\mu\nu}_{\text{SM}} & A^{\mu\nu}_{\text{SM}} & K^{\mu\nu} \end{pmatrix} (U_1^T)^{-1} (U_2^T)^{-1} \mathbb{I}_3 U_2^{-1} U_1^{-1} \begin{pmatrix} Z_{\mu\nu,\ \text{SM}} \\ A_{\mu\nu,\ \text{SM}} \\ K_{\mu\nu} \end{pmatrix}$$

$$+ \frac{1}{2} \begin{pmatrix} Z^{\mu}_{\text{SM}} & A^{\mu}_{\text{SM}} & K^{\mu} \end{pmatrix} (U_1^T)^{-1} (U_2^T)^{-1} \begin{pmatrix} \dfrac{m^2_{Z,\ \text{SM}}(1-\epsilon^2)^2 + m_K^2 \epsilon^2 t_W^2}{(1-\epsilon^2)(1-\epsilon^2 c_W^{-2})} & 0 & \dfrac{-m_K^2 \epsilon t_W}{(1-\epsilon^2)\sqrt{1-\epsilon^2 c_W^{-2}}} \\ 0 & 0 & 0 \\ \dfrac{-m_K^2 \epsilon t_W}{(1-\epsilon^2)\sqrt{1-\epsilon^2 c_W^{-2}}} & 0 & \dfrac{m_K^2}{1-\epsilon^2} \end{pmatrix}$$

$$\times U_2^{-1} U_1^{-1} \begin{pmatrix} Z_{\mu,\ \text{SM}} \\ A_{\mu,\ \text{SM}} \\ K_{\mu} \end{pmatrix}$$

Double Dark Portal model

- Steps for solving the neutral vector Lagrangian (pedagogical)
  - Rediagonalize mass matrix via Jacobi rotation (exact)
  - To O(ε³), masses and fields are

$$m_{\tilde{K}}^2 = m_K^2 + \frac{m_K^2 c_W^{-2} \epsilon^2 (m_{Z,\text{ SM}}^2 c_W^2 - m_K^2)}{m_{Z,\text{ SM}}^2 - m_K^2}\,, \quad m_{\tilde{Z}}^2 = m_{Z,\text{ SM}}^2 + \frac{m_{Z,\text{ SM}}^4 t_W^2 \epsilon^2}{m_{Z,\text{ SM}}^2 - m_K^2}$$

$$\begin{pmatrix} \tilde{Z}_\mu \\ \tilde{A}_\mu \\ \tilde{K}_\mu \end{pmatrix} = \begin{pmatrix} Z_{\mu,\text{ SM}} - \dfrac{t_W m_K^2}{m_{Z,\text{ SM}}^2 - m_K^2} \epsilon K_\mu - \dfrac{m_{Z,\text{ SM}}^4 t_W^2}{2(m_{Z,\text{ SM}}^2 - m_K^2)^2} \epsilon^2 Z_{\mu,\text{ SM}} \\ A_{\mu,\text{ SM}} - \epsilon K_\mu \\ K_\mu + \dfrac{t_W m_{Z,\text{ SM}}^2}{m_{Z,\text{ SM}}^2 - m_K^2} \epsilon Z_{\mu,\text{ SM}} - \left(\dfrac{1}{2} + \dfrac{m_K^4 t_W^2}{2(m_{Z,\text{ SM}}^2 - m_K^2)^2} \right) \epsilon^2 K_\mu \end{pmatrix}$$

    - Singular behavior at $m_K = m_{Z,\text{ SM}}$ is maximal mixing limit
  - Effects from field redefinitions seen in dark, SM currents

Double Dark Portal model

- Fermion bilinears experience the new currents

$$\begin{aligned}\mathcal{L} &\supset g Z_{\mu,\,\text{SM}} J_Z^\mu + e A_{\mu,\,\text{SM}} J_{\text{em}}^\mu + g_D K_\mu J_D^\mu \\ &= \tilde{Z}_\mu \left(g J_Z^\mu - g_D \frac{m_{Z,\,\text{SM}}^2 t_W}{m_{Z,\,\text{SM}}^2 - m_K^2} \epsilon J_D^\mu \right) \\ &+ \tilde{K}_\mu \left(g_D J_D^\mu + g \frac{m_K^2 t_W}{m_{Z,\,\text{SM}}^2 - m_K^2} \epsilon J_Z^\mu + e \epsilon J_{\text{em}}^\mu \right) \\ &+ \tilde{A}_\mu e J_{\text{em}}^\mu + \mathcal{O}(\epsilon^2)\end{aligned}$$

  - $U(1)_D$- charged fermions pick up ε weak charge mediated by *Z*
  - SM charged fermions pick up ε weak charge and ε electric charge mediated by dark photon
  - Photon remains massless, long-range
    - (Singular behavior at $m_K = m_{Z,\,\text{SM}}$ is maximal mixing limit)

Double Dark Portal model

- Scalar boson mixing
  - Higgs portal coupling leads to mass mixing between dark Higgs and SM Higgs
    - Mixing angle

$$\tan 2\alpha = \frac{\lambda_{HP} v_H v_D}{\lambda_D v_D^2 - \lambda_H v_H^2}$$

    - Masses

$$m^2_{S,\, H_0} = \lambda_H v_H^2 + \lambda_D v_D^2 \pm \sqrt{(\lambda_H v_H^2 - \lambda_D v_D^2)^2 + \lambda_{HP} v_H^2 v_D^2}$$

  - Dominant effect is cos α-suppression of Higgs couplings to fermions, dark Higgs mass eigenstate *S* picks up sin α-suppressed couplings to SM fermions

Double Dark Portal model

- Scalar-vector-vector interactions
  - Plays a key role in e^+e^- Higgs studies

$$
\begin{aligned}
\mathcal{L} \supset\ & m_{Z,\mathrm{SM}}^2 \left(\frac{\cos\alpha}{v_H}\right) \tilde{Z}_\mu \tilde{Z}^\mu H_0 \\
& + 2\epsilon t_W \frac{m_K^2 m_{Z,\ \mathrm{SM}}^2}{(m_{Z,\ \mathrm{SM}}^2 - m_K^2)} \left(\frac{\cos\alpha}{v_H} + \frac{\sin\alpha}{v_D}\right) \tilde{Z}_\mu \tilde{K}^\mu H_0 \\
& + m_K^2 \left(-\frac{\sin\alpha}{v_D}\right) \tilde{K}_\mu \tilde{K}^\mu H_0 \\
& + m_{Z,\mathrm{SM}}^2 \left(\frac{\sin\alpha}{v_H}\right) \tilde{Z}_\mu \tilde{Z}^\mu S \\
& + 2\epsilon t_W \frac{m_K^2 m_{Z,\ \mathrm{SM}}^2}{(m_{Z,\ \mathrm{SM}}^2 - m_K^2)} \left(-\frac{\cos\alpha}{v_D} + \frac{\sin\alpha}{v_H}\right) \tilde{Z}_\mu \tilde{K}^\mu S \\
& + m_K^2 \left(\frac{\cos\alpha}{v_D}\right) \tilde{K}_\mu \tilde{K}^\mu S + \mathcal{O}(\epsilon^2)
\end{aligned}
$$

Phenomenology

- Three new states $\tilde{K}$, S , χ
- Many new interactions
  - Deviations in *Z* couplings
  - Deviations in Higgs couplings
  - Exotic Higgs decays (invisible, semi-visible, fully visible)
  - Interactions with dark matter mediated by dark photon
- Rich phenomenology for DM physics and colliders
  - Double Dark Portal model ties together two marginal couplings simultaneously
  - Attractive framework for marrying Higgs deviations and direct coupling to light, very-weakly coupled particles

