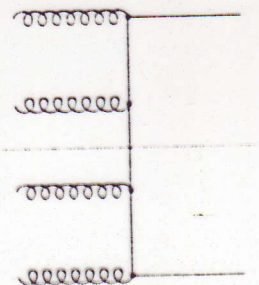


LO Pomeron-to-parton splitting functions



- LO Pomeron-to-parton splitting functions calculated in Eur. Phys. J. C **44** (2005) 69.
- **Notation:** ' $\mathbb{P} = G$ ' means gluonic Pomeron, ' $\mathbb{P} = S$ ' means sea-quark Pomeron, ' $\mathbb{P} = GS$ ' means interference between these.

$$z \Sigma^{\mathbb{P}=G}(z, \mu^2; \mu^2) = P_{q, \mathbb{P}=G}(z) = z^3 (1-z),$$

$$z g^{\mathbb{P}=G}(z, \mu^2; \mu^2) = P_{g, \mathbb{P}=G}(z) = \frac{9}{16} (1+z)^2 (1-z)^2,$$

$$z \Sigma^{\mathbb{P}=S}(z, \mu^2; \mu^2) = P_{q, \mathbb{P}=S}(z) = \frac{4}{81} z (1-z),$$

$$z g^{\mathbb{P}=S}(z, \mu^2; \mu^2) = P_{g, \mathbb{P}=S}(z) = \frac{1}{9} (1-z)^2,$$

$$z \Sigma^{\mathbb{P}=GS}(z, \mu^2; \mu^2) = P_{q, \mathbb{P}=GS}(z) = \frac{2}{9} z^2 (1-z),$$

$$z g^{\mathbb{P}=GS}(z, \mu^2; \mu^2) = P_{g, \mathbb{P}=GS}(z) = \frac{1}{4} (1+2z) (1-z)^2.$$

Evolve these input Pomeron PDFs from μ^2 up to Q^2 using NLO DGLAP evolution.

$$F_2^D = \sum_{a=g,q} C_{2,a} \cdot a^D + \underbrace{C_{2,P}}_{\text{direct IP}}$$

$$\frac{\partial a^D}{\partial \ln Q^2} = \underbrace{\sum_a P_{a,a'} \cdot a'^D}_{\text{DGLAP}} + \underbrace{P_{a|P}(z) f_P(x_P, Q^2)}_{\text{inhomogeneous term}}$$

• No free parameters
 $\hookrightarrow$ less g at $\beta > 0.4$

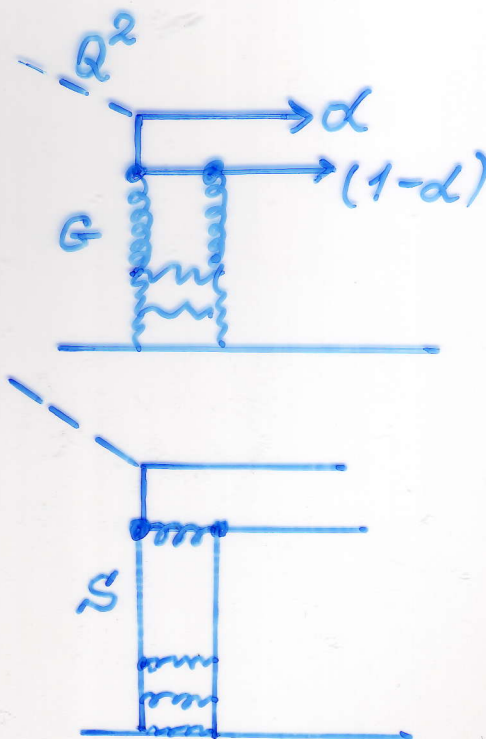
$$P = G + S$$

$\hookrightarrow$ sea quarks in t -channel

• $\frac{dG}{dK_t^2} = \frac{(K_t^2)^{2\gamma}}{K_t^4} \Rightarrow$ large $K_{t, \text{intr.}}$ in P -fragment.

