

GPDs from exclusive meson leptonproduction

P. Kroll

Fachbereich Physik, Univ. Wuppertal

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Outline:

- Exclusive processes, GPDs, parametrizations
- What did we learn about GPDs?
- Transverse target asymmetries for DVCS
- π^0 cross section and/or η/π^0 cross section ratio
- Beam spin asymmetry for π production
- Summary

Hard exclusive scattering - GPDs

DVCS and meson electroproduction

rigorous proofs of collinear factorization in generalized Bjorken regime:

Radyushkin, Collins et al, Ji et al

$(Q^2, W \rightarrow \infty, x_{Bj} \text{ fixed})$

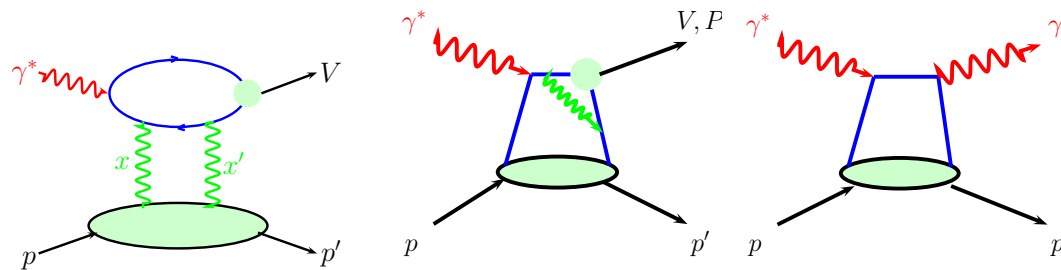
possible power corrections not under control \implies

unknown at which Q^2 asymptotic result can be applied

hard subprocesses

$$\gamma^* g \rightarrow V g,$$

$$\gamma^* q \rightarrow V(P, \gamma) q$$



and GPDs and meson w.f.

(encode the soft physics)

$$\langle K \rangle = \int_{-1}^1 d\bar{x} \mathcal{H}(\bar{x}, \xi, t) K(\bar{x}, \xi, t)$$

dominant transitions $\gamma_L^* \rightarrow V_L(P), \gamma_T^* \rightarrow \gamma_T$

other transitions power suppressed but often non-negligible (e.g. $\gamma_T^* \rightarrow V_T, \pi^+$)

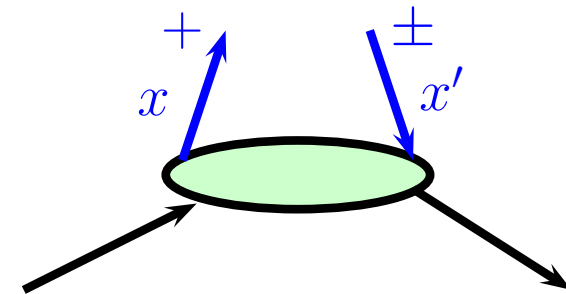
GPDs

D. Müller et al (94), Ji(97), Radyushkin (97)

GPDs: $K = K(\bar{x}, \xi, t)$ $x = \frac{\bar{x} + \xi}{1 + \xi}$ $x' = \frac{\bar{x} - \xi}{1 - \xi}$

defined by FT of $\langle p' | \bar{\Psi}(-z/2) \Gamma \Psi(z/2) | p \rangle$

$\Gamma = \gamma^+, \gamma^+ \gamma_5$ $K = H, E, \tilde{H}, \tilde{E}$ (non-flip)
 $= \sigma^{+j}$ $H_T, E_T, \tilde{H}_T, \tilde{E}_T$ (flip)



for quarks ($\xi < \bar{x} < 1$) and gluons

(antiquarks for $-1 < \bar{x} < -\xi$, $q\bar{q}$ pairs $-\xi < \bar{x} < \xi$)

reduction formula $H^q(\bar{x}, \xi = t = 0) = q(\bar{x})$, $\tilde{H}^q \rightarrow \Delta q(\bar{x})$, $H_T^q \rightarrow \delta^q(\bar{x})$

sum rules (proton form factors): $F_1^q(t) = \int d\bar{x} H^q(\bar{x}, \xi, t)$, $F_1 = \sum e_q F_1^q$

