

Electroweak Physics with HERA Data

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Goals in DIS $e p$ scattering

- High-precision measurements of the **proton structure**
- Search for **new physics**
- Electroweak physics:

Tests of the standard model, and after LHC discoveries

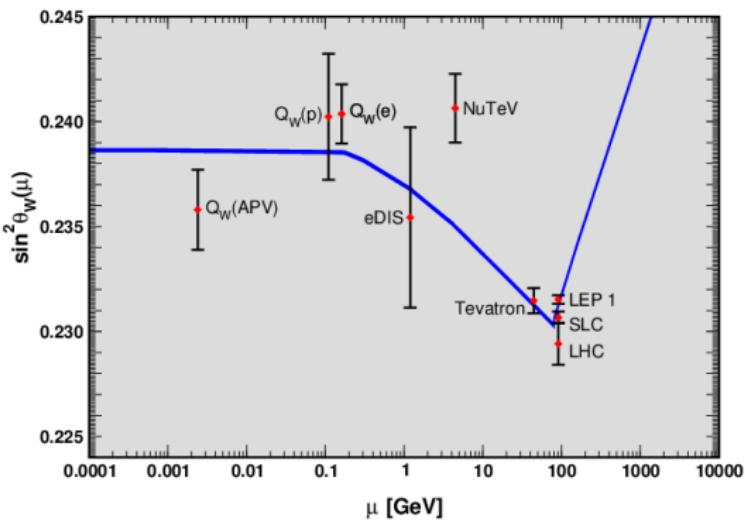
Tests of the standard model extended with new physics

This talk:

- Some remarks about the present status of $\sin^2 \theta_W$
- Electroweak parameters in DIS
- Expectations 20 years ago and present status

(not a complete review)

The Running Weak Mixing Angle: Present Measurements of $\sin^2 \hat{\theta}_W$ (\overline{MS})



PDG 2014

- $Q_W(APV)$: atomic parity violation (Cs)
- $Q_W(e)$: Moller scattering
- NuTeV: Neutrino scattering (re-analysis needed)

Z-pole measurements:

- LEP1 and SLC
- Tevatron
- First results from LHC

The Weak Mixing Angle: Present Situation

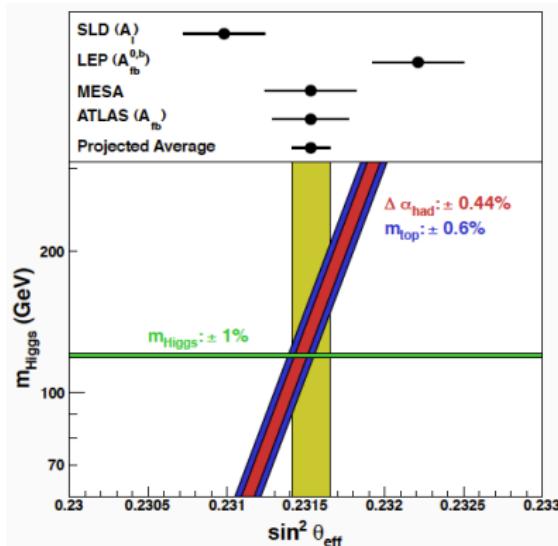
Present situation

- $\sin^2 \hat{\theta}_w(m_Z) = 0.23070 \pm 0.00026$ from A_{LR} , SLD
- $\sin^2 \hat{\theta}_w(m_Z) = 0.23193 \pm 0.00029$ from $A_{FB}^{b\bar{b}}$, LEP1
 - 3σ difference !
- $\sin^2 \hat{\theta}_w(m_Z) = 0.23125 \pm 0.00016$ world average
- $\sin^2 \hat{\theta}_w(m_Z) = 0.23104 \pm 0.00015$ from α , G_μ , m_Z and m_W

Very different implications for new physics:
look at S , T , U parameters, e.g.,

- from A_{LR} → $S = -0.18 \pm 0.15$ → Susy?
- from $A_{FB}^{b\bar{b}}$ → $S = +0.46 \pm 0.17$ → heavy Higgs? KK at 1 - 2 TeV?
- from average → $S = +0.11 \pm 0.11$ → new heavy doublets? KK above 3 TeV?

