

η and η' Mixing and the Witten-Veneziano Formula

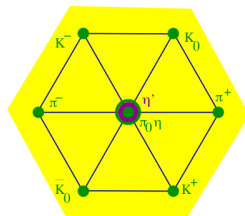
C. Urbach

HISKP, University of Bonn
Project C7 SFB/TR16

SFB/TR 9 Meeting, Durbach, September 2014

The $U_A(1)$ Problem

- nine lightest pseudo-scalar mesons show a peculiar spectrum:
 - 3 very light pions (140 MeV)
 - kaons and the η around 500 MeV
 - η' around 1 GeV



- this was called the $U_A(1)$ -problem
- $M_\pi \ll M_\eta \ll M_{\eta'}$
- in contrast: $\rho - \omega - \phi$ appear ideally mixed as expected from OZI rule
- spontaneously broken chiral symmetry: nine Goldstone bosons?

[Glashow (1968)]

- the $U(1)$ axial current is anomalous at quantum level

[Adler (1969), Jackiw and Bell (1969)]

$$\partial^\mu A_\mu^{0q} = \partial^\mu (\bar{q} \gamma_\mu \gamma_5 q) = 2m_q (\bar{q} i \gamma_5 q) + \frac{\alpha_s}{4\pi} F \tilde{F}$$

and not spontaneously broken

- anomaly vanishes to all orders in perturbation theory at zero momentum
- ⇒ need additional mechanism

- instantons with non-trivial topology allow to explain η' mass

[Kogut, Susskind (1975), Belavin et al. (1975), 't Hooft (1976)]

- ⇒ suggests consistency of QCD with nature

- π, η, η' masses can be computed using lattice QCD
- ... but how to establish the relation to the $U_A(1)$ anomaly?
- in the 't Hooft limit the Witten-Veneziano formula can be derived ($m_q = 0, N_c \rightarrow \infty, g^2 N_c$ and N_f fixed)

$$M_{\eta'}^2 = \frac{4N_f}{f_\pi^2} \chi_\infty$$

[Witten (1979), Veneziano (1979)]

- where χ_∞ is the topological susceptibility in pure Yang-Mills theory
- including $m_s > 0$ effects

$$M_\eta^2 + M_{\eta'}^2 - 2M_K^2 = \frac{4N_f}{f_\pi^2} \chi_\infty$$

- for exact SU(3) flavour symmetry ($m_u = m_d = m_s$) define (considering only $\bar{q}q$ states ...)
- the flavour octet state η_8

$$\eta_8 \equiv \frac{1}{\sqrt{6}}(\bar{u}i\gamma_5 u + \bar{d}i\gamma_5 d - 2\bar{s}i\gamma_5 s)$$

⇒ a (pseudo) Goldstone boson

- the flavour singlet state η_0

$$\eta_0 \equiv \frac{1}{\sqrt{3}}(\bar{u}i\gamma_5 u + \bar{d}i\gamma_5 d + \bar{s}i\gamma_5 s)$$

⇒ related to $U_A(1)$ anomaly

- SU(3) flavour symmetry broken by larger m_s : $m_s \gg m_u = m_d = m_\ell$
 \Rightarrow physical states will be a mixture, e.g.

$$\begin{pmatrix} |\eta\rangle \\ |\eta'\rangle \end{pmatrix} = \begin{pmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{pmatrix} \cdot \begin{pmatrix} |\eta_\ell\rangle \\ |\eta_s\rangle \end{pmatrix}$$

in the quark flavour basis

$$\eta_\ell \equiv \frac{1}{\sqrt{2}}(\bar{u}i\gamma_5 u + \bar{d}i\gamma_5 d), \quad \eta_s = \bar{s}i\gamma_5 s$$

- in nature $m_u \neq m_d \quad \Rightarrow \quad$ also π_0 mixes
- how did nature arrange the mixing pattern?
- can one determine the mixing angle(s)?