$E \rightarrow F_2$, $\tilde{H} \rightarrow F_A$, $\tilde{E} \rightarrow F_P$

polynomiality, universality, evolution, positivity constraints

Ji's sum rule $J_q = \frac{1}{2} \int_{-1}^1 d\bar{x} \bar{x} [H^q(\bar{x}, \xi, t = 0) + E^q(\bar{x}, \xi, t = 0)]$

Parameterizing the GPDs

needed in order to extract GPDs from data

double distribution ansatz (Mueller *et al* (94), Radyushkin (99))

$$K_i(\bar{x}, \xi, t) = \int_{-1}^1 d\rho \int_{-1+|\rho|}^{1-|\rho|} d\eta \delta(\beta + \xi\alpha - \bar{x}) K_i(\rho, \xi = 0, t) w_i(\rho, \eta) + D_i \Theta(\xi^2 - \bar{x}^2)$$

DD: zero-skewness GPD \times weight fct w_i (generates ξ dep.)

$$K_i(\rho, \xi = 0, t) = k_i(\rho) \exp [(b_i + \alpha'_i \ln(1/\rho))t]$$

$$k_i = q, \Delta q, \delta^q \text{ for } H, \tilde{H}, H_T \text{ or } c\bar{x}^{-\alpha_f(0)}(1 - \bar{x})^{\beta_f} \text{ for } E, \tilde{E}, \bar{E}_T$$

Regge-like t dep. (for small ξ and small $-t$ reasonable appr.)

advantage: polynomiality and reduction formulas automatically satisfied

D -term neglected

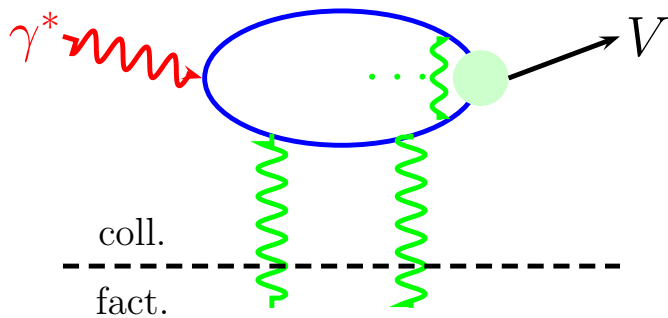
The subprocess amplitude for DVMP

mod. pert. approach - quark trans. momenta in subprocess

(emission and absorption of partons from proton collinear to proton momenta)

transverse separation of color sources \implies gluon radiation

(Sudakov factor Sterman et al(93))



LO pQCD

+ quark trans. mom.

+ Sudakov supp.

\implies asymp. fact. formula (lead. twist)

for $Q^2 \rightarrow \infty$

gluon part bears similarity to color dipole model

What has been done?

- analysis of FF ([Diehl et al, \(04\), update: Diehl-K 1302.4604](#))
using [ABM11](#) PDFs, fixes $H, E, (\tilde{H})$ for valence quarks
- analysis of $d\sigma_L/dt$ for ρ^0 and ϕ production [Goloskokov-K, hep-ph/0611290](#)
for $Q^2 \gtrsim 3 \text{ GeV}^2$ and $W \gtrsim 4 \text{ GeV}$ ($\xi \lesssim 0.1, -t \lesssim 0.5 \text{ GeV}^2$)
data from [H1, ZEUS, E665, HERMES](#)
fixes H for sea quarks and gluons for given H^{val} (E negligible)
update with [ABM11](#) required
- analysis of π^+ production, [Goloskokov-K, 0906.0460](#)
 $d\sigma/dt$ and A_{UT} data from [HERMES](#)
evidence for strong contr. from γ_T^* (H_T)
fixes \tilde{H} , pion pole and H_T (no clear signal for $\tilde{E}_{\text{non-pole}}$)
- SDME and A_{UT} for ρ^0 production,
 π^0 cross section and η/π^0 from CLAS (large skewness!),
lattice QCD [QCDSF and UKQCD, hep-lat/0612032](#)
hints at strong contributions from $\bar{E}_T = 2\tilde{H}_T + E_T$ **needs confirmation**

