

# Improved Jet Algorithms and the Underlying Event

### Matteo Cacciari<sup>1</sup>, Juan Rojo<sup>1</sup>, Gavin Salam<sup>1</sup>, Gregory Soyez<sup>2</sup>

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QCD/EW KEK-D Workshop, ITH Zurich, 03/06/2008

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# Outline

- Jets at the LHC
- Recent developments
- The underlying event
- Performance
- Heavy ions
- Conclusions

### Extras

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# JETS AT THE LHC

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# Counting jets



Three jets? Four jets? ...

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### Jet algorithms

Jet finding not an easy eyeball task ...

A jet algorithm is a non-ambiguous mapping between observed hadrons (or calorimeter towers) and the original hard scattering partons.

It is a systematic procedure that projects away the multi-particle dynamics, so as to leave a simple picture of the event.



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Jets are as close as we can get to a physical single quark or gluon.

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Image: A math a math

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# Analysis flowchart

#### Jet algorithms present in almost all analysis at the LHC



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#### Mainstream jet-algorithms

Iterative cone algorithms (top-down) → (JetClu, ILCA/Midpoint, LHC cones,...) Searches for cones centered on regions of energy flow Dominant at hadron colliders. With split-merge: CDF JetClu, Atlas Cone, PxCone, MidPoint Without split-merge: UA1 cone, CellJet, GetJet, CMS cone?

Sequential recombination algorithms (bottom-up) → (k<sub>t</sub>, Cambridge/Aachen, Jade) Recombine closest pair of particles, next closest, etc. Dominant at e<sup>+</sup>e<sup>-</sup> and ep colliders, and successfully used at Tevatron and now included in LHC analysis software.

#### What is needed of a jet algorithm

- Must be infrared and collinear (IRC) safe
- Must be identical procedure at parton level, hadron-level and experimental level So that theory calculations can be compared to measurements

#### What is **nice** for a jet algorithm

- Shouldn't be too sensitive to hadronization, underlying event and pileup, while being sensitive to perturbative radiation.
- Should be realistically applicable at detector level (i.e. allow fast implementations, to cope with the large multiplicities at hadronic colliders).

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# RECENT DEVELOPEMENTS IN JET ALGORITHMS

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References:

- M. Cacciari and G. Salam, hep-ph/0512210
- G. Salam and M. Cacciari, arXiv: arXiv:0704.0292
- M. Cacciari, G. Salam and G. Soyez, arXiv:0802.1189

# Progress in jet algorithms

The FastJet package: a fast implementation of the  $k_T$  algorithm (based on computational geometry) with background subtraction methods and external jet finders as plug-ins (Cacciari, Salam and Soyez 05-08).

### http://www.lpthe.jussieu.fr/~ salam/fastjet/

New jet algorithms have appeared in the last years:

- SISCone: A Seedless and Infrared Safe Cone algorithm (Salam and Soyez 07). IRC safe at all orders of pert. theory. Finds all stable cones using geometrical techniques.
- The anti-k<sub>T</sub> algorithm: A generalization of the k<sub>T</sub> algorithm with measure:

$$d_{ij} = \min(k_{ti}^{2n}, k_{tj}^{2n}) \Delta R_{ij}^2 / R^2$$
,  $n = -1$ ,  $d_{iB} = k_{ti}^2$ 

which acts like a perfect cone algorithm with area  $A_{jet} = \pi R^2$  (Cacciari, Salam, Soyez 08). For  $k_T$  one has n = 1 and for Cambridge n = 0.

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## Infrared safety

# Cone algorithms have been known to suffer from Infrared and Collinear unsafety for many years.

For the CDF MidPoint cone algorithm (even worse for the Iterative Cone!):

Table 2: Summary of the order  $(\alpha_s^4 \text{ or } \alpha_s^3 \alpha_{EW})$  at which stable cones are missed in various processes with a midpoint algorithm, and the corresponding last order that can be meaningfully calculated. Infrared unsafety first becomes visible one order beyond that at which one misses stable cones.

Theory investment in NLO computations:  $\sim 50$  people  $\times$  10 years  $\sim 30-50$  million \$  $\rightarrow$  Lost if IRC unsafe jet algorithms used at the LHC!