- mixing angle definition via Fock states theoretically difficult

⇒ use matrix elements instead

$$\begin{pmatrix} \mathbf{P}_{\eta}^{\ell} & \mathbf{P}_{\eta}^s \\ \mathbf{P}_{\eta'}^{\ell} & \mathbf{P}_{\eta'}^s \end{pmatrix} = \begin{pmatrix} c_{\ell} \cos \phi_{\ell} & -c_s \sin \phi_s \\ c_{\ell} \sin \phi_{\ell} & c_s \cos \phi_s \end{pmatrix}$$

defined via pseudo-scalar matrix elements

$$\mathbf{P}_P^q = \langle 0 | \bar{q} i \gamma_5 q | P \rangle, \quad P = \eta, \eta', \quad q = \ell, s$$

- note: usually defined via decay constants obtained from

$$\langle 0 | A_{\mu}^q | P(p) \rangle = i f_P^q p_{\mu}$$

with

$$A_{\mu}^{\ell} = \frac{1}{\sqrt{2}} (\bar{u} \gamma_{\mu} \gamma_5 u + \bar{d} \gamma_{\mu} \gamma_5 d), \quad A_{\mu}^s = \bar{s} \gamma_{\mu} \gamma_5 s$$

- angles are then obtained from

$$\frac{\mathbf{P}_{\eta'}^{\ell}}{\mathbf{P}_{\eta}^{\ell}} = \tan \phi_{\ell}, \quad \frac{\mathbf{P}_{\eta}^s}{\mathbf{P}_{\eta'}^s} = -\tan \phi_s$$

- from χ PT and $1/N_c$ arguments one expects

$$|\phi_{\ell} - \phi_s| \ll 1$$

- if correct, define a single angle ϕ obtained from

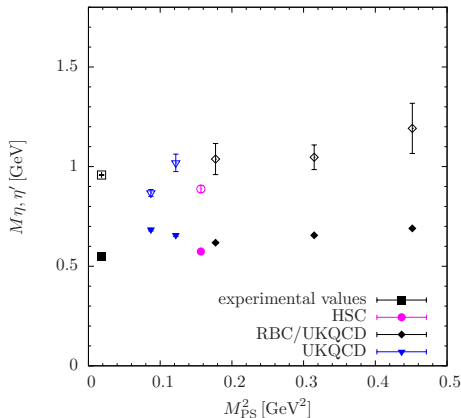
$$\tan^2 \phi = -\frac{\mathbf{P}_{\eta'}^{\ell} \mathbf{P}_{\eta}^s}{\mathbf{P}_{\eta}^{\ell} \mathbf{P}_{\eta'}^s}$$

- note: all renormalisation constants cancel in the ratios

fundamental questions:

- does QCD reproduce the experimental mass pattern?
- how does nature arrange for the mixing?
- can we relate the η' mass to the $U_A(1)$ anomaly

- η, η' involve OZI-rule violating diagrams
- ⇒ so-called fermionic disconnected diagrams
- due to noise computationally challenging
- ⇒ only few lattice results available
- lattice definition of χ_t
 - must be cleanly defined
 - must be properly renormalised



filled symbols: η open: η'

[HSC, J. J. Dudek et al., Phys. Rev. D83 (2011)]
 [RBC/UKQCD, N. Christ et al., Phys. Rev. Lett. 105 (2010)]
 [UKQCD, E. B. Gregory et al., Phys.Rev. D86 (2012)]

- let $O(x, t)$ be an operator with quantum numbers of a given state
- for instance for the pion it is given by

$$O(x, t) = \bar{u}i\gamma_5 d(x, t), \quad O(t) = \sum_x O(x, t)$$

projected to zero momentum

- create pion at time 0 and annihilate at t

$$C_\pi(t) = \langle O(0) O^\dagger(t) \rangle \propto \sum_n \langle 0 | O(0) | n \rangle \langle n | e^{-Ht} O(0) e^{Ht} | 0 \rangle$$

inserting a complete set of states $|n\rangle$ and using the time evolution operator $\exp -Ht$

- which leads to

$$C_\pi(t) \propto \sum_n |\langle 0 | O | n \rangle|^2 e^{-(E_n - E_0)t}$$

- and for large times the lowest state dominates

$$\lim_{t \rightarrow \infty} C_\pi(t) \propto e^{-(E_1 - E_0)t}$$

- $E_1 - E_0$ corresponds at zero momentum to the mass M_π
- and matrix elements of the pion for $n = 1$

$$\langle 0 | O | \pi \rangle$$

- also define effective masses

$$M = -\frac{d}{dt} \log C_\pi(t)$$

- $2 + 1 + 1$ quark flavour ensembles from ETM Collaboration
 $m_u = m_d < m_s < m_c$

