

# C-TCT measurements on scCVD diamond and its use at CNGS

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# Outline

- Motivation and Concept

Why and how to investigate charge transport/collection  
at cold temperatures?

- The Set-up

Why diamond?  
The experimental technique

- Results and Model

Current pulses and what we can learn from them

- Diamonds at CNGS

# Motivation and Concept

- Understand charge transport over a wide temp. range:
  - Beam Loss Monitors within the LHC magnets
    - Operate detector at 1.9 K
  - Electrical Devices in space crafts
    - Huge temperature range
- Perform transient current technique (TCT) at various controlled temperatures and field strengths:
  - Provides a powerful technique for study of transport characteristics
- Analyse current pulses in terms of transit time, area, ...
  - compare data to model



# Why Diamond

Pros:

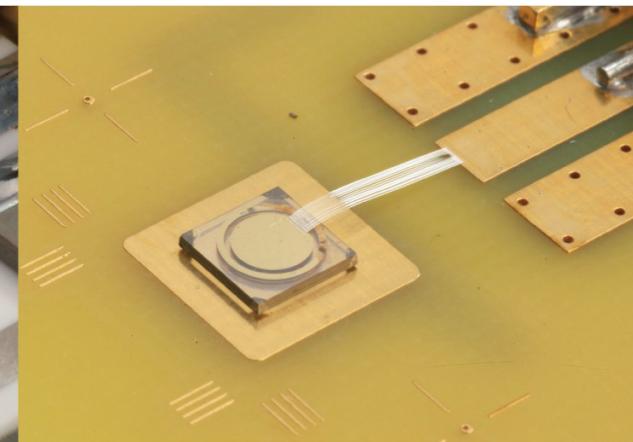
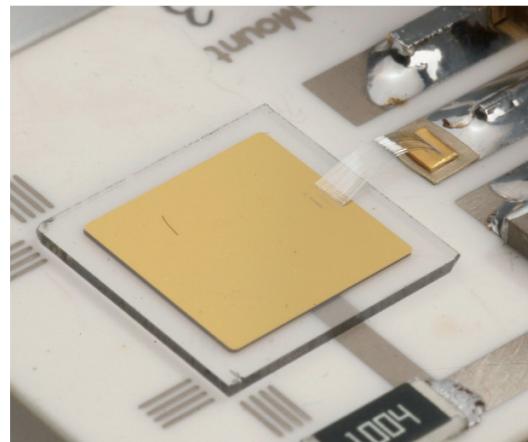
- High band gap (5.5 eV)  
→ Very high breakdown field  $> 1\text{e}7 \text{ V/cm}$
- Very high resistivity  $> 1\text{e}11 \Omega\text{cm}$ 
  - Large area
- Very low leakage current  $\sim \text{few pA}$
- Low dielectric constant (5.7)
  - Low capacitance → Low noise
- High displacement energy (43 eV/atom)
  - Radiation hard → No replacement
- High mobility ( $\sim 2000 \text{ cm}^2/\text{Vs}$ )
  - Fast signals
    - High collision rate
- Very wide sensitivity range
  - Single MIP to few THz tested
- Wide operational temperature range



Cons:



- High  $E_{\text{pair-creation}}$  (13.5 eV)
  - Less signal, but S2N-ratio comparable to Si
- Rather high costs
- Not as well understood as Si
  - More R&D efforts needed



# The Experimental Technique

- The Transient-Current Technique (TCT) measurement:

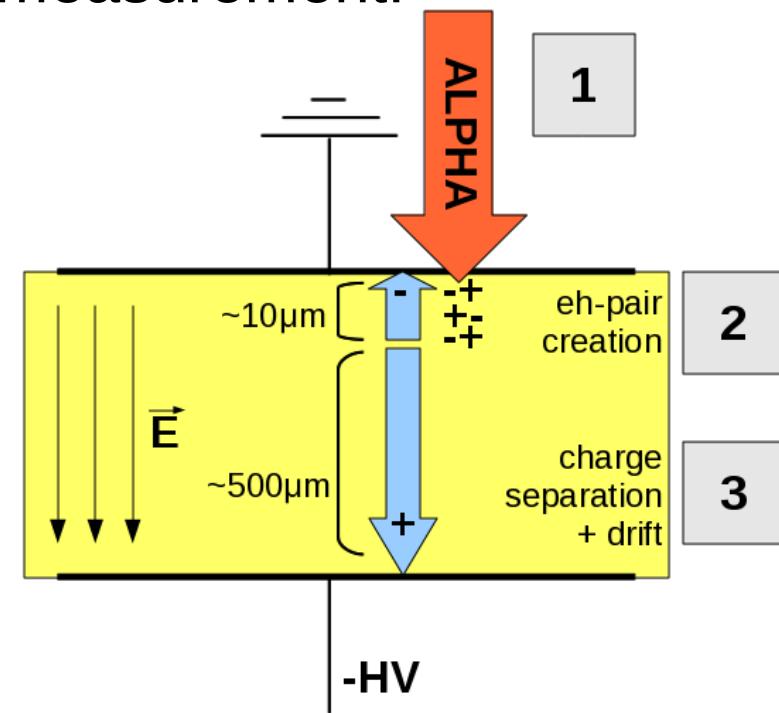
- 1)  $\alpha$  particles impinge on top side
- 2) Create eh-pairs **close** to electrode
- 3) Electric field separates charges
- 4) Drifting charges induce current

→ measure the **transient current**

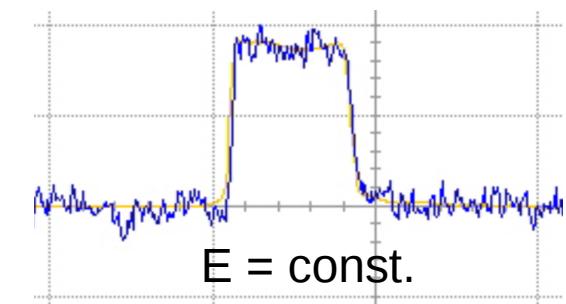
→ **Pos.** (**neg.**) bias → Measure  $e^-$  ( $h^+$ )

→ Use fast (2 GHz), low noise (3 mV)  
current amplifier, 40 dB

→ record signal with oscilloscope

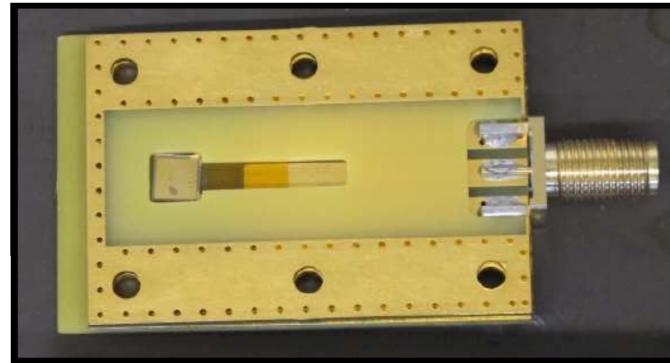


$$4 \quad i(t) = Q(t)v(t)/d$$

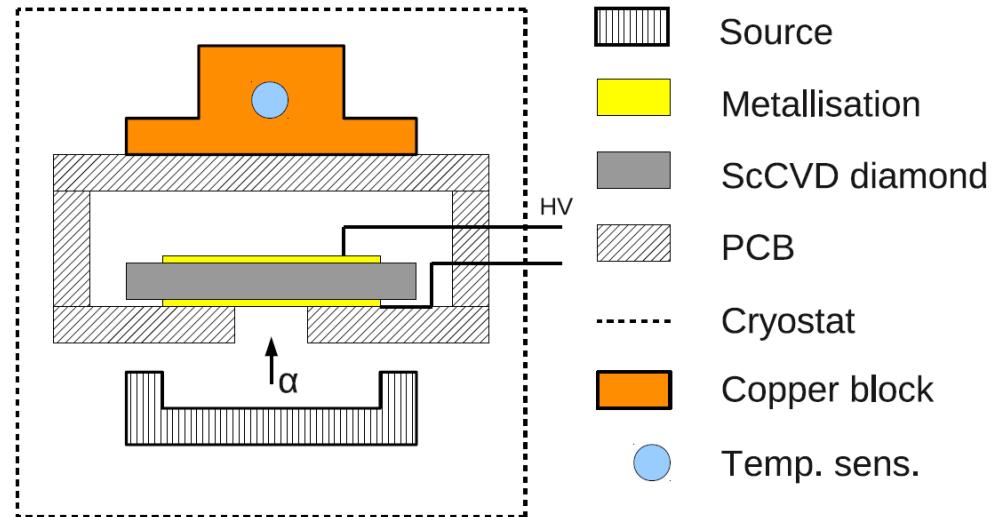


# The Set-up I

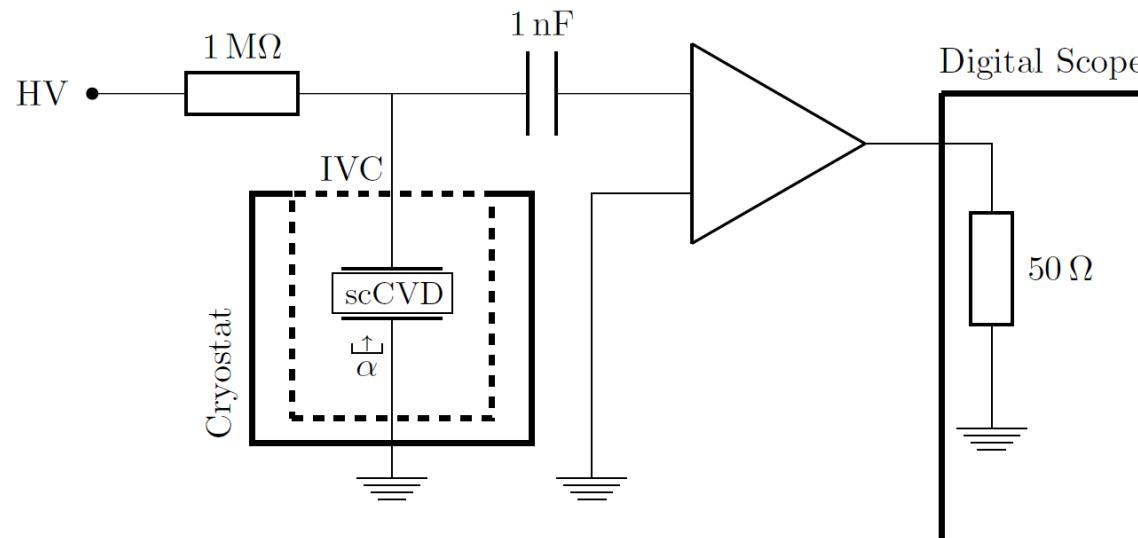
- Sensor + Sensor holder:



Sample size:  $5 \times 5 \text{ mm}^2$ ,  $500 \mu\text{m}$

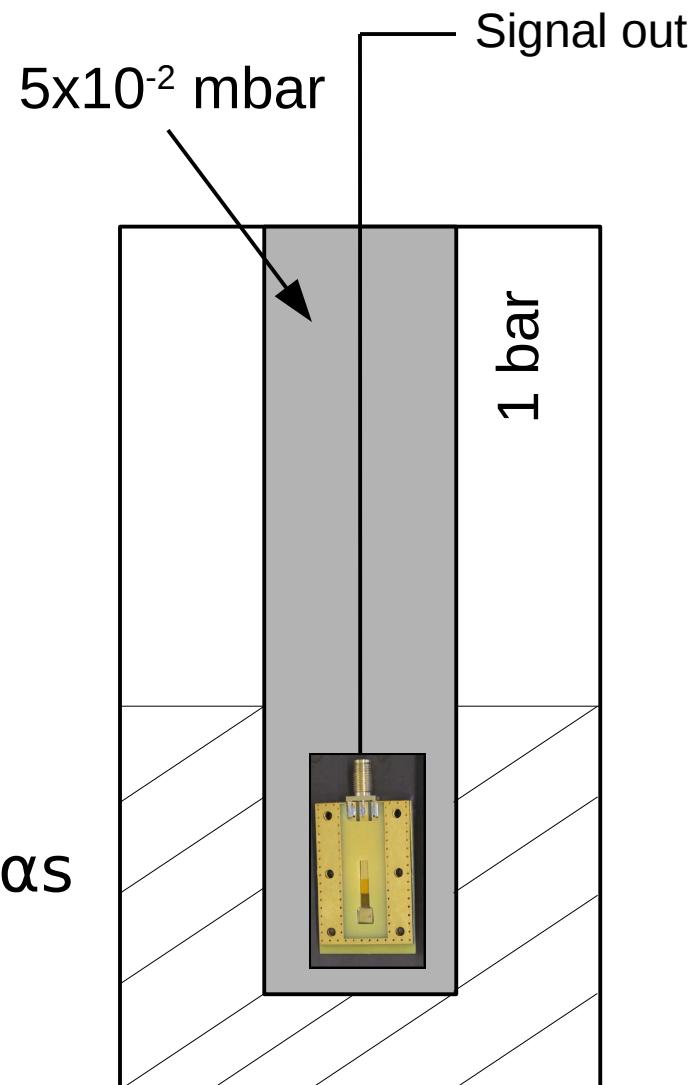


- Circuit schematics:

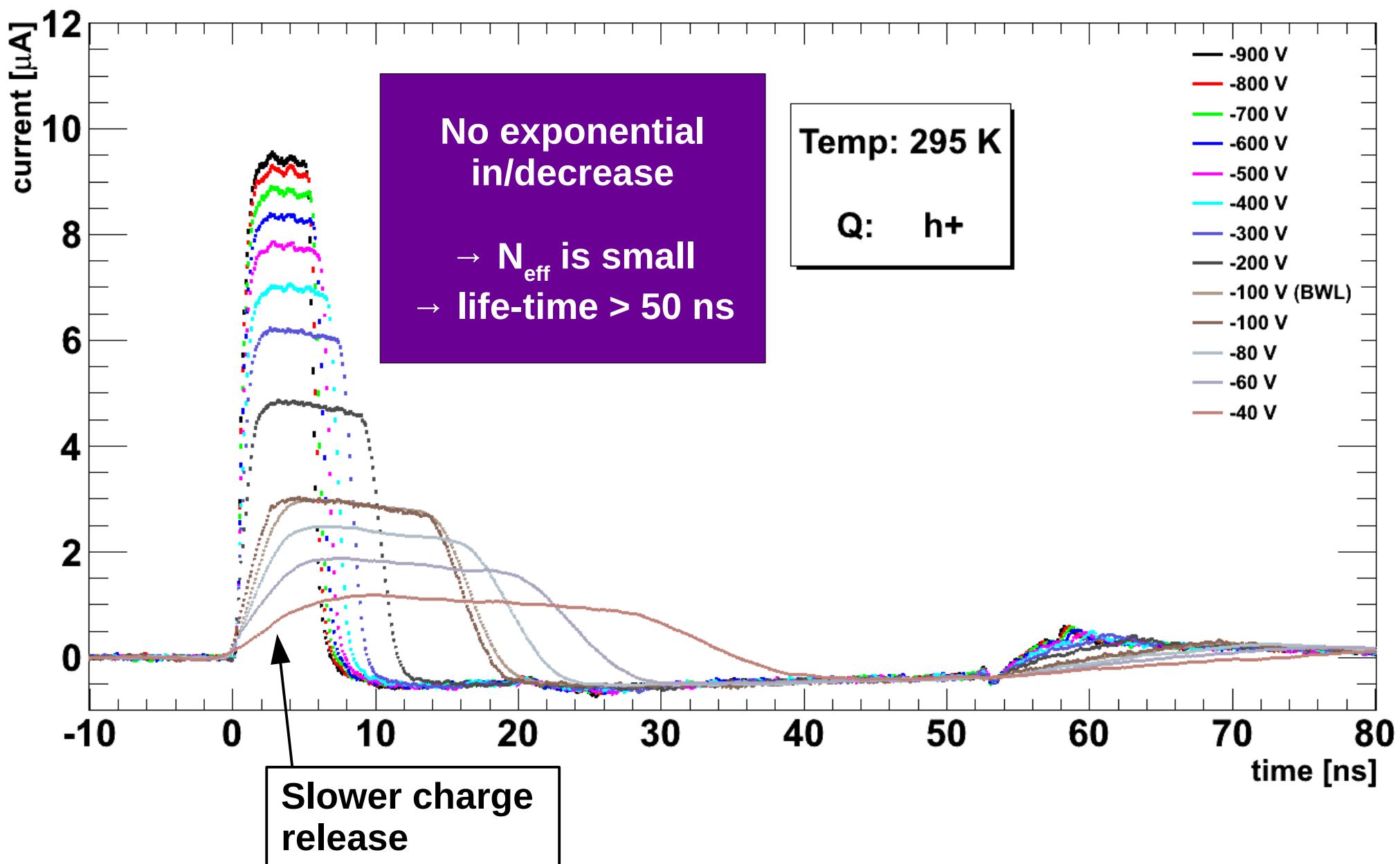


# The Cryostat

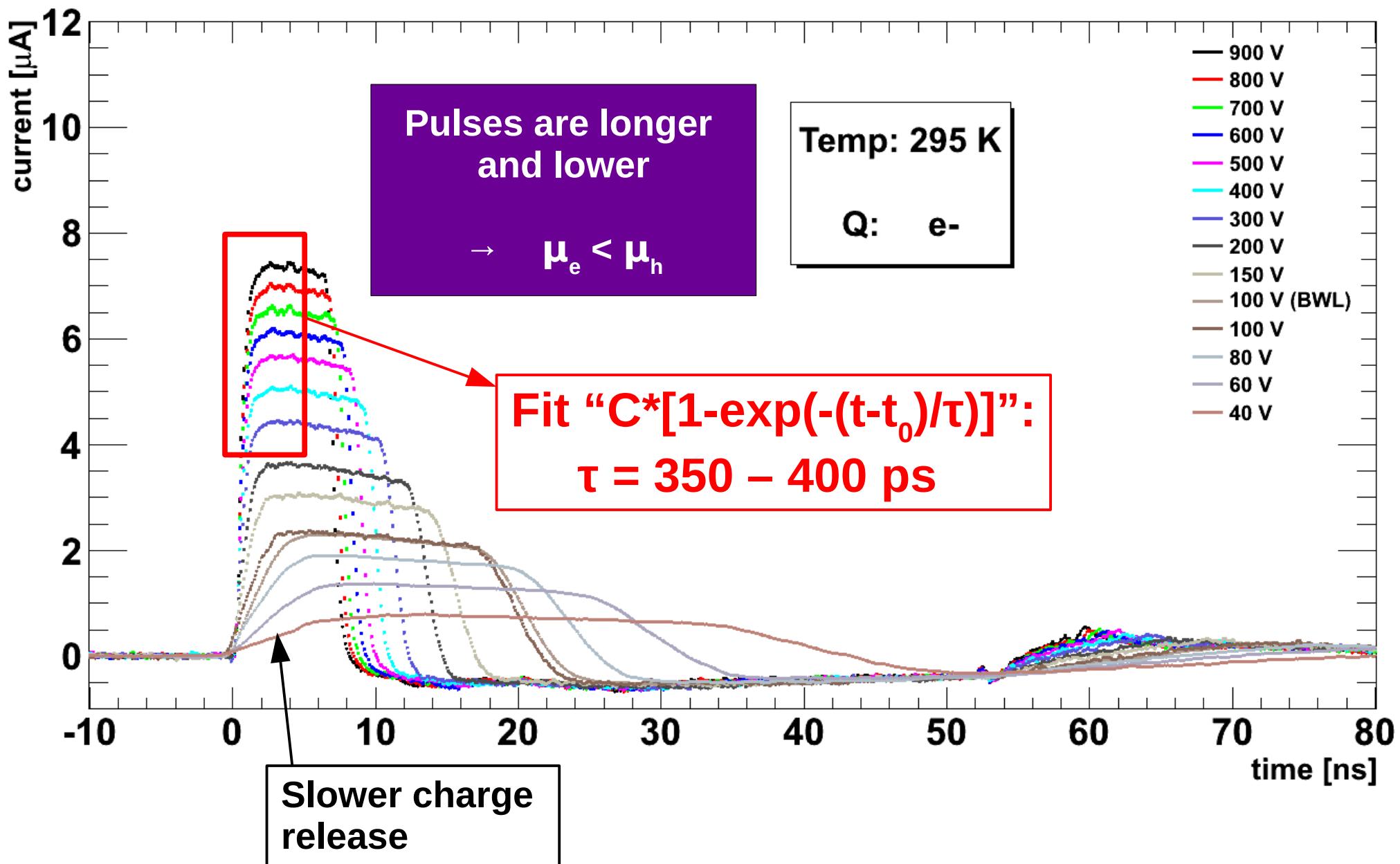
- Outer chamber:
  - liquid helium
  - pressure = 1 bar for  $T \geq 4.2 \text{ K}$
  - pressure few mbar for  $T = 1.6 \text{ K}$
- inner vacuum chamber:
  - helium gas @ few  $10^{-2} \text{ mbar}$
  - where the sensor is
    - almost no energy loss of the  $\alpha$ s
    - high breakdown voltage
    - good thermalisation



