

# Dissecting Soft Radiation with Factorization.

Frank Tackmann

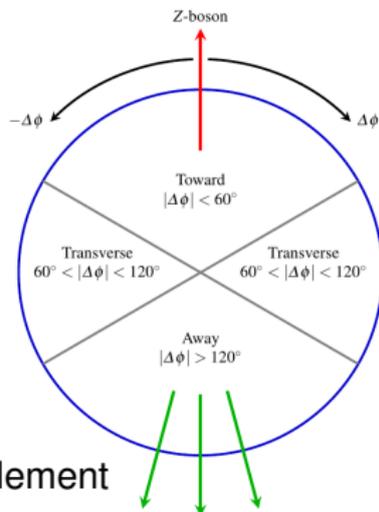
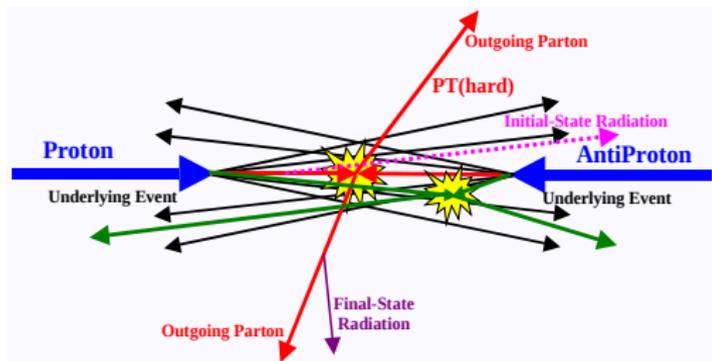
Deutsches Elektronen-Synchrotron

LHC Physics Discussion  
November 10, 2014

work with Iain Stewart and Wouter Waalewijn  
(arXiv:1405.6722)



# What is the “Underlying Event”?



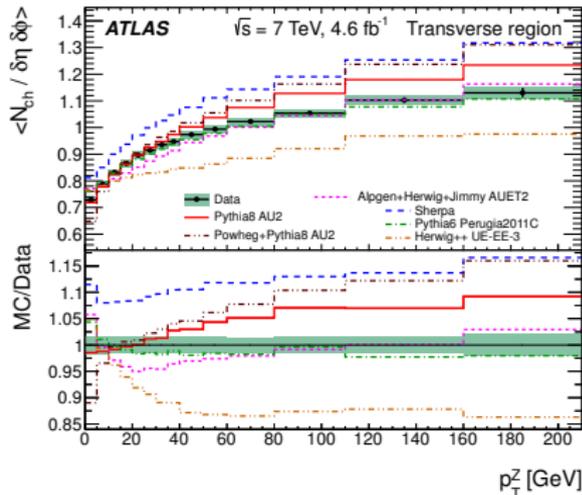
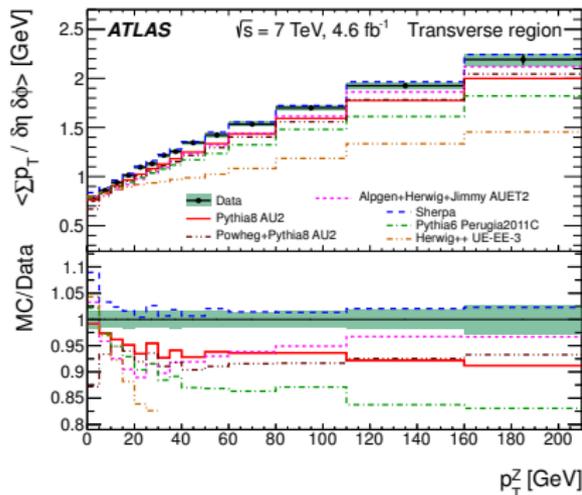
## Answers you might get, depending on who you ask

- Everything which is not the hard partonic matrix element
- All soft radiation underlying the hard scattering
- All particles not coming from primary hard process
- Monte Carlo: (soft) multiparton interactions (MPI)

## Current approach

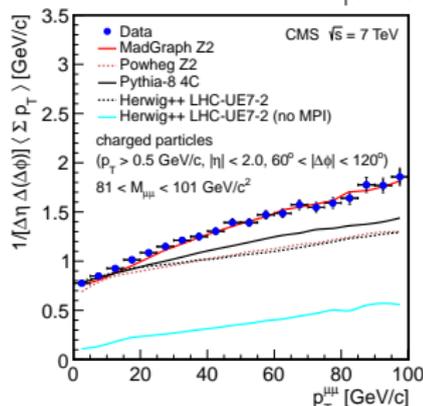
- Measure charged particle and energy densities in transverse regions
- Use these to tune Monte Carlo models (ISR, FSR, MPI, ....)