Standard Model Relation: Higgs Boson Mass versus $\sin^2 \theta_W$



Combination of precision measurements at the Z -pole

→ $M_{\text{Higgs}} - \sin^2 \hat{\theta}_W(\mu)$ SM relation (red-blue band)

Precision measurement of $\sin^2 \hat{\theta}_W(\mu)$ had provided indirect evidence for the allowed range of M_{Higgs}

Combination of measurements provide strong tests of the SM,
... and maybe evidence for new physics

MESA at Mainz:

Parity violating polarization asymmetry in elastic $e p$ scattering, $Q^2 \simeq 0.0025 \text{ GeV}^2$
expected precision: $\Delta \sin^2 \theta_W = 0.00037$

DIS at Large Q^2 : Neutral Current

Neutral current at tree level, polarized e^\pm scattering

$$\frac{d^2\sigma_{NC}}{dx dQ^2} = \frac{2\pi\alpha^2}{Q^4 x} \left(Y_+ \mathbf{F}_2 + Y_- x \mathbf{F}_3 - y^2 \mathbf{F}_L \right)$$

with

$$\mathbf{F}_2 = F_2^\gamma + \kappa_Z (-v_e \mp Pa_e) F_2^{\gamma Z} + \kappa_Z^2 (v_e^2 + a_e^2 \pm 2Pv_e a_e) F_2^Z$$

$$x \mathbf{F}_3 = +\kappa_Z (\pm a_e + Pv_e) x F_3^{\gamma Z} + \kappa_Z^2 (\mp 2v_e a_e - P(v_e^2 + a_e^2)) x F_3^Z$$

$$\kappa_Z(Q^2) = \frac{Q^2}{Q^2 + m_Z^2} \frac{1}{4 \sin^2 \theta_w \cos^2 \theta_w}$$

$$(F_2^\gamma, F_2^{\gamma Z}, F_2^Z) = x \sum (Q_q^2, 2Q_q v_q, v_q^2 + a_q^2) (q + \bar{q})$$

$$x(F_3^{\gamma Z}, F_3^Z) = 2x \sum (Q_q a_q, v_q a_q) (q - \bar{q})$$

$$v_f = l_f^{(3)} - 2Q_f \sin^2 \theta_w, \quad a_f = l_f^{(3)} \quad (f = e, u, d, \dots)$$

Parameters: $\alpha, m_Z, \sin^2 \theta_w$; more general: also $v_e, a_e; v_q, a_q$

DIS at Large Q^2 : Charged Current

Charged current at tree level

$$\frac{d^2\sigma_{CC}}{dx dQ^2} = \frac{1 \pm P}{2} \frac{2\pi\alpha^2}{Q^4 x} \kappa_W^2 \left(Y_+ \mathbf{W}_2 \pm Y_- x \mathbf{W}_3 - y^2 \mathbf{W}_L \right)$$

with

$$\kappa_W(Q^2) = \frac{Q^2}{Q^2 + m_W^2} \frac{1}{4 \sin^2 \theta_W}$$

Parameters: $\alpha, m_W, \sin^2 \theta_W$

Parameter Relations in the Standard Model

Relations at tree-level, e.g.,

- $\cos \theta_W = m_W/m_Z$ (by definition in the on-mass-shell scheme)

- Muon decay constant: $G_F = \frac{\pi\alpha}{\sqrt{2}\sin^2\theta_W m_W^2}$

→ $\frac{d^2\sigma_{CC}}{dx dQ^2} = \frac{1 \pm P}{2} \frac{G_F^2}{2\pi x} \left(\frac{m_W^2}{m_W^2 + Q^2} \right)^2 \sigma_{r,CC}$

and

→ $\frac{d^2\sigma_{NC}}{dx dQ^2}$ using $\kappa_Z(Q^2) = Q^2 \frac{G_F}{2\pi\alpha} \frac{m_Z^2}{Q^2 + m_Z^2}$

- One-loop corrections to the muon decay: Δr

$$G_F = \frac{\pi\alpha}{\sqrt{2}\sin^2\theta_W m_W^2} (1 + \Delta r) \quad \text{with} \quad \Delta r = \Delta r(\alpha, m_W, \sin\theta_W, m_{top}, M_{Higgs}, \dots)$$