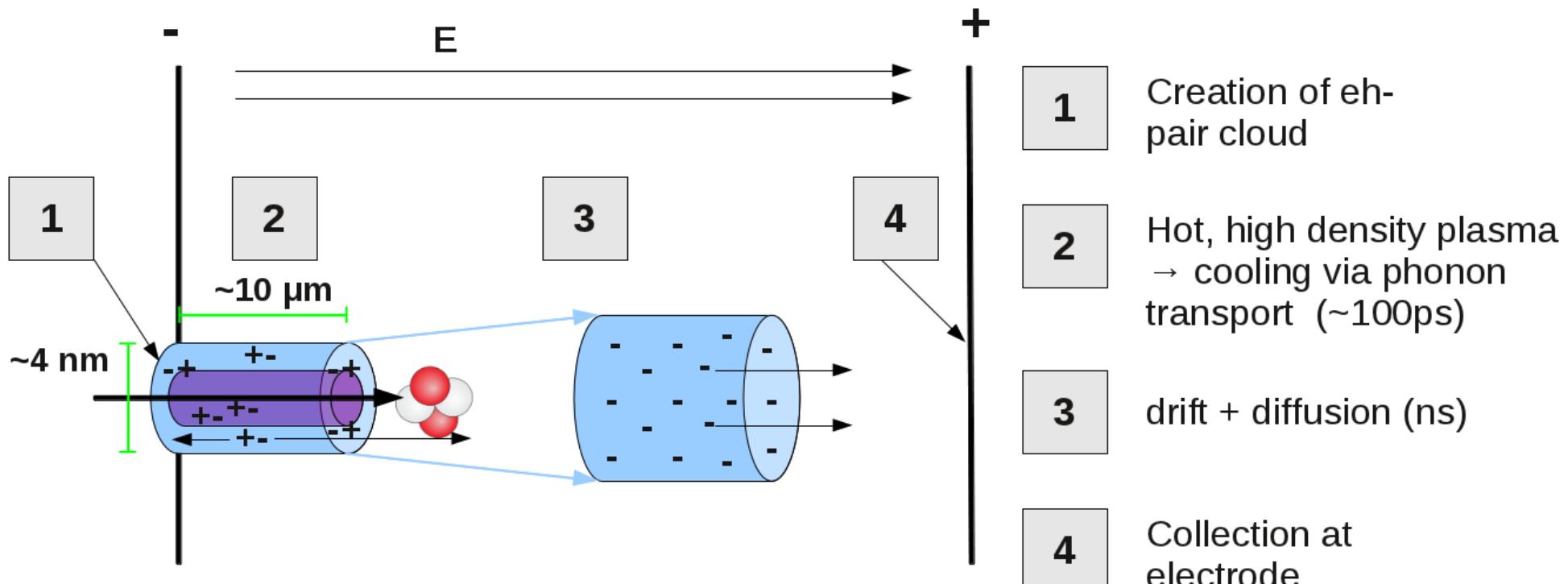
# TCT Pulses at RT: Holes



# TCT Pulses at RT: Electrons



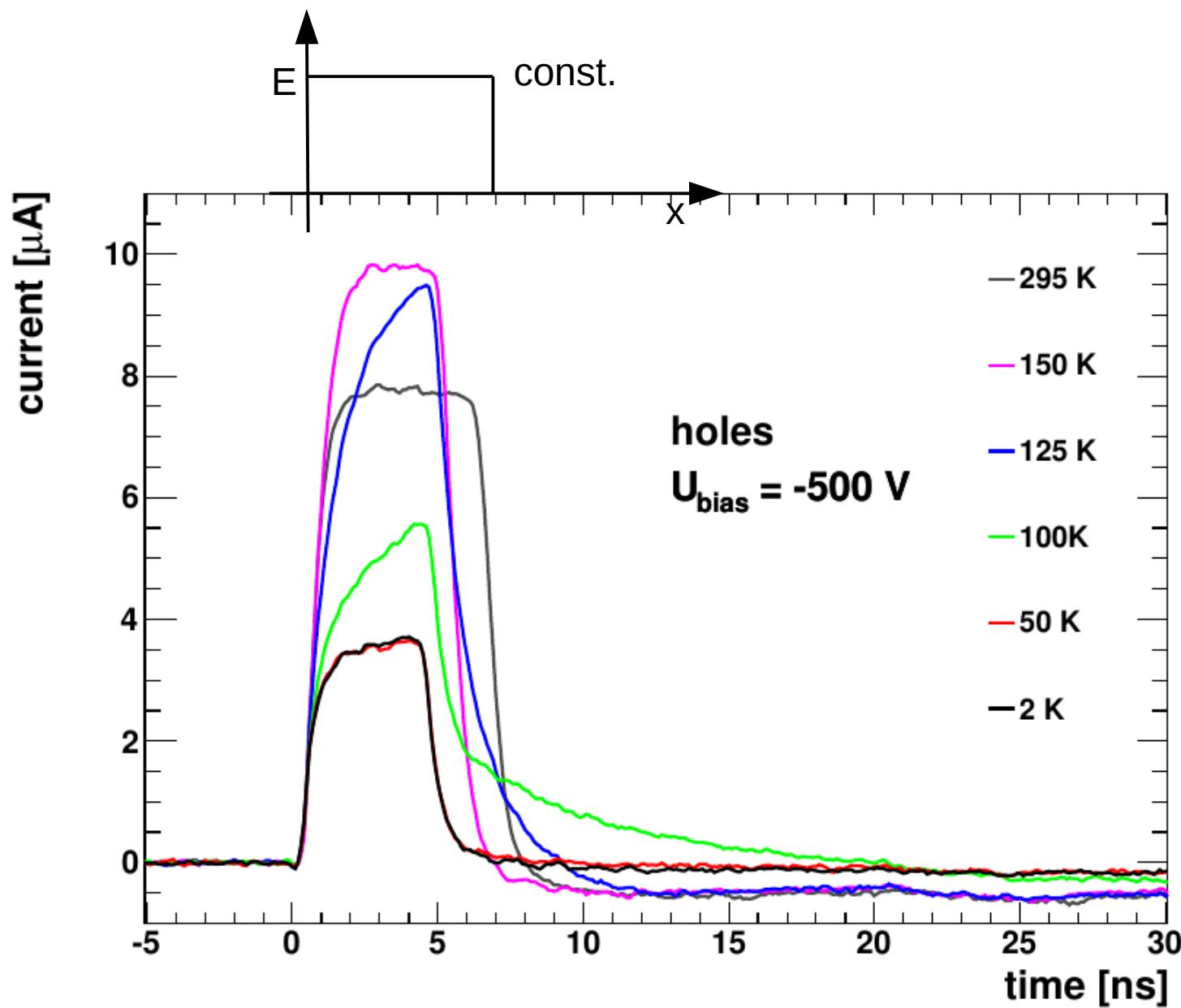
# The Plasma



- as produce **high density** charge cloud
- Outer charges **screen** inner ones only in very early phase of the plasma  
→ E-Field **penetrates plasma after 100 ps or so**
- E-Field influences “**inside**” to “**outside**” fraction

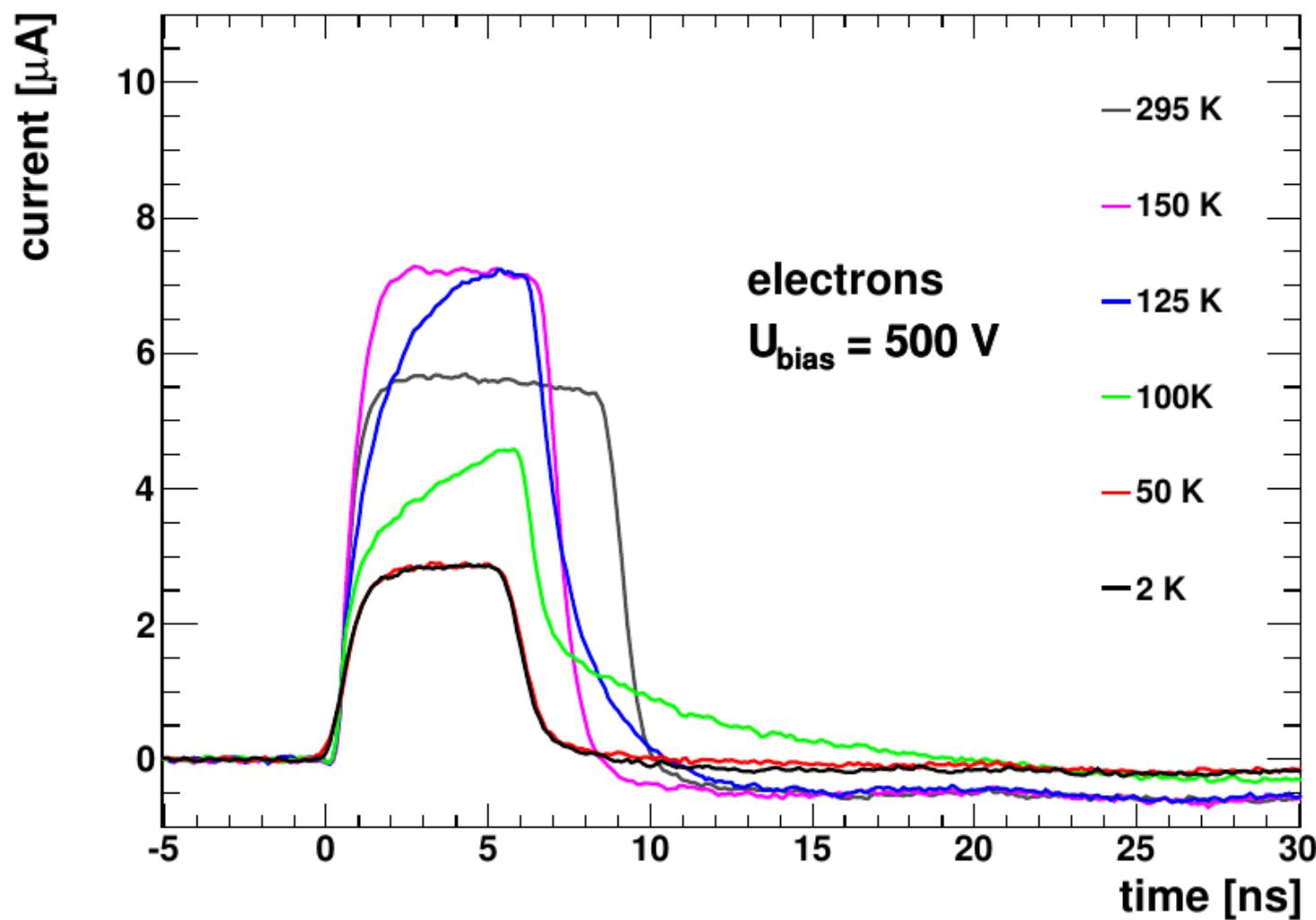
$$\rho_{cloud} \approx \frac{3.5 \cdot 10^5 \text{ pairs}}{(2 \text{ nm})^2 \pi 10 \mu m} \approx 10^{21} \text{ cm}^{-3}$$

# TCT Pulses at Various Temperatures

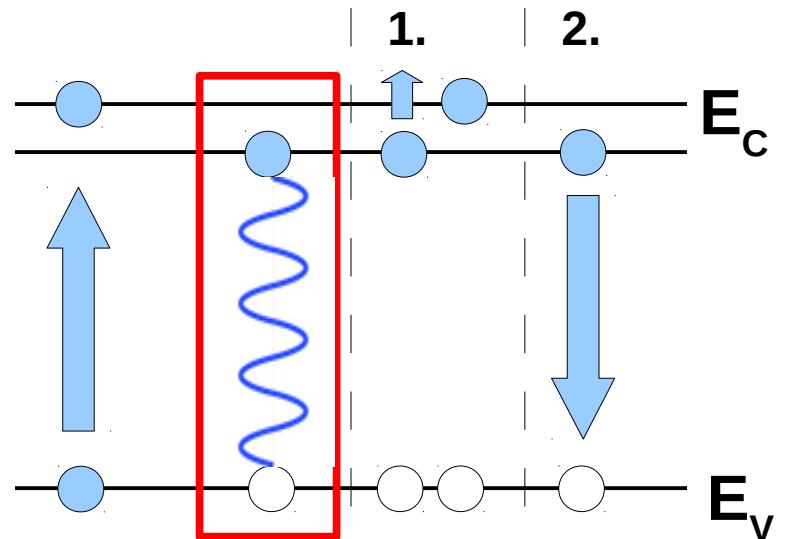
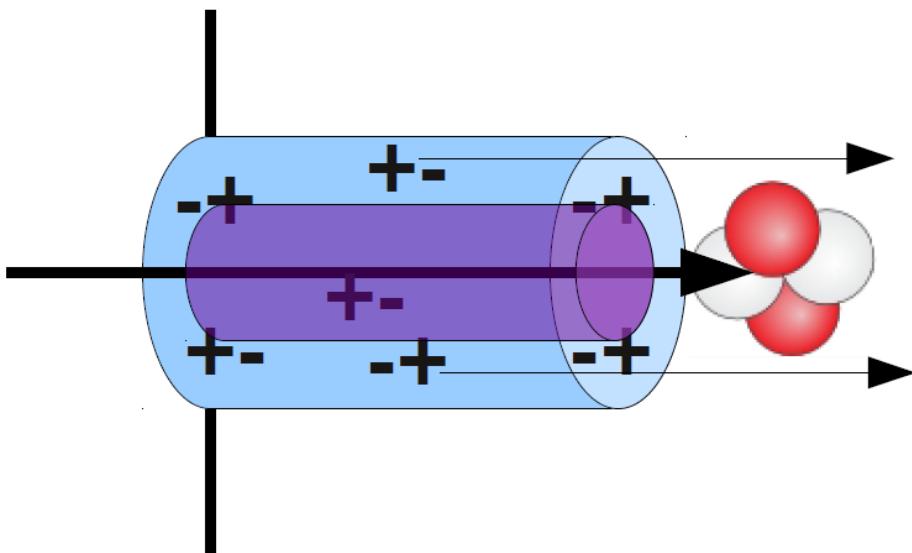


Verified for three samples, two metallisations !

# TCT Pulses at Various Temperatures



# The Plasma again

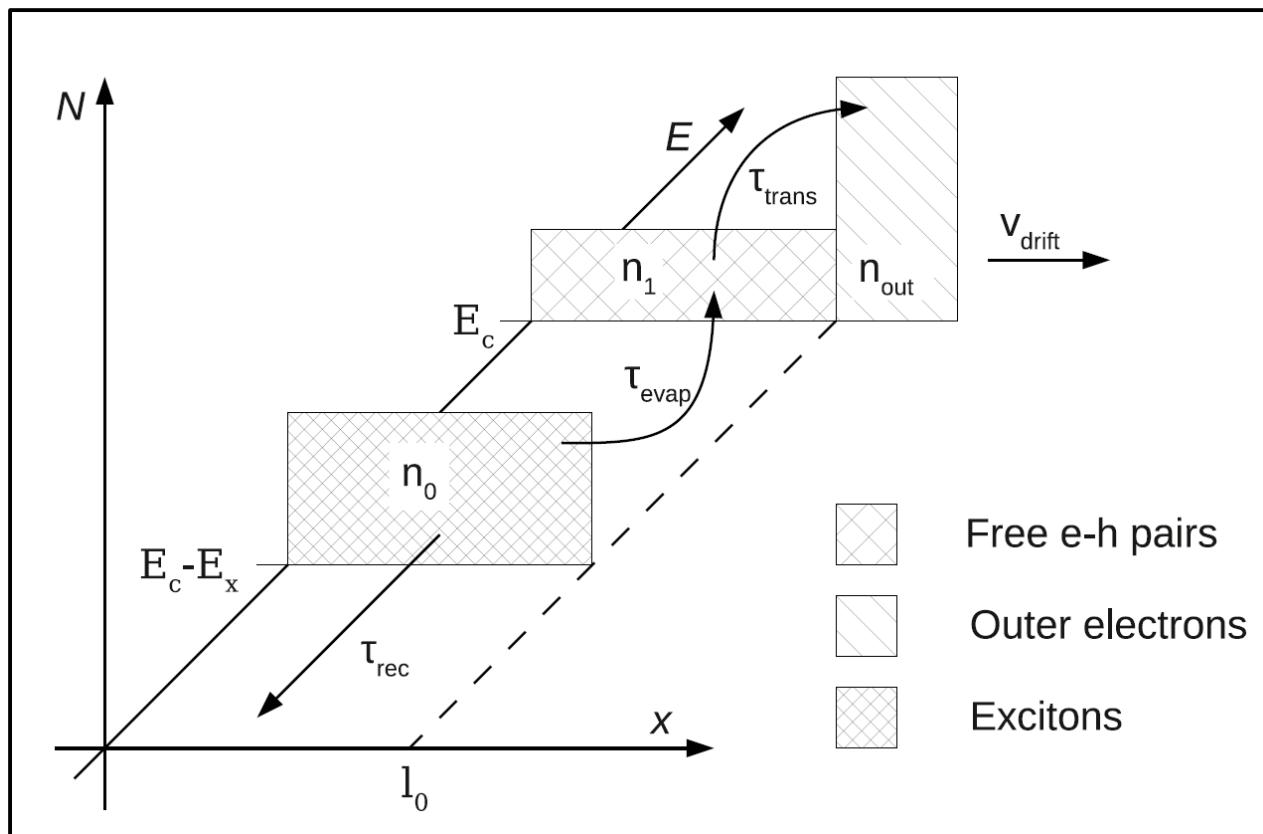


- Always collect “outer” charge
- Fraction inner to outer  
→ E-Field dependent
- “Inside”: Thermal equilibrium between excitons and free e-h pairs
- Charges are released with  $\tau_{\text{evap}} \sim \exp(E_x / kT)$   
→ charges being retained, thermal lifetime 100 ps (RT) to 100 us  
→ recombination time ( $\sim 2\text{ns}$ ) important for  $T < 150\text{ K}$