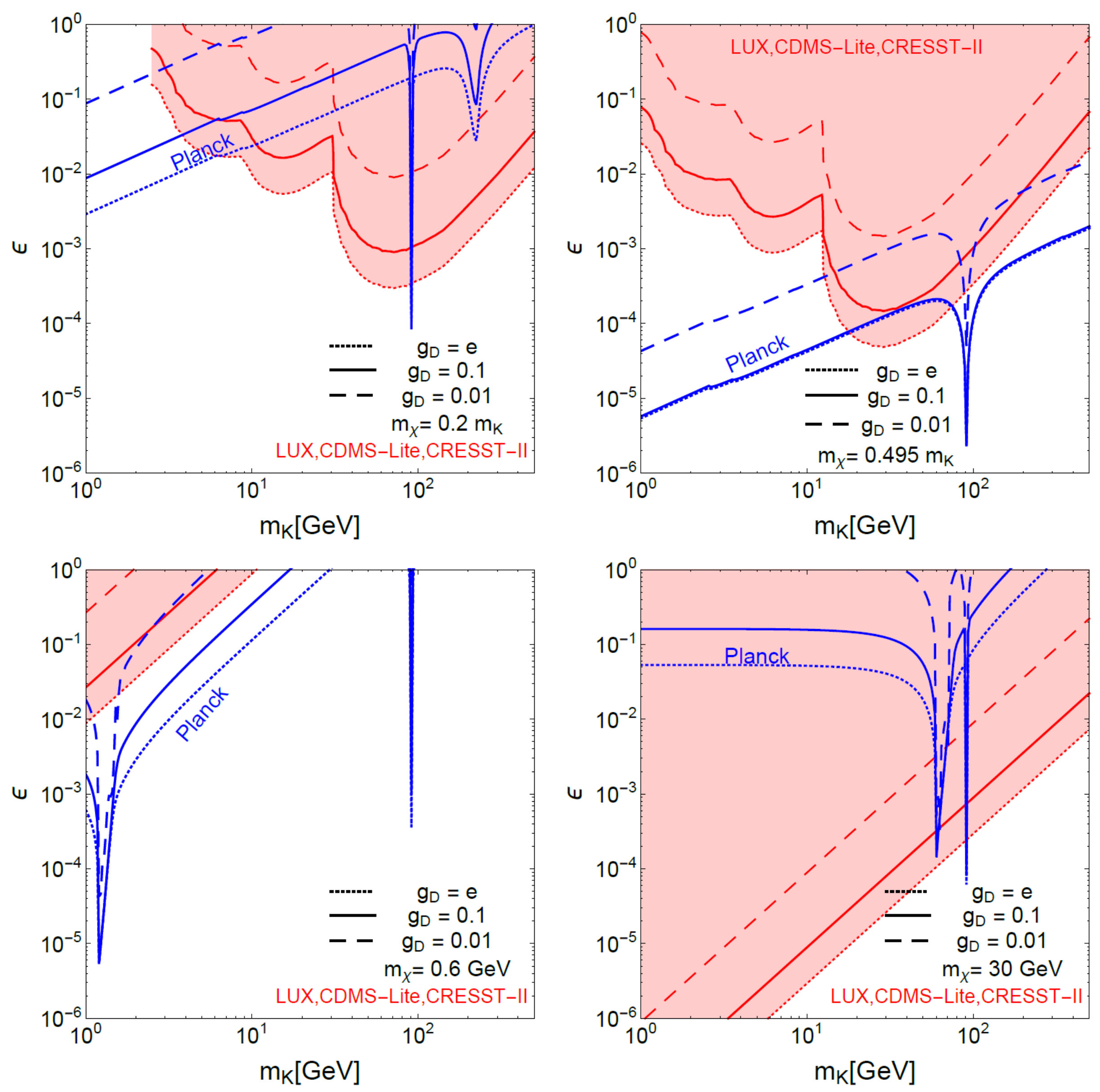
Dark matter direct detection

- Dark matter scattering off protons dominantly from dark photon exchange, suppressed by $(\varepsilon e)^2$
  - Intrinsic cancellation between weak charged currents mediated by massive Z and K vectors (at this order in ε)
  - Dark matter does not interact with photon, hence only protons contribute to direct detection

$$\sigma_p \simeq \frac{\epsilon^2 g_D^2 e^2}{\pi} \frac{\mu_{\chi p}^2}{m_{\tilde{K}}^4} \approx 10^{-44} \text{ cm}^2 \left(\frac{g_D}{e}\right)^2 \left(\frac{\epsilon}{10^{-5}}\right)^2 \left(\frac{10 \text{ GeV}}{m_{\tilde{K}}}\right)^2$$

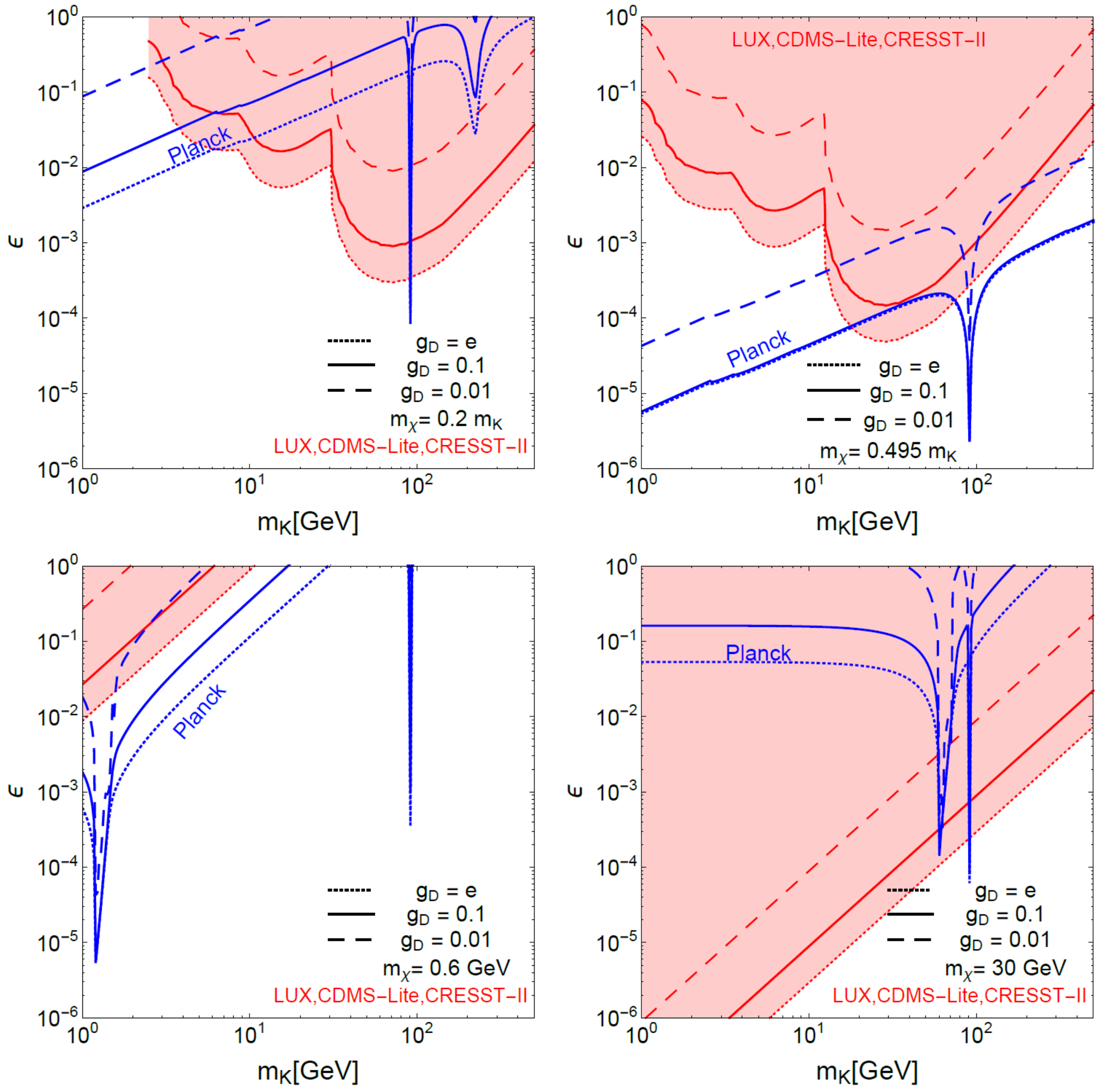
Dark matter direct detection

- Exclusion limits are highly sensitive to the dark matter mass
  - Nuclear recoil energy threshold becomes too soft for light dark matter (about 5 GeV)