- On-mass-shell scheme versus modified on-mass-shell scheme:
 (α, m_W, m_Z) or $(\alpha, m_W, \sin^2\theta_W)$ versus (α, G_F, m_Z)

Fits based on different parametrizations may lead to very different results

Effective Couplings

Including quantum loop corrections, universal contributions:

Self energy diagrams of the exchanged boson (γ , Z , W)

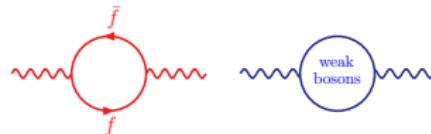
Feynman diagram showing the addition of a self-energy correction to the photon propagator. The left part shows a bare photon line with a wavy arrow. The right part shows the same line with a wavy arrow plus a loop correction enclosed in a red circle labeled Π_γ . A red arrow points to the equation $\alpha \rightarrow \alpha + \Delta\alpha = \alpha(Q^2)$.

Feynman diagram showing the addition of a self-energy correction to the Z boson propagator. The left part shows a bare Z boson line with a wavy arrow. The right part shows the same line with a wavy arrow plus a loop correction enclosed in a red circle labeled $\Pi_{\gamma Z}$. A red arrow points to the equation $\sin^2 \theta_W \rightarrow \sin^2 \theta_W + \Delta \sin^2 \theta_W = \sin^2 \theta_{\text{eff}}(Q^2)$.

Feynman diagram showing the addition of a self-energy correction to the W boson propagator. The left part shows a bare W boson line with a wavy arrow. The right part shows the same line with a wavy arrow plus a loop correction enclosed in a red circle labeled Π_W . A red arrow points to the equation $G_F \rightarrow G_F(1 + \Delta r)$.

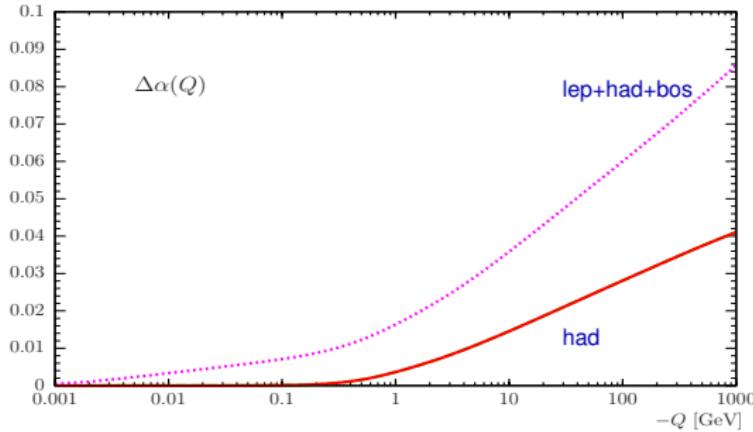
One-Loop Corrections: $\Delta\alpha$

- Self energy diagrams:



$$\propto \log \frac{Q^2}{m_f^2} \rightarrow O(10\%) \quad \text{small, } O(1\%)$$

- Photon self energy = vacuum polarization, absorbed and resummed in the running fine structure constant: $\alpha \rightarrow \alpha(Q^2) = \frac{\alpha}{1 - \Pi_\gamma(Q^2)}$

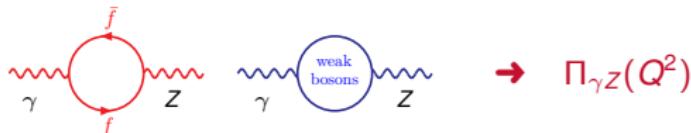


- Z-boson self energy: a small correction if written in terms of:

$$\frac{\alpha}{\sin \theta_W^2 \cos \theta_W^2} \rightarrow \frac{m_Z^2 G_\mu \sqrt{2}}{\pi} \frac{1 - \Delta r}{1 - \Pi_Z(Q^2)}$$

One-Loop Corrections: Scale-Dependent $\sin^2 \theta_W$

- Photon-Z mixing



can be absorbed into **effective**, running, **scale-dependent weak mixing angle**

Definitions of the weak mixing angle:

- On-shell definition: $\cos \theta_W = \frac{m_W}{m_Z}$ ($\rightarrow \Delta r$ contains large contribution from m_{top})
 - $\sin^2 \theta_{\text{eff}}(Q^2)$ absorbs $\Pi_{\gamma Z}(Q^2)$
 - **$\overline{\text{MS}}$ scheme:** $\sin^2 \hat{\theta}_W(\mu)$ (via $\tan \hat{\theta}_W(\mu) = g_1/g_2$)
less sensitive to m_{top} , suited for comparisons with extensions of the SM
- Relation

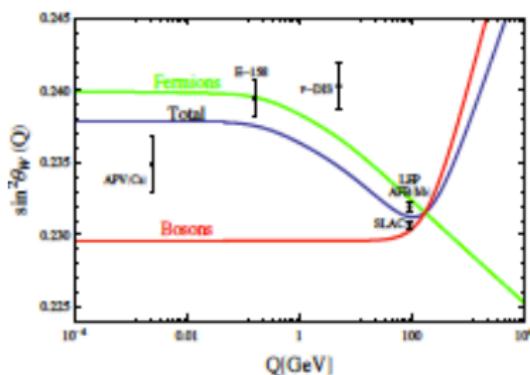
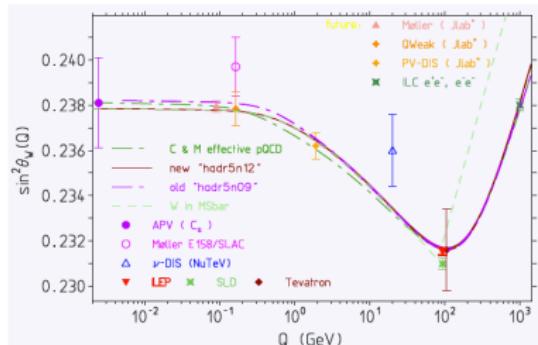
$$\sin^2 \hat{\theta}_W(\mu) = \left(1 + \frac{\rho_t}{\tan^2 \theta_W} + \dots \right) \sin^2 \theta_W$$

$$\text{with } \rho_t = 3G_\mu m_{top}^2 / 8\sqrt{2}\pi^2 = 0.00939 (m_{top}/173 \text{ GeV})^2$$

- Additional uncertainty: $\Delta \sin^2 \hat{\theta}_W \simeq \pm 0.0001$ for a 1% error on m_{top}

Scheme Dependence of $\sin^2 \hat{\theta}_W$

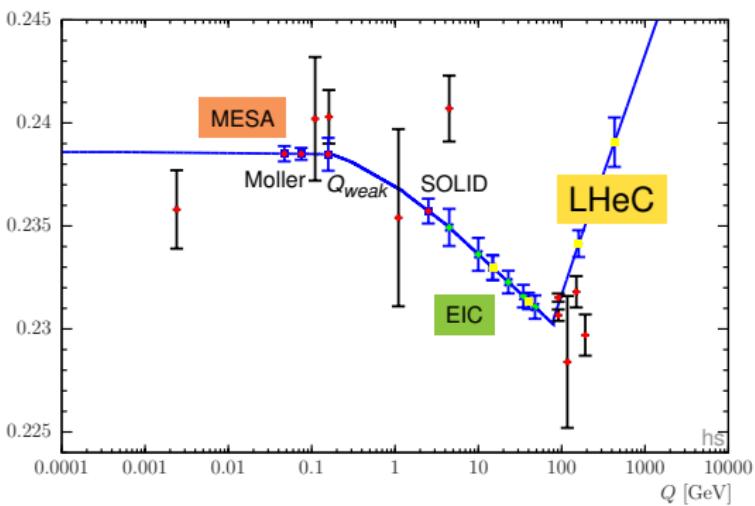
Definition of $\sin^2 \theta_W(Q)$: absorbing $\Pi_{\gamma Z}$ and part of vertex+box corrections



- C&M: Czarnecki, Marciano: pQCD with $m_{q,\text{eff}}$ and bosonic contributions
- 'old' and 'new' program hadr5 by Jegerlehner 2010-12, hadronic VP from data (different $SU(3)$ flavor splitting)
- Scheme dependence partly compensated by remaining $\delta_{1-\text{loop}}$
- Include higher orders
- Parameter dependence m_t , m_H , ...
- Prescriptions known up to small uncertainties

$\sin^2 \hat{\theta}_W$: Present Status and Possible Future

- Use running $\sin^2 \hat{\theta}_W(Q)$ to compare different measurements
- Match definition of $\sin^2 \hat{\theta}_W(Q)$ with the complete 1-loop corrections