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Observable	1st miss cones at	Last meaningful order
Inclusive jet cross section	NNLO	NLO
W/Z/H + 1 jet cross section	NNLO	NLO
3 jet cross section	NLO	LO
W/Z/H + 2 jet cross section	NLO	LO
jet masses in 3 jets, $W/Z/H + 2$ jets	LO	none

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### Infrared safety

IRC unsafety affects a large fraction of events!

For JetClu (similar to Atlas cone), half of events fails IRC safety tests.

Even for the MidPoint cone algorithm, 15% of events fail the test!



CMS recommends the use of SISCone as reference cone algorithm, and a solution of the second s

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## Progress in jet algorithms - Status 2006

Algorithm	Туре	IRC status	Notes	
exclusive $k_t$	$SR_{p=1}$	OK		
inclusive $k_t$	$SR_{p=1}$	OK	widespread: QCD-th, HERA	
Cambridge/Aachen	$SR_{\rho=0}$	OK		
Run II Seedless cone	SC-SM	OK	slow: N2 <sup>N</sup> !!	
CDF JetClu	IC <sub>r</sub> -SM	IR <sub>2+1</sub>	for top physics, searches	
CDF MidPoint cone	IC <sub>mp</sub> -SM	IR <sub>3+1</sub>	$\simeq$ Tev Run II recommend <sup><math>n</math></sup>	
CDF MidPoint searchcone	IC <sub>se,mp</sub> -SM	IR <sub>2+1</sub>		
D0 Run II cone	IC <sub>mp</sub> -SM	IR <sub>3+1</sub>	Tev Run II + cut on cone $p_t$	
ATLAS Cone	IC-SM	IR <sub>2+1</sub>		
PxCone	IC <sub>mp</sub> -SD	IR <sub>3+1</sub>	has cut on cone $p_t$ ,	
CMS Iterative Cone	IC-PR	Coll <sub>3+1</sub>		
PyCell/CellJet (from Pythia)	FC-PR	Coll <sub>3+1</sub>	widespread in BSM theory	
GetJet (from ISAJET)	FC-PR	Coll <sub>3+1</sub>	likewise	

SR = seq.rec.; IC = it.cone; FC = fixed cone;

SM = split-merge; SD = split-drop; PR = progressive removal

 $R_{n+1}$ : for *n* nearby hard partons, 1 soft emitted gluon can change hard jets  $Coll_{n+1}$ : for *n* nearby hard partons, 1 collinear splitting can change hard jets

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## Progress in jet algorithms - Status 2008

Algorithm	Туре	IRC status	Notes	
exclusive $k_t$	$SR_{p=1}$	OK		
inclusive $k_t$	$SR_{p=1}$	OK	widespread: QCD-th, HERA	
Cambridge/Aachen	$SR_{p=0}$	OK		
Run II Seedless cone	SC-SM	OK	slow: N2 <sup>N</sup> !!	
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CDF MidPoint searchcone	$IC_{se,mp}$ -SM	IR <sub>2+1</sub>		
D0 Run II cone	IC <sub>mp</sub> -SM	IR <sub>3+1</sub>	Tev Run II + cut on cone $p_t$	
ATLAS Cone	IC-SM	IR <sub>2+1</sub>		
PxCone	IC <sub>mp</sub> -SD	IR <sub>3+1</sub>	has cut on cone $p_t$ ,	
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## Progress in jet algorithms - Status 2008

Algorithm	Туре	IRC status	Evolution	
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inclusive $k_t$	$SR_{p=1}$	OK	$N^3 \rightarrow N \ln N$	
Cambridge/Aachen	$SR_{p=0}$	OK	$N^3 \rightarrow N \ln N$	
Run II Seedless cone	SC-SM	OK	$\rightarrow$ SISCone	
CDF JetClu	IC <sub>r</sub> -SM	IR <sub>2+1</sub>	$[\rightarrow SISCone]$	
CDF MidPoint cone	IC <sub>mp</sub> -SM	IR <sub>3+1</sub>	$\rightarrow$ SISCone	
CDF MidPoint searchcone	IC <sub>se,mp</sub> -SM	IR <sub>2+1</sub>	$[\rightarrow SISCone]$	
D0 Run II cone	IC <sub>mp</sub> -SM	IR <sub>3+1</sub>	$\rightarrow$ SISCone [with $p_t$ cut?]	
ATLAS Cone	IC-SM	IR <sub>2+1</sub>	$\rightarrow$ SISCone	
PxCone	IC <sub>mp</sub> -SD	IR <sub>3+1</sub>	[little used]	
CMS Iterative Cone	IC-PR	Coll <sub>3+1</sub>	$\rightarrow$ anti- $k_t$	
PyCell/CellJet (from Pythia)	FC-PR	Coll <sub>3+1</sub>	$ ightarrow$ anti- $k_t$	
GetJet (from ISAJET)	FC-PR	Coll <sub>3+1</sub>	$\rightarrow$ anti- $k_t$	