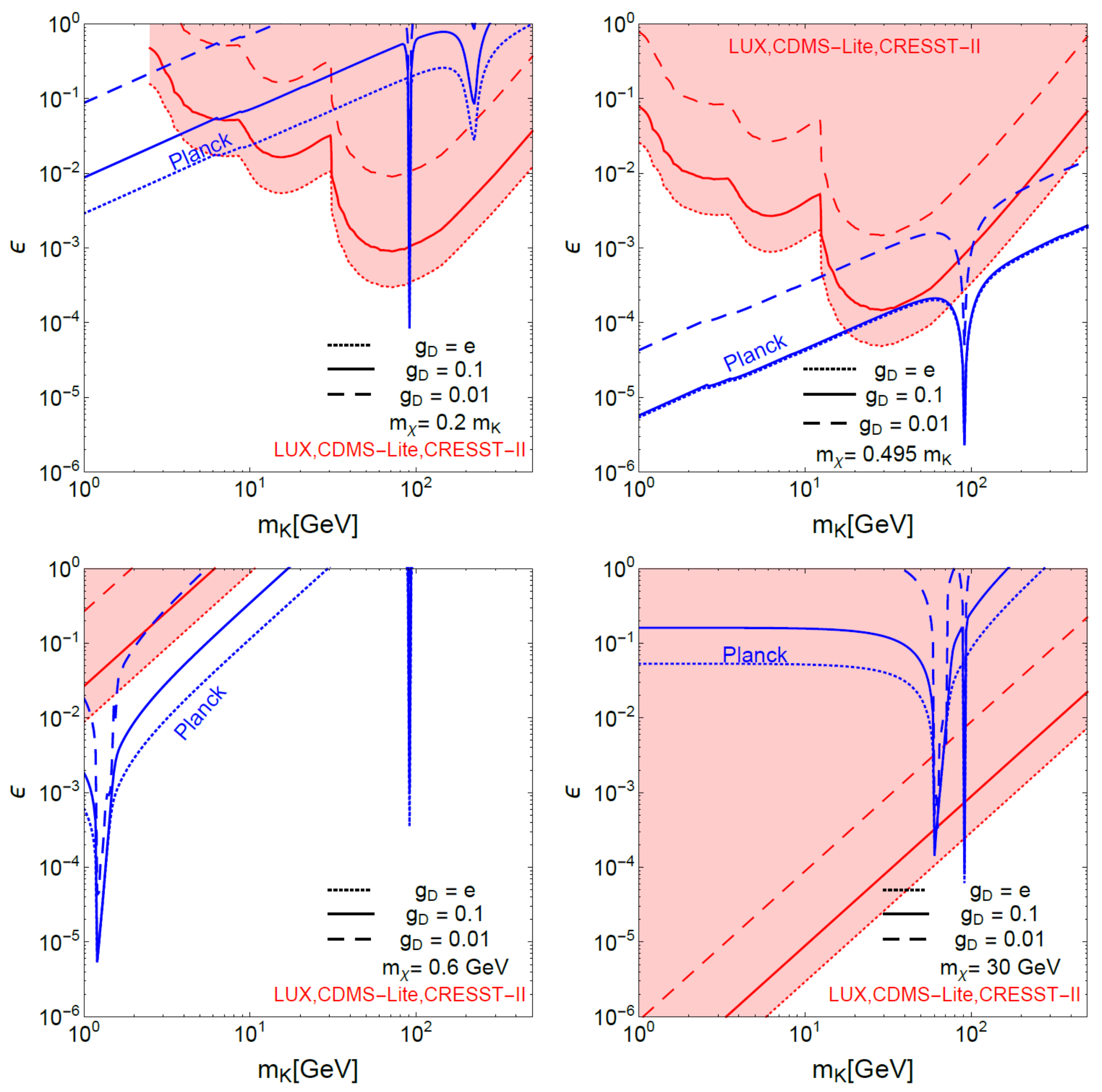
Dark matter direct detection

- Relic abundance (blue line) shows resonances at dark photon and *Z* masses
- DM is underabundant above blue line, overabundant below blue line



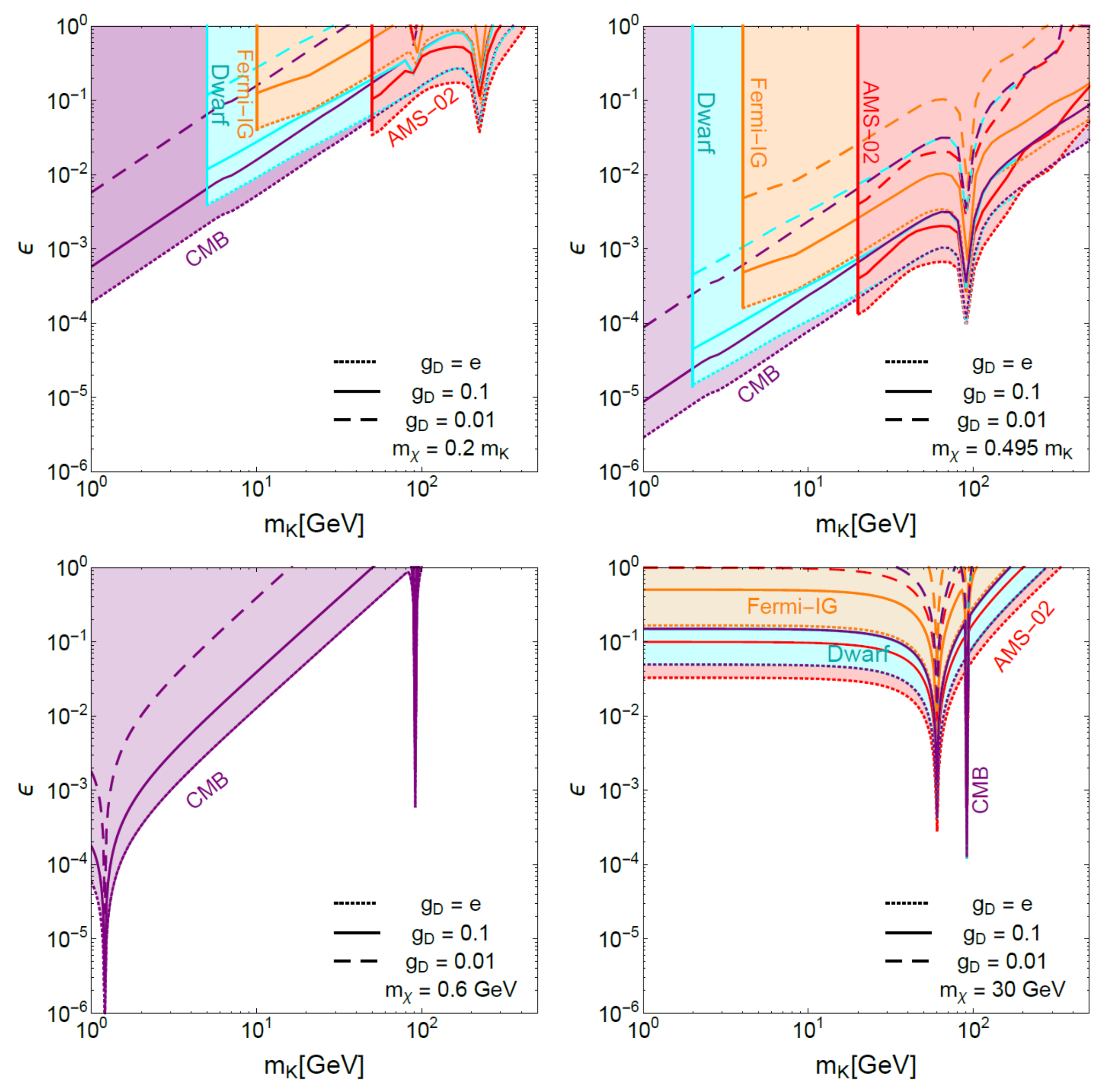
Dark matter direct detection

- Dark matter experiments fix the local relic abundance to 0.3 GeV/cm^3
  - On the other hand, the predicted dark matter relic abundance scales as ε^{-2}, while the scattering rate scales as ε^2
- Ratio of DD limits to relic abundance curve (for fixed m_K) gives the limit on local abundance



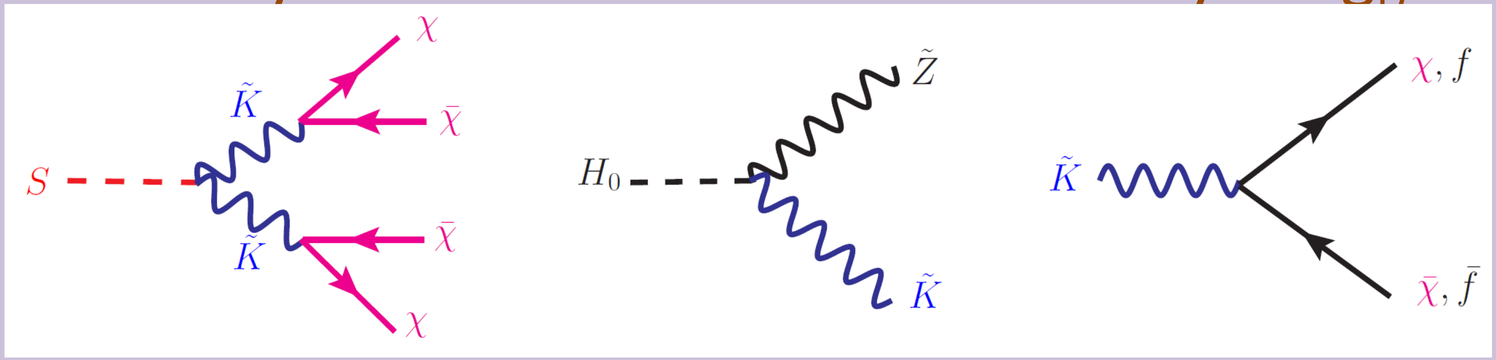
Dark matter indirect detection

- Present day annihilation constrained by observations of gamma ray spectra
- Early universe annihilation constrained by energy injection in CMB
- Strongest limits when DM mass is close to *Z* or dark photon resonance