[ETMC, R. Baron et. al., JHEP 06 111 (2010)]

- three lattice spacings (A , B and D ensembles):
 $a_A = 0.086$ fm, $a_B = 0.078$ fm and $a_D = 0.061$ fm
- charged pion masses range from ≈ 230 MeV to ≈ 500 MeV
- $L \geq 3$ fm and $M_\pi \cdot L \geq 3.5$ for most ensembles
- bare m_s and m_c fixed for each lattice spacing
- use $r_0 = 0.45(2)$ fm (from f_π) for η, η'
 r_0 Sommer scale determined from static $\bar{q}q$ potential
 $\Rightarrow r_0 M$ is dimensionless (with M a mass)

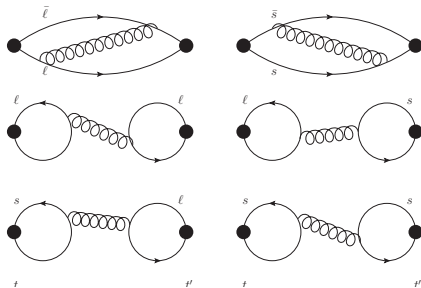
- need to estimate correlator matrix

$$\mathcal{C} = \begin{pmatrix} \eta_{\ell\ell} & \eta_{\ell s} \\ \eta_{s\ell} & \eta_{ss} \end{pmatrix}$$

- η_{XY} correlator of appropriate operators, e.g.

$$\eta_{ss}(t) \equiv \langle \bar{s} i\gamma_5 s(t) \bar{s} i\gamma_5 s(0) \rangle$$

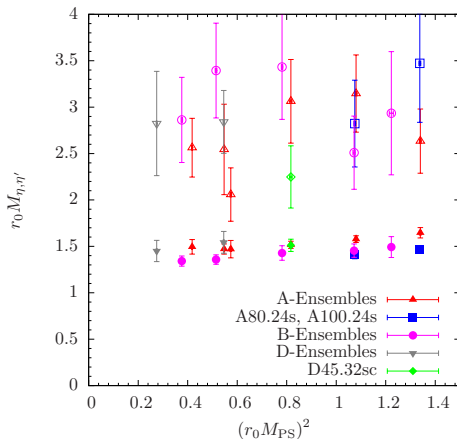
- connected and disconnected contributions



- diagonalise matrix \Rightarrow masses M_P and matrix elements \mathbf{P}_P^q
- $P = \eta$: lowest state, $P = \eta'$: second state, $P = \eta_c \dots$

- η' mainly the flavour singlet
- disconnected contributions significant
- ⇒ very noisy
- chiral extrapolation uncertain
- ⇒ no clear picture
- need for improvement

filled symbols: η open: η'

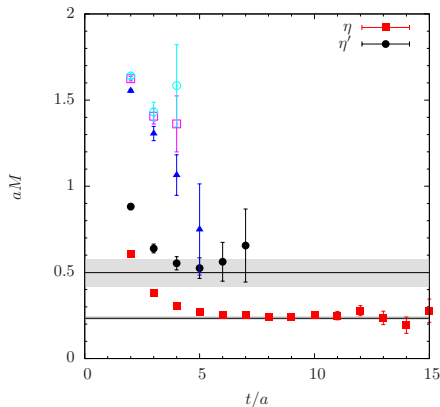


[C. Michael, K. Otnad, S. Reker, C.U., JHEP 1211 (2012) 048]

- effective masses for $\eta, \eta', \eta_c, \dots$
 - η mass well determined
 - η' :
 - at small t/a excited states dominant
 - disconnected contributions noisy
- ⇒ signal lost in noise before plateau reached

- large systematics in η'

⇒ need to reduce noise/increase statistics tremendously
or remove excited states to fit at smaller t/a

