

Light Higgs and Vector-like Quarks without Prejudice

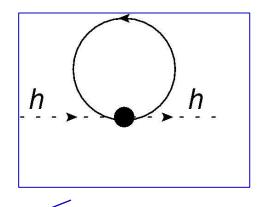
Based on JHEP 1307 (2013) 155 and references therein

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Why Vector-like Quarks?

- Vector-like Quarks (VLQ)?
 - Color triplet spin 1/2 fermions
 - Left and right spinors in the same representation
- Why VLQ?
 - No gauge anomalies
 - Models addressing the EW hierarchy problem
 - Expected to be light: $(m_f \lesssim 1 \text{ TeV})$



After canceling Λ^2 term

$$\delta m_h^2 \approx \frac{3m_t^2}{4\pi^2 v^2} m_f^2 \log \frac{\Lambda^2}{m_f^2}$$

Scope

- SM + single VLQ representation
- Consider VLQ representations of the SM chiral quarks:
 - Singlet up-like VLQ: (3,1,2/3)
 - Singlet down-like VLQ: (3,1,-1/3)
 - Doublet VLQ: (3,2,1/6)
- Phenomenology (Flavor, EWPO)
 - Renormalizable level
 - Leading higher dimensional terms: $H^{\dagger}H\overline{q}_{i}Q, H^{\dagger}H\overline{Q}Q$
- Implications for the Higgs boson

Renormalizable VLQ models

- Mass term in the weak (chiral) basis $-\mathcal{L}_{\text{mass}} = \bar{u}_{L}^{i} \mathcal{M}_{u}^{ij} u_{R}^{j} + \bar{d}_{L}^{i} \mathcal{M}_{d}^{ij} d_{R}^{j} + \text{h.c.}$
 - **Bi-unitary rotations**

•

$$\mathcal{M}_{u,d,\text{diag}} = U_L^{u,d} \mathcal{M}_{u,d} U_R^{u,d\dagger}$$

• Interactions in the physical (mass) eigenbasis

$$\begin{split} V_{ij}^{L} &\equiv (U_{L}^{d})_{jk}^{*}(U_{L}^{u})_{ik} \\ \mathcal{L}_{W} &= -\frac{g}{\sqrt{2}} (V_{ij}^{L} \bar{u}^{i} \gamma^{\mu} P_{L} d^{j} + V_{ij}^{R} \bar{u}^{i} \gamma^{\mu} P_{R} d^{j}) W_{\mu}^{+} + \text{h.c.} , \\ \mathcal{L}_{Z} &= -\frac{g}{2c_{W}} \left(X_{ij}^{u} \bar{u}^{i} \gamma^{\mu} P_{L} u^{j} - X_{ij}^{d} \bar{d}^{i} \gamma^{\mu} P_{L} d^{j} + Y_{ij}^{u} \bar{u}^{i} \gamma^{\mu} P_{R} u^{j} - Y_{ij}^{d} \bar{d}^{i} \gamma^{\mu} P_{R} d^{j} - 2s_{W}^{2} J_{\text{EM}}^{\mu} \right) Z_{\mu} , \\ \mathcal{L}_{h}^{(0)} &= -(X_{ij}^{u} - Y_{ij}^{u}) \frac{m_{j}}{v} \bar{u}^{i} P_{R} u^{j} h - (X_{ij}^{d} - Y_{ij}^{d}) \frac{m_{j}}{v} \bar{d}^{i} P_{R} d^{j} h + \text{h.c.} , \\ &\swarrow &\downarrow &\downarrow &\downarrow &\downarrow &\downarrow &\downarrow &\downarrow \\ X^{u} &\equiv V^{L} V^{L\dagger} &\downarrow &\downarrow &\downarrow &\downarrow &\downarrow \\ Y^{u} &\equiv V^{R} V^{R\dagger} &\downarrow &X^{d} &\equiv V^{L\dagger} V^{L} \end{split}$$

✓ Higgs interactions fixed by the neutral current interactions

Renormalizable Lagrangian

 $-\mathcal{L}_{U}^{(0)} = y_{d}^{ij}\bar{q}_{L}^{i}Hd_{R}^{j} + y_{u}^{ij}\bar{q}_{L}^{i}\tilde{H}u_{R}^{j} + y_{U}^{i}\bar{q}_{L}^{i}\tilde{H}U_{R} + M_{U}\bar{U}_{L}U_{R} + \text{h.c.}$

Mass matrices (V^R=0, Y^{u,d}=0, X^d=1)

$$\mathcal{M}_u = \begin{pmatrix} y_u v / \sqrt{2} & y_U v / \sqrt{2} \\ 0 & M_U \end{pmatrix}, \quad \mathcal{M}_d = (y_d v / \sqrt{2})$$