SR = seq.rec.; IC = it.cone; FC = fixed cone;

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### Progress in jet algorithms

Disclaimer: topics which cannot be covered include:

- Jet flavour and b-jets (Banfi, Salam and Zanderighi, hep-ph/0601139, arXiv:0704.2999 [hep-ph])
- Analytical studies of jet algorithms (Dasgupta, Magnea and Salam, arXiv:0712.3014 [hep-ph])
- Strategies to asses jet algorithms performance, Rabbertz et al., Campanelli et al. in arXiv:0803.0678
- Jet substructure as a new Higgs search channel at the LHC, Butterworth, Davison, Rubin, Salam, arXiv:0802.2470

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# JET AREAS, THE UNDERLYING EVENT AND ITS MEASURE

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References:

M. Cacciari talk at Les Houches 2007 Cacciari and Salam, arXiv:0707.1378 Cacciari, Salam and Soyez, arxiv:0802.1188

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### Jet areas

Jet areas provide a quantitative definition of the concept of the *size* of a jet (Cacciari, Salam and Soyez, arxiv:0802.1188) Various definitions (active, passive, Voronoi), coincide for high multiplicity Active area  $\rightarrow$  add a large number of ghosts (particles with infinitesimal  $p_T$ ) and perform jet clustering:



Area computations only meaningful for IRC safe algorithms and practical with fast implementations  $\rightarrow$  Both conditions fulfilled in FastJet.

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### Jet areas



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 $A_{\rm jet} \neq \pi R^2$  in general (but note anti- $k_T$ )

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### Jet areas



SISCone smallest active area  $\rightarrow$  Reduced sensitivity to diffuse soft radiation

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# Standard approach to the UE



Marchesini-Webber idea: look at transverse region to measure underlying event



### **Topological selection**

The jets are classified as belonging to the noise on the ground of their **position** 

But: difficult to use in busy events, misses point-to-point and event-by event fluctuations

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### Area-based UE measure and subtraction

We propose a dynamical selection: jets assigned to background due to its characteristics: UE/PU can be measured and subtracted based on jet areas Basic idea: Jets formed of soft stuff only have a  $p_T$  proportional to its area Take at  $p_T \sim 200$  GeV dijet event with 20 min. bias collisions:



### Area-based UE measure and subtraction

We propose a dynamical selection: jets assigned to background due to its characteristics: UE/PU can be measured and subtracted based on jet areas

► Determine the noise density per unit area  $\rho$  in an event-by-event basis with the  $k_T$ -algorithm and R = 0.5

$$\rho \equiv \text{median}\left[\left\{\frac{p_{tj}}{A_j}\right\}\right] \tag{1}$$

Subtract the PU contribution to the jet using its area

$$\boldsymbol{p}_{\mu j}^{(\mathrm{sub})} = \boldsymbol{p}_{\mu j} - \boldsymbol{A}_{\mu j} \, \boldsymbol{\rho} \pm \sigma_{\rho} \sqrt{\boldsymbol{A}_{\mathrm{j}}} \tag{2}$$

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Note that no cut in the  $p_T$  of input particles required N.B. : Intrinsic uncertainty in subtraction due to background fluctuations  $\sigma_{\rho}$ Other limitations include backreaction (effect of min. bias particles in jet clustering)

### Subtraction in practice

First example: Z' invariant mass distribution



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### Subtraction in practice

Second example: Z' Z' invariant mass distribution



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### Subtraction in practice

Second example: Z' Z' invariant mass distribution



Note that subtraction not only corrects the peak position, it also improves sizably the mass resolution

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## Measuring PileUp and Underlying event

The area-based subtraction strategy can also be used to measure UE and PU directly from experimental data (or MC models).