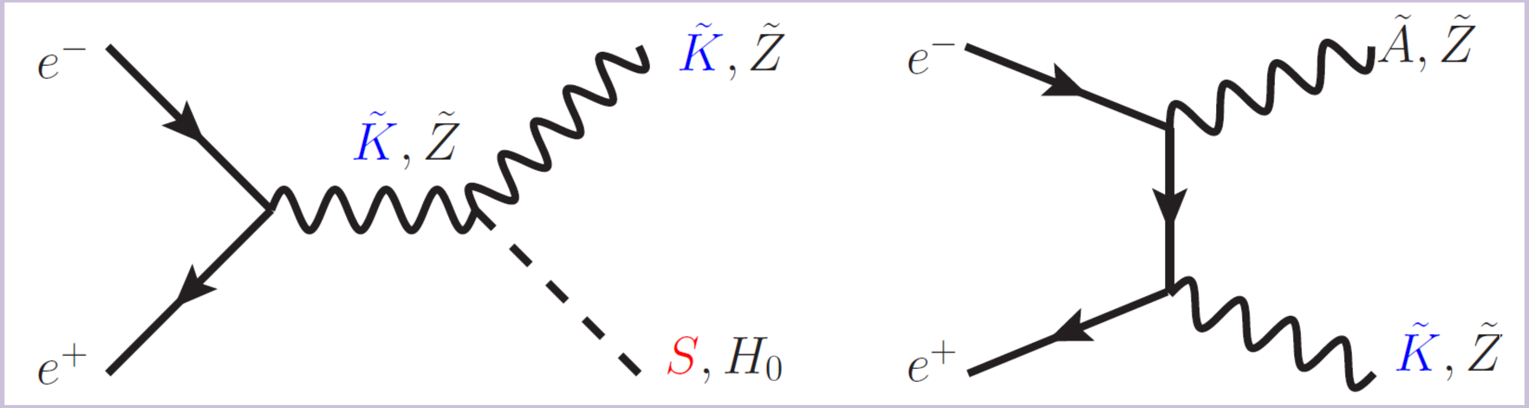
Collider phenomenology

- Modifications to *Z* couplings probed in precision electroweak observables
- Modifications to Higgs couplings tested by LHC and can be seen at a future Higgs factory
  - Also induce invisible and semi-visible exotic Higgs decays
- Will assume dark decays of *S* and *K* are on-shell
  - Ensured by kinematics and mild hierarchy for g_D and ε



Going beyond κ-framework, Higgs EFT

- New light states cause deviations in Higgs physics and can be directly produced



- Exploit radiative return process for hidden photon production
  - Recoil mass technique adapted to monophoton events and other SM candles as recoil taggers

Exploiting radiative return and recoil mass techniques at e$^+$e$^-$ machines

- Radiative return – use ISR photon to make 2-2 production on-shell
    - At LHC, "radiative return" is better known as "mono-jet"
- Recoil mass method – use four-momentum conservation in 2-2 process
    - In case of invisible decay and radiative return, equivalent to searching for a monophoton peak
        - Design driver for e$^+$e$^-$ electromagnatic calorimeter

$$E_{\text{vis}} = \frac{\sqrt{s}}{2} + \frac{m_{\text{vis}}^2 - m_X^2}{2\sqrt{s}}$$

$$m_{\text{recoil}} = m_X = \sqrt{s + m_{\text{vis}}^2 - 2E_{\text{vis}}\sqrt{s}}$$

Exotic invisible decay of Higgs

- Familiar case: Higgs recoiling against Z for invisible Higgs decays
  - Invisible decay combines sensitivity to sin α and ε, overall rate driven by g_D

$$\Gamma(H_0 \to \text{inv}) \approx \Gamma(H_0 \to SS) + \Gamma(H_0 \to \tilde{K}\tilde{K}) + 0.2 \times \Gamma(H_0 \to \tilde{K}\tilde{Z})$$

- Individual rates are

$$\Gamma(H_0 \to SS) = g_D^2 \sin^2\alpha \frac{m_{H_0}}{32\pi} \sqrt{1 - \frac{4m_S^2}{m_{H_0}^2}} \frac{(m_{H_0}^2 + 2m_S^2)^2}{m_{H_0}^2 m_K^2} \ ,$$

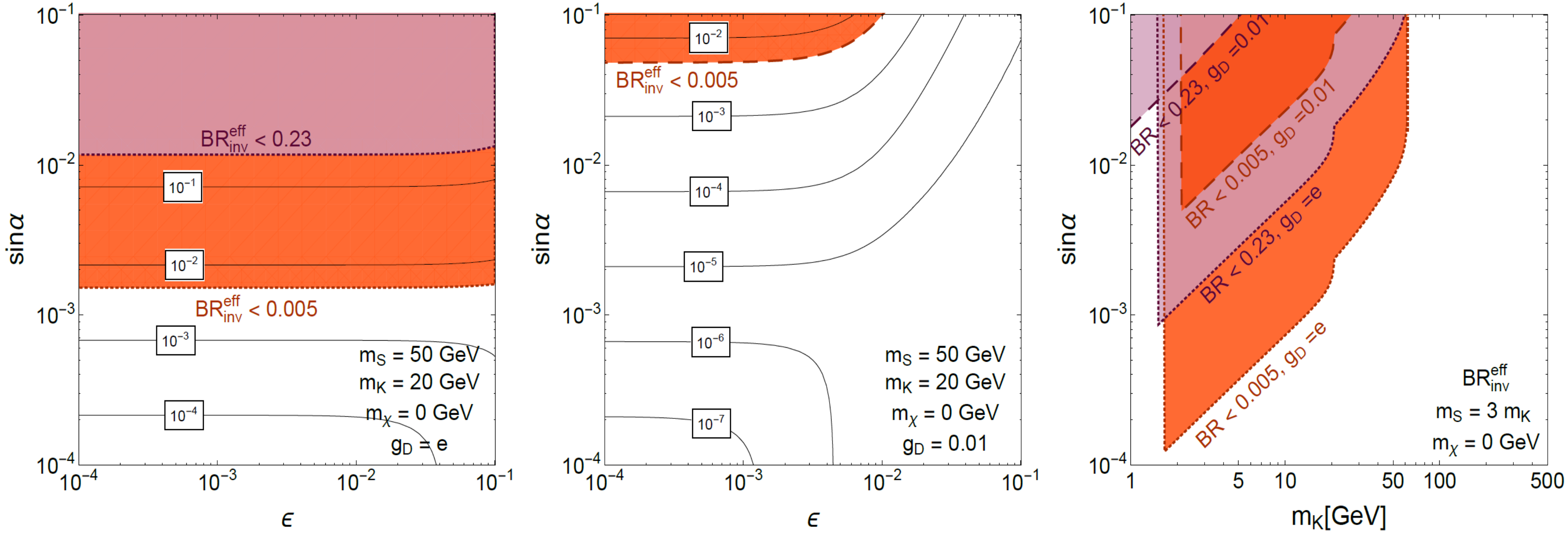
$$\Gamma(H_0 \to \tilde{K}\tilde{K}) = g_D^2 \sin^2\alpha \frac{m_{H_0}}{32\pi} \sqrt{1 - \frac{4m_{\tilde{K}}^2}{m_{H_0}^2}} \frac{m_{H_0}^4 - 4m_{H_0}^2 m_{\tilde{K}}^2 + 12 m_{\tilde{K}}^4}{m_{H_0}^2 m_{\tilde{K}}^2} \frac{m_K^2}{m_{\tilde{K}}^2} \ ,$$

$$\Gamma(H_0 \to \tilde{K}\tilde{Z}) = \frac{\epsilon^2 t_W^2 \left(\frac{\cos\alpha}{v_H} + \frac{\sin\alpha}{v_D}\right)^2}{16\pi m_{H_0}^3 \left(m_K^2 - m_{Z,\text{ SM}}^2\right)^2} \frac{m_K^4 m_{Z,\text{ SM}}^4}{m_{\tilde{K}}^2 m_{\tilde{Z}}^2} \sqrt{m_{H_0}^4 + \left(m_{\tilde{K}}^2 - m_{\tilde{Z}}^2\right)^2 - 2m_{H_0}^2 \left(m_{\tilde{K}}^2 + m_{\tilde{Z}}^2\right)}$$
$$\times \left((m_{H_0}^2 - m_{\tilde{K}}^2 - m_{\tilde{Z}}^2)^2 + 8 m_{\tilde{K}}^2 m_{\tilde{Z}}^2\right)$$