• Tree level constraints $|X^{u} - \mathbb{I}|_{3 \times 3} < \begin{bmatrix} 0.001 & 2.1 \times 10^{-4} & 0.14 \\ 0.0026 & 0.14 \\ 0.13 \end{bmatrix} \xrightarrow{7} t \rightarrow Z q$ CKM non-unitarity, $Z \rightarrow qq$ • Neglecting the mixing with the first two generations

$$\tan(2\theta_{tU}) = \frac{\sqrt{2}vy_U^t M_U}{M_U^2 - [(y_u^t)^2 + (y_U^t)^2]v^2/2},$$

$$m_t m_{u'} = M_U y_u^t \frac{v}{\sqrt{2}}, \qquad m_t^2 + m_{u'}^2 = M_U^2 + \frac{v^2}{2} [(y_u^t)^2 + (y_U^t)^2].$$

 t – T mixing constrained from the ρ parameter

$$\Delta \rho = \frac{\alpha N_C}{16\pi s_W^2} \frac{m_t^2}{m_W^2} s_{tU}^2 \left[-(1+c_{tU}^2) + s_{tU}^2 r + 2c_{tU}^2 \frac{r}{r-1} \log(r) \right] \quad \gtrsim \$$

• Higgs (and Z) couplings:

 $X_{tt}^{u} = c_{tU}^{2}, X_{tu'}^{u} = c_{tU}s_{tU}$ and $X_{u'u'}^{u} = s_{tU}^{2}$

S: $M X_{u'u'}^{u} = s_{tU}^{2}$ $M u'u' = s_{tU}^{2}$ $M u'u' = s_{tU}^{2}$ $M u'u' = s_{tU}^{2}$ $M u'u' = s_{tU}^{2}$

0.5

0.4

0.3

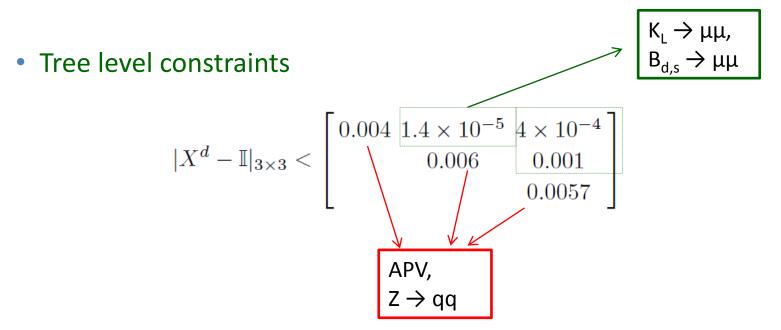
0.2

✓ Negligible deviations in the Higgs production and decays

Renormalizable Lagrangian

 $-\mathcal{L}_{D}^{(0)} = y_{d}^{ij}\bar{q}_{L}^{i}Hd_{R}^{j} + y_{u}^{ij}\bar{q}_{L}^{i}\tilde{H}u_{R}^{j} + y_{D}^{i}\bar{q}_{L}^{i}HD_{R} + M_{D}\bar{D}_{L}D_{R} + \text{h.c.}$

• V^R=0, Y^{u,d}=0, X^u=1



✓ Higgs properties remain SM-like!

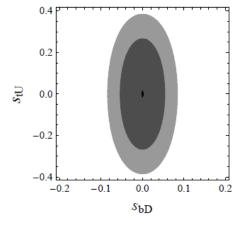
Mass terms (V^L unitary, X^{u,d}=1)

$$\mathcal{M}_u = \begin{pmatrix} y_u v / \sqrt{2} & 0\\ y_U v / \sqrt{2} & M_Q \end{pmatrix}, \quad \mathcal{M}_d = \begin{pmatrix} y_d v / \sqrt{2} & 0\\ y_D v / \sqrt{2} & M_Q \end{pmatrix}$$