### Model independent UE/PU measurement

Note the size of the background fluctuations (basic limit on the accuracy of method)

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## Measuring PileUp and Underlying event

The area-based subtraction strategy can also be used to measure UE and PU directly from experimental data (or MC models). Compare measured UE with known amount in MC:



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### Underlying event distributions

Measure distribution of  $\rho_{UE}$  from Pythia and Herwig at TeV and LHC:



Proposal: Repeat the comparison with modern event generators  $\rightarrow$  Last call for predictions before LHC start-up

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### Underlying event distributions

Measure distribution of  $\rho_{UE}$  from Pythia and Herwig at TeV and LHC:



Proposal: Repeat the comparison with modern event generators  $\rightarrow$  Last call for predictions before LHC start-up

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# QUANTIFYING THE PERFORMANCE OF JET ALGORITHMS AT THE LHC

References:

M. Cacciari, J. Rojo, G. Soyez and G. Salam, Les Houches 2007, Tools and Jets Summary Report, arXiv:0803.0678

+ Work in progress

Juan Rojo <u>Improved Je</u>t Algorithms and the Underlying Event



### The performance of jet algorithms

How to determine which jet algorithm is optimal for a given process? We need quality measures that:

- Do not depend on ill-defined quantities (like original parton p<sub>T</sub> or direction, a parton is an ambiguous concept in pQCD)
- Do not assume any underlying probability distribution of (MC) data, *i.e.*, whether the reconstructed mass distribution is gaussian, asymmetric, ...

Processes studied:

- Fictitious narrow width  $Z' \rightarrow q\bar{q}$  decays for various  $M_{Z'}$  (100 GeV 4 TeV): Source of mono-energetic quark jets
- Fictitious narrow width  $H \rightarrow gg$  decays for various  $M_H$  (100 GeV 4 TeV): Source of mono-energetic gluon jets

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1.  $Q_{f=z}^{w}(R) \rightarrow$  The width of the smallest (reconstructed) mass window that contains a fraction f = z of the generated massive objects:

 $f = \left(\frac{\# \text{ reco. massive objects in window of width } w}{\text{Total } \# \text{ generated massive objects}}\right)$ 

2. The max. fraction of evs. f in window of width  $w = x\sqrt{M}$ :

$$Q_{w=x\sqrt{M}}^{f}(R) \equiv \left(\frac{\text{Max $\#$ reco. mass. obj. in width $w = x\sqrt{M}$}}{\text{Total $\#$ generated massive objects}}\right)^{-1}$$

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Consistent results obtained with both measures

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### Quality measures

 Q<sup>w</sup><sub>f=z</sub>(R) → The width of the smallest (reconstructed) mass window that contains a fraction f = z of the generated massive objects:

$$f = \left(\frac{\# \text{ reconstructed massive objects in window of width } w}{\text{Total } \# \text{ generated massive objects}}\right)$$

2.  $Q_{w=x\sqrt{M}}^{f}(R) \rightarrow$  The max. fraction of events f in window of width  $w = x\sqrt{M}$ :

 $Q^f_{w=x\sqrt{M}}(R) \equiv \left(\frac{\mathrm{Max}\ \#\ \mathrm{reconstructed\ massive\ objects\ in\ window\ of\ width\ w=x\sqrt{M}}}{\mathrm{Total}\ \#\ \mathrm{generated\ massive\ objects}}\right)^{-1}\ ,$ 



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Improved Jet Algorithms and the Underlying Event

### Effective luminosity ratio

Mapping variations of quality measures and of effective luminosity ratio  $\rho_{\mathcal{L}}$  (Crude) Assumption: flat background, unaffected by jet clustering. Define effective power to discriminate signal from background:

$$\begin{split} \boldsymbol{\Sigma}^{\text{eff}}\left(\text{JA}, \boldsymbol{R}\right) &\equiv \frac{N_{\text{signal}}}{\sqrt{N_{\text{back}}}}\\ \boldsymbol{r_w} &\equiv \frac{Q_{f=z}^w\left(\text{JA}_2, R_2\right)}{Q_{f=z}^w\left(\text{JA}_1, R_1\right)} = \frac{N_{\text{back}}\left(\text{JA}_2, R_2\right)}{N_{\text{back}}\left(\text{JA}_1, R_1\right)},\\ &\frac{\boldsymbol{\Sigma}^{\text{eff}}\left(\text{JA}_2, R_2\right)}{\boldsymbol{\Sigma}^{\text{eff}}\left(\text{JA}_2, R_2\right)} = \sqrt{r_w}, \quad \boldsymbol{\rho_{\mathcal{L}}} = \frac{1}{r_w}.\\ \boldsymbol{r_f} &\equiv \frac{Q_{w=x\sqrt{M}}^f\left(\text{JA}_2, R_2\right)}{Q_{w=x\sqrt{M}}^f\left(\text{JA}_1, R_1\right)} = \frac{N_{\text{signal}}\left(\text{JA}_1, R_1\right)}{N_{\text{signal}}\left(\text{JA}_2, R_2\right)},\\ &\frac{\boldsymbol{\Sigma}^{\text{eff}}\left(\text{JA}_1, R_1\right)}{\boldsymbol{\Sigma}^{\text{eff}}\left(\text{JA}_2, R_2\right)} = r_f, \quad \boldsymbol{\rho_{\mathcal{L}}} = \frac{1}{r_f^2}. \end{split}$$

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### Jet algorithms

We study the performance of modern, Infrared and Collinear safe jet algorithms:

- 1.  $k_T$  algorithm
- 2. Cambridge/Aachen (Cam/Aa) algorithm
- 3. SISCone algorithm
- 4. Anti- $k_T$  algorithm
- 5.  $Cam/Aa(filt) \rightarrow novel jet finding strategy based on Cambridge/Aachen with subjet filtering$

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We would like a to have a jet algorithm which keeps most perturbative radiation while *filtering away* most of the UE/PileUp noise  $\rightarrow$ 

- 1. Cluster all the particles in the event with a given jet definition  $(JA_1, R_1)$ .
- 2. Take each of the jets of the event and cluster its constituents only with another jet definition  $(JA_2, R_2)$  with  $R_2 < R_1$ . This operation will provide us with a set of subjets of the original jet.
- 3. Keep the *n*<sub>sj</sub> subjets of a jet with largest *p*<sub>T</sub> and throw way the remaining subjets.
- 4. Original jets are replaced merging the selected subjets

This procedures defines the subjet filtering strategy. Natural choice  $\rightarrow JA_1=JA_2 \rightarrow Cambridge/Aachen algorithm (distance <math>d_{ij}$  based on angular separation).

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Does the subjet filtering really improve jet finding? In J. Butterworth et al., (arXiv:0802.2470 [hep-ph]), it was shown that using a SJA-type algorithm, the Higgs boson search channels  $Z/WH(\rightarrow b\bar{b})$  at  $M_H \sim 120$  GeV transform from background-swamped channels to promising discovery channels.

Jet definition	$\sigma_S/{\rm fb}$	$\sigma_B/{\rm fb}$	$S/\sqrt{B \cdot fb}$
C/A, R = 1.2, MD-F	0.57	0.51	0.80
$K_{\perp}, R = 1.0, y_{cut}$	0.19	0.74	0.22
SISCone, $R = 0.8$	0.49	1.33	0.42

TABLE I: Cross section for signal and the Z+jets background in the leptonic Z channel for 200 <  $p_{TZ}/\text{GeV}$  < 600 and 110 <  $m_J/\text{GeV}$  < 125, with perfect b-tagging; shown for our jet definition, and other standard ones at near optimal R values.



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### Motivation for study of the subjet filtering in a general context.

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The performance of jet algorithms - Narrow  $H \rightarrow gg$ 



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The performance of jet algorithms - Narrow  $H \rightarrow gg$ 



Less favored choices for the  $M_H = 2$  TeV case:

1. Use SISCone, but  $R_{\text{best}}^{100 \text{ GeV}} = 0.6$  instead of  $R_{\text{best}}^{2 \text{ TeV}} = 1.1 \rightarrow \rho_{\mathcal{L}} \sim 0.55$ 2. Use  $R_{\text{best}}^{2 \text{ TeV}}$ , choose not SISCone, SubJet/Filtering but  $k_{\mathcal{T}} \rightarrow \rho_{\mathcal{L}} \sim 0.6$ In both cases  $\rightarrow$  Lose almost half effective discriminating power  $\Sigma^{\text{eff}}$ !