• direct diffr. dijets

$$LO \Rightarrow G/S \propto \left(\frac{Q^2 L(1-d)}{E_T^2} \right)^2$$



$$x_P f_P^G = \frac{1}{B_D} \left[R_g \frac{\alpha_s(\mu^2)}{\mu} x g(x, \mu^2) \right]^2$$

$$x g \propto (\mu^2)^\gamma$$

$$x_P f_P \propto (\mu^2)^{2\gamma} (x_P)^{-1}$$

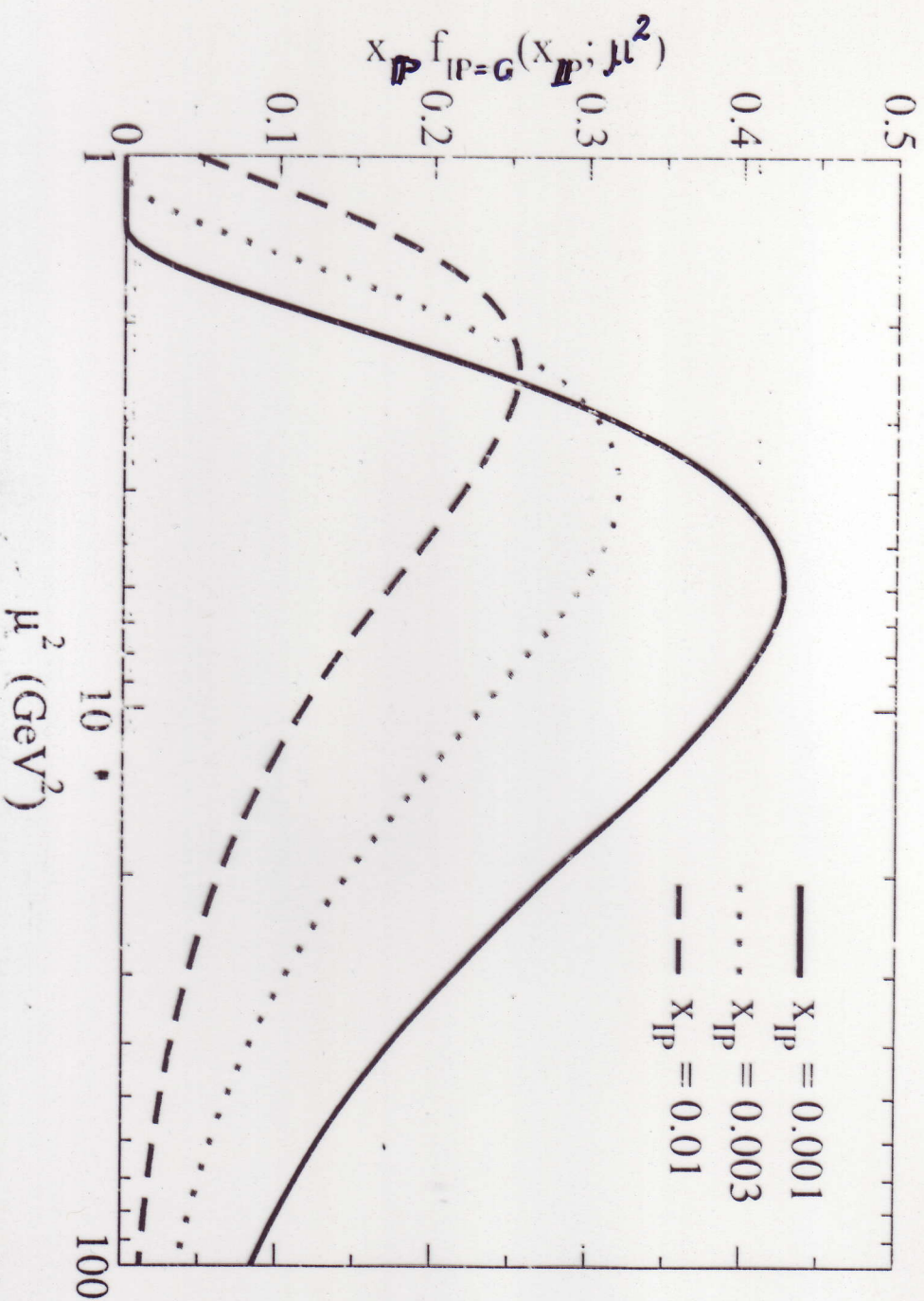


Figure 4: The μ^2 dependence of the flux factor f_{IP} , given by (7) with $N = 1$, for three different values of x_P , and using the MRST2001 NLO gluon distribution of the proton [16].