# Transverse Region in Drell-Yan.



[ATLAS 1409.3433, CMS 1204.1411]

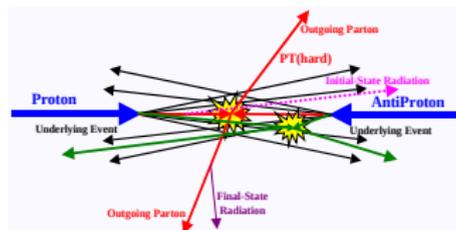
- MPI models in MC seem to produce too soft/too many particles
- Part of “UE activity” is due to primary and not-so-soft radiation (well known, same is seen in jet production)



# What is the “Underlying Event”?

Better question to ask:

What are the contributions from various soft effects to a given observable?



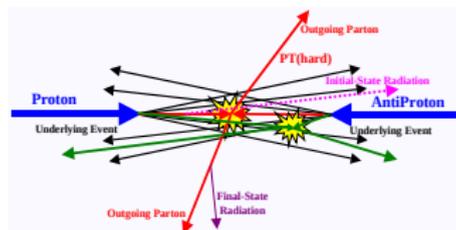
We want to distinguish contributions from

- Soft ISR: Perturbative soft radiation from primary partons
- Hadronization: Nonperturbative soft radiation
- MPI: soft radiation not associated with primary interaction

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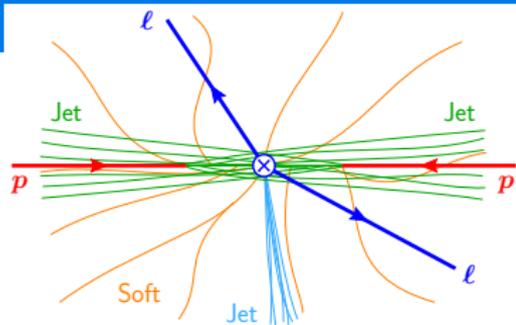
Consider an observable which

- is directly sensitive to soft effects
- corresponds to a well-defined hard process
- we understand well enough with factorization to make predictions

⇒ Consider jet mass spectrum in exclusive  $Z$ +jet and  $H$ +jet events

# Jet Mass Spectrum in $V+Jet$ .

- Require jet with  $p_T^J \geq 200 \text{ GeV}$
- Veto additional jets with  $p_T^J > 50 \text{ GeV}$   
(jet-veto effect essentially drops out for normalized spectrum)



Jet mass spectrum can be factorized for  $m_J \ll p_T^J$  (without MPI)

[Jouttenus, Stewart, FT, Waalewijn]

$$\frac{d\sigma}{dm_J^2} = \sum_{\kappa} H_{\kappa} (I \otimes f)^2 \int dk_s J_{q,g}(m_J^2 - 2p_T^J k_s) S_{\kappa}(k_s)$$

where each function has a field-theoretic definition (in SCET)

Hard function  $H_{\kappa}$ : partonic process  $\kappa: qq \rightarrow Zq, q\bar{q} \rightarrow Zg, gg \rightarrow Hg$

Beam function  $I \otimes f$ : collinear energetic ISR plus PDFs

Jet function  $J_{q,g}$ : collinear energetic FSR

Soft function  $S_{\kappa}(k_s)$ : primary soft radiation (ISR, FSR, pert. and nonpert.)