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### The performance of jet algorithms - $R_{\text{best}}$

Gluon jets require larger values of  $R_{\text{best}}$  than quark jets. For quark jets with  $p_T \ge 250 \text{ GeV} \rightarrow R_{\text{best}} \ge 0.7$ For gluon jets with  $p_T > 250 \text{ GeV} \rightarrow R_{\text{best}} > 0.9$ Good jet resolution requires large values of R

Approximate scaling  $R_{\text{best}} \sim \ln M$  (pert. radiation)



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### The PileUp case

Check robustness of results when PileUp present ( $\mathcal{L}_{high} = 0.25 \text{mb}^{-1}$ )



After subtraction  $R_{
m best}^{
m PU} \sim R_{
m best}^{
m no~PU}$  for  $p_T^{
m jet} \ge 150 \ {
m GeV}$ 

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Improved Jet Algorithms and the Underlying Event

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### The PileUp case

Check robustness of results when PileUp present ( $\mathcal{L}_{high} = 0.25 \text{mb}^{-1}$ )



For  $p_T^{\text{jet}} \ge 250 \text{ GeV}$  jets  $R_{\text{best}}^{\text{PU}} \ge 0.9$  also in the presence of high lumi PU Cam/Aa(filt) has  $R_{\text{best} \ge 1}$  due to its reduced jet area

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### The PileUp case

Check robustness of results when PileUp present ( $\mathcal{L}_{high} = 0.25 \text{mb}^{-1}$ )



Same hierarchy of jet algorithms that in the no PU case

SISCone and Cam/Aa(filt) are the best jet algorithms for all values of the jet p<sub>T</sub> (p<sub>T</sub><sup>jet</sup> ∼ M<sub>H</sub>/2)

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# JET RECONSTRUCTION IN HEAVY ION COLLISIONS

References:

M. Cacciari, G. Salam, arXiv:0707.1378

N. Armesto, M. Cacciari, J. Rojo, G. Salam, C. Salgado, G. Soyez, work in

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progress

#### Jets in heavy ion collisions

Jets will be of paramount importance to fully exploit the potential of the HIC program at the LHC

Jets will be most abundant hard probes in HIC at the LHC From CMS HIC TDR (J. Phys. G: Nucl. Part. Phys. 34 2307)

	PbPb $\sqrt{s_{NN}} = 5.5 \text{ TeV}$ $\mathcal{L} = 5 \times 10^{26} \text{ cm}^{-2} \text{ s}^{-1}$		pPb $\sqrt{s_{NN}} = 8.8 \text{ TeV}$ $\mathcal{L} = 1.4 \times 10^{30} \text{ cm}^{-2} \text{ s}^{-1}$	
Process	Yield/10 <sup>6</sup> s	Ref.	Yield/10 <sup>6</sup> s	Ref.
	$ \eta $ :	<b>≦</b> 2.4		
$jet (p_T > 50 \text{GeV/c})$	$2.2  imes 10^7$	[47]	$1.5 imes10^{10}$	[48]
jet $(p_{\rm T} > 250 {\rm GeV/c})$	$2.2 \times 10^3$	[47]	$5.2 \times 10^6$	[48]
$Z^0$	$3.2 \times 10^5$	[49]	$6.8  imes 10^6$	[48]

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Table 1.1. The expected yield of several hard probes in 10<sup>6</sup> s PbPb and pPb LHC runs

#### Jets in heavy ion collisions

Jets will be of paramount importance to fully exploit the potential of the HIC program at the LHC

- Jets will be most abundant hard probes in HIC at the LHC
- Jets free of inclusive particle measurements biases
- A solid pQCD baseline is required to detect and quantify medium effects
- Need to asses to which extend can reconstructed QCD jets be disentangled from background, and which is the minimum size of medium effects which could then be probed

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All the successful jet technology from *pp* can be transferred to the extremely dense environment of heavy ion collisions

#### Event characterization

#### Wide range of predictions for UE in HIC at the LHC

Process	$\langle N_{\rm particles} \rangle$	$\left  \left\langle \frac{dN}{d\eta} \right _{\eta=0} \right\rangle$	$\left\langle \left. \frac{dN_{\rm ch}}{d\eta} \right _{\eta=0} \right\rangle$	$\left<  ho_{(\eta,\phi)=(0,0)} \right>$	T [s]
pp  ightarrow gg	160	30	15	0.5 GeV	$2 \cdot 10^{-4}$
$pp \rightarrow gg(+PbPb/SH)$	$4.7 \cdot 10^{4}$	5350	3020	450 GeV	1.2
$pp \rightarrow gg(+PbPb/NSH)$	$2.7 \cdot 10^{4}$	2230	1230	150 GeV	0.2