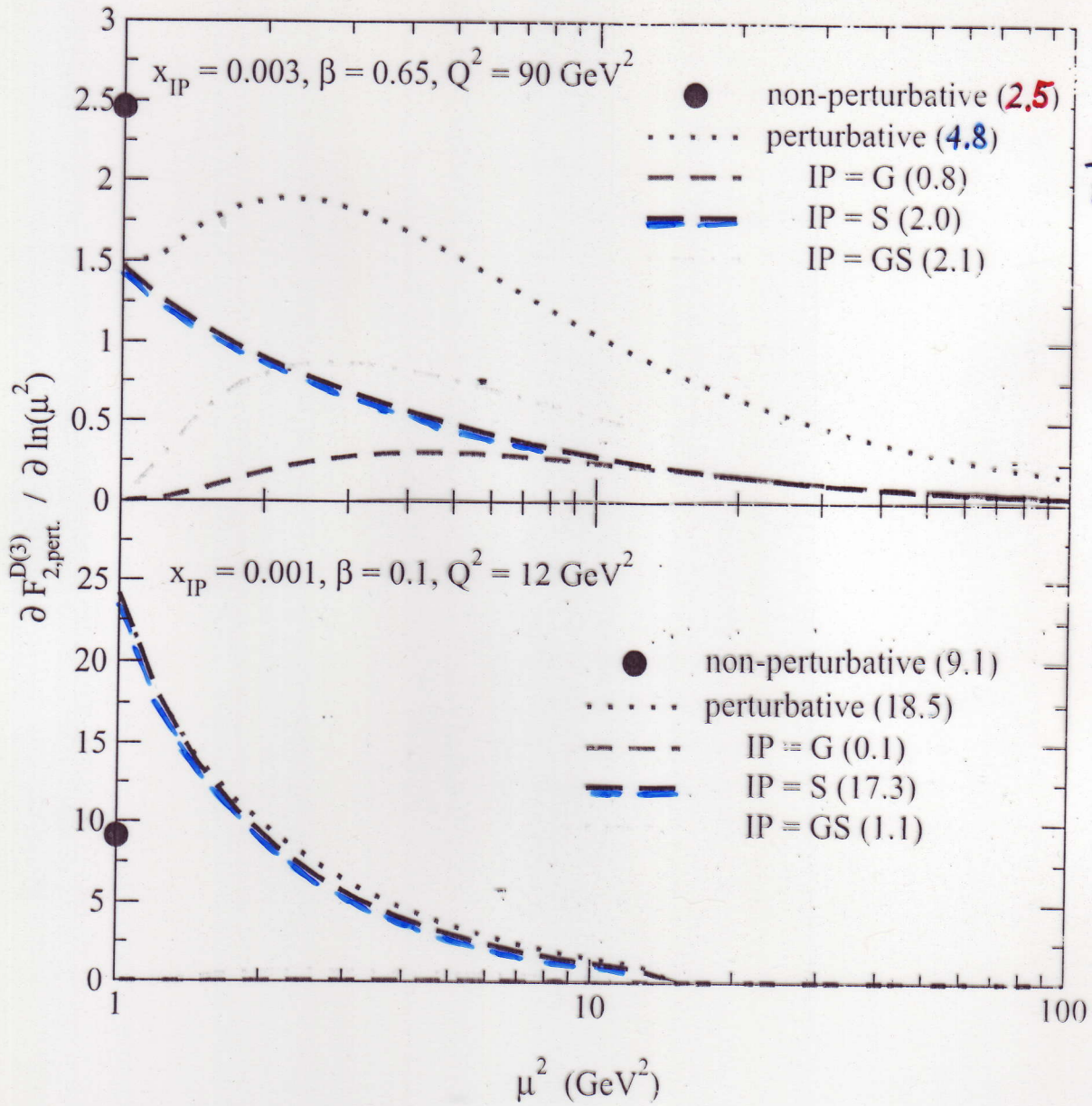


Figure 6: The perturbative ($IP = G, S$, and the GS interference) and non-perturbative contributions to $F_2^{D(3)}$, for two sets of x_{IP} , β , and Q^2 values, found in the analysis of HERA-DDIE3 data in Ref. [7]. The plots show the μ^2 dependence of the perturbative contributions: their integral over μ^2 is shown by the numbers in parentheses in the legend.