# Soft Function Factorization.

Convolution of  $J$  and  $S$  determines shape of jet mass spectrum

$$\frac{d\sigma}{dm_J^2} \sim \int dk_s J_{q,g}(m_J^2 - 2p_T^J k_s) S_\kappa(k_s)$$

Soft function can be factorized further

$$S_\kappa(k_s) = \int dk \underbrace{S_\kappa^{\text{pert}}(k_s - k)}_{\text{pert. soft ISR \& FSR}} \underbrace{F_\kappa(k)}_{\text{nonpert. hadronization}}$$

For  $k_s \gg k \sim \Lambda_{\text{QCD}}$  can expand in moments

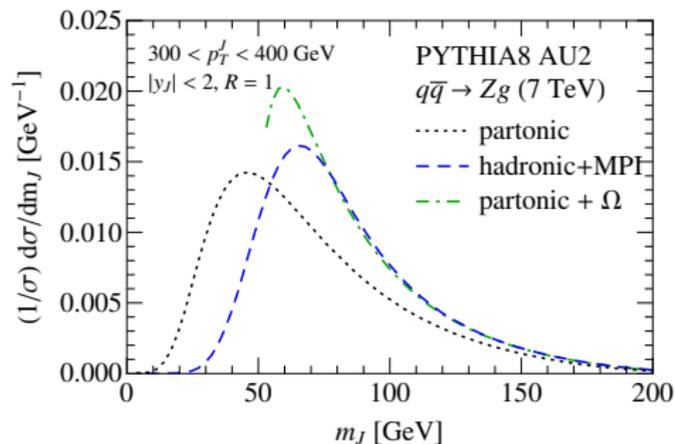
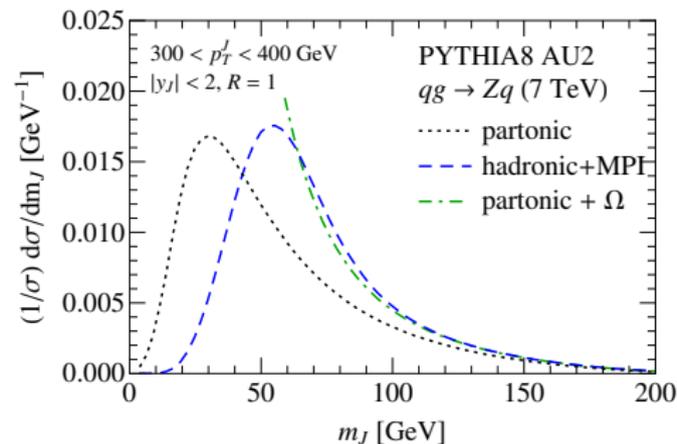
$$F_\kappa(k) = \delta(k) + \Omega_\kappa \delta'(k) + \dots$$

$$\Rightarrow S_\kappa(k_s) = S_\kappa^{\text{pert}}(k_s - \Omega_\kappa) + \dots$$

For  $m_J^2 \gg p_T^J \Lambda_{\text{QCD}}$  (tail of the spectrum) expect nonpert. corrections to shift the jet mass

$$m_J^2 = (m_J^2)^{\text{pert}} + 2p_T^J \Omega_\kappa$$

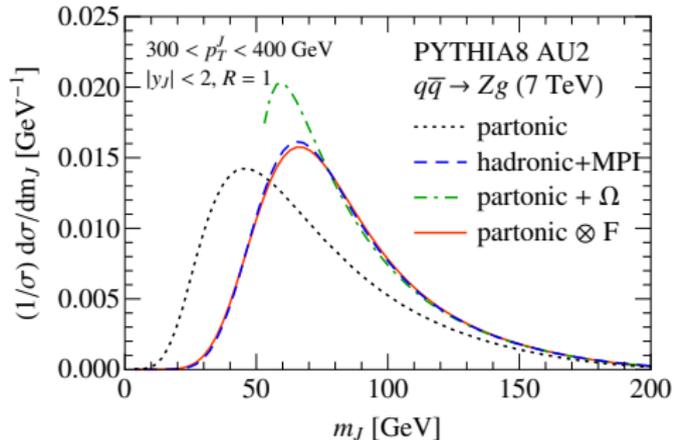
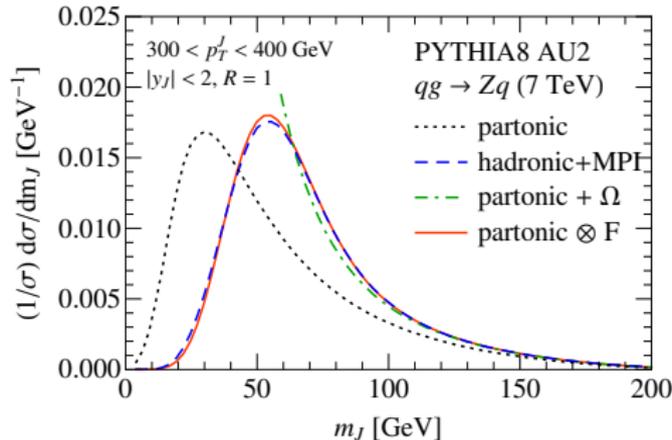
# Jet Mass Spectrum in Pythia8.