Note that estimations for background density  $\rho \in (100 - 500)$  GeV satisfy tipically  $\rho \ge p_T^{\text{jet}} \rightarrow \text{New regime for jet finding}$ Note also clustering times (without cuts in  $p_T$ )

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#### A typical dijet event in HIC at LHC

pp 
ightarrow gg events with  $p_T^{
m jet} \sim 100$  GeV and R = 0.4 - No PbPb



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 $pp \rightarrow gg$  events with  $p_{T}^{\rm jet} \sim 100~{\rm GeV}$  and R=0.4 -  ${\rm PbPb}$  model NSH



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#### A typical dijet event in HIC at LHC

 $pp \rightarrow gg$  events with  $p_T^{\rm jet} \sim 100$  GeV and R=0.4 - PbPb model SH



The large UE in HIC can be measured and subtracted with similar techniques as in the pp case

Caveat: background level depends on jet position in  $(\eta, \phi)$  plane First option: Extract parabolic dependence  $\rho(y) = \rho_0 + \rho_1 y^2$ 



The large UE in HIC can be measured and subtracted with similar techniques as in the pp case

Second option: Circular range of D = kR centered on jet axis (reduce

sensitivity to background structure details)



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#### Background subtraction in practice

Example: Generate Pythia events with  $p_T^{\min} = 100 \text{ GeV}$ Hardest jet distribution from  $k_T$  jet algorithm and R = 0.4



#### Hardest jet distribution

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#### Background subtraction in practice

Example: Generate Pythia events with  $p_T^{\min} = 100 \text{ GeV}$ Hardest jet distribution from  $k_T$  jet algorithm and R = 0.4



#### Hardest jet distribution

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### Background subtraction in practice

Example: Generate Pythia events with  $p_T^{\min} = 100 \text{ GeV}$ Hardest jet distribution from  $k_T$  jet algorithm and R = 0.4



#### Hardest jet distribution

Jets down to  $p_T \sim 50$  GeV can be reconstructed even in worst HIC scenario = 2000

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#### Jet shape

Jet substructure  $\phi(r)$  useful discriminator of medium effects Cluster jet constituents with  $R_{\rm sj}$  ( $r\equiv R_{\rm sj}/R_{\rm jet}<1$ ) and keep hardest subjet with  $p_T^{\rm sj}(\leq p_T^{\rm jet})$ .  $R_{\rm jet}=0.5,~0.15\leq R_{\rm sj}\leq 0.5$ 



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#### Jet shape

#### $R_{AA}(\phi(r)) \equiv \phi_{\rm med}(r)/\phi_{\rm vac}(r)$ , with $\phi_{\rm vac}(r)$ from pp jets



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#### Jet shape - LL pQCD

If medium effects parametrized by  $(1 + f_{med})$  in the singular part of the splitting functions (Borghini et al. 05) then:

$$\phi_{\mathrm{med}}(r, f_{\mathrm{med}}) = 1 - \frac{\alpha_s}{\pi} \ln \frac{R_{\mathrm{j}}}{R_{\mathrm{sj}}} \left[ C_A \left( \frac{3f_{\mathrm{med}}}{8} + 2\ln 2 - \frac{43}{96} \right) + \frac{7N_f T_R}{48} \right]$$

for  $\alpha_s = 0.2$  and  $f_{med} \sim 3 \ [\phi_{vac}(r) = \phi_{med}(r, f_{med} = 0)] \rightarrow \text{Agreement with}$  $\phi(r)$  results from MC simulations + subjets ( $L = 2 \text{ fm}, \ \hat{q}L = 20 \text{ GeV}^2$ )



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#### Conclusions and outlook

- Sizable progress in jet algorithms in the recent years
- Area-based method provide an alternative strategy for measurement/subtraction of Underlying Event and Pile-Up
- It is necessary now to quantify the performance of jet algorithms and define optimal choices
- Modern jet finding technology has also important potential applications in heavy ion collisions
- Still room for further improvements, both in theory and in experiment: close collaboration essential to fully exploit the LHC potential!

# Thanks for your attention!

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#### Conclusions and outlook

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# Thanks for your attention!