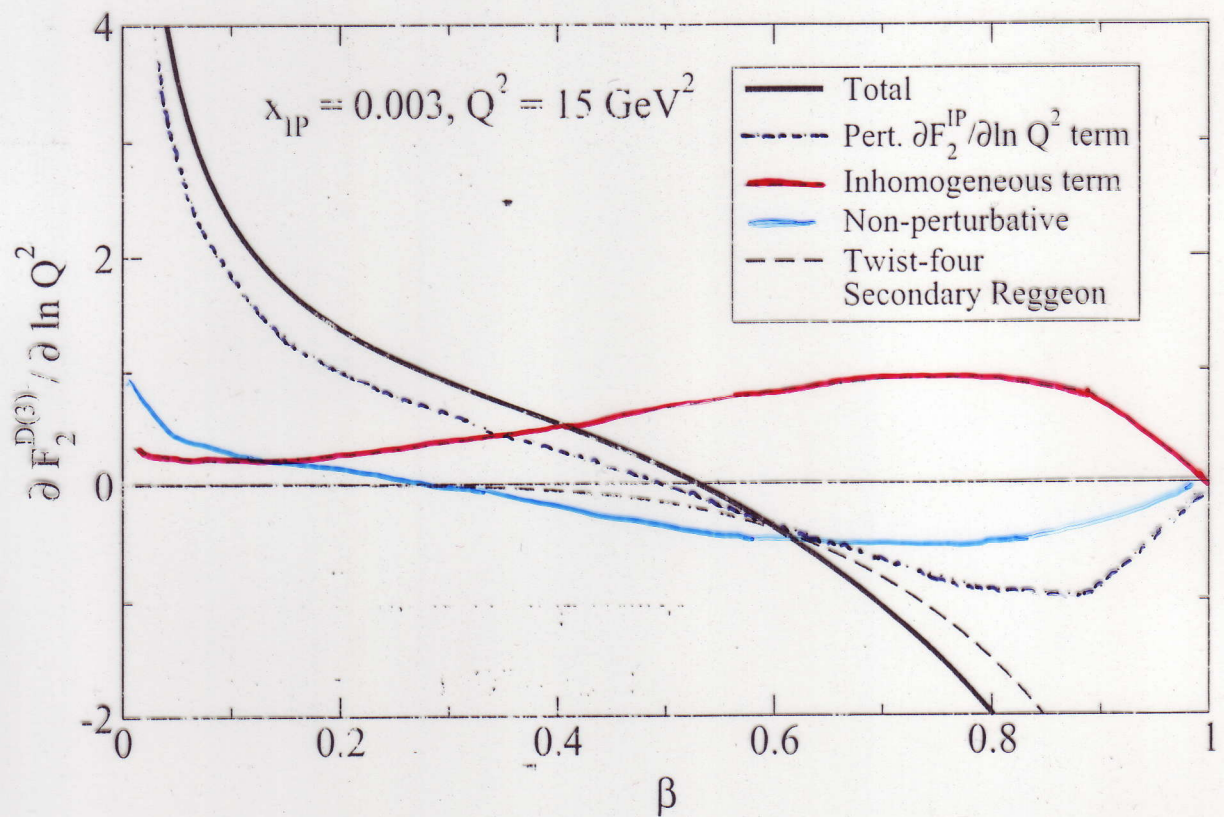


Figure 7: The breakdown of the contributions to the description of the slope $\partial F_2^{D(3)}/\partial \ln Q^2$ as a function of β , for $Q^2 = 15 \text{ GeV}^2$ and $\lambda_{IP} = 0.003$ obtained in the MRW analysis [7] of the combined ZEUS and H1 DDIS data, which uses MRST2001 NLO partons [16] to calculate the perturbative Pomeron flux.