Tree level constraints

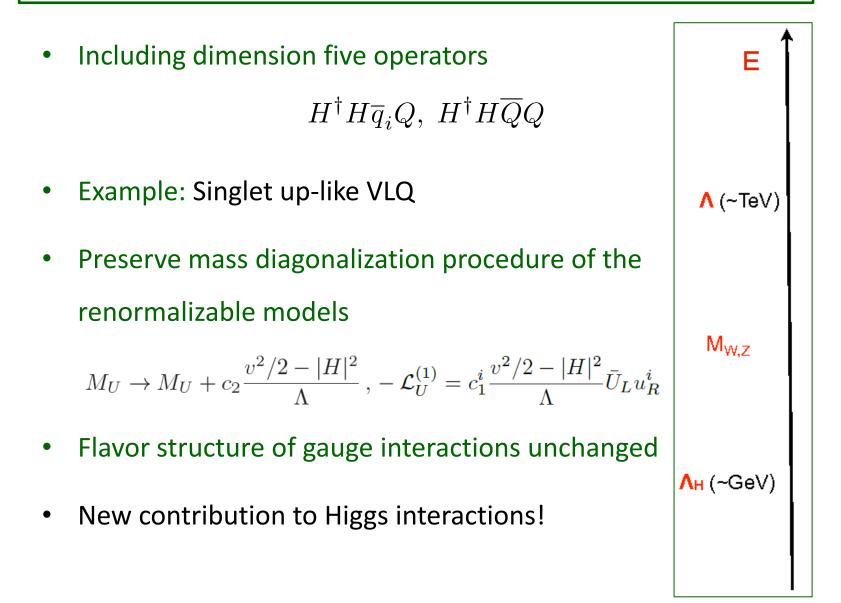
$$|Y^{u}|_{3\times 3} < \begin{bmatrix} 0.11 \ 2.1 \times 10^{-4} \ 0.14 \\ 0.018 \ 0.14 \\ - \end{bmatrix}, \quad |Y^{d}|_{3\times 3} < \begin{bmatrix} 0.1 \ 1.4 \times 10^{-5} \ 4 \times 10^{-4} \\ 0.21 \ 0.001 \\ 0.03 \end{bmatrix}$$

- t-T and b-B system
 - Mass splitting and mixing angles related
 - Mixing angles constrained from Z->bb
 - Mass splitting from ρ parameter



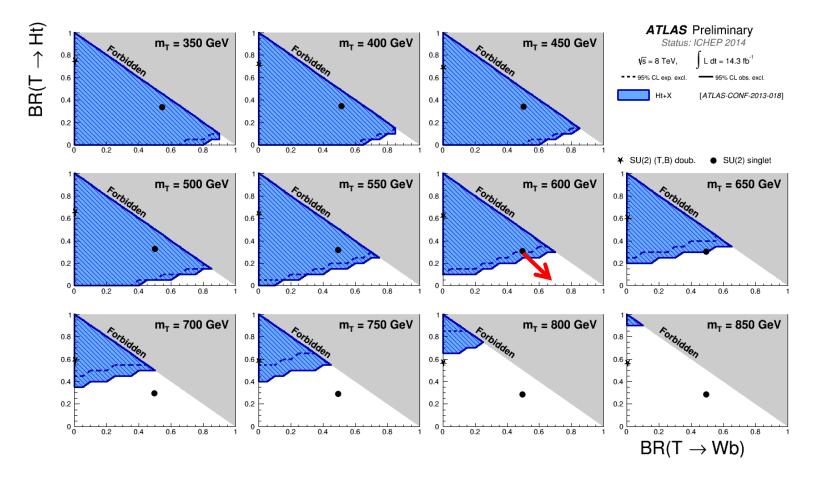
✓ Higgs properties remain SM-like!

Non-renormalizable VLQ models



Implications for VLQ direct searches

- Renormalizable models: fixed branching ratios for $T \rightarrow th$, Wb, Zt
- New contribution to $T \rightarrow ht$



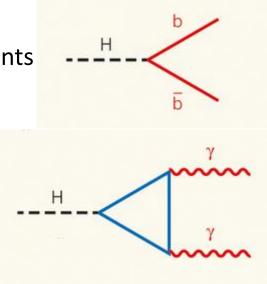
Implications for Higgs physics

Including dimension five operators (mass basis)

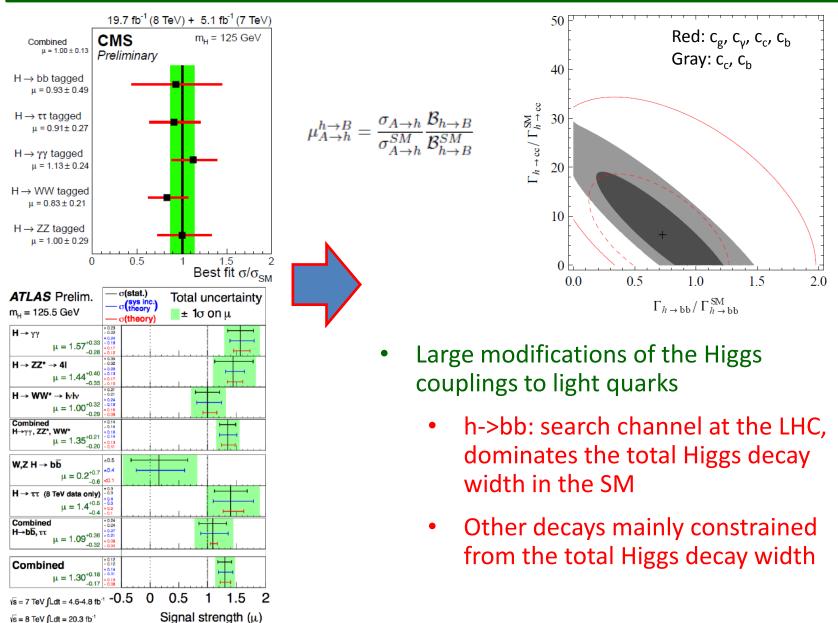
$$\mathcal{L}_{h}^{(1)} = \left(\frac{X_{ij}^{u\prime}}{\Lambda} \bar{u}_{L}^{i} u_{R}^{j} + \frac{X_{ij}^{d\prime}}{\Lambda} \bar{d}_{L}^{i} d_{R}^{j}\right) \left[vh + \frac{h^{2}}{2}\right] + \text{h.c.}$$