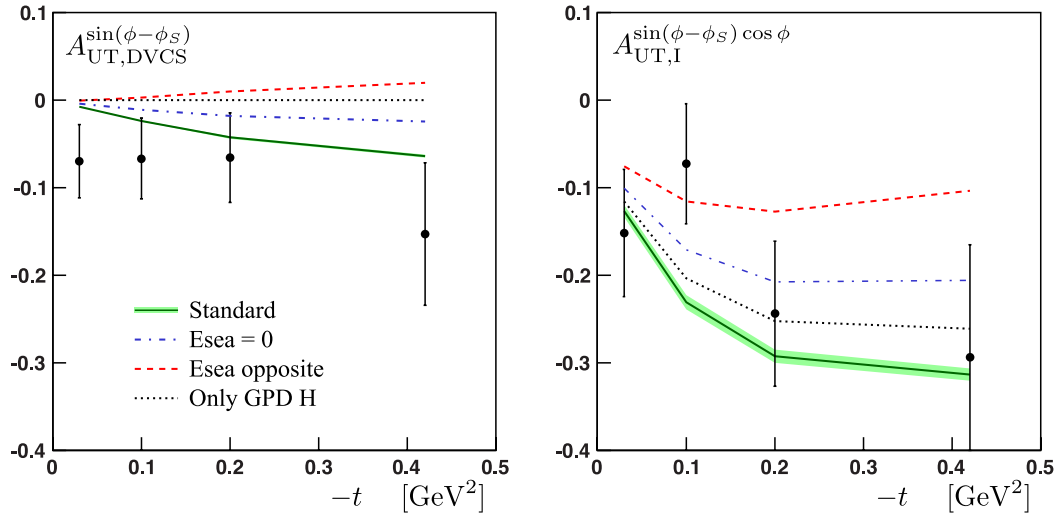
Applications

our set of GPDs allows for [parameter free calculations](#) of other hard exclusive reactions

- $\gamma^* p \rightarrow \omega p$ [Goloskokov-K, 1407.1141](#)
compared to SDMEs from [HERMES\(14\)](#) good agreement
prominent role of pion pole
- DVCS [K-Moutarde-Sabatie, 1210.6975](#)
compared to data from [Jlab, HERMES, H1, ZEUS](#)
good agreement with small skewness data
less good agreement with Jlab data

What future data from HERA and HERMES
will provide additional constraints and tests
of this set of GPDs?

Target asymmetry in DVCS



data: HERMES 06

$$\langle Q^2 \rangle \simeq 2.7 \text{ GeV}^2$$

$$\langle x_B \rangle \simeq 0.1$$

K.-Moutarde-Sabatie

$$A_{UT,DVCS}^{\sin(\phi-\phi_s)} \sim \text{Im} \left[\langle E \rangle^* \langle H \rangle \right]$$

$$A_{UT,I} \sim \text{BH} - \text{DVCS interference}$$

negative $\langle E^s \rangle$ favored

(flavor symmetric sea for E assumed)

E^s restricted by positivity bound for FT of zero-skewness GPDs Burkardt

$$(\vec{\Delta}_\perp \rightarrow \vec{b}_\perp; t = -\Delta_\perp^2) \quad \frac{b_\perp^2}{m^2} \left(\frac{\partial e_s(\bar{x}, \vec{b}_\perp)}{\partial b_\perp^2} \right)^2 \leq s^2(\bar{x}, \vec{b}_\perp) - \Delta s^2(\bar{x}, \vec{b}_\perp)$$

Parton angular momenta

Ji's sum rule (for $\xi = 0, t = 0$): ($q = u, \bar{u}, d, \dots$)

$$J^q = \frac{1}{2} [q_{20} + e_{20}^q] \quad J^g = \frac{1}{2} [g_{20} + e_{20}^g]$$