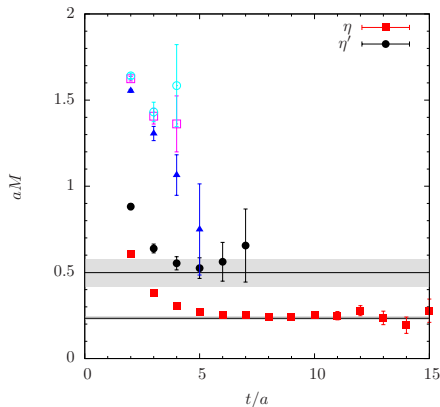


- effective masses for $\eta, \eta', \eta_c, \dots$
- η mass well determined
- η' :
 - at small t/a excited states dominant
 - disconnected contributions noisy \Rightarrow signal lost in noise before plateau reached

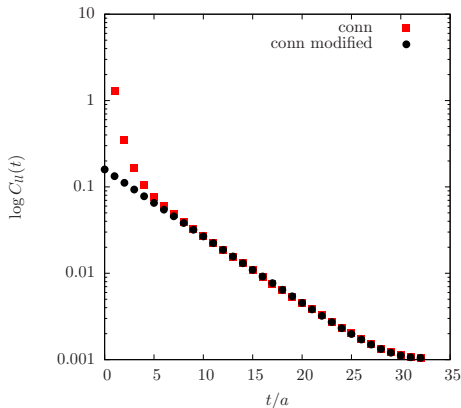
- large systematics in η'

\Rightarrow need to reduce noise/increase statistics tremendously

or remove excited states to fit at smaller t/a

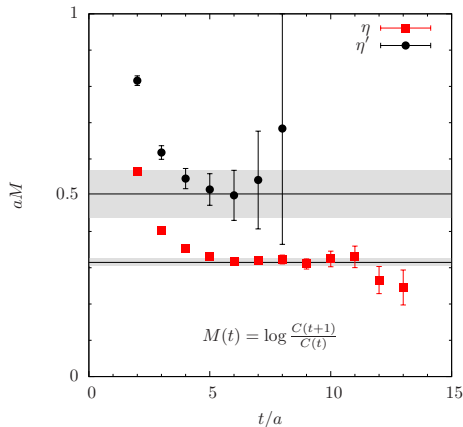


- recall: disconnected contributions noisy
- **lets make an assumption:**
disconnected contributions couple only to η and η' states, not to higher states
[Neff et al. (2001), K.Jansen, C.Michael, C.U. (2008)]
- replace connected contributions by only the ground states
- if assumption justified:
there should be a plateau in the effective masses from very low times on!

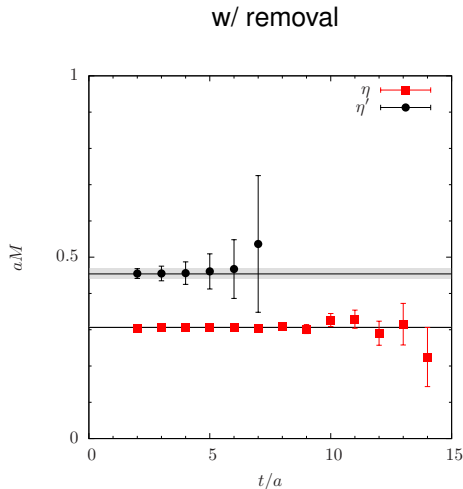


- we see a plateau from $t/a = 2$ on
- for both η and η'
- η : good agreement with previous results
- η' : possibly much better determination
- assumption justified?
- systematic uncertainties?

w/o removal



- we see a plateau from $t/a = 2$ on
- for both η and η'
- η : good agreement with previous results
- η' : possibly much better determination
- assumption justified?
- systematic uncertainties?