Remarkably well satisfied by Pythia's hadronization *and also* MPI

- Tail reproduced by  $m_J^2 \rightarrow m_J^2 + 2p_T^J \Omega$  (with  $\Omega \simeq 2.5$  GeV)

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- Tail reproduced by  $m_J^2 \rightarrow m_J^2 + 2p_T^J \Omega$  (with  $\Omega \simeq 2.5$  GeV)

- Even more general  $\frac{d\sigma}{dm_J^2} \rightarrow \int dk \frac{d\sigma^{\text{partonic}}}{dm_J^2} (m_J^2 - 2p_T^J k) F(k)$

⇒ Apparent degeneracy between hadronization and MPI  
(not too surprising: MPI populates jet with constant background of soft particles)

# Jet Mass Moment.

Consider the 1st moment of the normalized jet mass spectrum, which tracks the shift in the spectrum

$$M_1 = \frac{1}{\sigma} \int dm_J^2 m_J^2 \frac{d\sigma}{dm_J^2}$$

- Factorization predicts

$$M_1 = M_{1\kappa}^{\text{pert}}(p_T^J, y_J, R) + 2p_T^J \Omega_\kappa(R) \quad (+ \text{MPI})$$

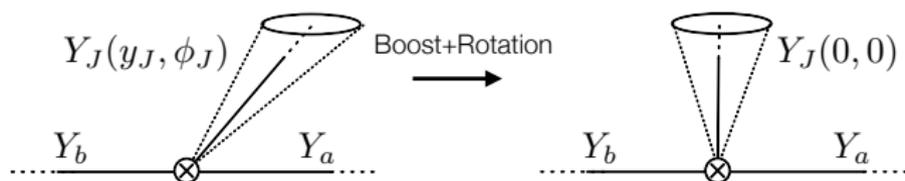
- Monte Carlo can be parametrized as

$$M_1 = M_{1\kappa}^{\text{partonic}}(p_T^J, y_J, R) + 2p_T^J \Omega_\kappa^{\text{had}}(R) + 2p_T^J \Upsilon_\kappa^{\text{MPI}}(R)$$

We will compare to

- ▶ Pythia8 with AU2-MSTW2008LO and 4C tunes
- ▶ Herwig++ 2.7 with default UE-EE-5-MRST tune

# Hadronization from Factorization.



Nonperturbative matrix element of Wilson lines  $Y_i$

$$\Omega_\kappa(R) = \langle 0 | Y_J^\dagger(y_J, \phi_J) Y_b^\dagger Y_a^\dagger M(y_J, \phi_J, R) Y_a Y_b Y_J(y_J, \phi_J) | 0 \rangle$$

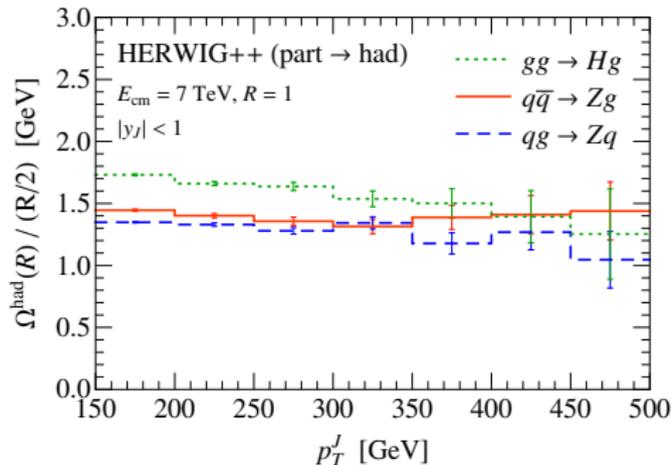
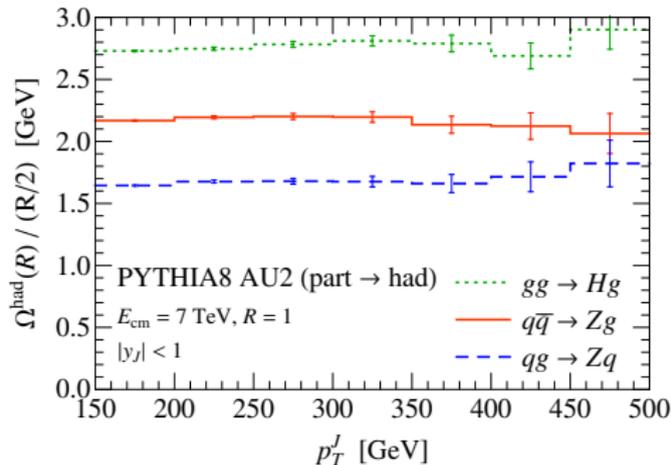
- Depends on partonic channel  $\kappa$  via color representation of  $Y_i$
- Independent of  $p_T^J$  (by definition) and  $y_J, \phi_J$  (shown by boosting)