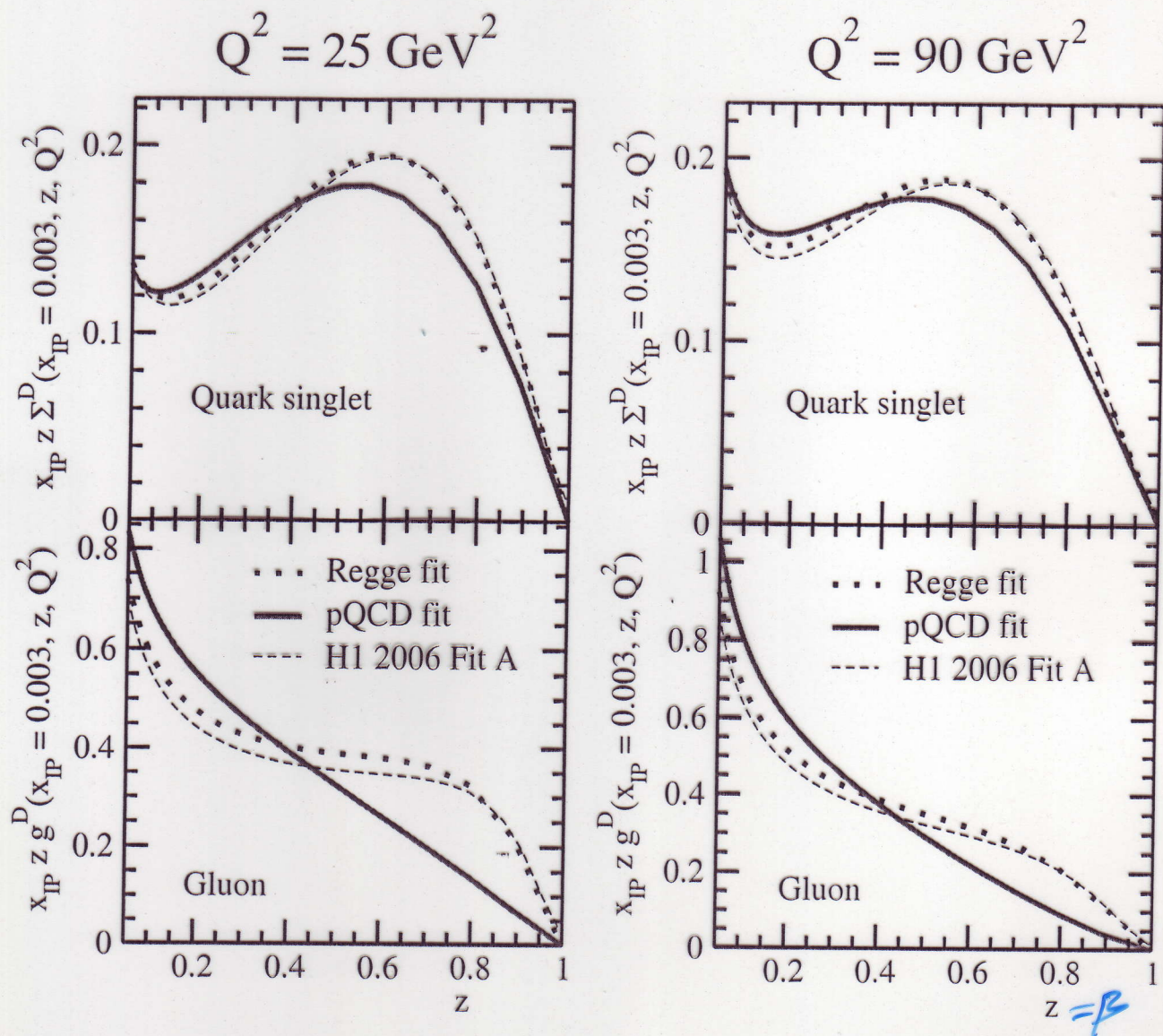


Figure 8: The "Regge" and "pQCD" DPDFs with $Q_{\min}^2 = 8.5 \text{ GeV}^2$ compared to the H1 2006 Fit A.