$q_{20} (= \int_0^1 dx x q(x))$, g_{20} from [ABM11](#)

e_{20} for valence quarks from form factor analysis

$e_{20}^s = 0 \div -0.026$ (from A_{UT} for DVCS and saturation of positivity bound)

e_{20}^g from sum rule: $e_{20}^g = -\sum e_{20}^{q_v} - 2\sum e_{20}^{\bar{q}}$

$$J^{u+\bar{u}} = 0.261 \div 0.235$$

$$J^{s+\bar{s}} = 0.017 \div -0.009$$

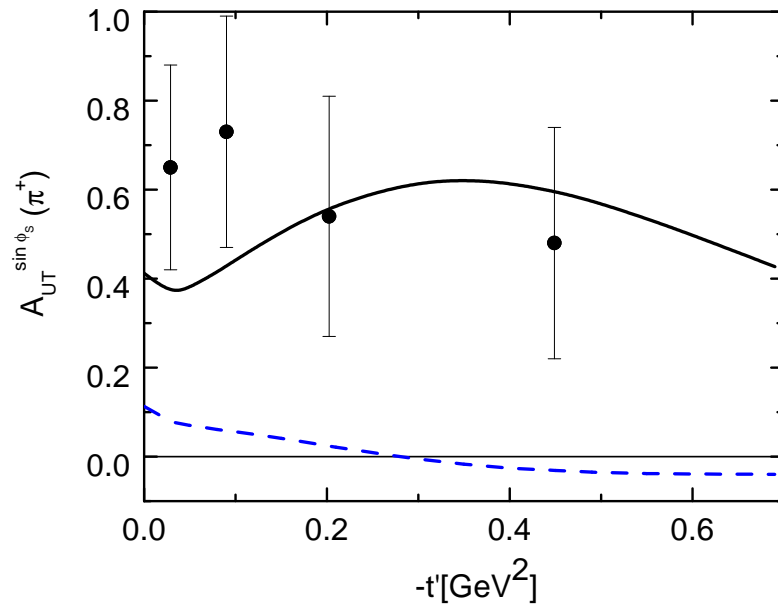
$$J^{d+\bar{d}} = 0.035 \div 0.009$$

$$J^g = 0.187 \div 0.265$$

at scale 2 GeV uncertainties due to E^s (l.h.s. $e_{20}^s = 0$, r.h.s. $e_{20}^s = -0.026$)

need for smaller errors of A_{UT} for DVCS

Transversity GPDs



HERMES(09) π^+

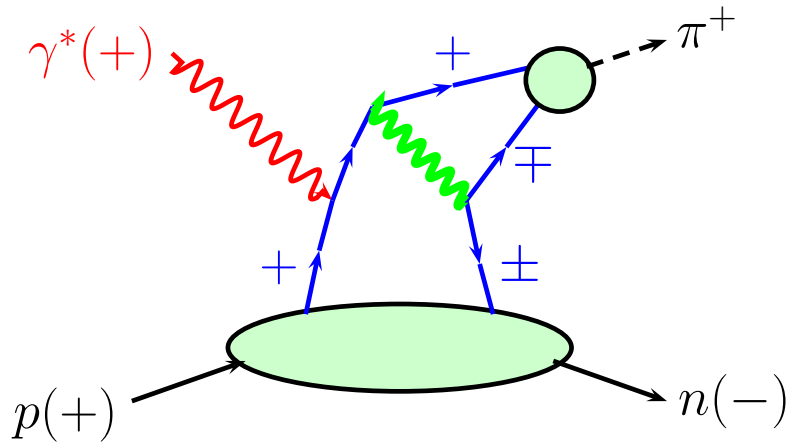
$Q^2 \simeq 2.5 \text{ GeV}^2$, $W = 3.99 \text{ GeV}$

ϕ_s orientation of target spin vector

large, does not vanish for $t' \rightarrow 0$

$$A_{UT}^{\sin \phi_s} \sim \text{Im} \left[\mathcal{M}_{0-,++}^* \mathcal{M}_{0+,0+} \right]$$

large n-f. ampl. $\mathcal{M}_{0-,++}$ required, non-vanishing in forward direction



usual GPDs: $\mathcal{M}_{0-,++} \sim t'$
 transversity GPDs: $\mathcal{M}_{0-,++} \rightarrow c$
 go along with twist-3 pion w.f.