Can be expanded in powers of  $R$

$$\Omega_\kappa(R) = \frac{R}{2} \Omega_{q,g}^{(1)} + \frac{R^3}{8} \Omega_\kappa^{(3)} + \frac{R^5}{32} \Omega_\kappa^{(5)} + \dots$$

- Only odd powers of  $R/2$  (to all orders)
- Leading  $\Omega_{q,g}^{(1)}$  is universal for quark and gluon initiated jets
  - ▶  $\Omega_q^{(1)}$  is the same that appears in 1-jettiness (thrust) in DIS

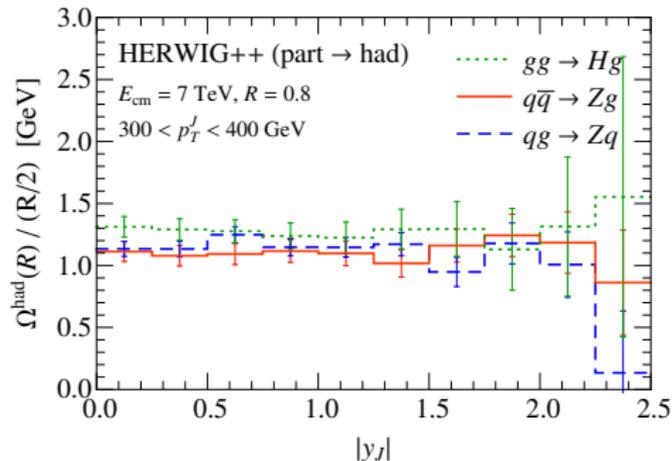
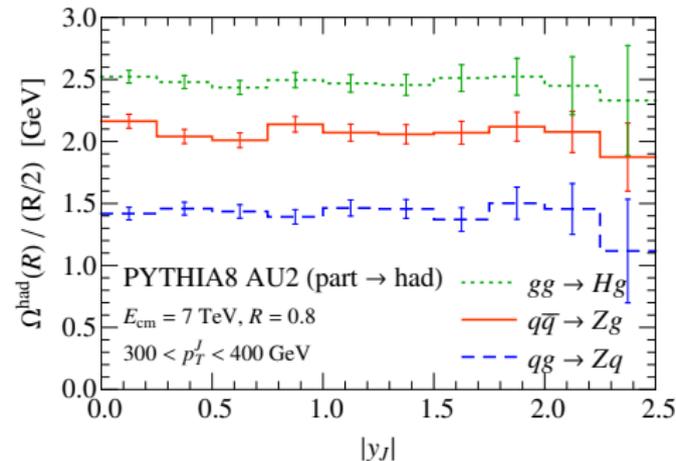
# Hadronization in Monte Carlo.