Exotic invisible decay of Higgs

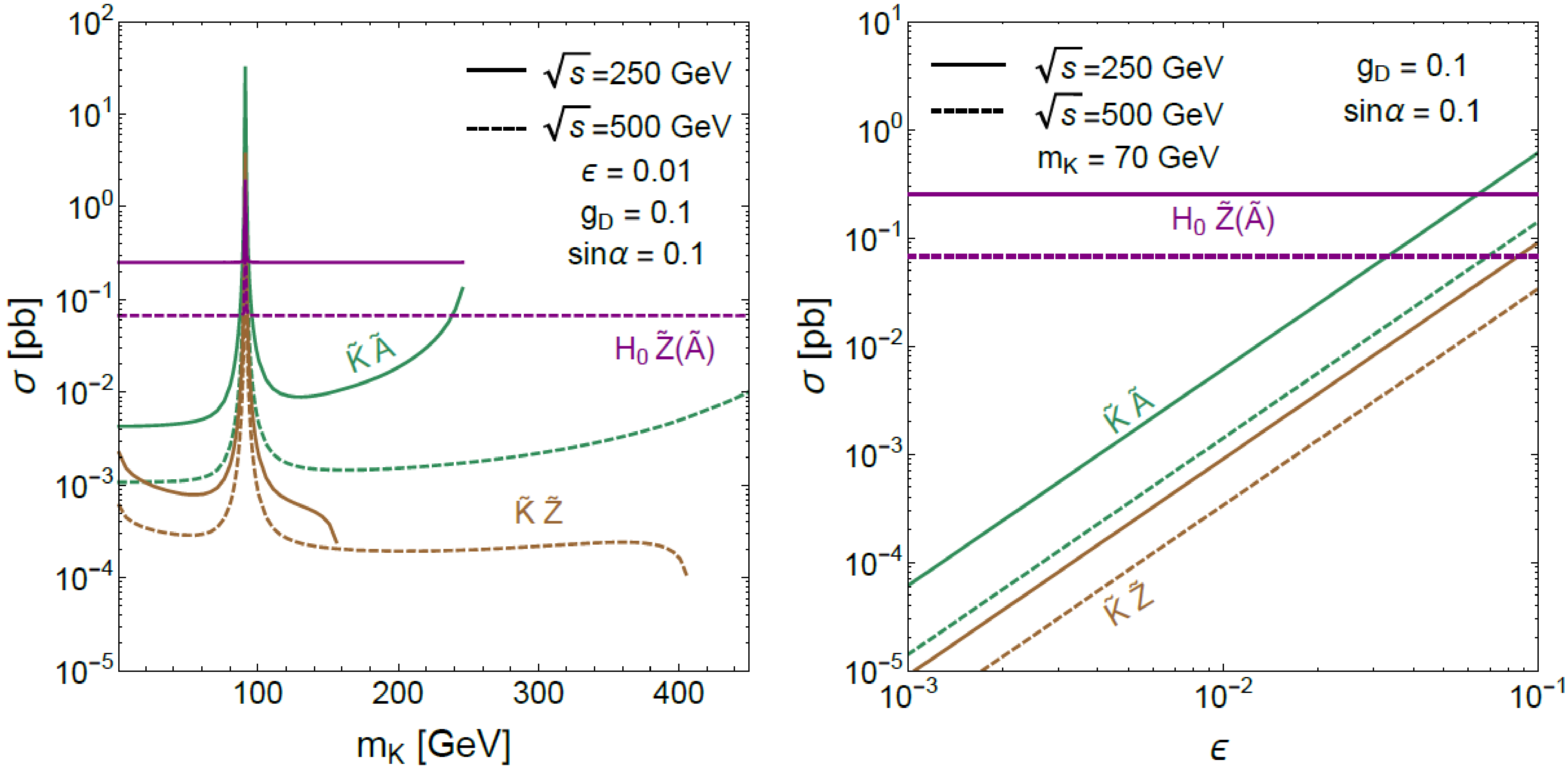
- Familiar case: Higgs recoiling against Z for invisible Higgs decays
  - Invisible decay combines sensitivity to sin α and ε, overall rate driven by g_D

$$\Gamma(H_0 \to \text{inv}) \approx \Gamma(H_0 \to SS) + \Gamma(H_0 \to \tilde{K}\tilde{K}) + 0.2 \times \Gamma(H_0 \to \tilde{K}\tilde{Z})$$



Direct production of new light states

- Possible new physics within kinematic reach
  - Signatures too difficult at LHC, exploit e^+e^- capabilities



Prospects for dark photon

- Many possible visible and invisible final states

  $e^+e^- \to \tilde{Z}H_0$ Study $\tilde{Z} \to \ell\ell$ and semi-visible $H_0 \to (\ell\ell)_Z\chi\chi$

  $e^+e^- \to \tilde{Z}\tilde{K}$ Study $\tilde{Z} \to \ell\ell$ and $\tilde{K} \to \bar{\chi}\chi$ or $\ell\ell$

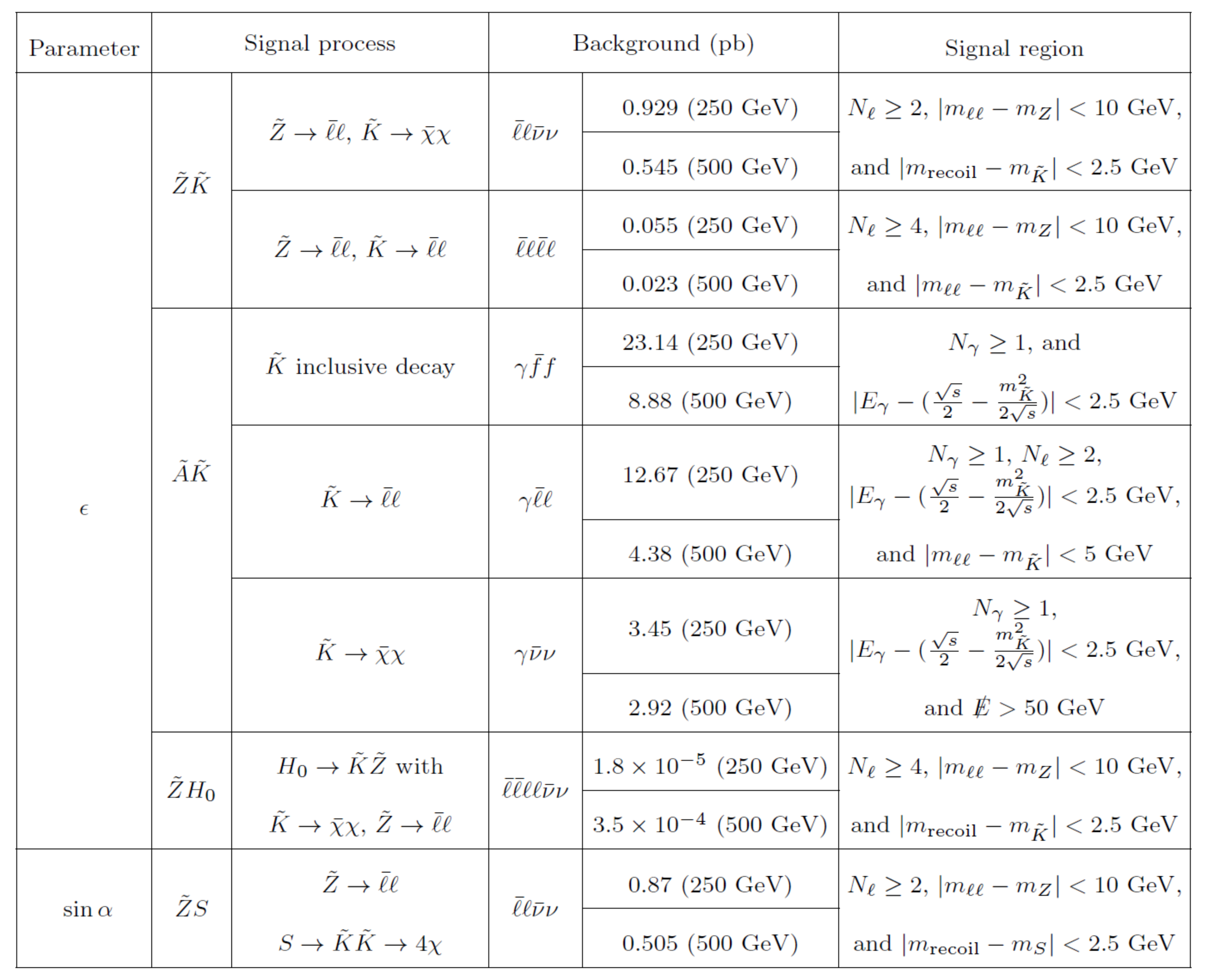
  $e^+e^- \to \gamma\tilde{K}$ Study $\tilde{K}$ inclusive decays, and exclusive $\tilde{K} \to \bar{\chi}\chi$ or $\ell\ell$

  $e^+e^- \to \tilde{Z}S$ Study $\tilde{Z} \to \ell\ell$ and $S \to 4\chi$

  - Event simulation using MG5+Pythia+Delphes
    - Use parametrized preliminary CEPC detector card
  - SM backgrounds and cuts driven by e^+e^- environment
- Rates for visible states are lower by $(\varepsilon/g_D)^2$, best sensitivity from requiring missing energy threshold
  - LEP direct constraints ($\varepsilon < 0.03$) not competitive