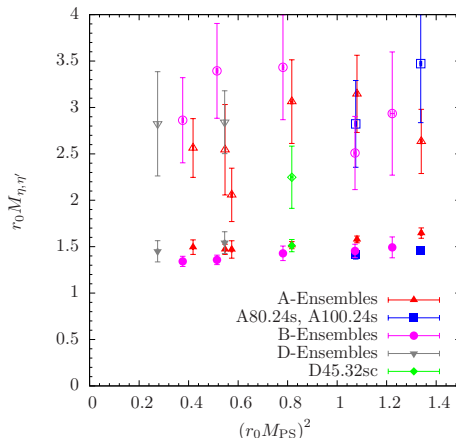
η :

- masses agree well
- improved precision

η' :

- masses determined much better
- always agreement within 2σ
- systematics hard to quantify

w/o removal



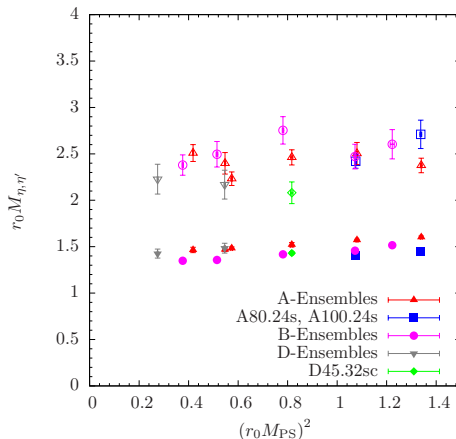
η :

- masses agree well
- improved precision

η' :

- masses determined much better
- always agreement within 2σ
- systematics hard to quantify

w/ removal



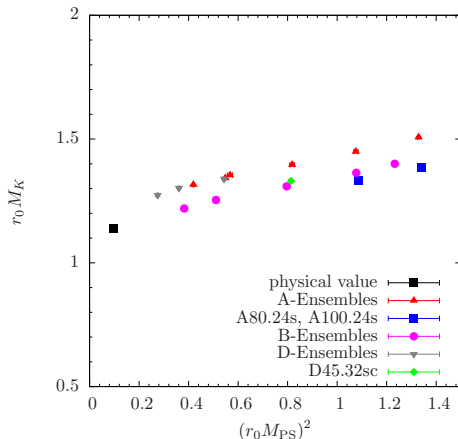
- can take the difference w/ and w/o to estimate systematics

- m_s not perfectly tuned to its physical value
 - two re-tuned ensembles for a_A
- ⇒ can estimate m_s dependence

- estimate

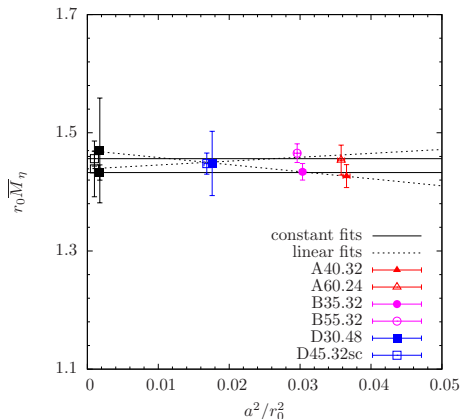
$$D_\eta \equiv \frac{d(aM_\eta)^2}{d(aM_K)^2} = 1.47(11)$$

- now assume:
 D_η independent of
 a, m_ℓ, m_s, m_c
- ...correct η masses



Continuum Limit: Scaling Test for M_η

- use two ensembles sets
(A60, B55, D45)
(A40, B35, D30)
with $r_0 M_{\text{PS}} \approx \text{const}$
 - correct M_η using D_η linearly in M_K^2
 $\Rightarrow r_0 M_K = 1.34$ fixed
 - compatible with both,
constant and linear continuum
extrapolation
- \Rightarrow lattice artifacts seem under control



[PRL 111 181602 (2013)]

- now correct all $(r_0 M_\eta)^2$ using D_η

$\Rightarrow (r_0 \bar{M}_\eta)^2 \propto \text{const} + (r_0 M_{\text{PS}})^2$

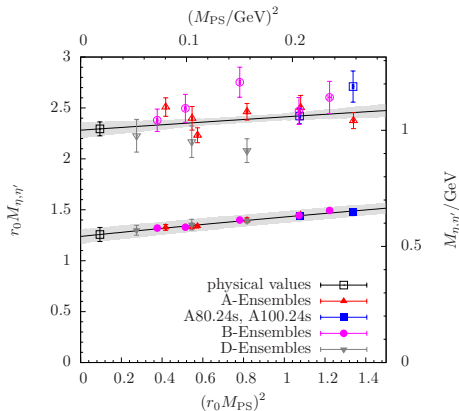
- all a -values fall on the same curve!