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# **ADDITIONAL MATERIAL**

Juan Rojo Improved Jet Algorithms and the Underlying Event ・ロ・・団・・ヨ・・ヨ・ 「田・今へで」

## The Anti- $k_T$ algorithm

The Anti- $k_T$  algorithm has a very reduced sensitivity to Back-Reaction (effect of UE particles in jet clustering):

$$\left. \Delta \rho_T^{\mathrm{BR}} \right|_{k_T} \sim 5 \mathrm{GeV} \ , \qquad \left. \Delta \rho_T^{\mathrm{BR}} \right|_{\mathrm{Anti}-k_T} \sim 1 \mathrm{GeV}$$

for  $p_T^{
m jet} \sim$  100 GeV, R = 0.5,  $ho \sim$  150 GeV.



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Improved Jet Algorithms and the Underlying Event

## The Anti- $k_T$ algorithm

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The performance of jet algorithms - Narrow Z' 
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The performance of jet algorithms - Narrow Z' 
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#### The performance of jet algorithms - Narrow Z' ightarrow q ar q



Less favored choices for the  $M_{Z'} = 2$  TeV case:

1. Use SISCone, but  $R_{\rm best}^{100 \text{ GeV}} = 0.5$  instead of  $R_{\rm best}^{2 \text{ TeV}} = 0.8 \rightarrow \rho_{\mathcal{L}} \sim 0.7$ 

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2. Use  $k_T$  instead of SISCone or SubJet/Filtering  $\rightarrow \rho_{\mathcal{L}} \sim 0.8$ 

In both cases  $\rightarrow$  Lose  $\sim$  25% of effective discrimination power  $\Sigma^{\rm eff}!$ 

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#### The performance of jet algorithms - Hadronic $t\bar{t}$

Probe jet algorithms performance in a complicated, multi-jet environment.

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#### The performance of jet algorithms - Hadronic $t\bar{t}$

Probe jet algorithms performance in a complicated, multi-jet environment.



All jet algs. blue similar performance near  $R_{\rm opt} \sim 0.3 - 0.4$ Use SISCone, but with R = 0.7 instead of  $R_{\rm best} \sim 0.4 \rightarrow \rho_{\mathcal{L}} \sim 0.5$ The effective lumi ratio  $\rho_{\mathcal{L}}$  worsens very rapidly for large values of  $R > R_{\rm best}$ .

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Juan Rojo

### SJA with Pile-Up

The SJA leads to better signal discrimination and is less sensitive to Pile-Up than Cam/Aa:

Juan Rojo Improved Jet Algorithms and the Underlying Event



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The SJA leads to better signal discrimination and is less sensitive to Pile-Up than Cam/Aa:



#### Example: the $k_t$ algorithm:

1. Calculate (or update) distances between all particles *i* and *j*, and between *i* and beam:

$$d_{ij} = \min(k_{ti}^2, k_{tj}^2) \frac{\Delta R_{ij}^2}{R^2}, \qquad d_{iB} = k_{ti}^2, \qquad \Delta R_{ij}^2 = \Delta y_{ij}^2 + \Delta \phi_{ij}^2$$

2. Find smallest of  $d_{ij}$  and  $d_{iB}$ 

If d<sub>ij</sub> is smallest, recombine i and j (add result to particle list, remove i, j)

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- if  $d_{iB}$  is smallest call *i* a jet (remove it from list of particles)
- 3. If any particles are left, repeat from step 1.

One parameter: R (like cone radius), whose natural value is 1  $k_t$  algorithm attempts approximate inversion of the QCD shower branching process  $\rightarrow$  Theoretical sound basis.

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# Sequential recombination algorithms

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The performance of jet algorithms -  $R_{\rm best}$ 

Compare the MC results with analytical estimations of  $R_{\rm best}$  from arXiv:0712.3014

Better for gluons that for quarks

(N.B. analytical results in the small-R approximation)



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The performance of jet algorithms -  $R_{\text{best}}$ 

Compare  $R_{\text{best}}$  after subtraction of pile-up For low lumi (0.05 nb<sup>-1</sup>) very small effects, at high lumi (0.25 nb<sup>-1</sup>) one finds moderate effects  $R_{\text{best}}^{\text{PU}} \sim R_{\text{best}}^{\text{no PU}}$ .

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## The performance of jet algorithms - Z'

Pile-up (at high LHC lumi) needs to be subtracted for a meaningful comparison (Use FastJet area-based tools):



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