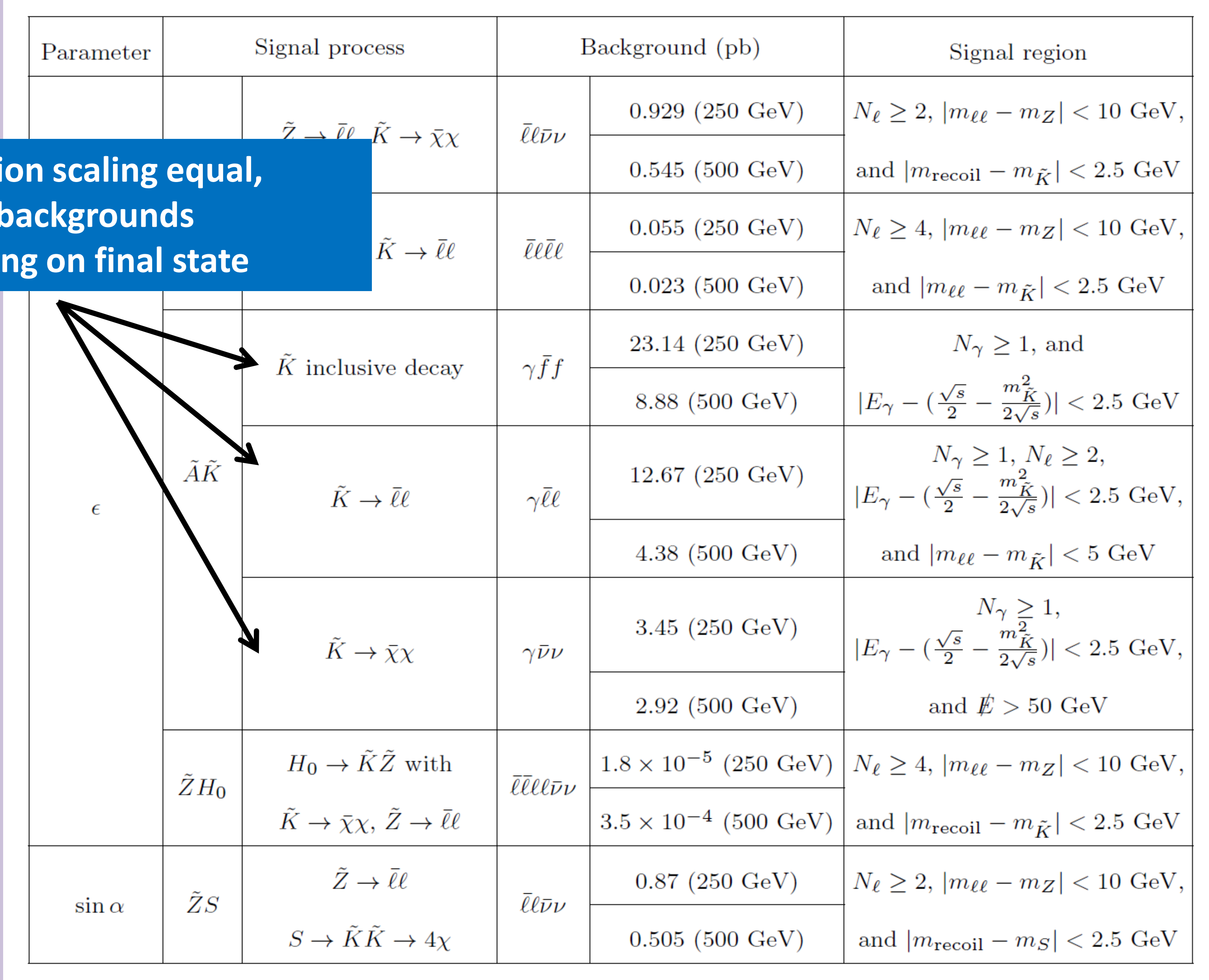
Hook, Izaguirre, Wacker [1006.0973]

Collider study cuts



Parameter	Signal process		Background (pb)		Signal region
ϵ	$\tilde{Z}\tilde{K}$	$\tilde{Z} \to \bar{\ell}\ell$, $\tilde{K} \to \bar{\chi}\chi$	$\bar{\ell}\ell\bar{\nu}\nu$	0.929 (250 GeV)	$N_\ell \geq 2$, $\lvert m_{\ell\ell} - m_Z \rvert < 10$ GeV,
				0.545 (500 GeV)	and $\lvert m_{\text{recoil}} - m_{\tilde{K}} \rvert < 2.5$ GeV
		$\tilde{Z} \to \bar{\ell}\ell$, $\tilde{K} \to \bar{\ell}\ell$	$\bar{\ell}\ell\bar{\ell}\ell$	0.055 (250 GeV)	$N_\ell \geq 4$, $\lvert m_{\ell\ell} - m_Z \rvert < 10$ GeV,
				0.023 (500 GeV)	and $\lvert m_{\ell\ell} - m_{\tilde{K}} \rvert < 2.5$ GeV
	$\tilde{A}\tilde{K}$	$\tilde{K}$ inclusive decay	$\gamma\bar{f}f$	23.14 (250 GeV)	$N_\gamma \geq 1$, and
				8.88 (500 GeV)	$\lvert E_\gamma - (\frac{\sqrt{s}}{2} - \frac{m_{\tilde{K}}^2}{2\sqrt{s}}) \rvert < 2.5$ GeV
		$\tilde{K} \to \bar{\ell}\ell$	$\gamma\bar{\ell}\ell$	12.67 (250 GeV)	$N_\gamma \geq 1$, $N_\ell \geq 2$, $\lvert E_\gamma - (\frac{\sqrt{s}}{2} - \frac{m_{\tilde{K}}^2}{2\sqrt{s}}) \rvert < 2.5$ GeV,
				4.38 (500 GeV)	and $\lvert m_{\ell\ell} - m_{\tilde{K}} \rvert < 5$ GeV
		$\tilde{K} \to \bar{\chi}\chi$	$\gamma\bar{\nu}\nu$	3.45 (250 GeV)	$N_\gamma \geq 1$, $\lvert E_\gamma - (\frac{\sqrt{s}}{2} - \frac{m_{\tilde{K}}^2}{2\sqrt{s}}) \rvert < 2.5$ GeV,
				2.92 (500 GeV)	and $\not{E} > 50$ GeV
	$\tilde{Z}H_0$	$H_0 \to \tilde{K}\tilde{Z}$ with	$\bar{\ell}\ell\bar{\ell}\ell\bar{\nu}\nu$	1.8×10^{-5} (250 GeV)	$N_\ell \geq 4$, $\lvert m_{\ell\ell} - m_Z \rvert < 10$ GeV,
		$\tilde{K} \to \bar{\chi}\chi$, $\tilde{Z} \to \bar{\ell}\ell$		3.5×10^{-4} (500 GeV)	and $\lvert m_{\text{recoil}} - m_{\tilde{K}} \rvert < 2.5$ GeV
$\sin\alpha$	$\tilde{Z}S$	$\tilde{Z} \to \bar{\ell}\ell$	$\bar{\ell}\ell\bar{\nu}\nu$	0.87 (250 GeV)	$N_\ell \geq 2$, $\lvert m_{\ell\ell} - m_Z \rvert < 10$ GeV,
		$S \to \tilde{K}\tilde{K} \to 4\chi$		0.505 (500 GeV)	and $\lvert m_{\text{recoil}} - m_S \rvert < 2.5$ GeV