suppressed by μ_π/Q
 as compared to leading-twist ampl.
 $\mu_\pi = m_\pi^2 / (m_u + m_d) \simeq 2 \text{ GeV}$
 at scale 2 GeV (Braun-Filianov (90))

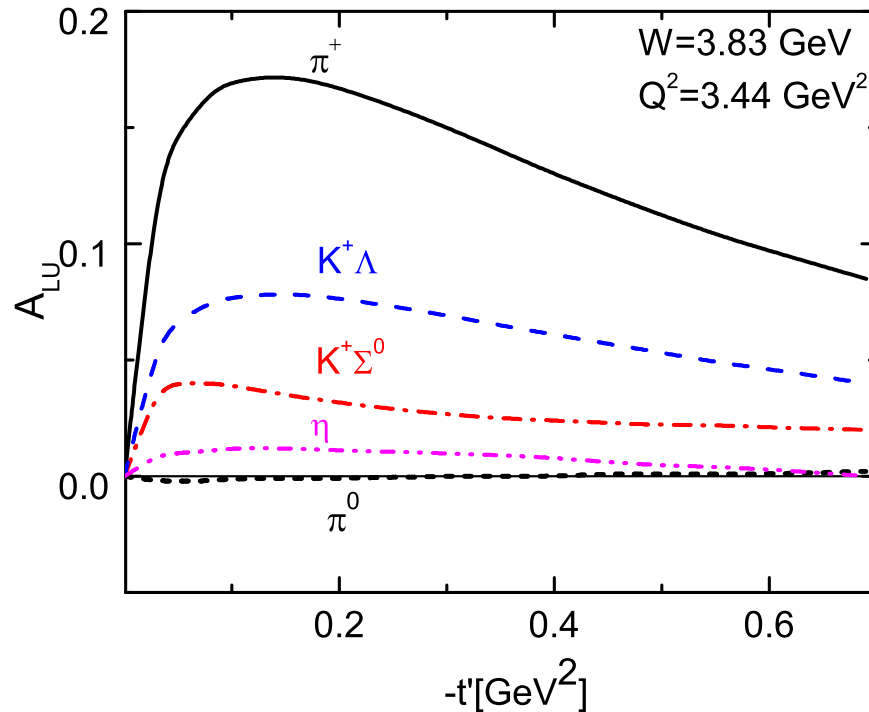
$$\mathcal{M}_{0-,++} = e_0 \sqrt{1 - \xi^2} \int d\bar{x} \mathcal{H}_{0-,++}^{\text{twist}-3} H_T \quad \mathcal{M}_{0+,\pm\pm} = -e_0 \frac{\sqrt{-t'}}{2m} \int d\bar{x} \mathcal{H}_{0-,++}^{\text{twist}-3} \bar{E}_T$$

H_T constrained by δ_q (Anselmino et al(09)) or lattice QCD

\bar{E}_T constrained by lattice QCD

QCDSF-UKQCD

The beam spin asymmetry for π production



Goloskokov-K, 1106.4897

$$A_{LU} \sim \text{Im} \left[(\mathcal{M}_{0+,++}^* - \mathcal{M}_{0+,-+}^*) \mathcal{M}_{0+,0+} + (\mathcal{M}_{0-,++}^* - \mathcal{M}_{0-,-+}^*) \mathcal{M}_{0-,0+} \right]$$

if H_T and \bar{E}_T are the dominant transversity GPDs:

$$A_{LU} \sim \text{Im} \left[\mathcal{M}_{0-,++}^* \mathcal{M}_{0-,0+} \right] \sim \text{Im} \left[\langle H_T \rangle^* \langle \tilde{E} \rangle \right]$$