Present and future of $\sin^2 \hat{\theta}_W(Q)$

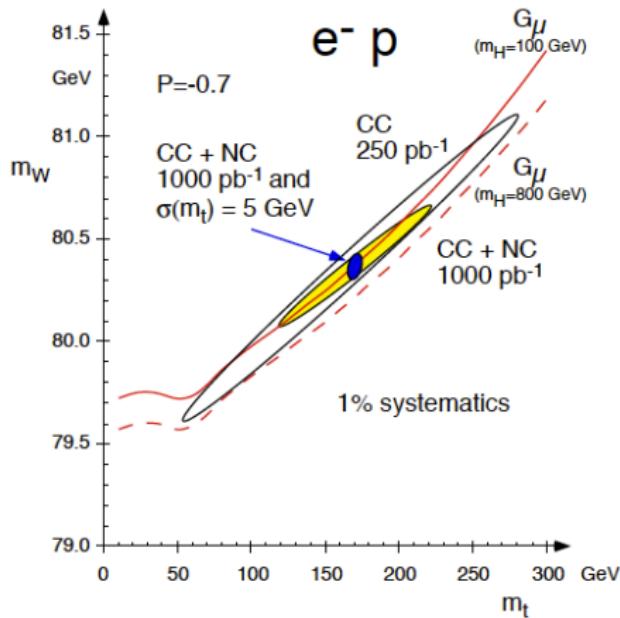
- Existing: APV, $Q_w(e)$, $Q_w(p)$, eDIS, NuTeV, LEP, SLD, Tevatron, CMS, ATLAS
 - Low energy: MESA, Moller, Q_{weak}
 - Medium energy: SOLID, EIC
 - High energy: LHC, LHeC, ...
- ... and HERA ?

The Optimistic View of 20 Years Ago

R. Beyer et al. in
Future Physics at HERA
1995/96

Sensitivity of NC+CC data
electrons with polarization
SM fit to $m_W - m_t$
(on-mass-shell scheme:
 $\alpha, m_Z, m_W, m_t, m_H$)

For 10^3 pb^{-1} , $\Delta m_t = \pm 5 \text{ GeV}$
→ $\Delta m_W = \pm 55 \text{ MeV}$



Present HERA I

H1 2006 with $\simeq 120 \text{ pb}^{-1}$
unpolarized $e^\pm \text{ NC+CC}$
(Phys. Lett. B632)

fit with fixed G_F :

$$m_W = 82.87 \pm 1.82^{+0.30}_{-0.16} \text{ GeV}$$

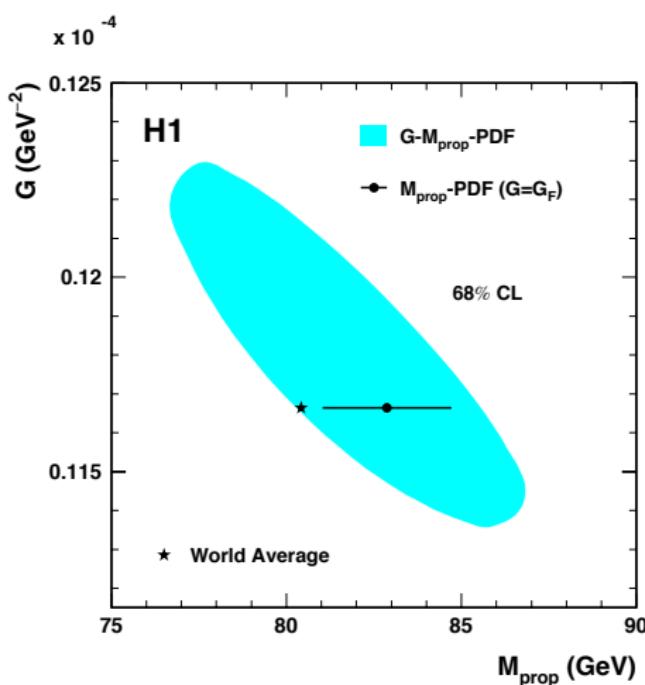
SM fit:

$$m_W = 80.786 \pm 0.205^{+0.063}_{-0.098} \text{ GeV}$$

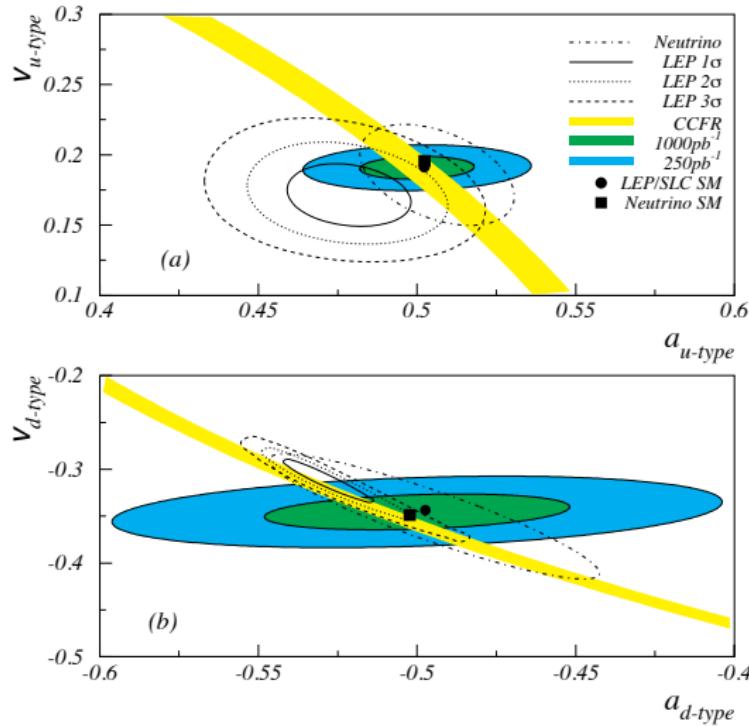
corresponds to

$$\sin^2 \theta_W = 0.2151 \pm 0.0040^{+0.0019}_{-0.0011}$$

- Use full statistic!
- Use polarization!



Quark Couplings: Expected in 1995



A more general approach
to test the Standard Model:

Quark coupling constants

v_q, a_q

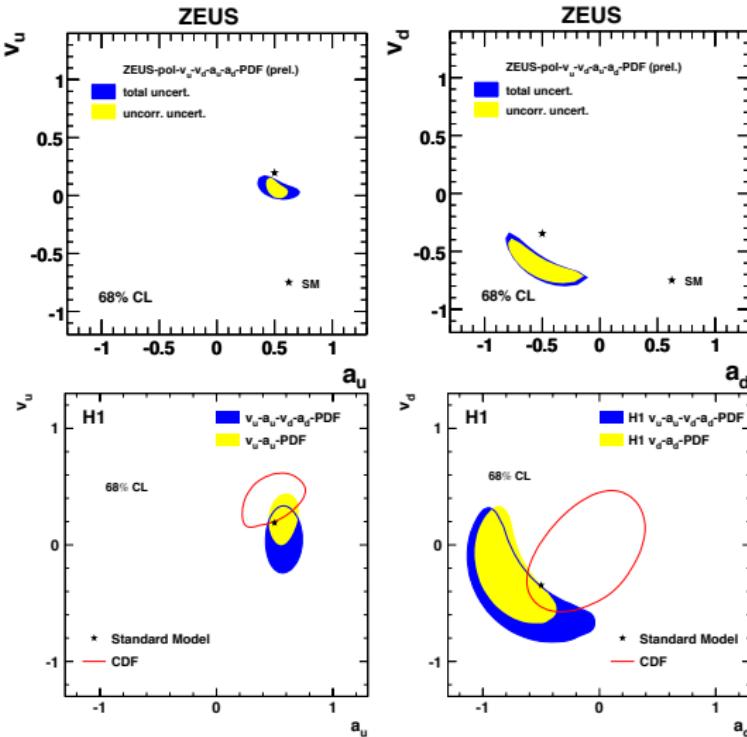
from fits to NC and CC
for e_L^- , e_R^- , e_L^+ , e_R^+