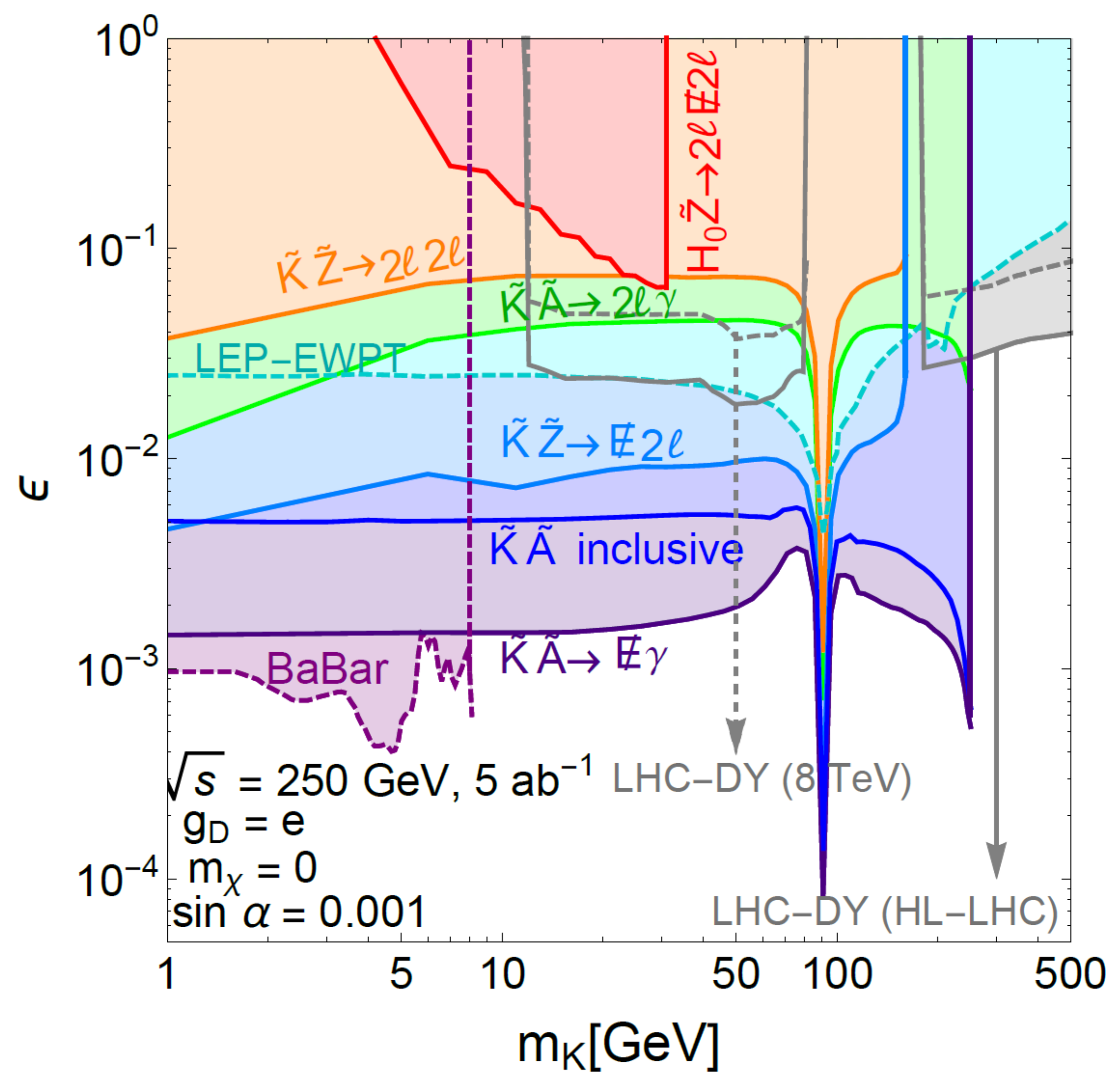
Collider study cuts



Production scaling equal, varying backgrounds depending on final state

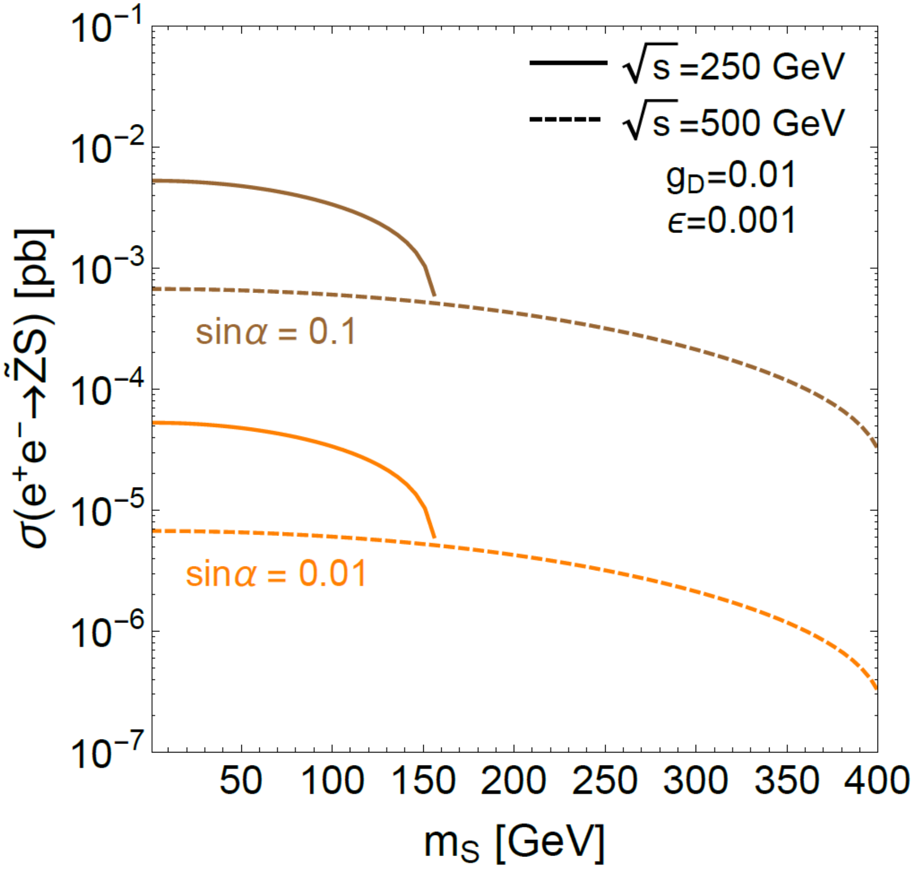
Parameter	Signal process		Background (pb)		Signal region
ϵ		$\tilde{Z} \to \bar{\ell}\ell$, $\tilde{K} \to \bar{\chi}\chi$	$\bar{\ell}\ell\bar{\nu}\nu$	0.929 (250 GeV) 0.545 (500 GeV)	$N_\ell \geq 2$, $\lvert m_{\ell\ell} - m_Z \rvert < 10$ GeV, and $\lvert m_{\text{recoil}} - m_{\tilde{K}} \rvert < 2.5$ GeV
		$\tilde{K} \to \bar{\ell}\ell$	$\bar{\ell}\ell\bar{\ell}\ell$	0.055 (250 GeV) 0.023 (500 GeV)	$N_\ell \geq 4$, $\lvert m_{\ell\ell} - m_Z \rvert < 10$ GeV, and $\lvert m_{\ell\ell} - m_{\tilde{K}} \rvert < 2.5$ GeV
	$\tilde{A}\tilde{K}$	$\tilde{K}$ inclusive decay	$\gamma\bar{f}f$	23.14 (250 GeV) 8.88 (500 GeV)	$N_\gamma \geq 1$, and $\lvert E_\gamma - (\frac{\sqrt{s}}{2} - \frac{m_{\tilde{K}}^2}{2\sqrt{s}}) \rvert < 2.5$ GeV
		$\tilde{K} \to \bar{\ell}\ell$	$\gamma\bar{\ell}\ell$	12.67 (250 GeV) 4.38 (500 GeV)	$N_\gamma \geq 1$, $N_\ell \geq 2$, $\lvert E_\gamma - (\frac{\sqrt{s}}{2} - \frac{m_{\tilde{K}}^2}{2\sqrt{s}}) \rvert < 2.5$ GeV, and $\lvert m_{\ell\ell} - m_{\tilde{K}} \rvert < 5$ GeV
		$\tilde{K} \to \bar{\chi}\chi$	$\gamma\bar{\nu}\nu$	3.45 (250 GeV) 2.92 (500 GeV)	$N_\gamma \geq 1$, $\lvert E_\gamma - (\frac{\sqrt{s}}{2} - \frac{m_{\tilde{K}}^2}{2\sqrt{s}}) \rvert < 2.5$ GeV, and $\not{E} > 50$ GeV
	$\tilde{Z}H_0$	$H_0 \to \tilde{K}\tilde{Z}$ with $\tilde{K} \to \bar{\chi}\chi$, $\tilde{Z} \to \bar{\ell}\ell$	$\bar{\ell}\ell\bar{\ell}\ell\bar{\nu}\nu$	1.8×10^{-5} (250 GeV) 3.5×10^{-4} (500 GeV)	$N_\ell \geq 4$, $\lvert m_{\ell\ell} - m_Z \rvert < 10$ GeV, and $\lvert m_{\text{recoil}} - m_{\tilde{K}} \rvert < 2.5$ GeV
$\sin\alpha$	$\tilde{Z}S$	$\tilde{Z} \to \bar{\ell}\ell$ $S \to \tilde{K}\tilde{K} \to 4\chi$	$\bar{\ell}\ell\bar{\nu}\nu$	0.87 (250 GeV) 0.505 (500 GeV)	$N_\ell \geq 2$, $\lvert m_{\ell\ell} - m_Z \rvert < 10$ GeV, and $\lvert m_{\text{recoil}} - m_S \rvert < 2.5$ GeV