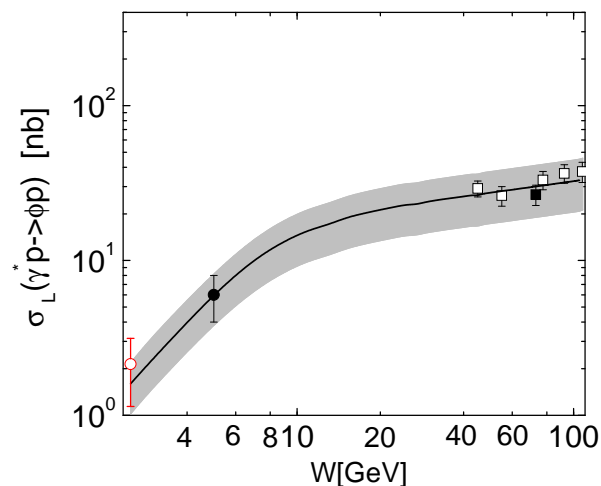
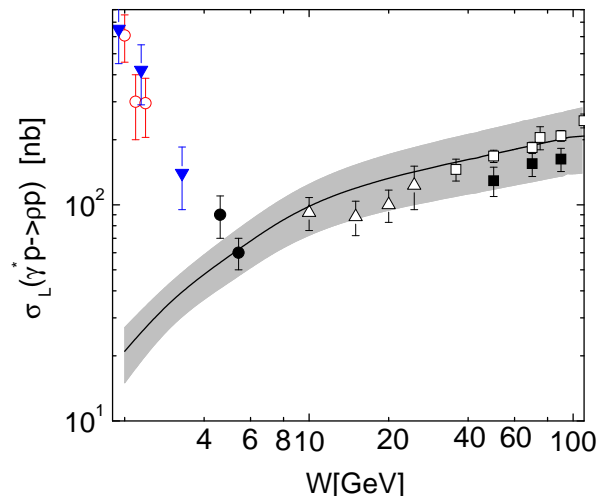
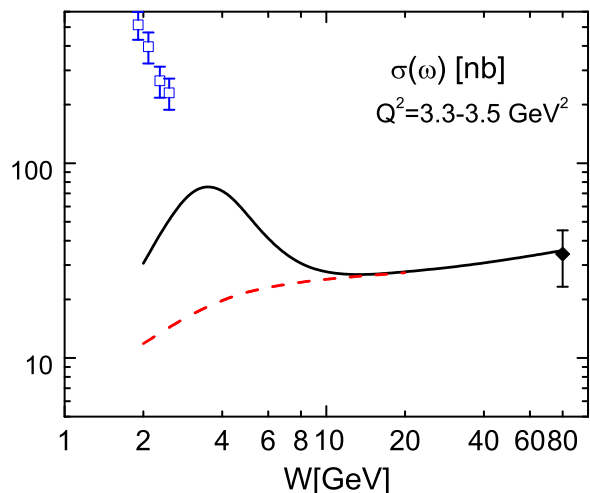
additional information on H_T and, for π^0 , on $\tilde{E}_{\text{non-pole}}$

Summary

- more accurate data on A_{UT} for DVCS and data on π^0 lepton production as well as the η/π^0 ratio would help to fix E^{sea} and the transversity GPDs
- of interest are also data on
 $A_{UL}(\pi)$: \bar{E}_T important
any information on strangeness production, e.g.
 $A_{UT}(K^+\Lambda, K^+\Sigma)$: sensitive to H_T^s
 $\sigma(\omega)$: may answer question whether $H^{\bar{u}} = H^{\bar{d}}$ or not

BACK-UP

The ω cross section



at $Q^2 = 3.5 \text{ GeV}^2$

4 GeV^2

3.8 GeV^2

data from H1, ZEUS, E665, HERMES, CORNELL, CLAS

Goloskokov-K, hep-ph/0611290, 1407.1141

decreasing W at fixed $Q^2 \implies$ increasing x_B, ξ and $-t_0 = 4m^2\xi^2/(1-\xi^2)$