Agrees with factorization predictions (very well in Pythia, not as well in Herwig)

- independence of  $p_T^J$

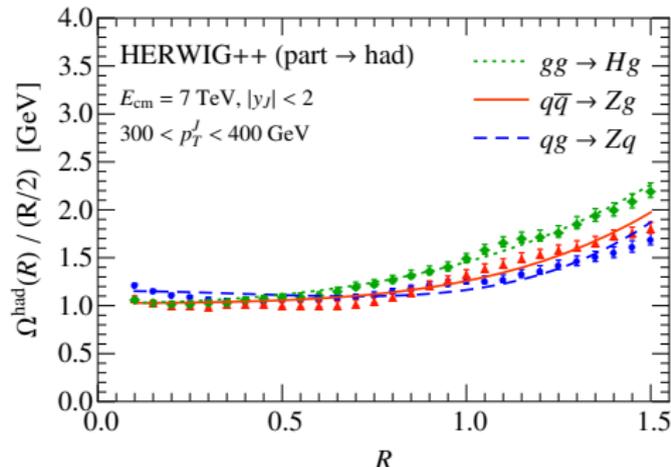
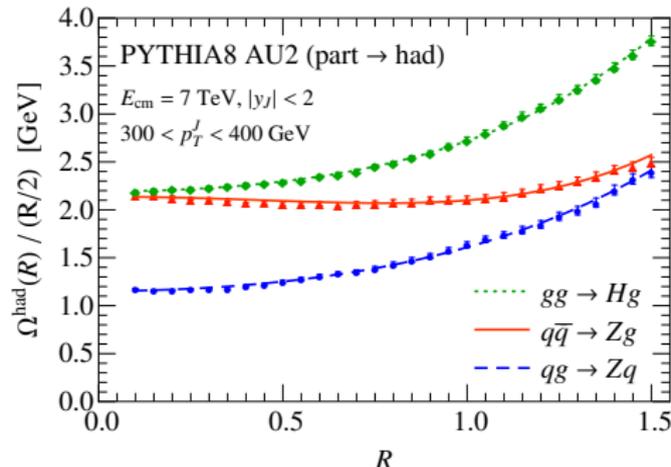
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- independence of  $p_T^J$
- independence of  $y_J$
- $R$  dependence starts linear (better fit with odd powers at large  $R$ )
- quark/gluon universality in limit of small  $R$

$\Rightarrow$  Channel dependence very different between Pythia and Herwig

# Primary Soft ISR.

Kinematics and dimensional analysis imply  $M_{1\kappa}^{\text{pert}} \sim (p_T^J R)^2$

$$\frac{M_{1\kappa}^{\text{pert}}}{2p_T^J} \sim p_T^J R^{2-\gamma_\kappa} + p_T^J R^4$$

- $\gamma_\kappa \propto \alpha_s > 0$  due to resummation effects

$R^4$  contribution arises from soft ISR

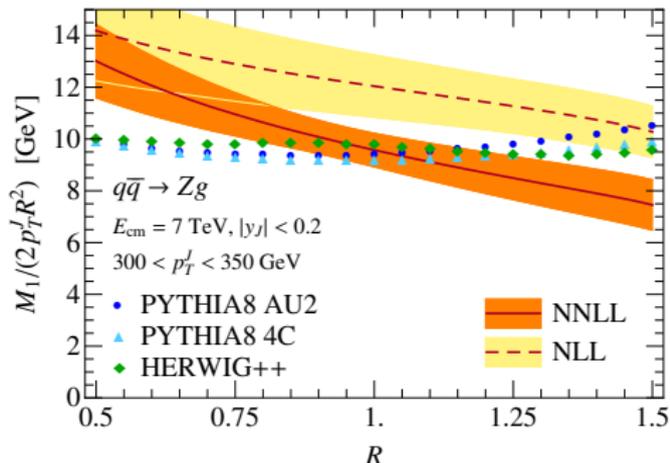
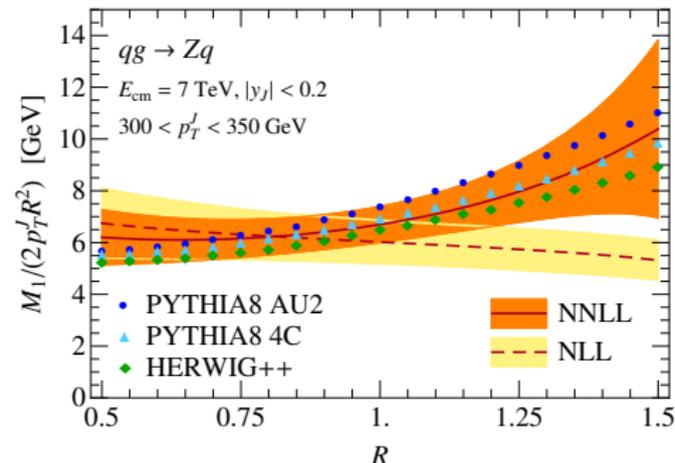
$$S_\kappa^{\text{pert}}(k_s) \supset \frac{\alpha_s C_\kappa}{\pi} R^2 \frac{1}{k_s}$$