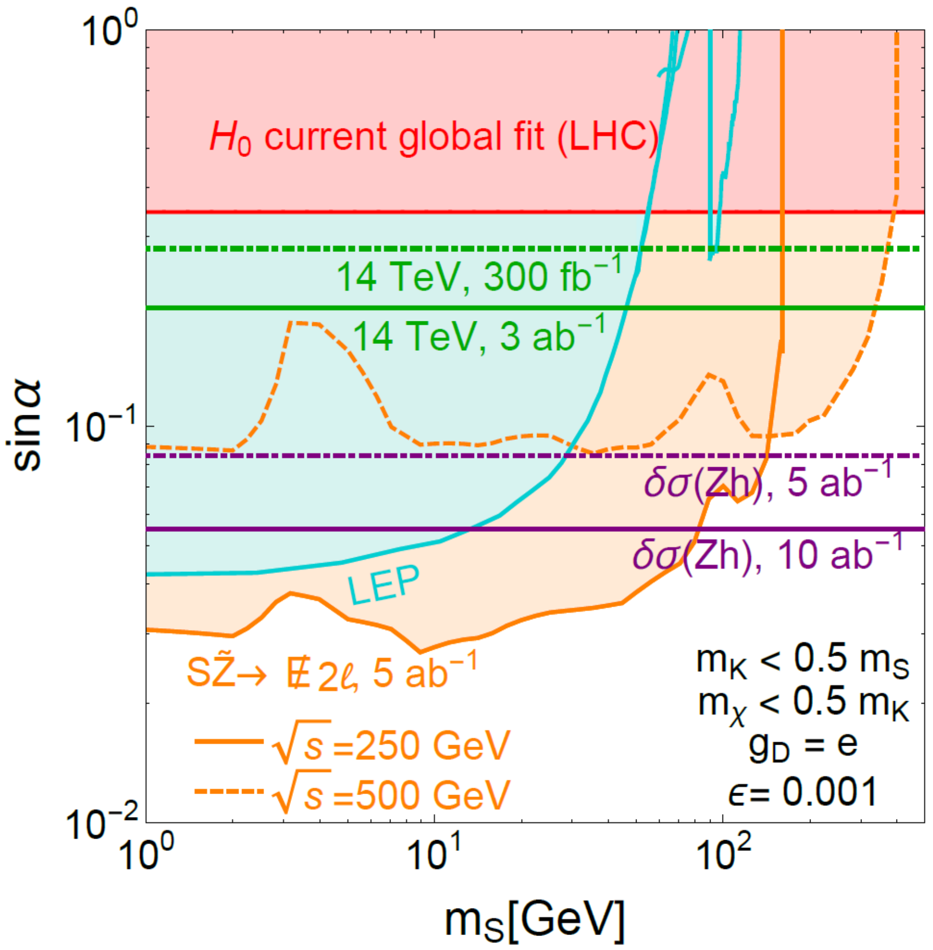
Dark photon sensitivity



Prospects for dark scalar

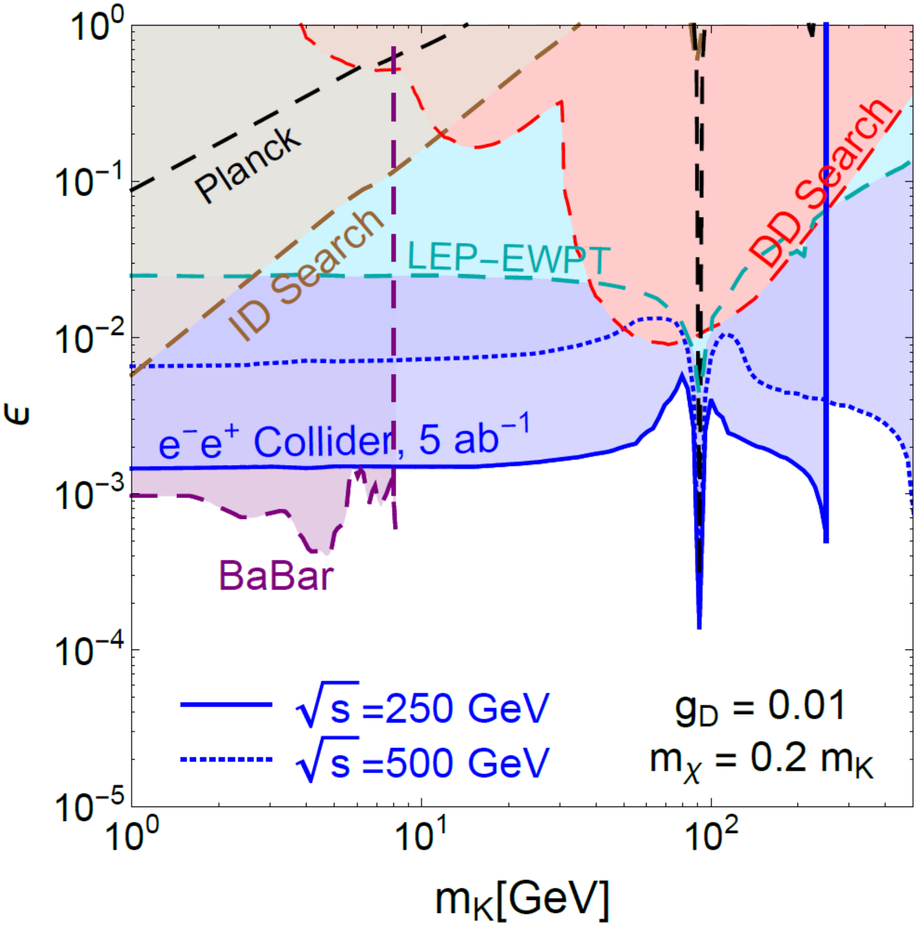
- Similarly, direct dark Higgs production and precision Higgs measurements

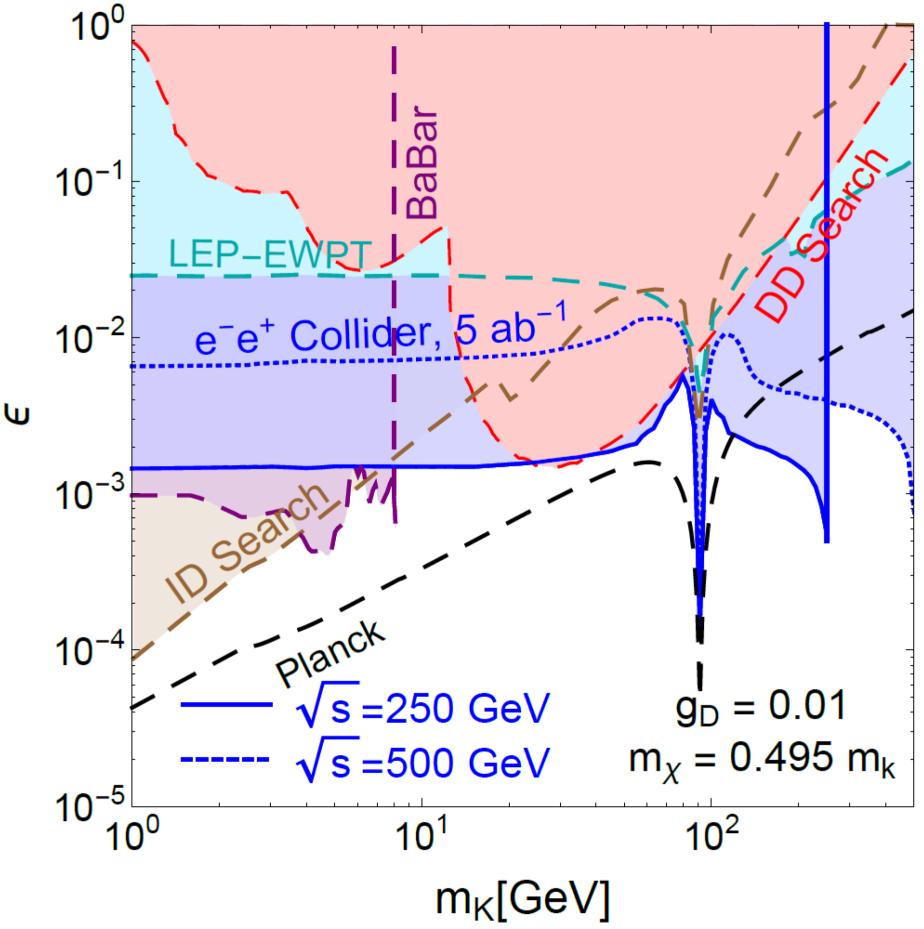




Comparing to complementary DM probes

- Dark matter discovery possible at e^+e^- machines



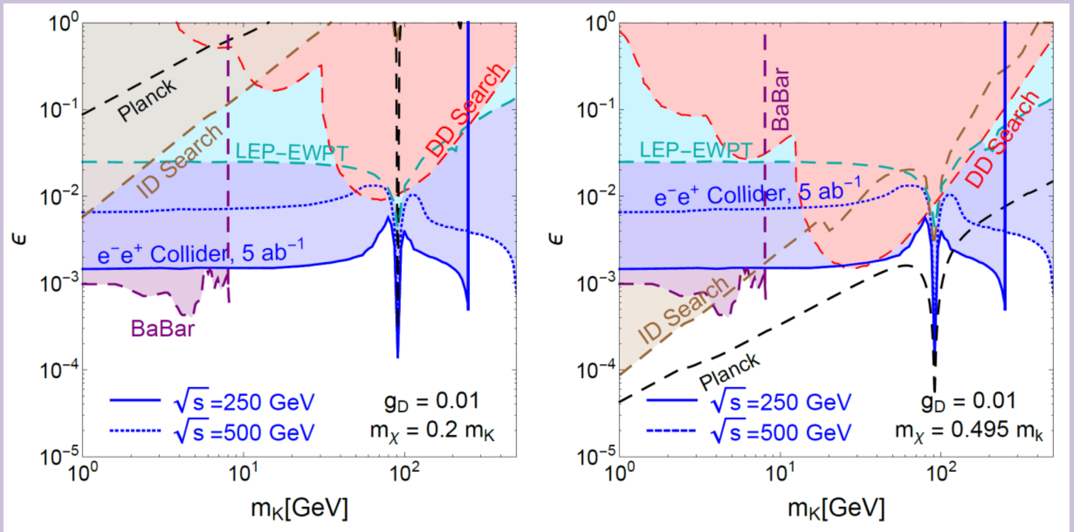


Outlook

- Several complementary studies on dark currents
  - Cui, D'Eramo – 1705.03897
  - Dror, Lasenby, Pospelov – 1705.06726, 1707.01503
  - Ismail, Katz, Racco – 1707.00709
  - Ismail, Katz – 1712.01840
- Also studies of light DM direct detection
  - e.g. Kahlhoefer, Kulkarni, Wild – 1707.08571
- Anomalon-induced operators (Dobrescu, Yu/Liu, Michaels, Wang, Yu [1801.upcomings]) provide new opportunities for dark current detection at LHC and beam dumps

Conclusions

- Physics potential of e^+e^- machine goes well beyond precision Standard Model program
- Direct production of new, light, very weakly-coupled hidden particles possible
- Double Dark Portal model is a concrete framework for studying two marginal couplings in tandem



Introduction and Motivation

- Era of exploratory particle physics
  - Possible NP models span decades in scale and couplings
  - Strong gains to come from e^+e^- precision Higgs program
  - ILC, FCC-ee, CEPC, CLIC machines under serious consideration
- Missing piece of story: e^+e^- collider production of new particles
  - More than a Higgs factory, but production of new, light states – especially when sensitivity exceeds possibilities at (HL-)LHC
  - Will discuss dark vector and dark scalar production and their SM and DM decays