**Exciton “Decay-modes”:**  
1. dissociation  
2. recombination

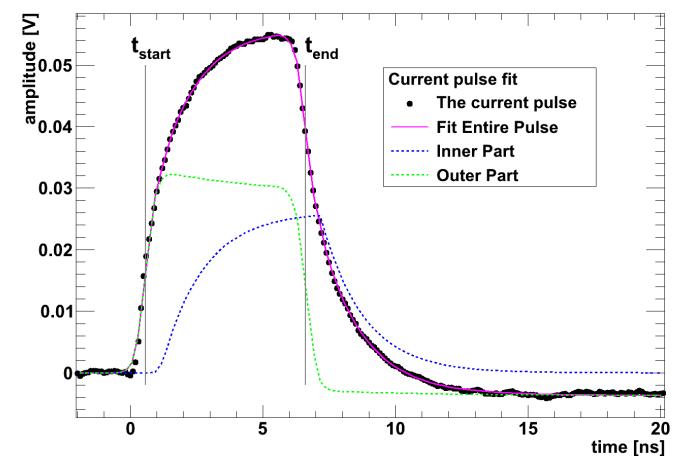
# The Model

- Creation of excitons and subsequent **ionisation** or **recombination**
- If ionised, charges are transported



Two competing processes:  $\tau_{\text{rec}}$  vs  $\tau_{\text{evap}}$

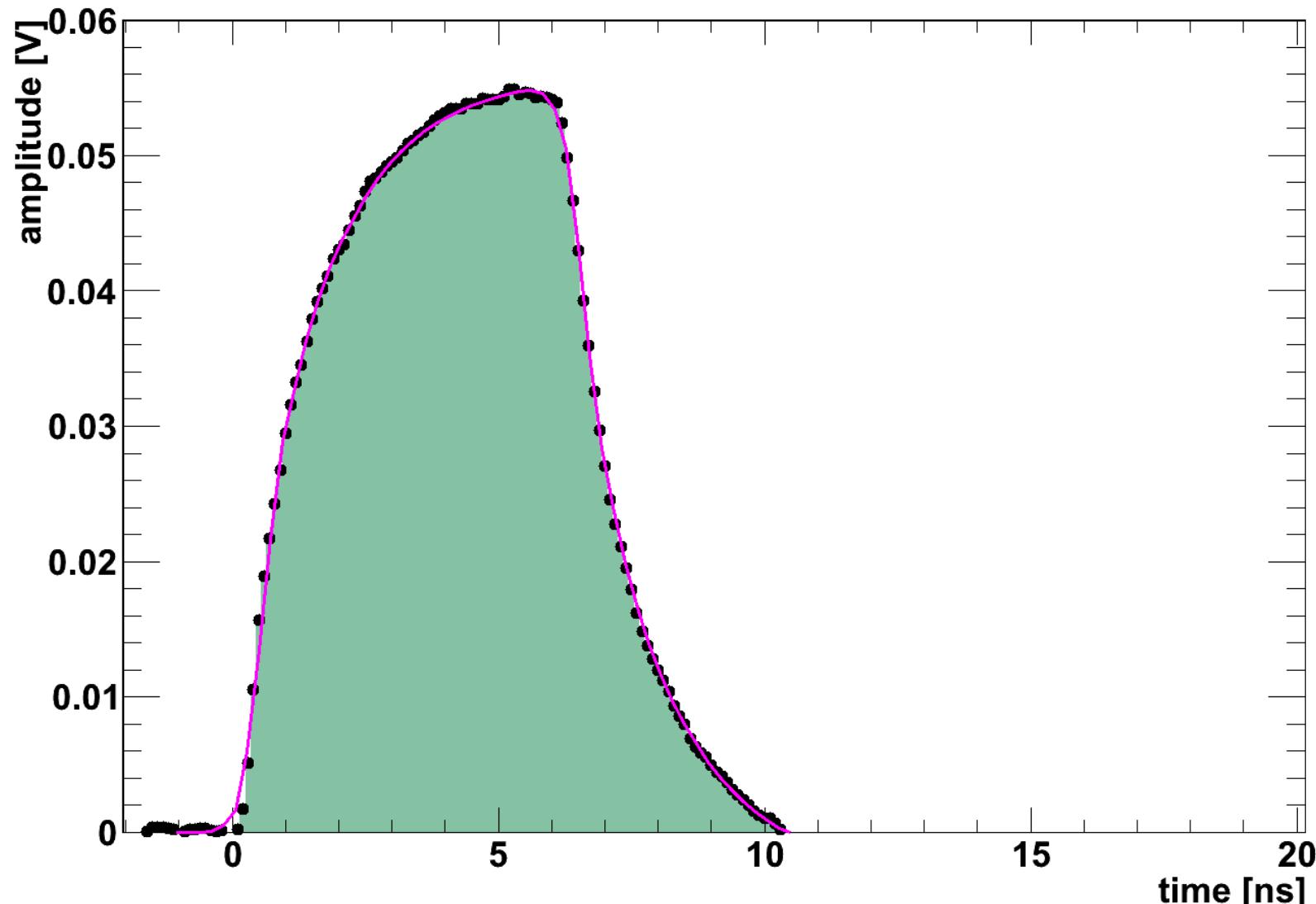
$$\frac{1}{\tau} = \frac{1}{\tau_{\text{rec}}} + \frac{1}{\tau_{\text{evap}}}$$



$$n_{\text{out}}(t) = n_{\text{out}}(0) + n_0(0) \frac{\tau}{\tau_{\text{evap}}} (1 - \exp(-t/\tau))$$

# Analysis of TCT Pulses

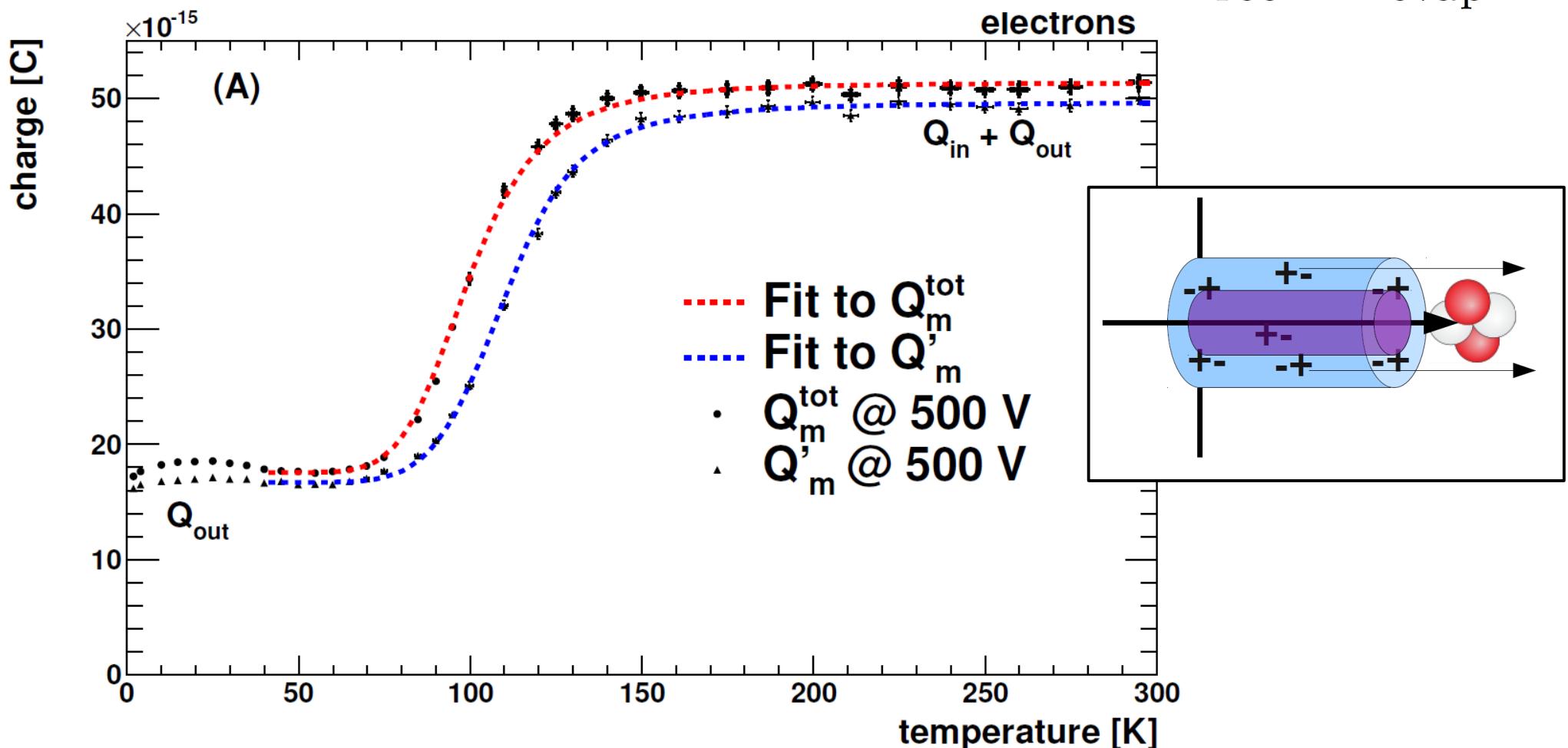
- Calculate integral of pulse:



# Integrated Charge

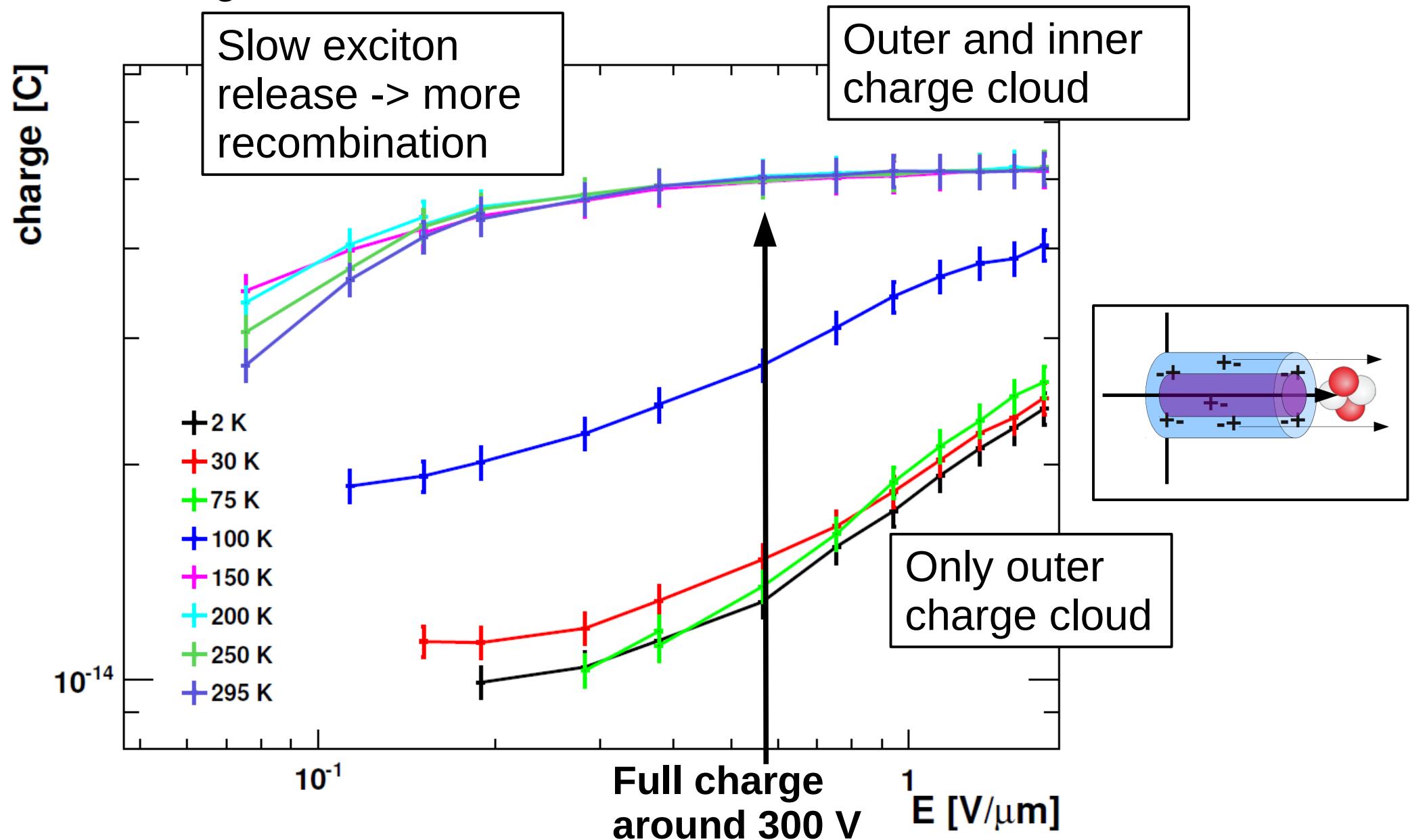
- Integrate  $n_{\text{out}}(t)$  over  $t$ , fit model to the integral of pulses over  $T$ :

$$Q_{\text{out}} + \frac{\tau_{\text{rec}}}{\tau_{\text{rec}} + \tau_{\text{evap}}} Q_{\text{in}}$$



# Integrated Charge

- Charge as function of E-field:

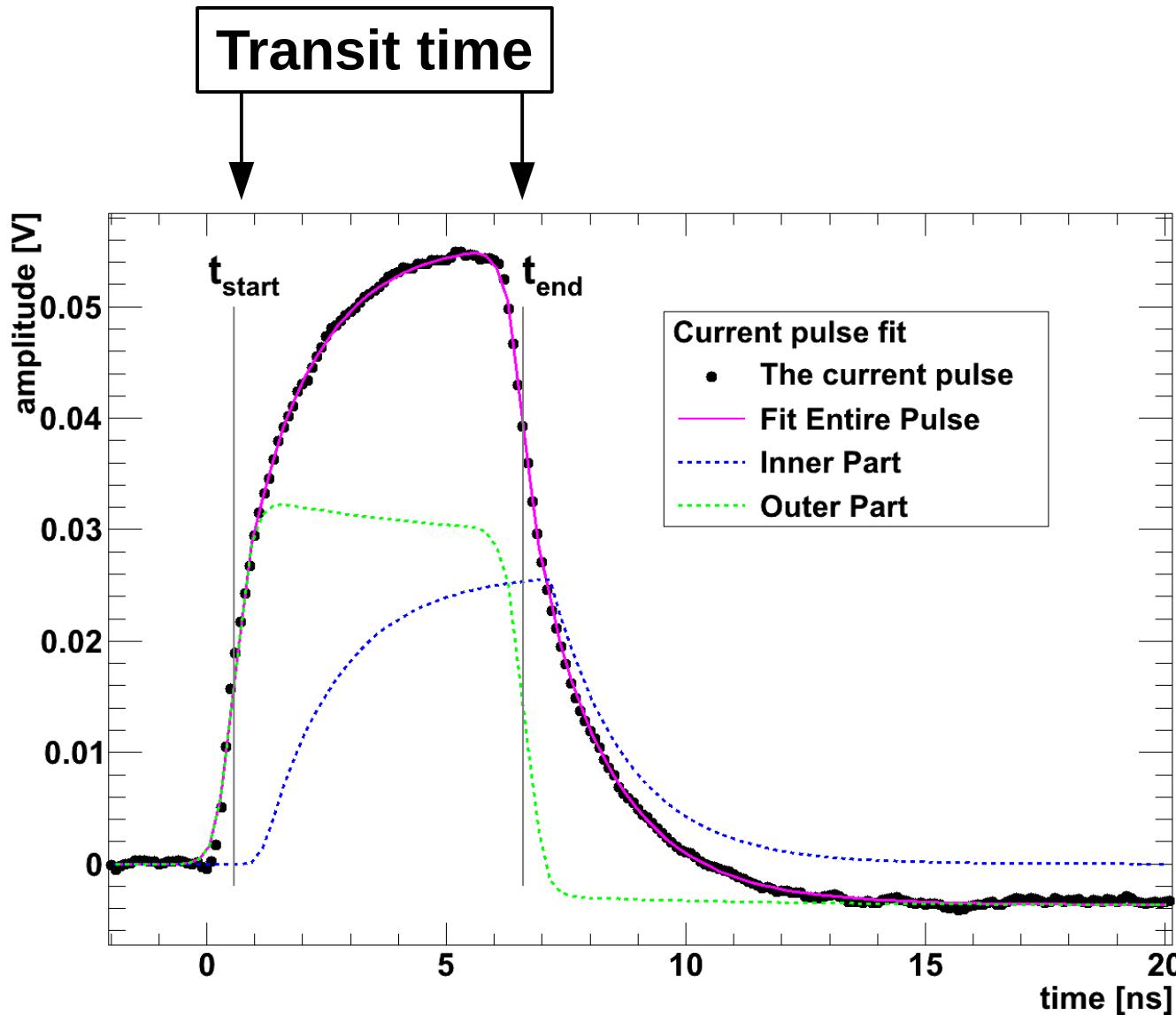


# Integrated Charge

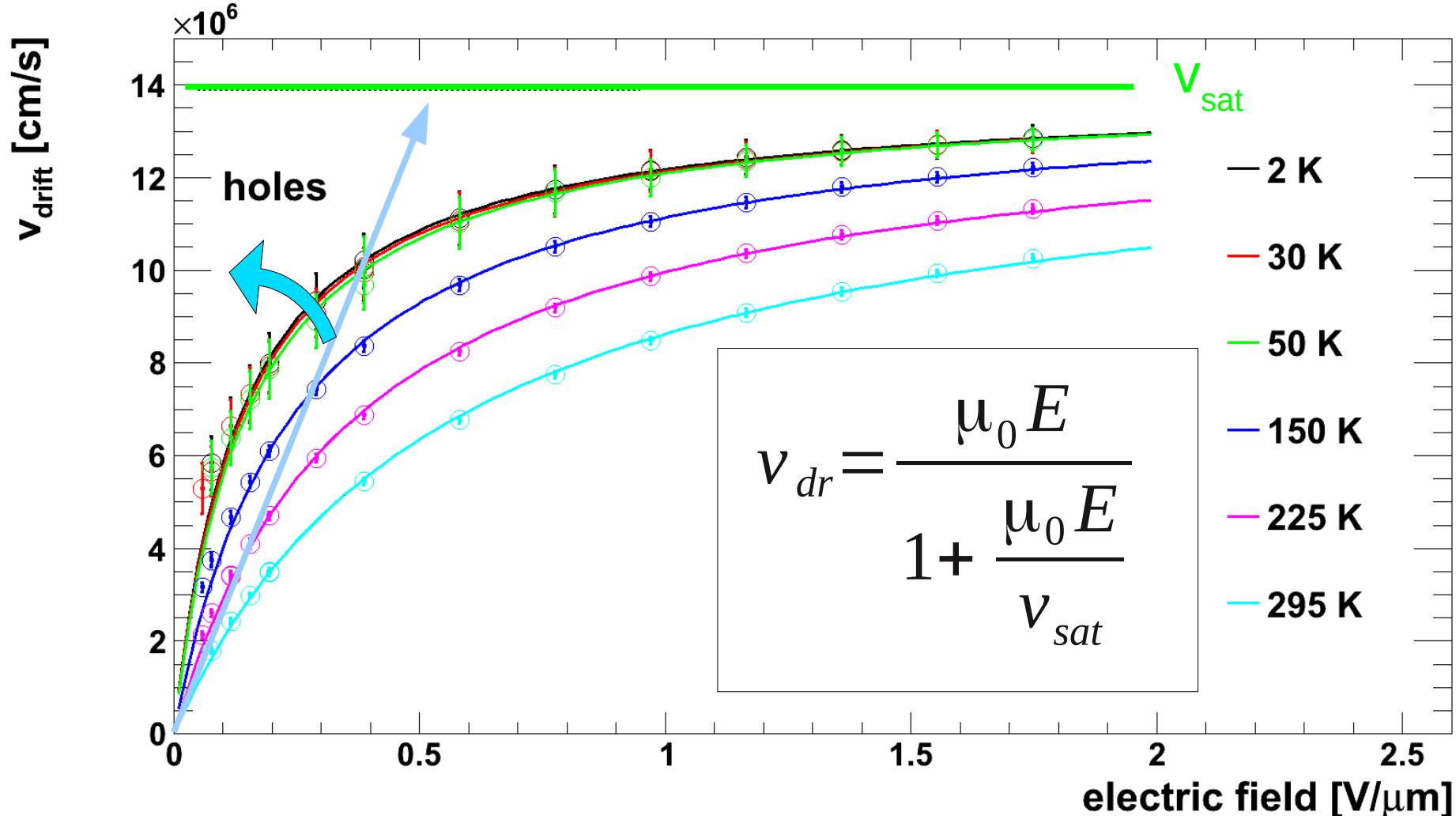
- For  $T > 150$  K:
  - Full charge at 300 V
  - Recombination leads to charge loss for  $|U| < 300V$ , why?
  - NO TRAPPING DURING! See flat pulses.  
Q-E behaviour not dominated by CCD ( $\tau_{\text{trap}} \approx 1\mu\text{s}$ )
  - This is charge recombination in the early plasma
  - > Dominant charge loss mechanism in high-quality, unirradiated scCVD diamond
- For  $T < 80$  K:
  - ALL inner charges recombine
  - only measure outer charges
  - See how “outer” part grows with field strength  
-> sensitive to physical processes within first  $\sim 100$  ps of the plasma!

# Analysis of TCT Pulses

- Fit rectangular  $(2 - \text{Erfc}(t)) - (2 - \text{Erfc}(t-t'))$   
+ exponential  $(1 - \exp(-t/\tau)) - (1 - \exp(-(t-t')/\tau))$

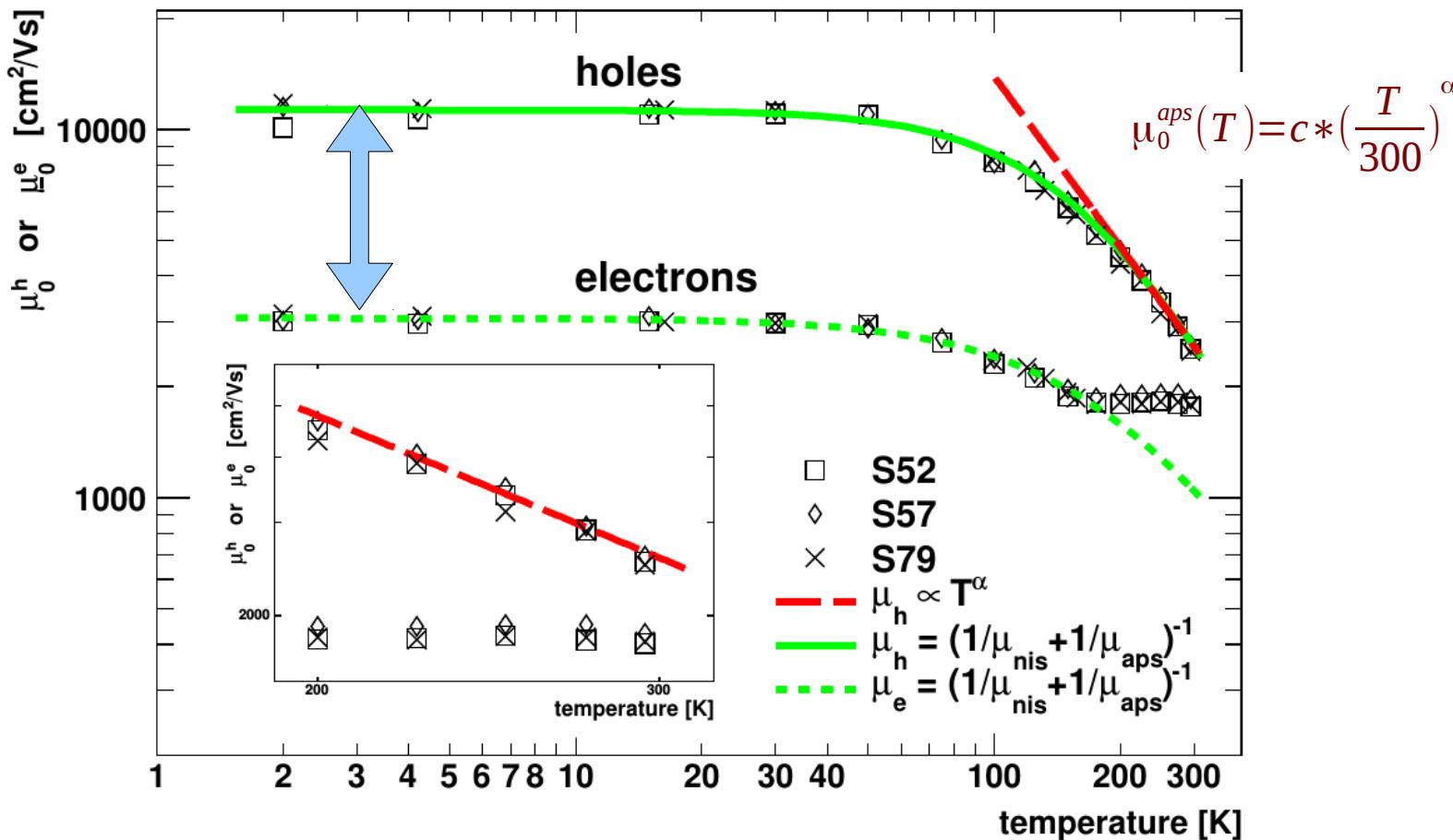


# Hole Drift Velocity



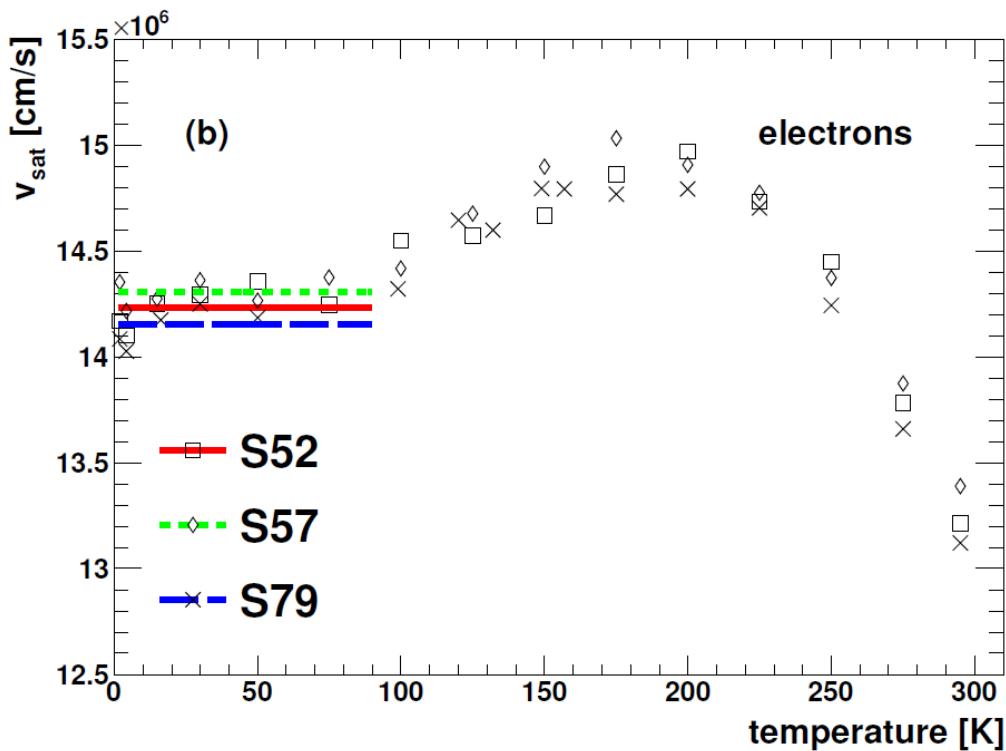
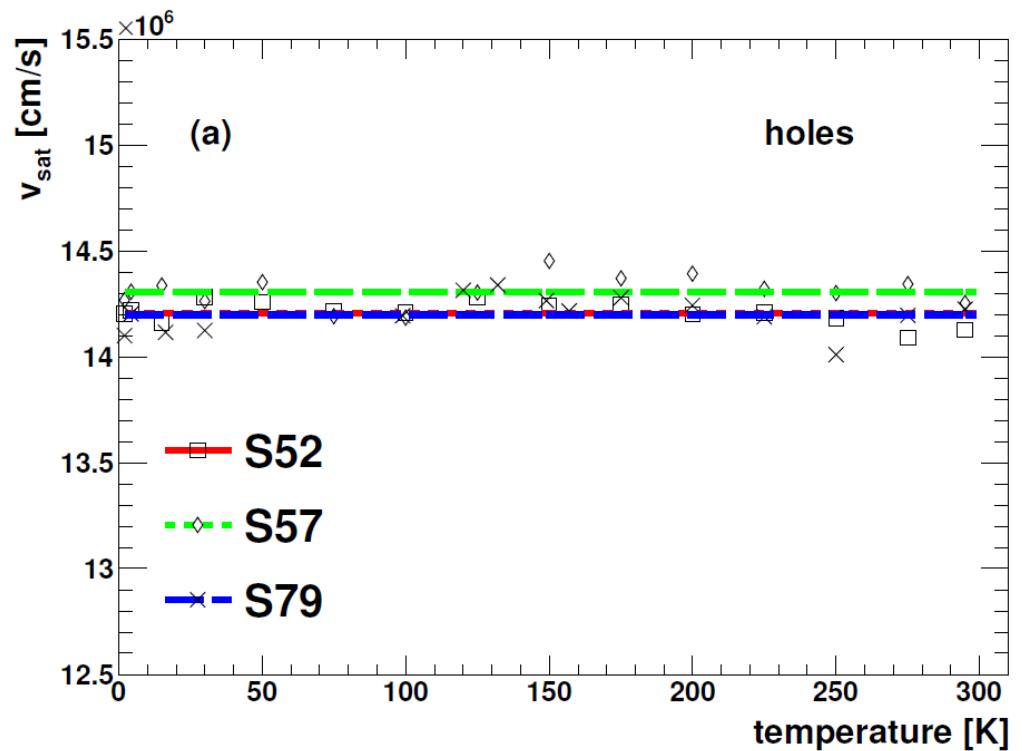
- $\mu_h$  increases with decreasing T down to 2 K
- $v_{sat}$  ~ constant with temperature @ 14e6 cm/s

# Mobility vs. Temperature



- Holes: For  $300\text{ K} < T < 200\text{ K}$  the mobility increases as acoustic phonon scattering decreases:  $\alpha = -1.49$ ,  $\mu_0 = 2400\text{ cm}^2/\text{Vs}$
- 50 K to 10 K: mobility stays  $\sim$ const.  
-> scattering dominated by neutral impurity scattering

# Saturation velocity



- Holes:  $v_{\text{sat}} \sim \text{constant with } T$ ,
- Electrons:  $v_{\text{sat}}$  has certain dependence,

low T limit: 14e6 cm/s

low T limit: 14e6 cm/s

# Mobility vs. Temperature

Journal of Applied Physics / Volume 113 / Issue 17 / ARTICLES / Electronic Structure and Transport

J. Appl. Phys. **113**, 173706 (2013); <http://dx.doi.org/10.1063/1.4802679> (9 pages)



FULL-TEX



## Temperature dependence of charge carrier mobility in single-crystal chemical vapour deposition diamond

Hendrik Jansen<sup>1</sup>, Daniel Dobos<sup>1</sup>, Thomas Eisel<sup>1</sup>, Heinz Pernegger<sup>1</sup>, Vladimir Eremin<sup>2</sup>, and Norbert Wermes<sup>3</sup>

<sup>1</sup>CERN, CH-1211 Genve 23, Switzerland

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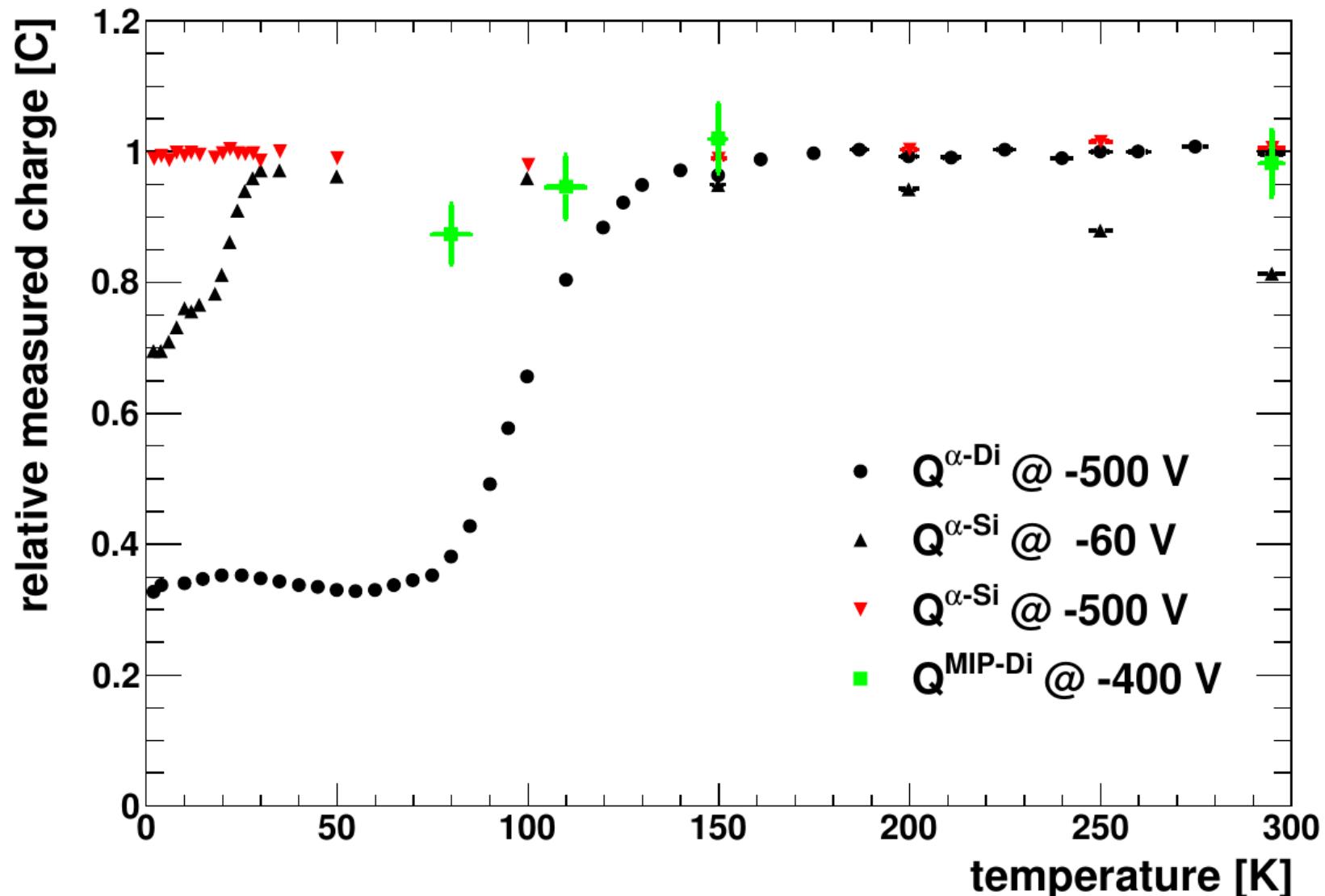
<sup>3</sup>University of Bonn, Nussallee 12, 53115 Bonn, Germany

[View Map](#)

(Received 25 February 2013; accepted 8 April 2013; published online 6 May 2013)