- Single Higgs couplings not related to weak currents!
- Di-Higgs couplings
- Flavor violating Higgs couplings
 - K, D, B meson mixings severe constraints
 - \checkmark t \rightarrow hq
- Flavor diagonal Higgs couplings
 - \checkmark h \rightarrow bb, cc, ss, dd, uu
- Loop induced Higgs processes

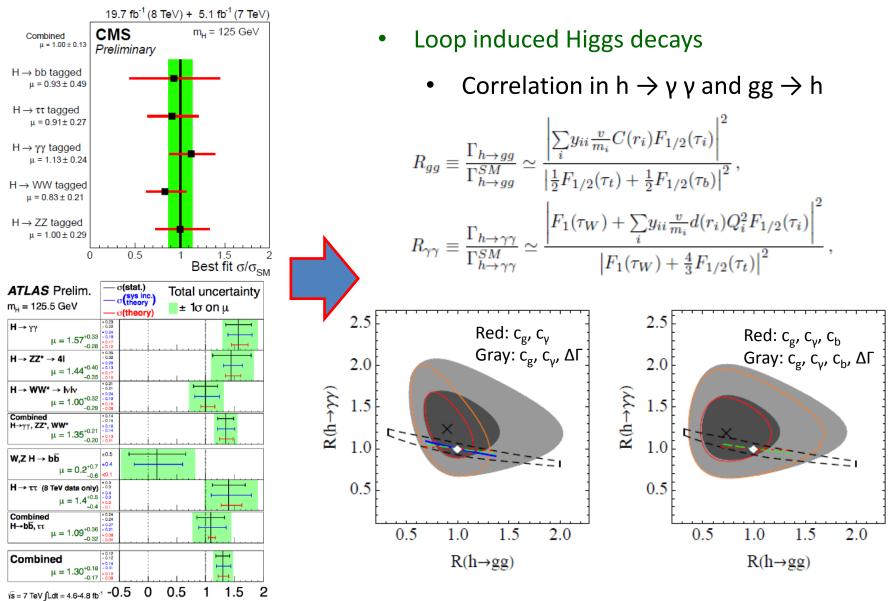
 \checkmark h \rightarrow $\gamma\gamma$, gg \rightarrow h



Impact of the existing LHC Higgs data



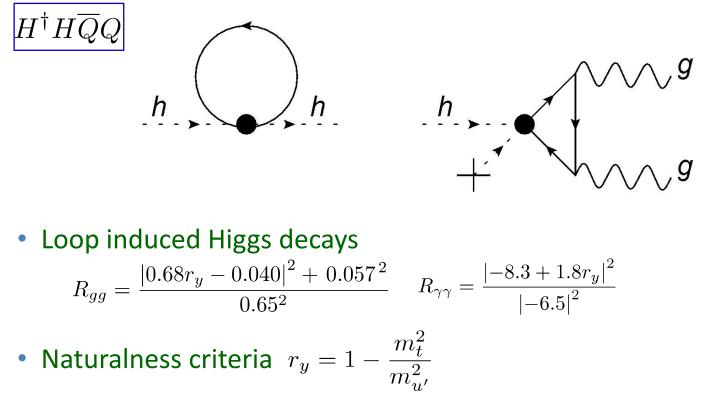
Impact of the existing LHC Higgs data



Signal strength (µ)

vs = 8 TeV (Ldt = 20.3 fb⁻¹

Singlet up-like VLQ: Naturalness and Higgs data



- From the Higgs data fit: m_u > 400 GeV
- Complementary to direct searches

Conclusions

- We study the impact of VLQs on Higgs physics
- Within renormalizable models Higgs properties remain SM-like
- Leading dimension five operators impact the LHC Higgs

data