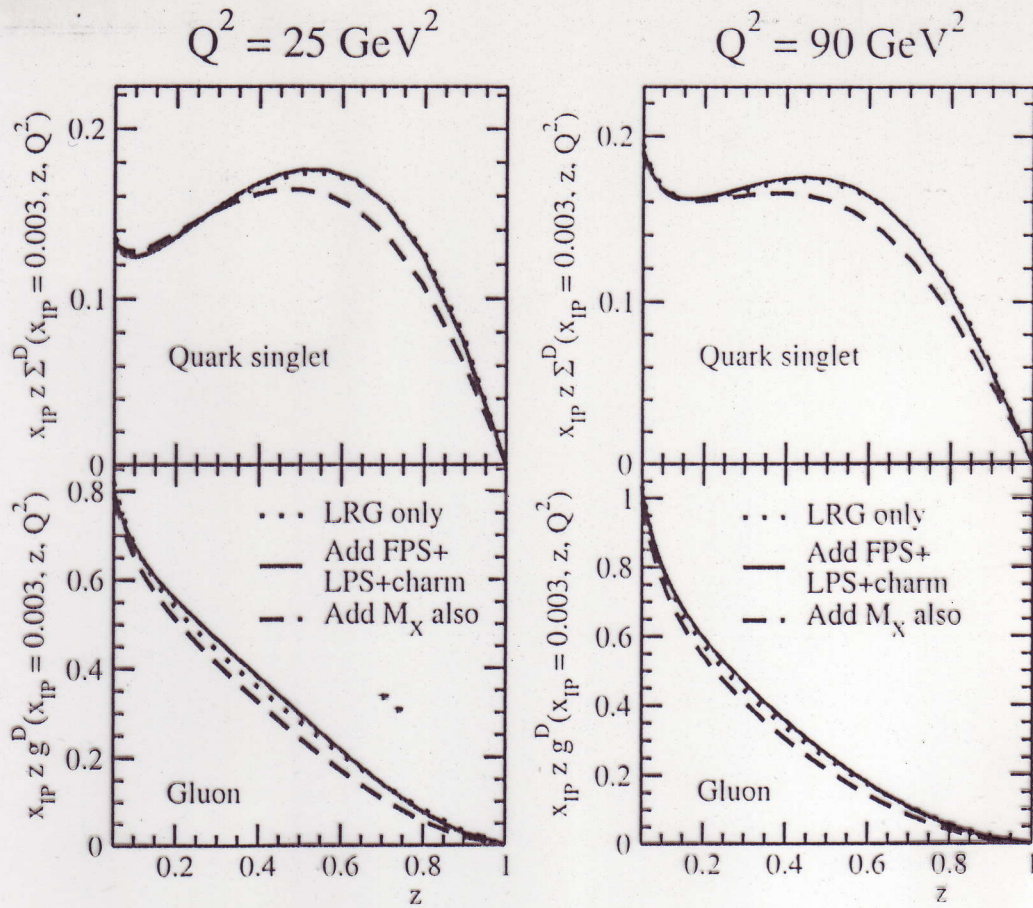


Figure 12: The effect of including additional data sets on the "pQCD" DPDFs.

- $\partial F_2^{D(3)}/\partial \ln Q^2$: Additional contributions to scaling violations apart from DGLAP contribution, important for $\beta \gtrsim 0.3$.
- $\partial F_2^{D(3)}/\partial \ln Q^2$: Peak due to threshold for $\gamma^* \mathbb{P} \rightarrow c\bar{c}$ at $\beta = Q^2/(Q^2 + 4m_c^2)$.

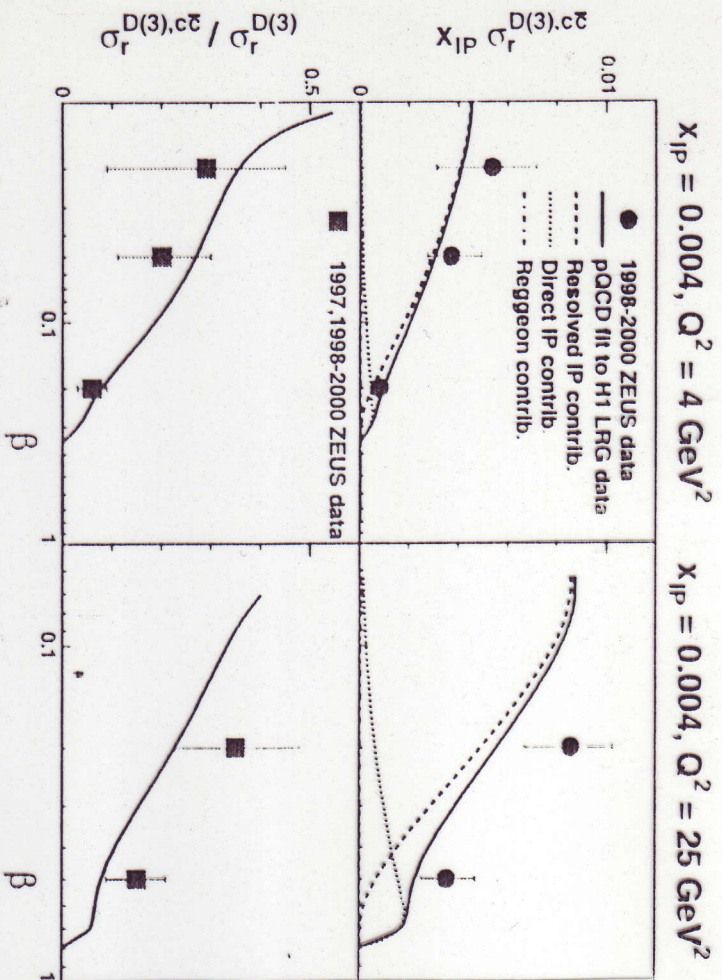


Figure 13: Description of ZEUS diffractive charm data using the “pQCD” DPDFs from a fit to H1 LRG data with $Q_{\min}^2 = 8.5 \text{ GeV}^2$.