- Interference between two beams, formally enters at  $\mathcal{O}(\alpha_s)$  and NNLL

- Color factor depends on channel

$\kappa$	$C_\kappa$
$gq \rightarrow Zq$	$C_A/2 = +3/2$
$q\bar{q} \rightarrow Zg$	$C_F - C_A/2 = -1/6$
$gg \rightarrow Hg$	$C_A/2 = +3/2$

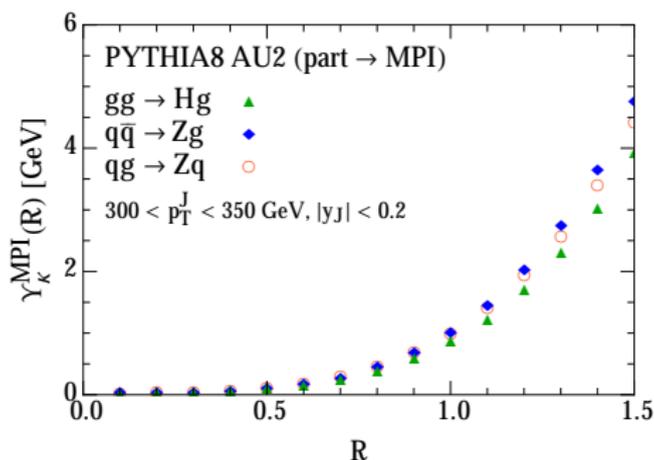
# Partonic Contribution in Monte Carlo.



Partonic contribution in MC very well described by

$$\frac{M_{1\kappa}^{\text{partonic}}}{2p_T^J} = c_2^\kappa(p_T^J, y_J) R^{2-\gamma_\kappa} + c_4^\kappa(p_T^J, y_J) R^4$$

- Soft ISR modelled by radiation pattern from both initial and final showers
  - ▶ Good agreement for  $qg \rightarrow Zq$  (not as good for  $gg \rightarrow Hg$ )
  - ▶ MCs do not reproduce negative interference for  $q\bar{q} \rightarrow Zg$



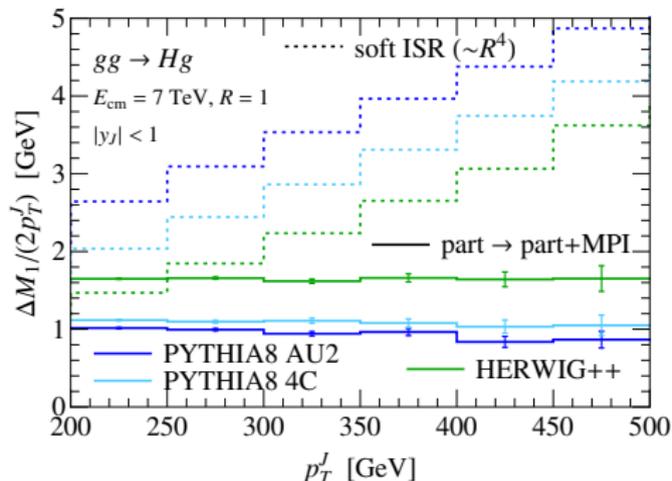
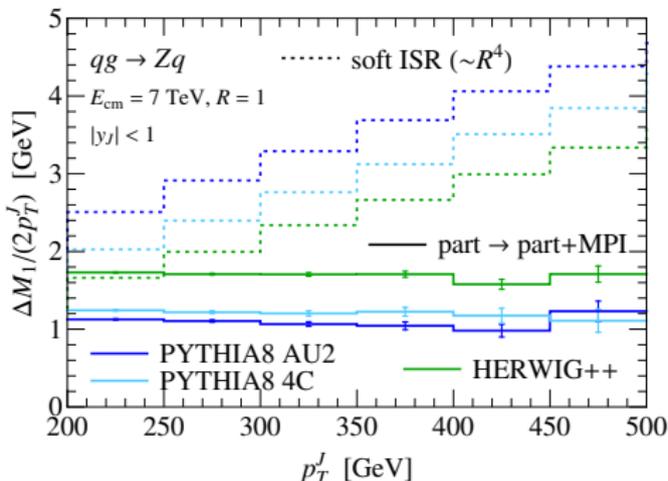
## MPI almost channel independent and with perfect $R^4$ behavior