- extrapolate $(r_0 \bar{M}_\eta)^2$ linearly in $(r_0 M_{\text{PS}})^2$ to $M_{\text{PS}} = M_\pi$

\Rightarrow result

$$M_\eta = 552(10)_{\text{stat}} \text{ MeV}$$

- similarly with $(\bar{M}_\eta / \bar{M}_K)^2$ or GMO relation



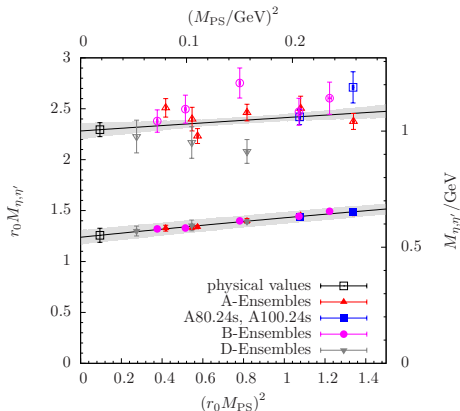
[PRL 111 181602 (2013)]

- no clear dependence on
 - lattice spacing
 - strange quark mass
- errors still significant
- include all data in extrapolation
- $(r_0 M_{\eta'})^2 \propto \text{const} + (r_0 M_{\text{PS}})^2$

⇒ result

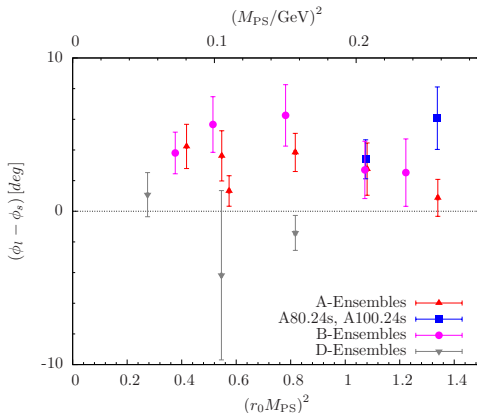
$$M_{\eta'} = 1006(54)_{\text{stat}} \text{ MeV}$$

- fitting A , B and D separately gives compatible results



Mixing Angles ϕ_ℓ and ϕ_s

- $\Delta\phi = 3(1)_{\text{stat}}(3)_{\text{sys}}^\circ$
confirms expectation,
smallness of OZI
corrections
- data well described by a
single angle
- but cannot exclude a finite
difference



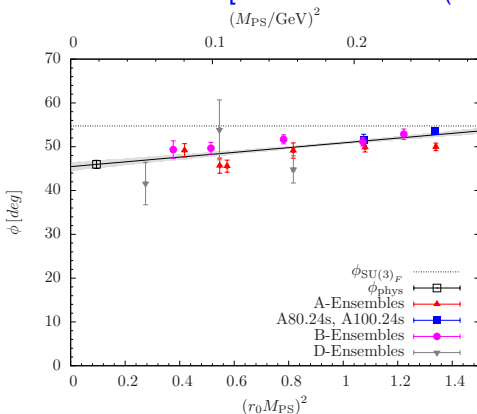
- systematic uncertainty from fits to data from different lattice spacings
separately

[PRL 111 181602 (2013)]

- extrapolate ϕ linearly in $(r_0 M_{\text{PS}})^2$

$$\Rightarrow \phi = 46(1)_{\text{stat}}(3)_{\text{sys}}^\circ$$

- systematic error from fitting different lattice spacings separately
- with $m_\ell = m_s$ we reproduce the SU(3) symmetric value 54.7°



- interpretation: η' meson mainly the flavour singlet state

\Rightarrow far from ideally mixed

K. Cichy, E. Garcia-Ramos, K. Jansen

- spectral projectors R_M allow for a clean definition of χ_t

[Giusti, Lüscher, (2009)]

$$\chi_t = \frac{Z_S^2}{Z_P^2} \frac{\langle \text{Tr}[\gamma_5 R_M^2] \text{Tr}[\gamma_5 R_M^2] \rangle}{V}$$