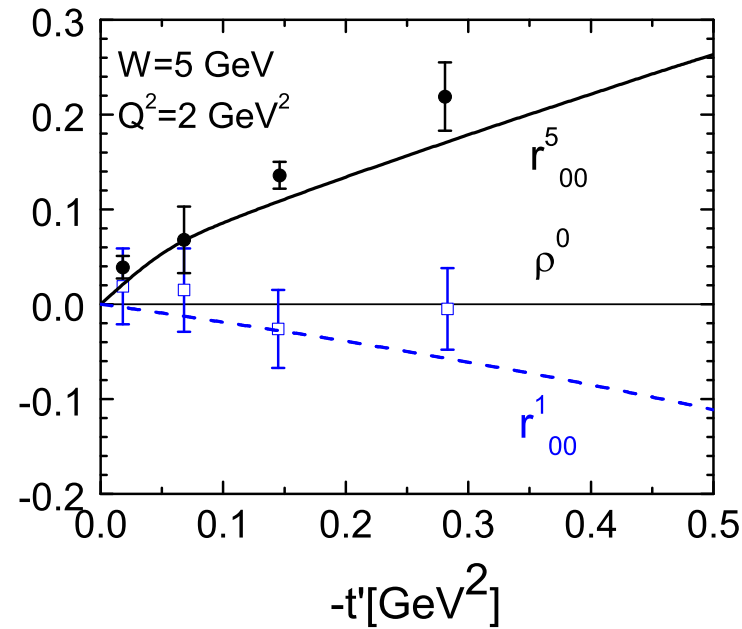
Handbag: GPDs decrease with increasing $-t$, decrease of c.s. related to $-t_0$

large ρ^0, ω cross sections at low W not understood, probably not hard physics

$\sigma \sim |gluon + sea|^2$: $\sigma(\omega)/\sigma(\rho^0) = 1/9$ for $H^{\bar{u}} = H^{\bar{d}}$ true?

(up to wave function effects, e.g. $f_\rho = 209 \text{ MeV}$, $f_\omega = 187 \text{ MeV}$)

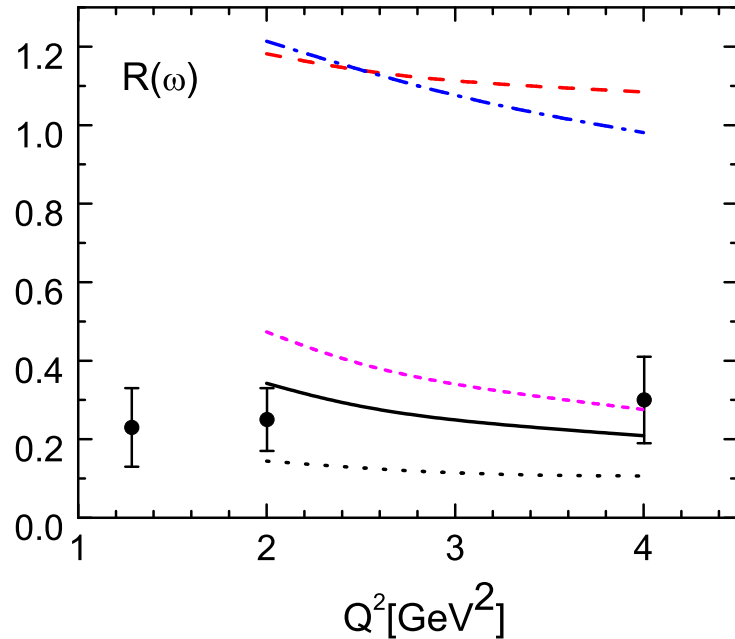
\bar{E}_T in ρ^0 production



$$r_{00}^1 \sim -|\langle \bar{E}_T \rangle|^2 \quad r_{00}^5 \sim \text{Re} \left[\langle \bar{E}_T \rangle^* \langle H \rangle \right]$$