- Like soft ISR, it populates jet with a constant background of soft particles
    - ▶ Expect same  $R^4$  scaling
  - Essentially independent of primary hard process
    - ▶ Expect channel independence, and also  $p_T^J, y_J$  independence
- $\Rightarrow$  Can in principle distinguish different “UE-like”  $R^4$  contributions

# Soft ISR vs. MPI.

## Compare

- Soft ISR: fitted  $R^4$  contribution in partonic results (dotted)
- Effect from turning on MPI (solid)



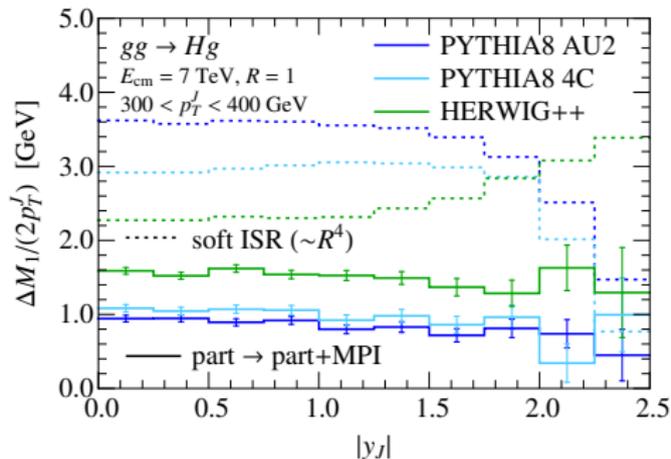
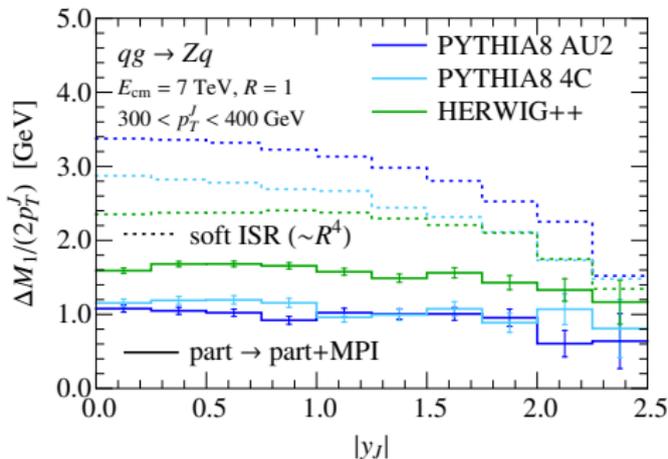
## Ambiguity between MPI and soft ISR in Monte Carlos

- Different MCs/tunes compensate more/less soft ISR with less/more MPI

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## Ambiguity between MPI and soft ISR in Monte Carlos

- Different MCs/tunes compensate more/less soft ISR with less/more MPI
- Different  $y_J$  dependence of soft ISR in Pythia and Herwig

# Summary.

$$\frac{M_1}{2p_T^J} \sim \underbrace{c_2^\kappa(p_T^J, y_J) R^{2-\gamma_\kappa}}_{\text{leading pert.}} + \underbrace{c_4^\kappa(p_T^J, y_J) R^4}_{\text{soft ISR}} + \underbrace{\frac{R}{2} \Omega_{q,g}^{(1)} + \frac{R^3}{8} \dots}_{\text{hadronization}} + \underbrace{R^4}_{\text{MPI}}$$

Jet-mass spectrum and in particular its 1st moment as a function of jet radius  $R$  allows to separate physically distinct sources of soft effects

- Hadronization (nonpert. soft)
- Primary soft ISR interference
- Soft MPI

Allows a 3-way comparison between

- Properties that can be predicted from QCD by factorization
- MC models
- Measurements

⇒ Try to measure  $M_1(R)$  for different  $p_T^J, y_J$ , and partonic channels