- spectral projectors estimated stochastically
- dedicated quenched ensembles with Iwasaki gauge action

see Cichy, Garcia-Ramos, Jansen (2013) for $N_f = 2$ and $2 + 1 + 1$ results

- four values of the lattice spacing from 0.07 fm to 0.14 fm
- Wilson twisted mass valence quarks

⇒ automatic $\mathcal{O}(a)$ improvement

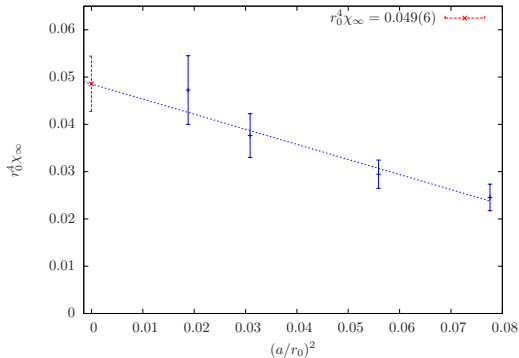
- box length fixed to

$$L \sim 2.8 \text{ fm}$$

- Z_S/Z_P computed with spectral projectors
- linear scaling in a^2 as expected

\Rightarrow continuum limit:

$$r_0^4 \chi_\infty = 0.049(6)$$



- using $r_0 = 0.5 \text{ fm} \quad \Rightarrow \quad \chi_\infty = (185(6) \text{ MeV})^4$

\Rightarrow WV formula well fulfilled with physical mass values

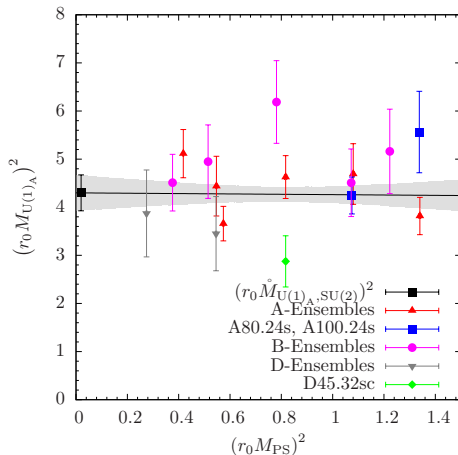
K. Cichy, E. Garcia-Ramos, K. Jansen, K. Ottnad, F. Zimmermann, C.U.

- define

$$M_{U_A(1)}^2 = M_\eta^2 + M_{\eta'}^2 - 2M_K^2$$

- extrapolates almost constant in M_{PS}^2
- systematics very similar to $M_{\eta'}$
- dominant statistical noise
- the extrapolated value agrees with separat mass extrapolations

⇒ to be investigated further



- study $U_A(1)$ problem in LQCD

- η, η' mixing angle

$$\phi = 46(1)_{\text{stat}}(3)_{\text{sys}}^\circ$$

- η, η' masses

$$M_\eta = 551(8)_{\text{stat}}(6)_{\text{sys}} \text{ MeV}$$

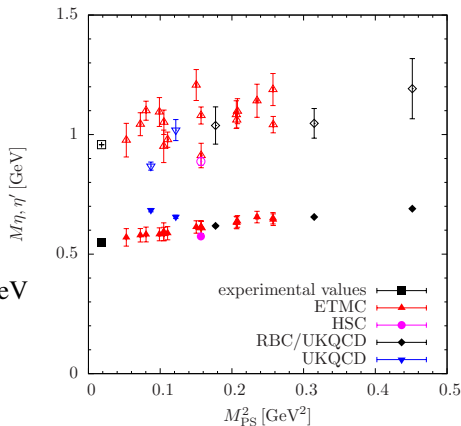
$$M_{\eta'} = 1006(54)_{\text{stat}}(38)_{\text{sys}}(+61)_{\text{ex}} \text{ MeV}$$

- Yang-Mills χ_∞

$$\chi_\infty = (185(6) \text{ MeV})^4$$

from spectral projectors

- test of Witten-Veneziano formula



[HSC, J. J. Dudek et al., Phys. Rev. D83 (2011)]

[RBC/UKQCD, N. Christ et al., Phys. Rev. Lett. 105 (2010)]

[UKQCD, E. B. Gregory et al., Phys.Rev. D86 (2012)]

[ETMC, C. Michael et al., PRL 111 (2013)]