

Singlet extensions of the MSSM with $\mathbb{Z}_4^R$ -symmetry

Patrick K. S. Vaudrevange

Universe Cluster Munich

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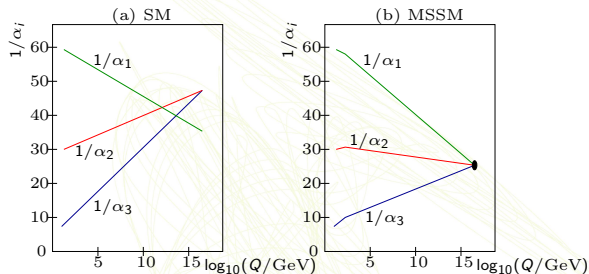
XXVII Workshop Beyond the Standard Model 2015

Based on:

► M. Ratz, P. V.: arXiv:1502.07207

Motivation for Supersymmetry

- Unification of gauge coupling constants α_i , $i = 1, 2, 3$



- Hierarchy stabilisation
- Dark matter candidate
- ...

⇒ Minimal Supersymmetric Extension of the Standard Model (MSSM)

Motivation against Supersymmetry

- ▶ Little hierarchy problem:
 - ▶ Higgs mass a bit large
 - ▶ SUSY not yet found

⇒ Singlet extension of the MSSM

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NMSSM

M. Dine, N. Seiberg, S. Thomas 2007

- ▶ Add SM singlet N coupled to $H_u H_d$

$$\mathcal{W} \supset \underbrace{\mu H_u H_d}_{\mu\text{-term}} + \lambda_N N H_u H_d + \frac{1}{2} \mu_N N^2$$

- ▶ Integrate out N

$$\Rightarrow \mathcal{W}_{\text{eff}} \supset \mu H_u H_d - \frac{\lambda_N}{2\mu_N} (H_u H_d)^2$$

- ▶ Raise Higgs mass, ameliorate the fine-tuning
- ▶ But: μ -problem & proton decay $\Rightarrow \mathbb{Z}_4^R$

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The $\mathbb{Z}_4^R$ symmetry

If one demands:

1. anomaly freedom
2. fermion masses (Yukawa couplings & neutrino mass operator)
3. consistency with SU(5)
4. gauge coupling unification (anomaly universality)

$\Rightarrow$ only R symmetries can forbid the μ -term

If additionally

5. consistency with SO(10)

$\Rightarrow$ only $\mathbb{Z}_4^R$ symmetry can forbid the μ -term, where

H.M. Lee, S. Raby, M. Ratz., G.G. Ross, R. Schieren, K. Schmidt-Hoberg, P.V. 2010/2011

$$R(\text{matter}) = 1, \quad R(\text{Higgs}) = 0, \quad R(\mathcal{W}) = 2.$$

K.S. Babu, I. Gogoladze, K. Wang 2002

The $\mathbb{Z}_4^R$ symmetry

- $\mathbb{Z}_4^R$ charges:

$$R(\text{matter}) = 1, \quad R(\text{Higgs}) = 0, \quad R(W) = 2.$$

- Most general superpotential:

$$\begin{aligned} \mathcal{W} = & \mu H_u H_d + Y^u Q \bar{U} H_u + Y^d Q \bar{D} H_d + Y^e L \bar{E} H_d \\ & + \kappa L H_u + \lambda L L \bar{E} + \lambda' L Q \bar{D} + \lambda'' \bar{U} \bar{D} \bar{D} \\ & + \frac{1}{M_P} (L H_u L H_u + Q Q Q L + \bar{U} \bar{U} \bar{D} \bar{E} + \dots) \end{aligned}$$

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$\mathbb{Z}_4^R$ broken to matter parity by anomaly $\Rightarrow \mu \sim m_{3/2}$

NMSSM with $\mathbb{Z}_4^R$ symmetry

- ▶ Add SM singlet N coupled to $H_u H_d$

$$\mathcal{W} \supset \lambda_N N H_u H_d \Rightarrow R(N) = 2$$

- ▶ Hence, linear term unconstrained, $\Lambda \sim M_P$

$$\mathcal{W} \supset \Lambda^2 N + \lambda_N N H_u H_d$$

- ▶ Forbid linear term $\Lambda^2 N$ by symmetry?

Such symmetry would forbid μ -term

$\Leftrightarrow$

Only $\mathbb{Z}_4^R$ can forbid μ -term

- ▶ No symmetry! Use holomorphic zeros!

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NMSSM with $\mathbb{Z}_4^R$ symmetry

- ▶ NMSSM with $\mathbb{Z}_4^R \times \text{U}(1)_{\text{anom}}$
- ▶ $\text{U}(1)_{\text{anom}}$ FI-term $\xi > 0$ cancelled by SM singlet ϕ , $R(\phi) = 0$:

$$D_{\text{anom}} = \xi - |\phi|^2 \stackrel{!}{=} 0 \quad \text{with} \quad \varepsilon = \frac{\langle \phi \rangle}{M_{\text{P}}} \sim \sin \vartheta_{\text{Cabbibo}} \sim 0.2$$

- ▶ Non-perturbative term e^{-bS} with $\mathbb{Z}_4^R \times \text{U}(1)_{\text{anom}}$ charges 2 and $s > 0$

$$\Rightarrow \mathcal{W}_{\text{non-pert.}} \supset \phi^s e^{-bS} \sim m_{3/2}$$

e.g. Arkani-Hamed, Dine & Martin 1998

NMSSM with $\mathbb{Z}_4^R$ symmetry

► Charges

	ϕ	$H_u H_d$	N	e^{-bS}
$U(1)_{\text{anom}}$	-1	$h > 0$	$n < 0$	$s > 0$
$\mathbb{Z}_4^R$	0	0	2	2

► Generalised NMSSM

M. Drees 1989

$$\mathcal{W}_{\text{eff}} \supset f^2 N + \mu H_u H_d + \lambda N H_u H_d + \mu_N N^2 + \kappa N^3$$

where

$$f \sim \varepsilon^{n/2} m_{3/2}, \quad \mu \sim \varepsilon^h m_{3/2}, \quad \lambda \sim \varepsilon^{n+h}$$

$$\mu_N \sim \varepsilon^{2n} m_{3/2}, \quad \kappa \sim \varepsilon^{3n} \frac{m_{3/2}^2}{M_{\text{P}}^2}$$

M. Ratz, P.V. 2015

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Conclusion

- ▶ The $\mathbb{Z}_4^R$ NMSSM
- ▶ Holomorphic zeros forbid linear term
- ▶ Effective superpotential

$$\mathcal{W}_{\text{eff}} \supset f^2 N + \mu H_u H_d + \lambda N H_u H_d + \mu_N N^2 + \kappa N^3$$

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