# Comparison with Silicon

- $\tau_{\text{evap}} \sim \exp(E_x / kT) !$



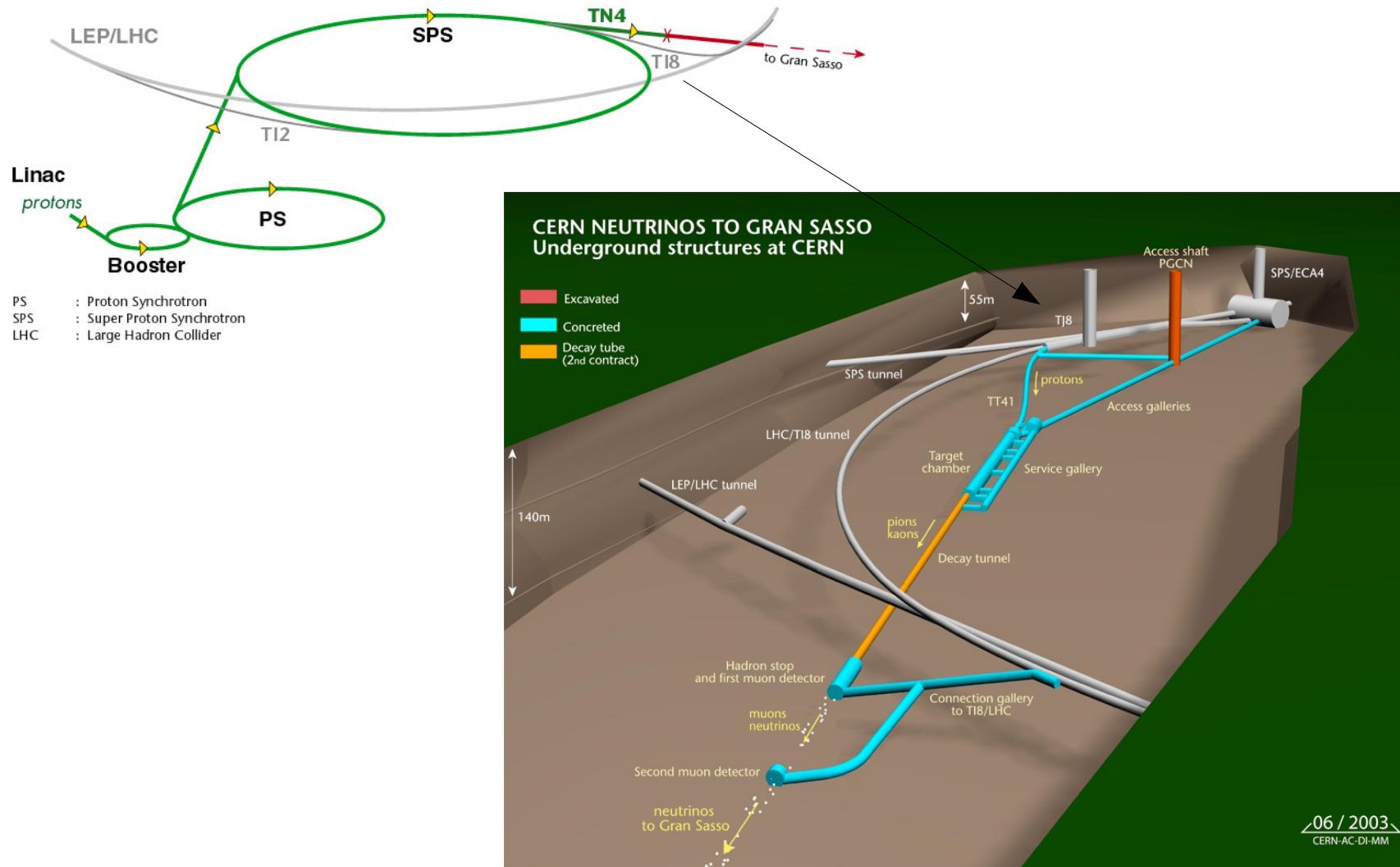
# Conclusion from cold measurements

- Diamond detects radiation even at 2 K
- Charge collection is function of the temperature:
  - full charge at  $T > 150$  K
  - collection of 'outer charge only' for  $T \leq 75$  K
- Creation of excitons by alpha particles in the diamond
  - maybe implications even at room temperature?  
Surface recombination, etc.
- First diamond mobility measurement down to 2 K
- Further reading:
  - H. Pernegger et al., JAP 97 (7) (2005) 073704.
  - C. Erginsoy, Phys. Rev. 79 (1950) 1013–1014
  - H. Jansen et al., Physics Procedia 37 (2012) 2005
  - H. Jansen et al., JAP 113 173706 (2013)

# Usage of CVD Diamond at CNGS

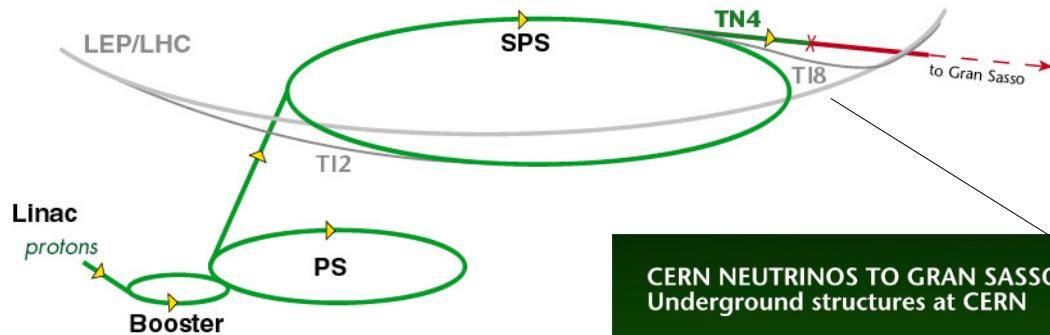
# Diamond Detectors at CNGS

- Where?



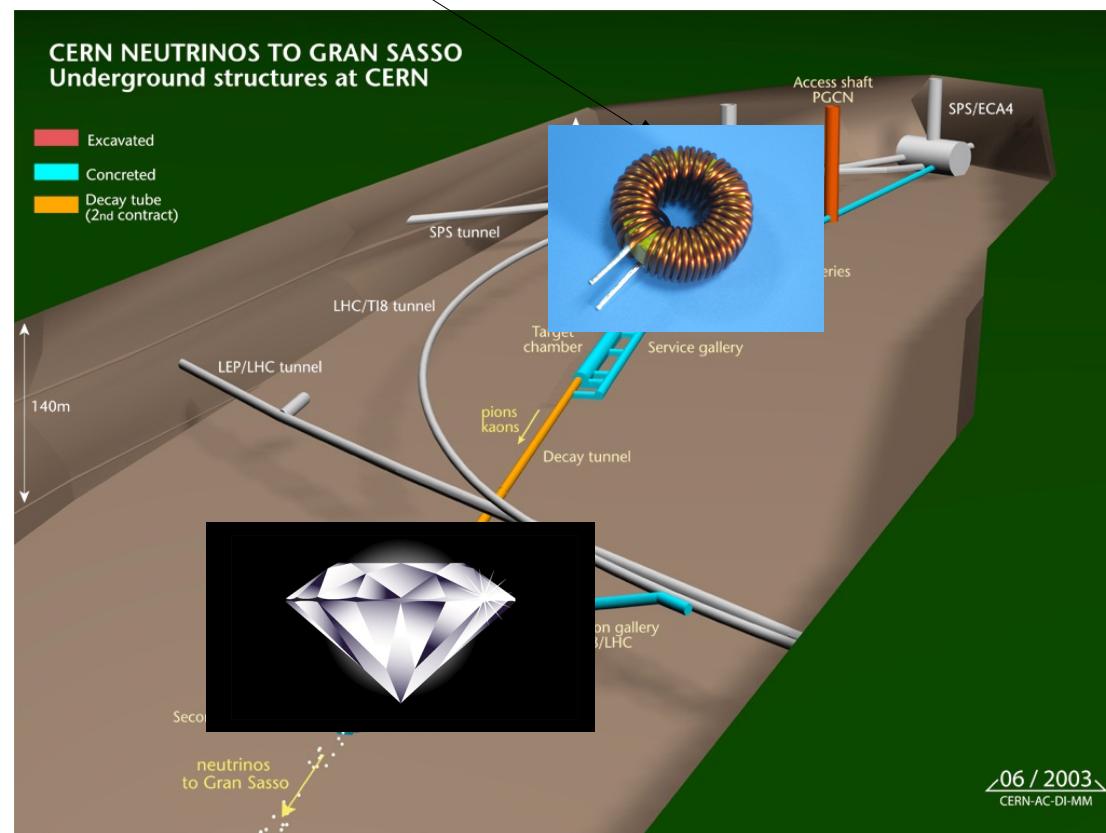
# Diamond Detectors at CNGS

- Where?



PS : Proton Synchrotron  
SPS : Super Proton Synchrotron  
LHC : Large Hadron Collider

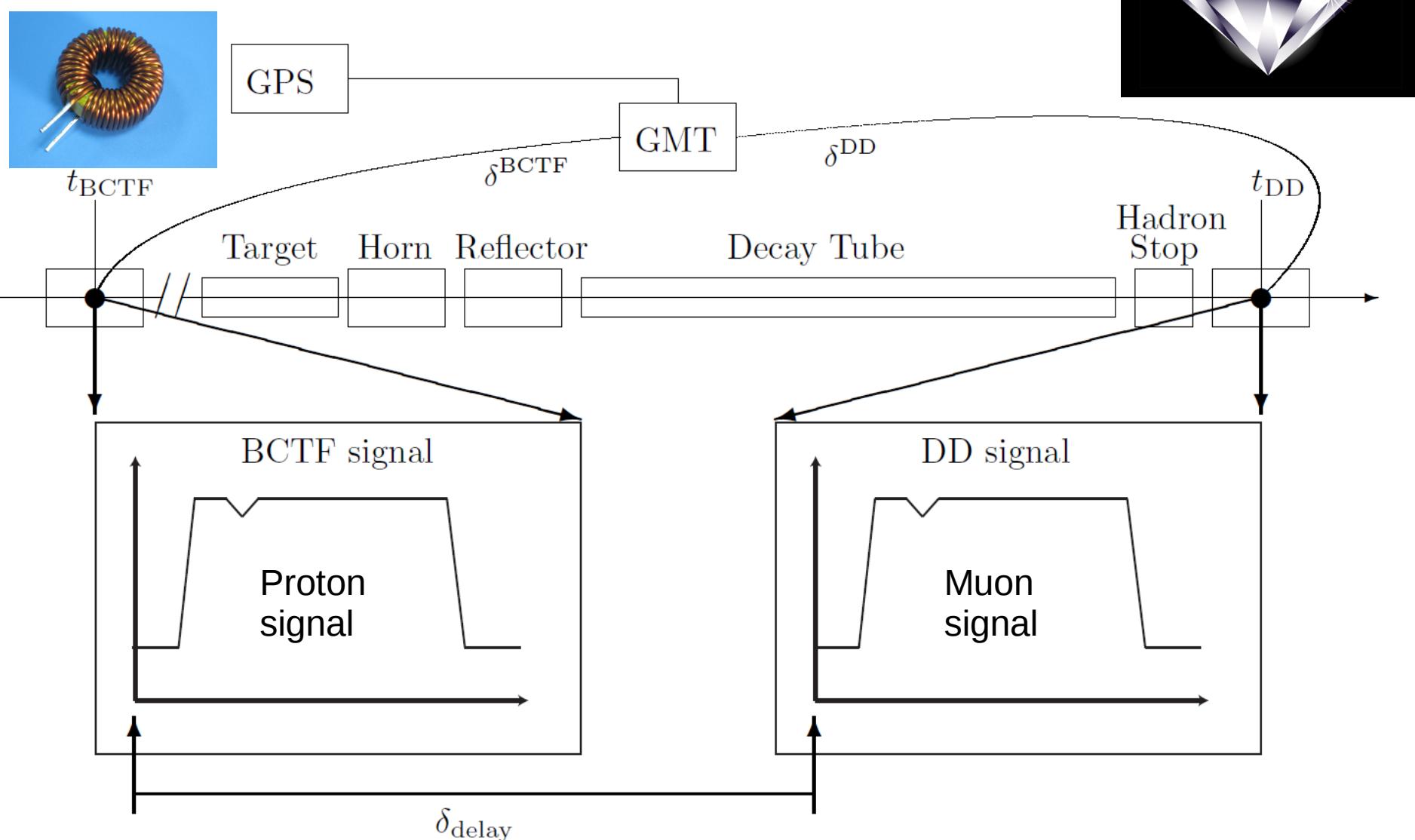
## Beam Current Transformer



## Diamond Detectors

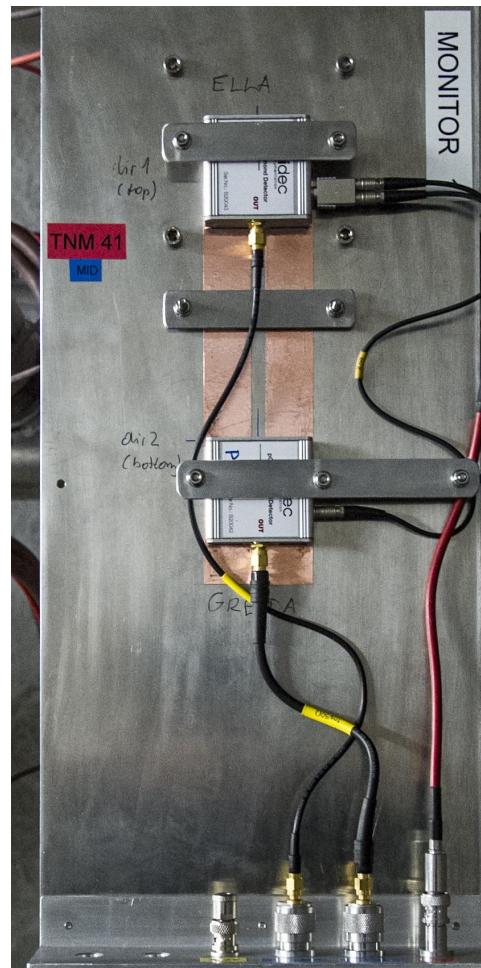
# Diamond Detectors at CNGS

- Connection from detectors to Control Room:



# Diamond Detectors at CNGS

- TNM 41:

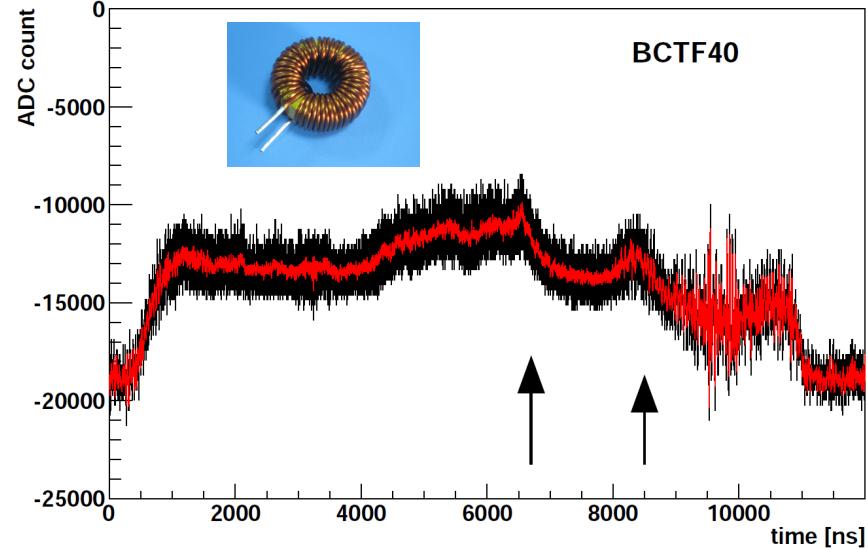
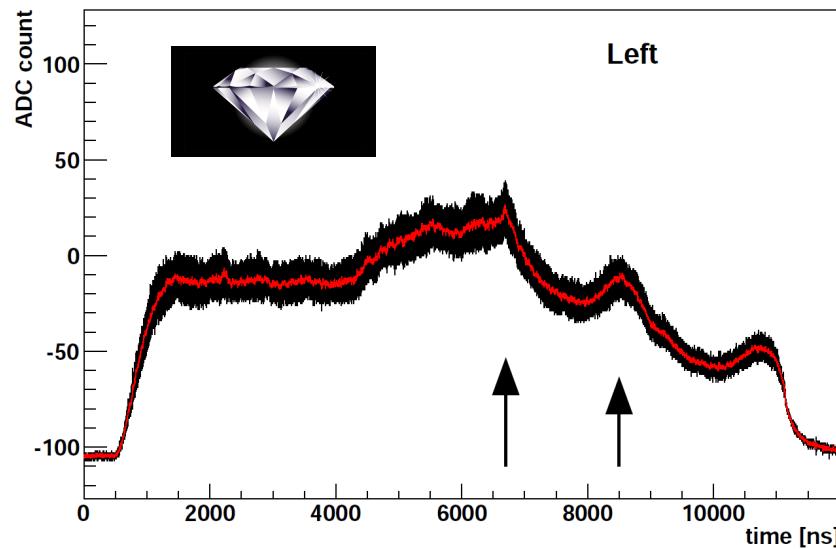
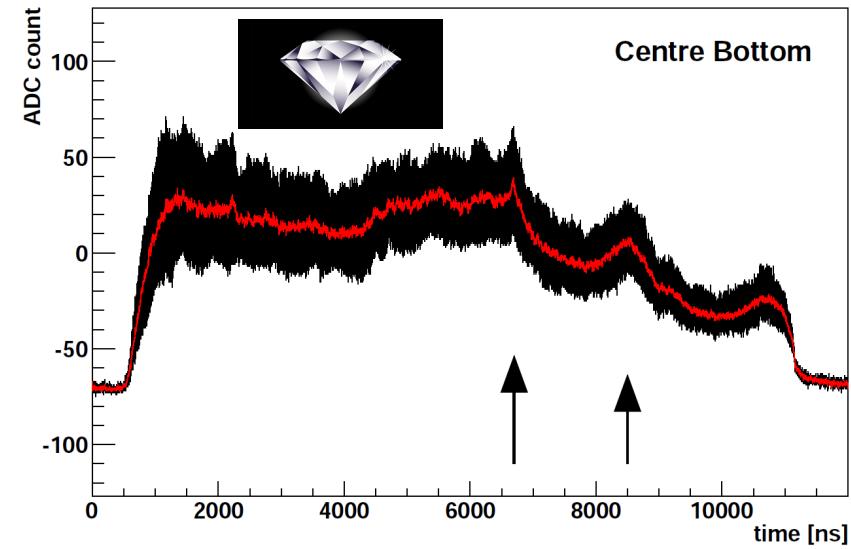
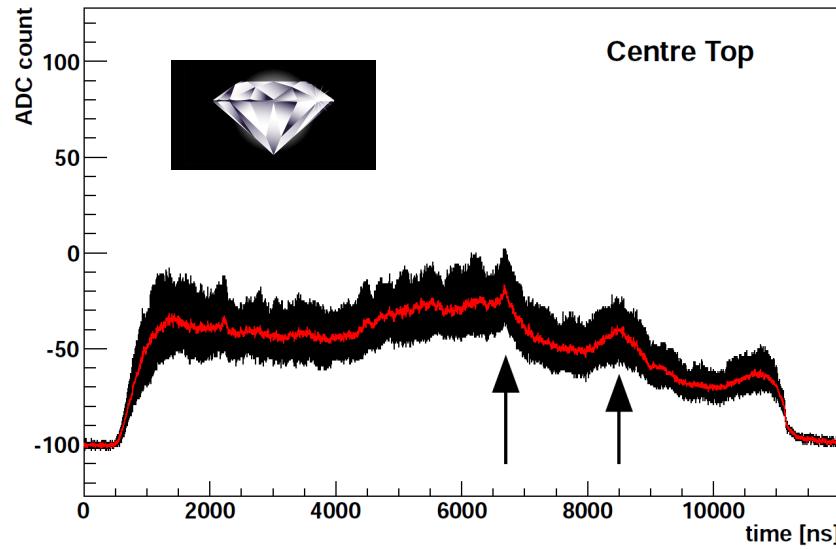


Poly-crystalline CVD diamond detectors:

- 8mm x 8mm
- 500 um thickness
- no amplification, direct scope read-out, 50 Ohm

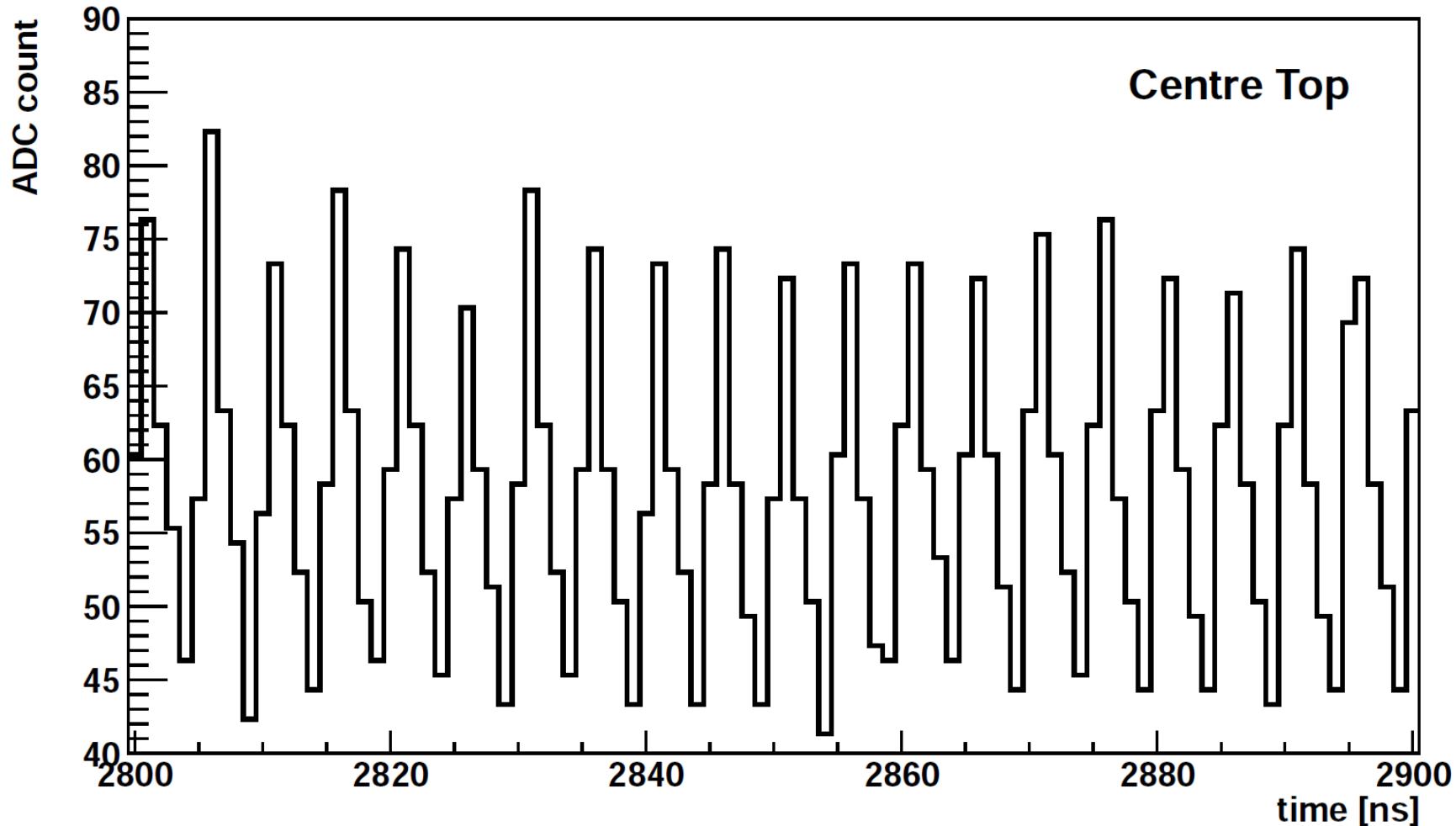
# Intensity/Profile Measurements

- From Single Extraction:



# Intensity/Profile Measurements

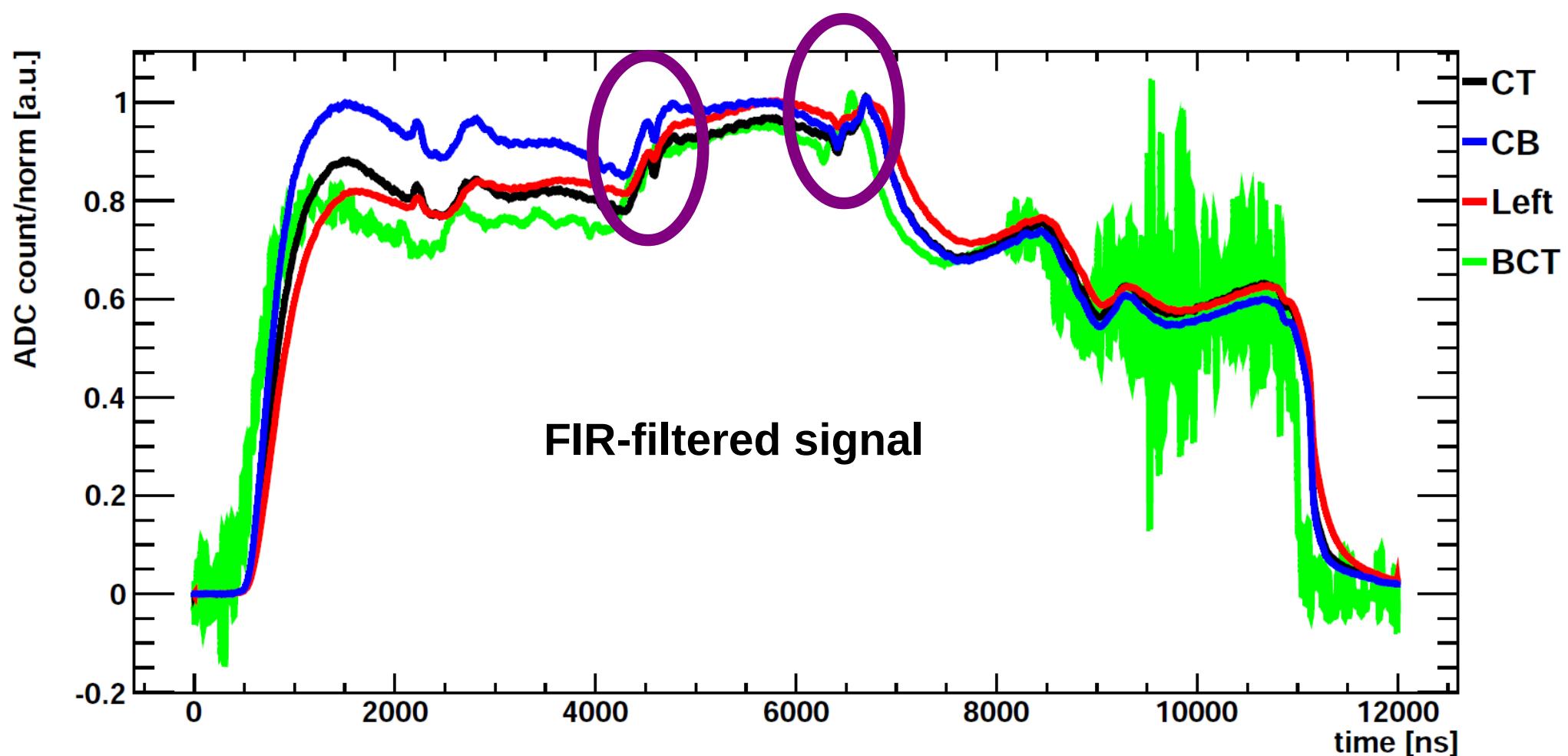
Zoom:



→ 5 ns bunch structure from SPS RF evident

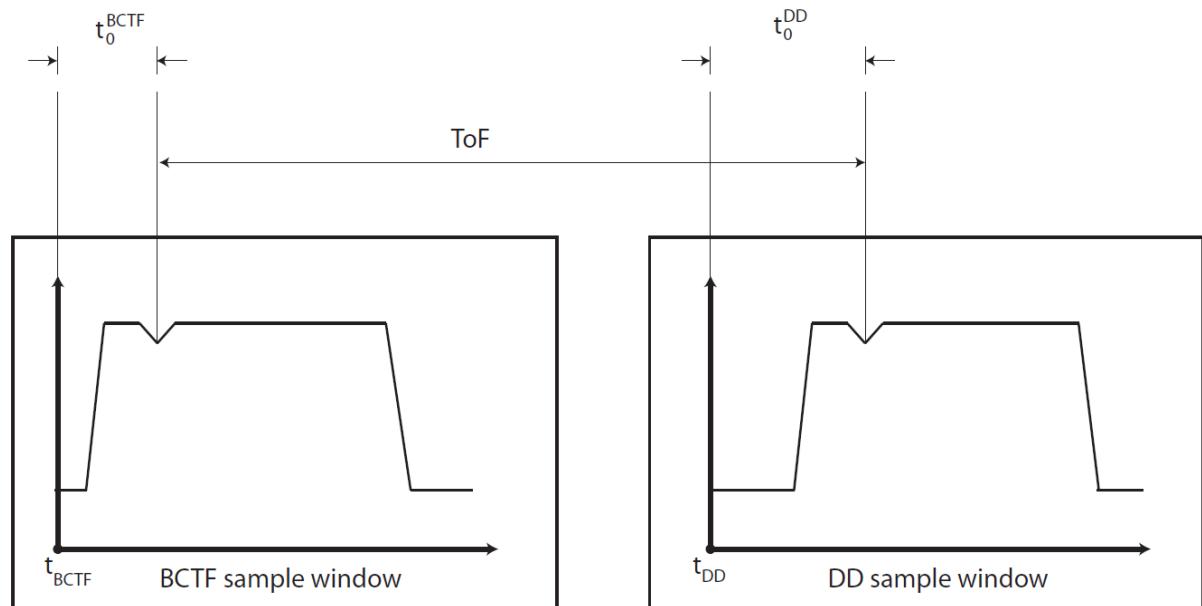
# Compare p- to $\mu$ -beam structure

- Signals from BCT and DDs agree nicely
- Even local features show up in both DDs and BCTF



# Verification of CNGS Timing

- Measure delays between read-out windows
- Measure timeline offset of signals within read-out window



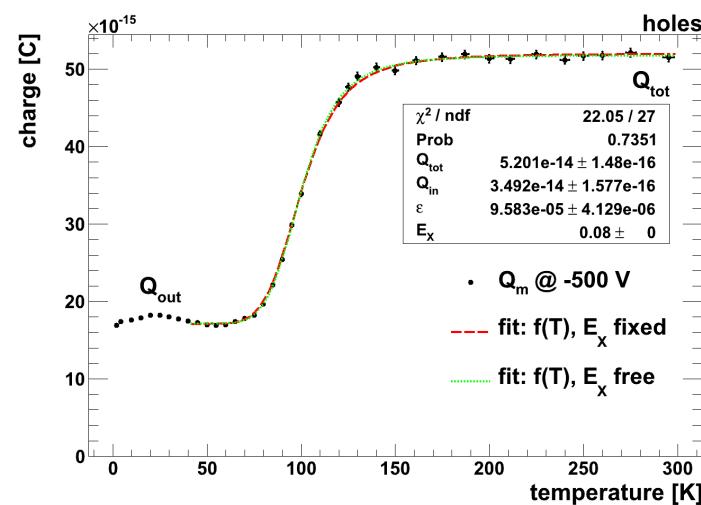
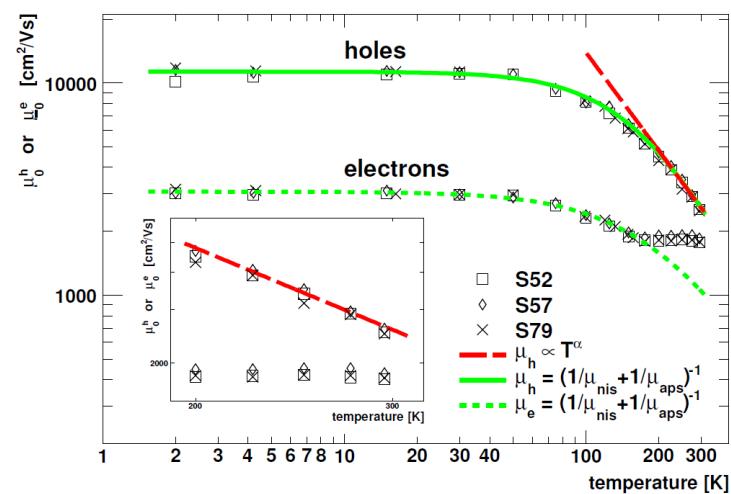
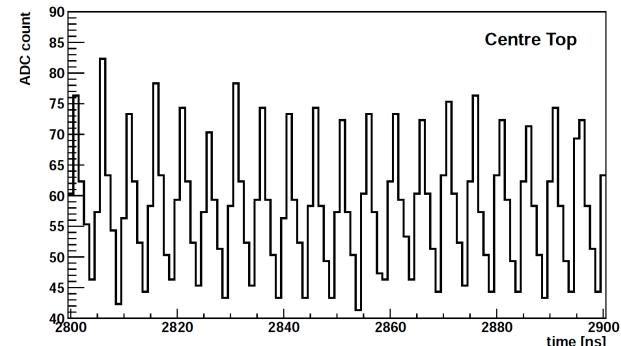
→ Calculate the Time-of-Flight:

$$\text{ToF}_{\hat{\Theta}_x} = (6205.4 \pm 3.6) \text{ ns}$$

$$\text{ToF}_{\text{nom}} = \frac{1859.95 \text{ m}}{v} = (6205.3 \pm 2.5) \text{ ns}$$

# Conclusion

- Diamond is fast:
  - Measure individual bunches up to  $\sim 200$  MHz.
  - Usable as intensity monitor.
- Diamond is usable as cryogenic BLM, just keep in mind T-dependence of mobility and charge yield



Thanks go to:

Erich Griesmayer

Jaako Haerkkoenen

Michael Pomorski

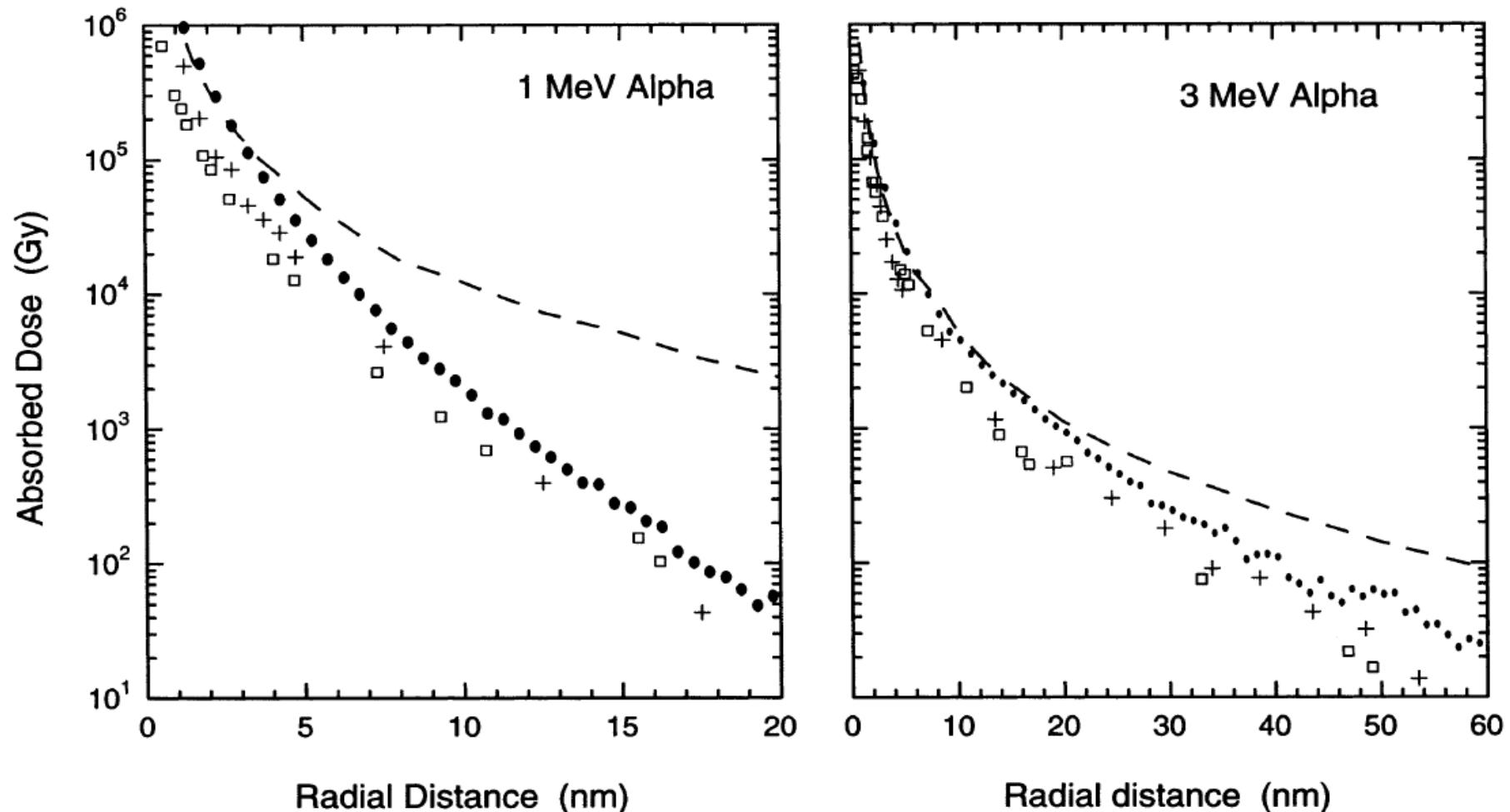
Thomas Eisel and the [CryoLab@CERN](#)

Contact:

[Hendrik.jansen@desy.ch](mailto:Hendrik.jansen@desy.ch)