Goloskokov-K. 1310.1472

ω production



data from [HERMES 1407.2119](#)
 $W = 4.8 \text{ GeV}, t' = -0.08 \text{ GeV}^2$

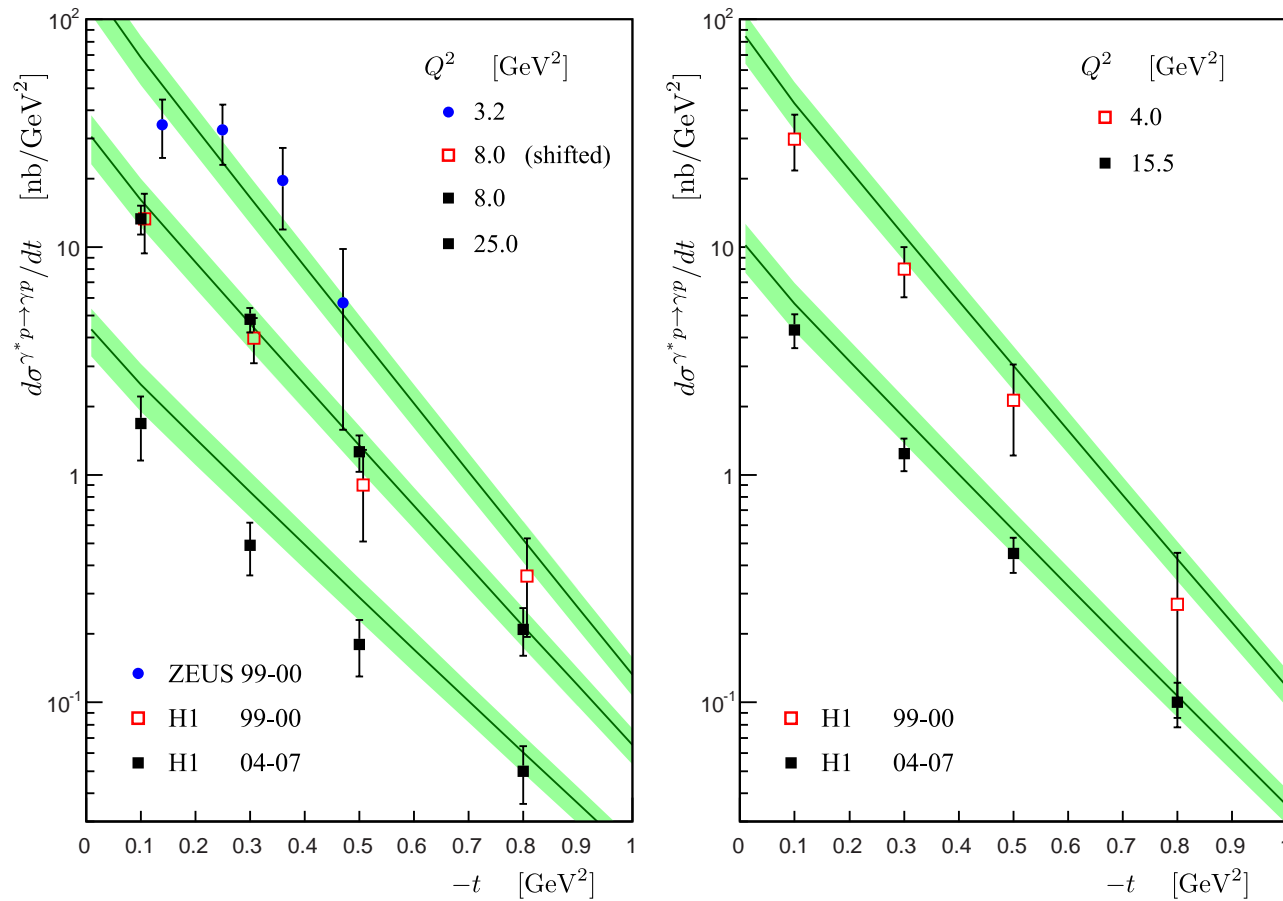
[Goloskokov-K, 1407.1141](#)

$$R = \frac{1}{\varepsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}} = d\sigma(\omega_L)/d\sigma(\omega_T) \quad (\text{full result: black line})$$

red line: without pion pole, magenta line: $R = d\sigma_L/d\sigma_T$

blue (dotted) line: predictions at $W = 8(3.5) \text{ GeV}$

DVCS predictions



K-Moutarde-Sabatie, 1210.6975