Polarization is essential

v_u, a_u, v_d, a_d :
13%, 6%, 17%, 17%

R. J. Cashmore et al. in
Future Physics at HERA
1995/96

Quark Couplings: Present Status



Present status

ZEUS-prel-07-027
H1, PLB632 (2006) 35

Unique at HERA:
flavor and v/a separation

See next talk by Katarzyna Wichmann

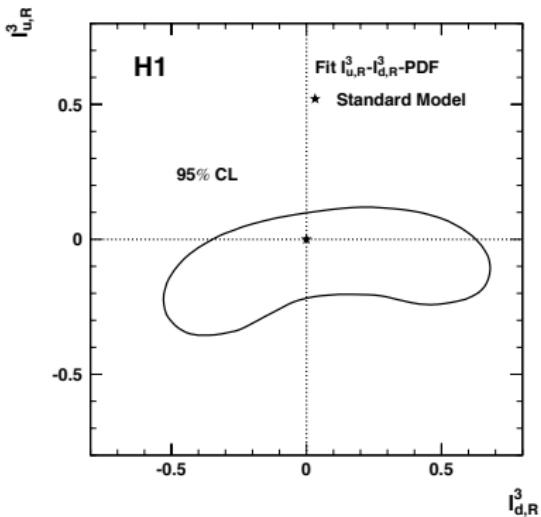
Electroweak Fits Restrict Physics Beyond the SM

H1 2006 with $\simeq 120 \text{ pb}^{-1}$
unpolarized e^\pm NC+CC
(Phys. Lett. B632)

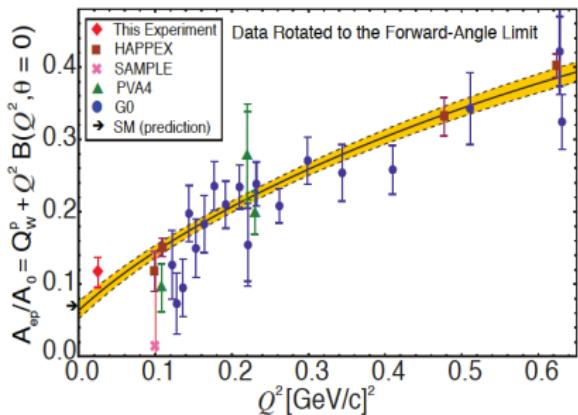
Light quarks with right-handed couplings:

$|I_{u,R}^3| = 0.5$ excluded at 95% CL

→ Match to low-energy fits
using couplings $C_{1u} \pm C_{1d}$, $C_{2u} \pm C_{2d}$
(see contact interactions fits)



First Results from Q_{weak}

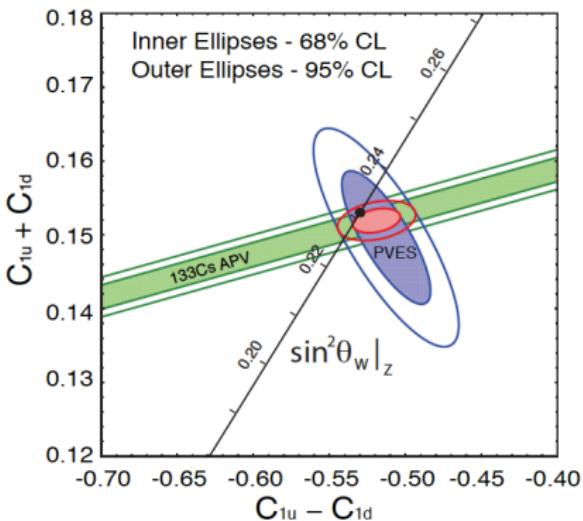


- 2-dim fit to quark couplings
- Best fit at $\sin^2 \theta_w = 0.23116$ (expect $\Delta \sin^2 \theta_w = 0.0008$)
- $Q_W^p = 0.064 \pm 0.012$ ($Q_W^p(SM) = 0.0710 \pm 0.0007$)

PRL111 (2013)

First results from Q_{weak}

- Analyzed 4% of data (expect factor 5 improved errors with full data taken)
- Used previous data to extrapolate to $Q^2 = 0$



Summary

- Present measurements of the weak mixing angle need improvement
 - Expect results from LHC, Q_{Weak}, MESA, SOLID, ...
 - Combined EW+QCD fit of HERA data is a 'must do'
 - Significant improvements for EW analysis of HERA data can be expected
- Complementarity:
only many measurements from different experiments allow conclusive tests
of the SM and of its extensions