# Radial Track Size

- Alpha-particles in WATER:  
avg. energy of secondary electrons is 60 eV

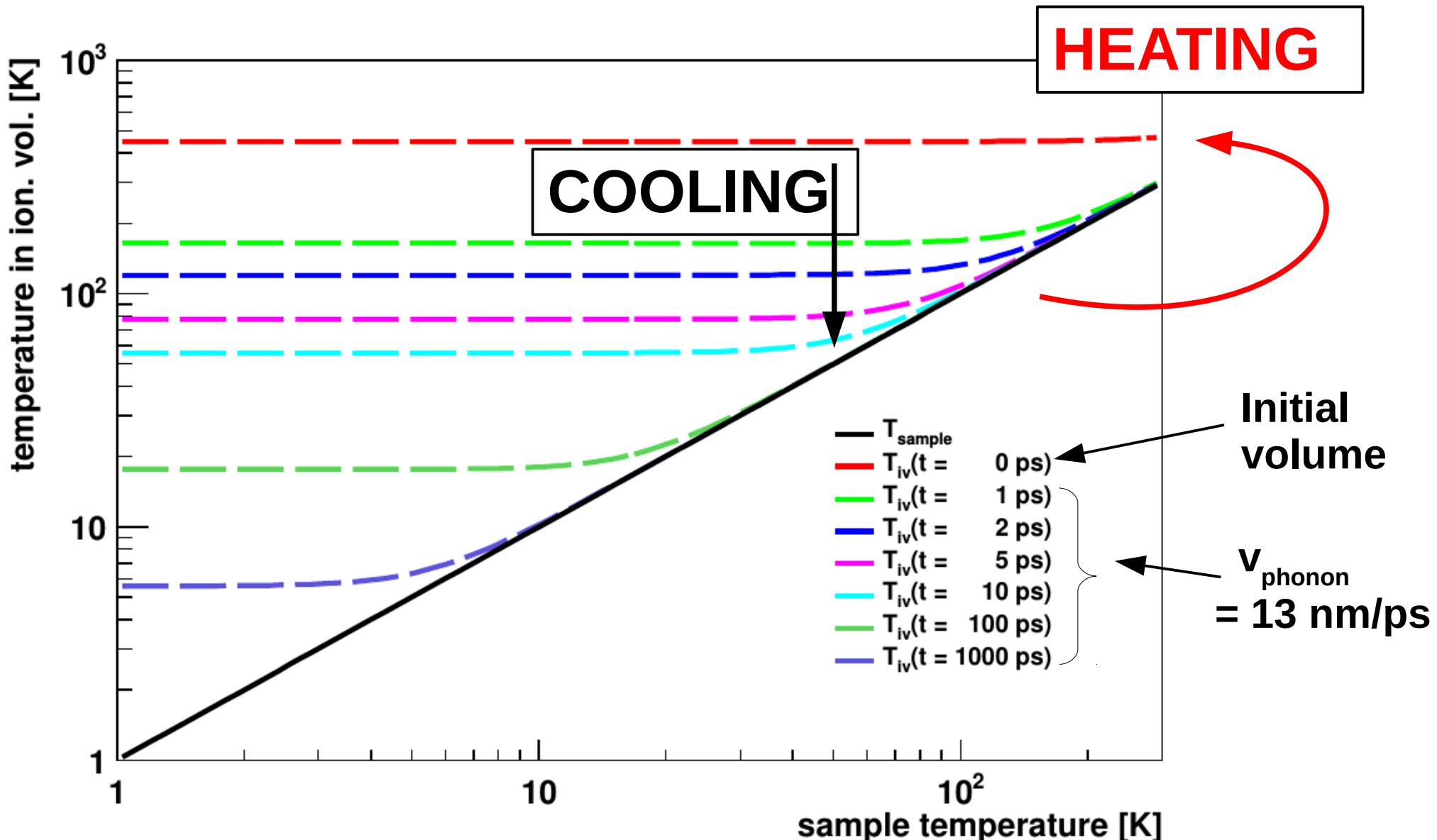


J Phys Chem B, 106 42 (2002) 11055

# The Plasma

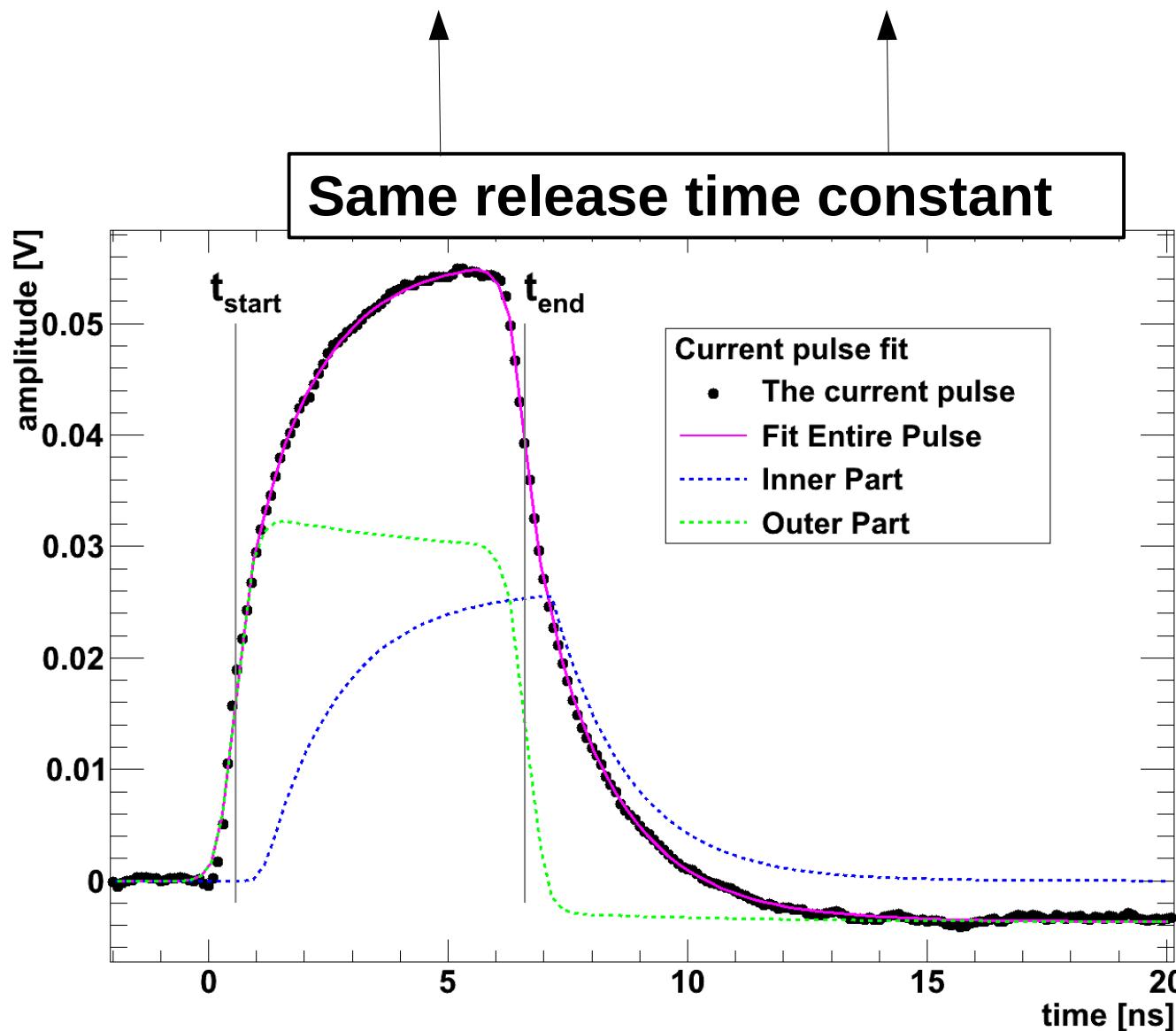
- The alpha-particle **heats** the ionisation volume!

$$c_v \sim T^3$$



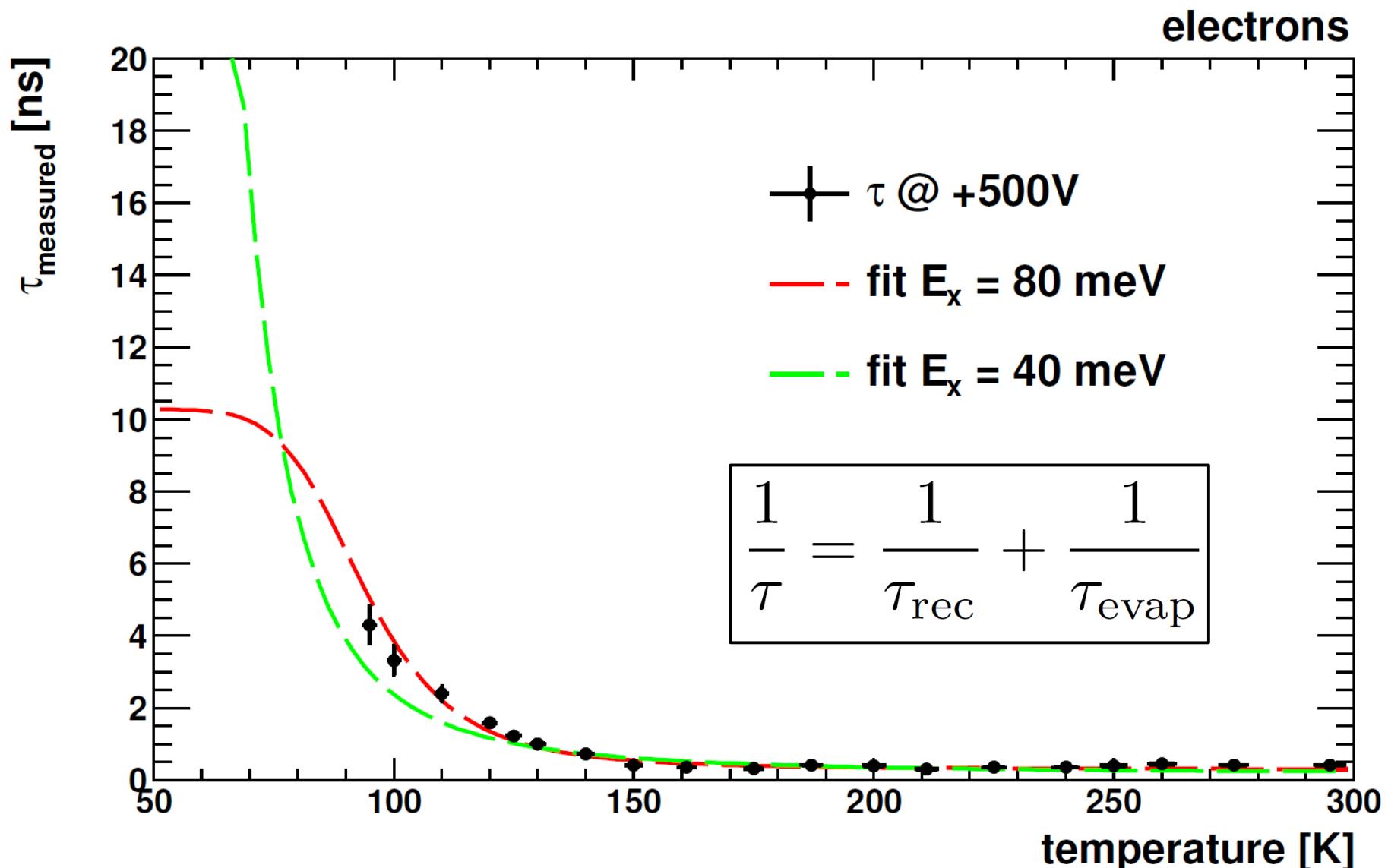
# Analysis of TCT Pulses

- Fit rectangular  $(2 - \text{Erfc}(t)) - (2 - \text{Erfc}(t-t'))$   
+ exponential  $(1 - \exp(-t/\tau)) - (1 - \exp(-(t-t')/\tau))$



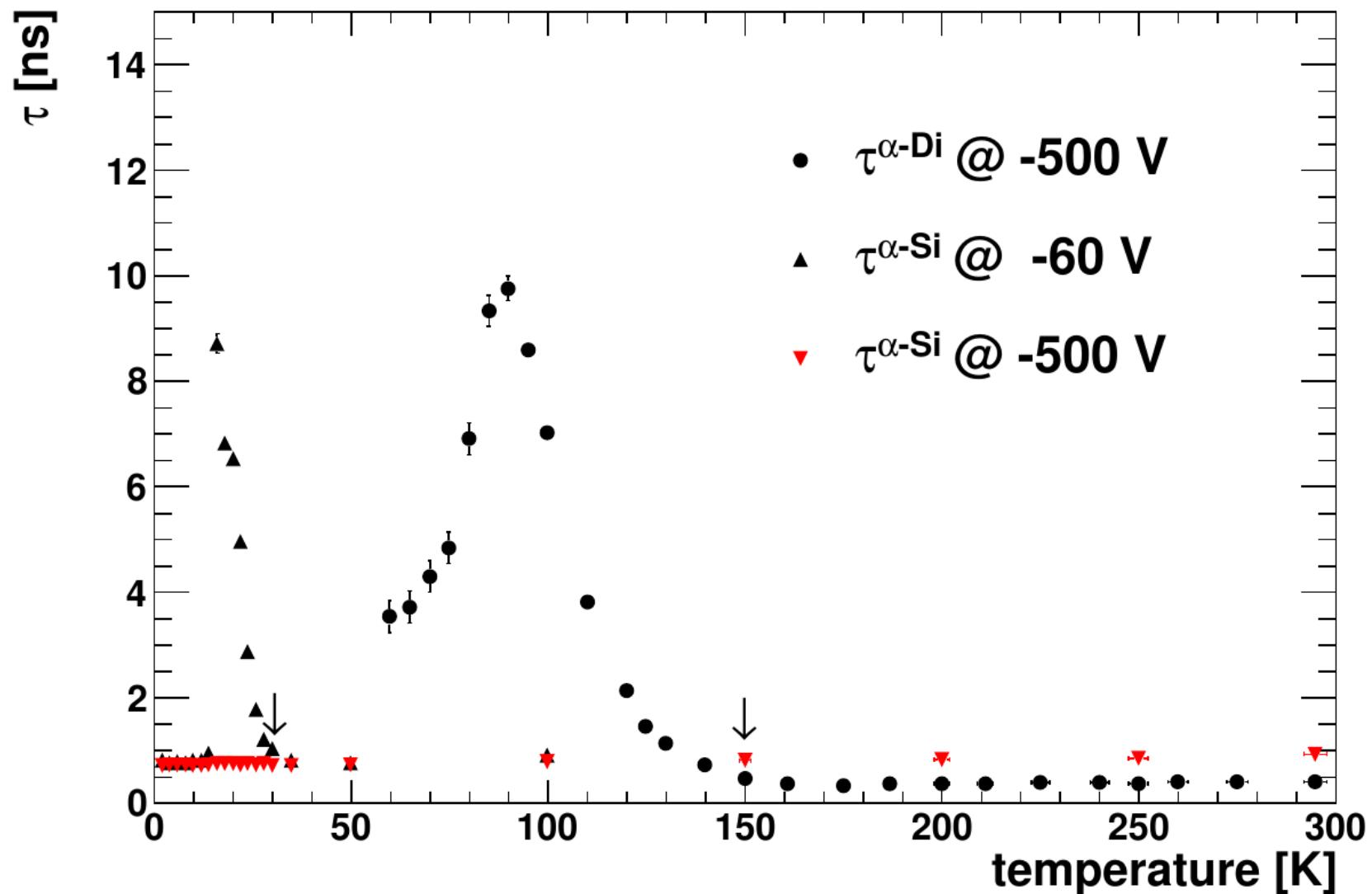
# Analysis of TCT Pulses

- Plot fitted  $\tau$  a.f.o. temperature, and fit T-dependent model:



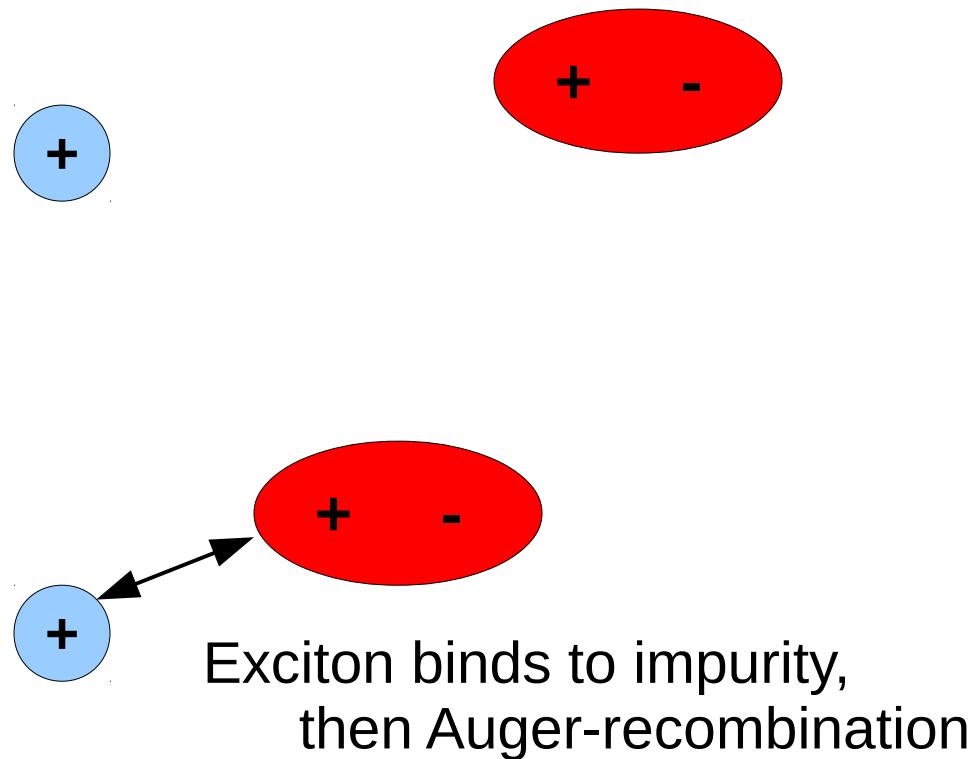
# Comparison with Silicon

- $\tau_{\text{evap}} \sim \exp(E_x / kT) !$

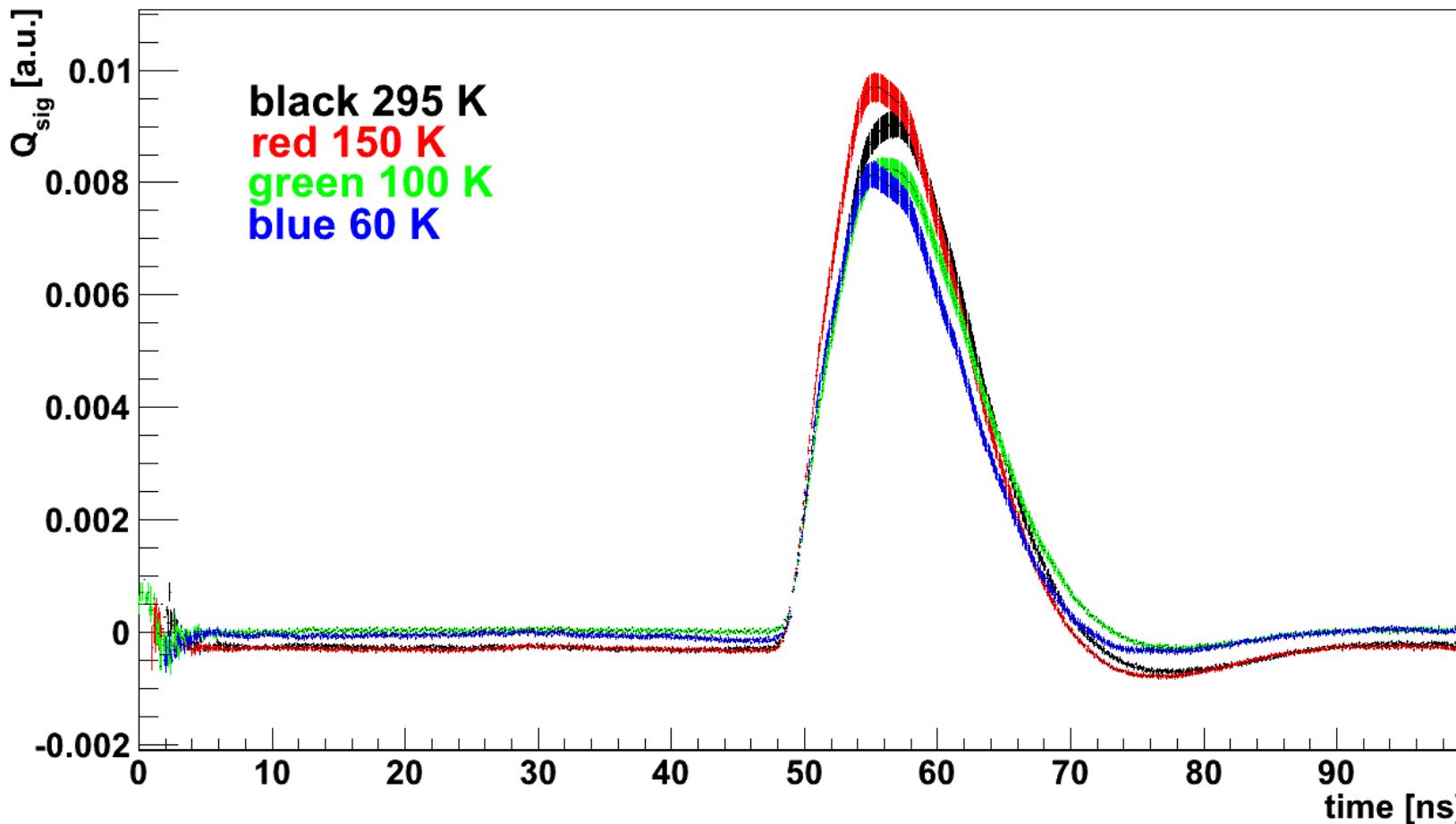


# Excitons

- Recombination via impurities



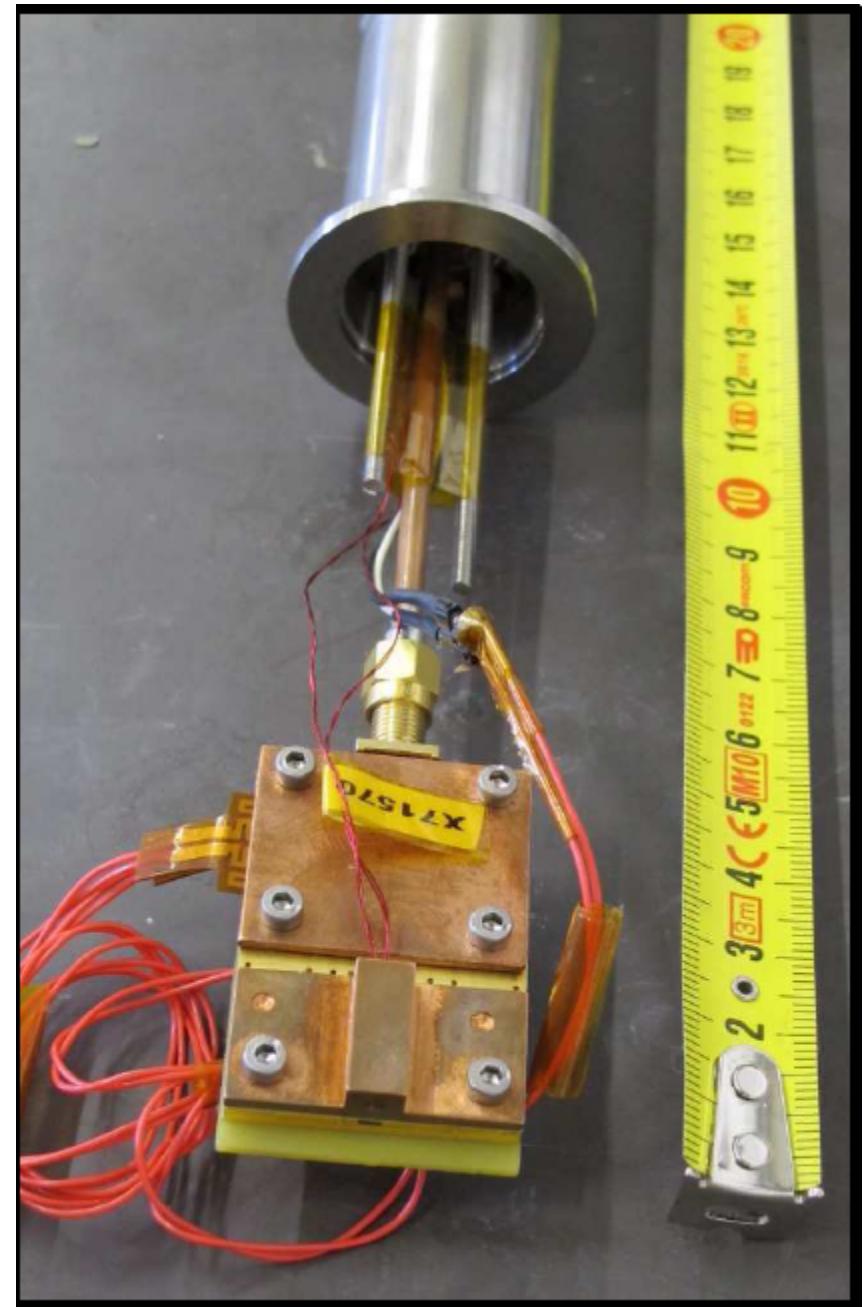
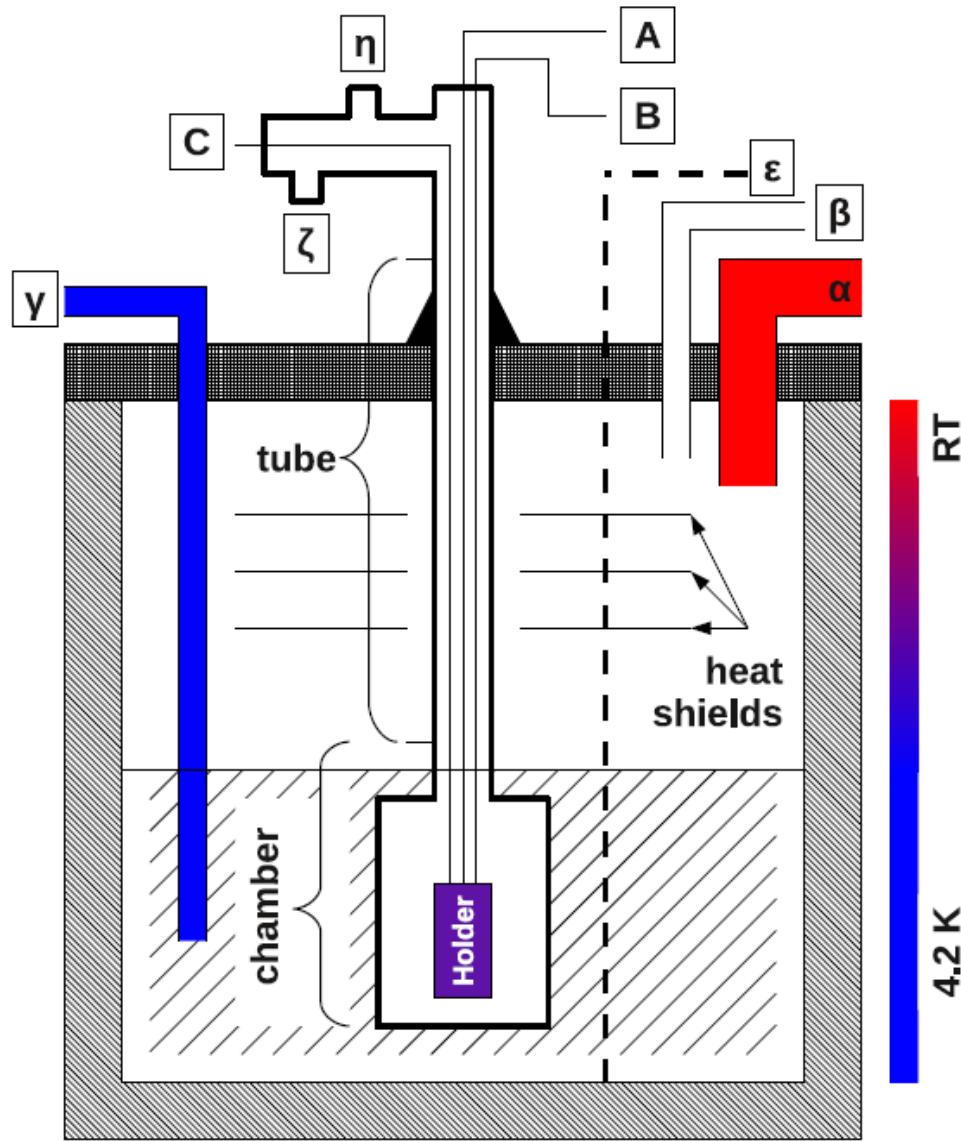
# Cosmic Muons



- Use **charge-sensitive amplifier** here
- Charge degradation much less with MIPs
  - expected due to lower charge density!

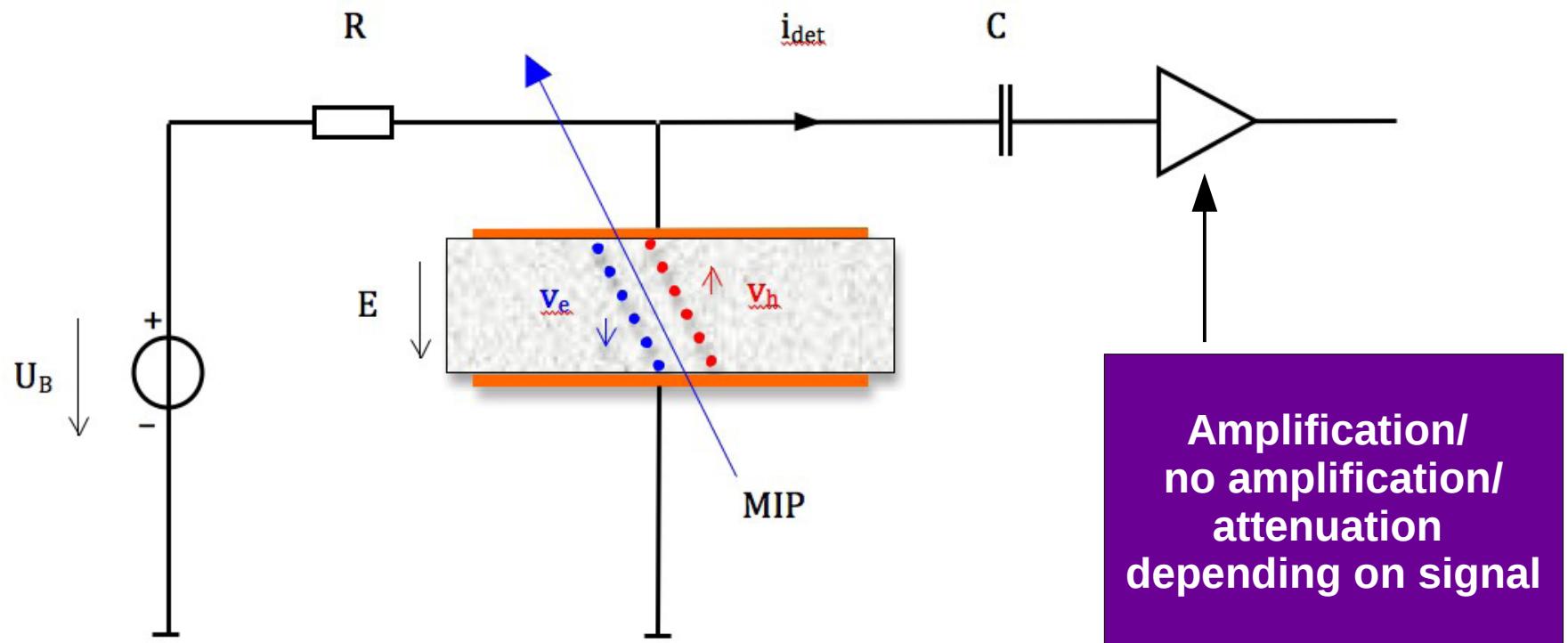
# Detailed set-up

## overview



# Common circuit for Diamond Detectors

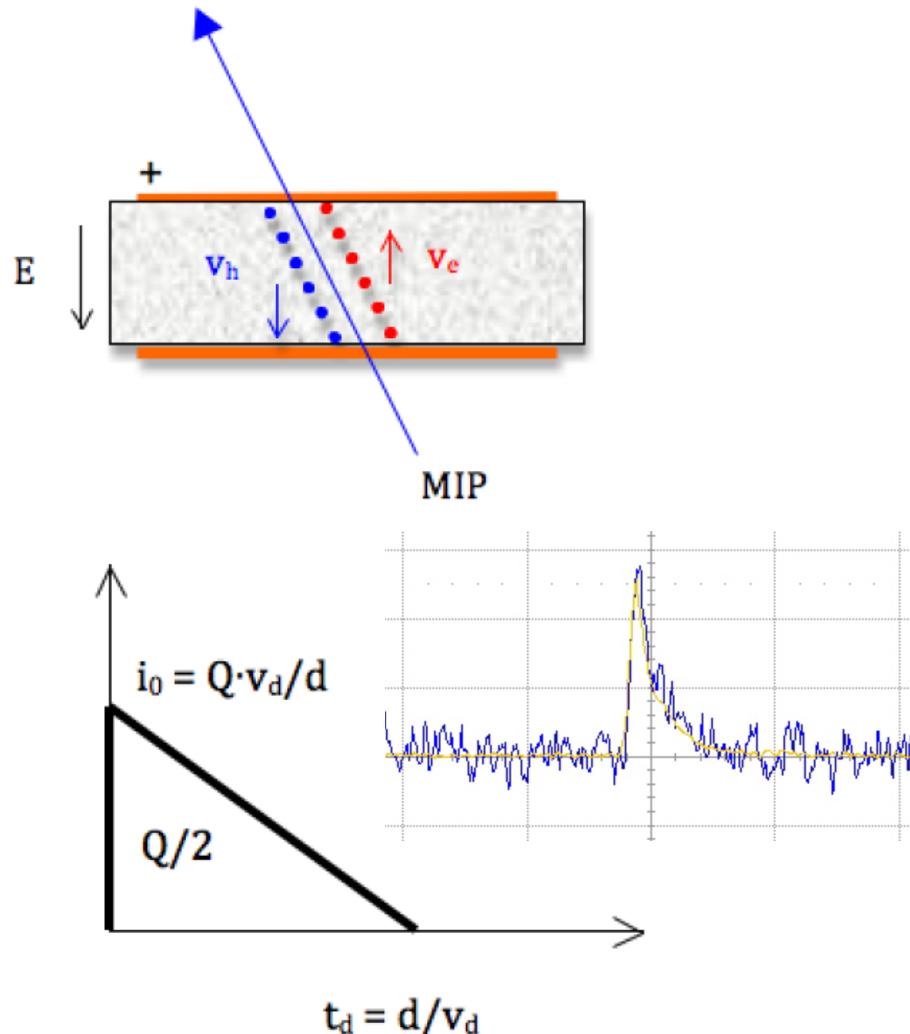
- Diamond is a semi-conductor based, solid state ionization chamber:



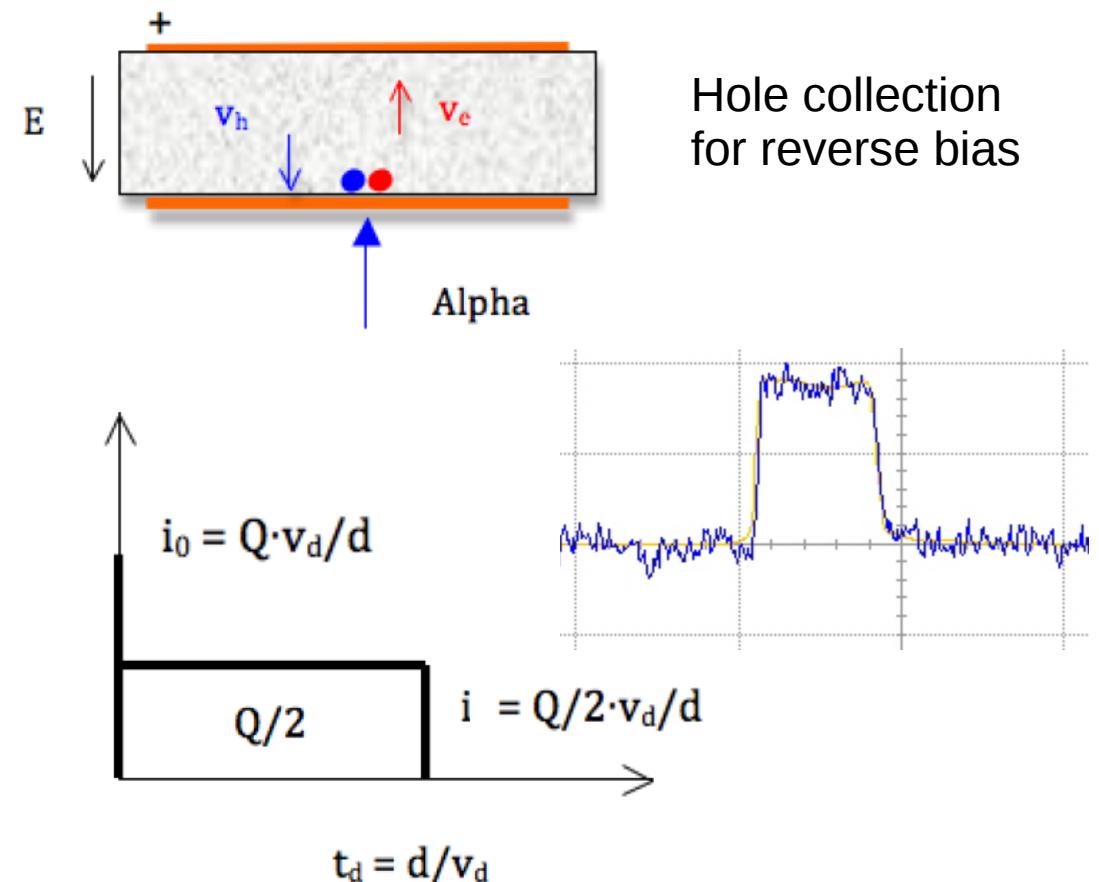
Courtesy  
CIVIDEC

# Modes of Operation

- Counting Mode



- Calorimetric Mode



Hole collection  
for reverse bias

Courtesy  
CIVIDEC

# pCVD vs scCVD

- pCVD Diamond:

- very short signal

- ~2ns FWHM @ 1V/um

- optimal double pulse resolution
  - charges lost at trapping centres

- scCVD Diamond:

- short signal

- ~5ns FWHM @ 1V/um

- optimal Signal-to-Noise ratio
  - lower trapping centre concentration

$$CCD = d * Q_{\text{sig}} / Q_0$$

