

s Channel Production of Heavy Vector Bosons in the Three-Site Higgsless Model [arXiv:0809.0023]

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HGF workshop — LHC-D session — SUSY and BSM
Aachen, 27.11.2008



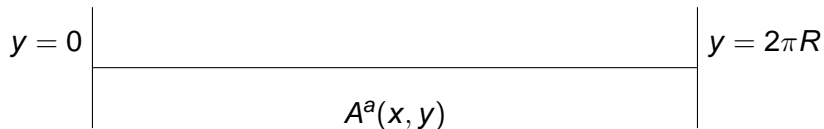
Outline:

- 1 Introduction
- 2 The model
- 3 Parameter space & Co.
- 4 Z' / W' production
- 5 Separating the resonances
- 6 Conclusions

Formalism:

5D gauge theory:

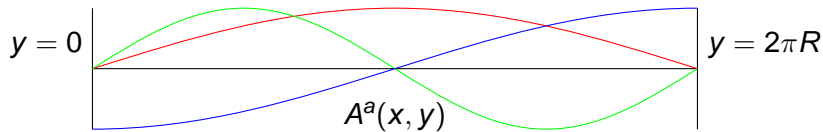
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- interpretation: “tower” of Kaluza-Klein (KK) modes, $m_i = |k_i^5|$



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Dimensional deconstruction:

- replace $y \in [0; 2\pi R]$ with discrete lattice sites y_i



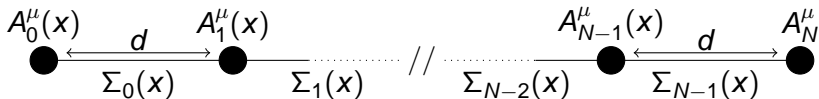
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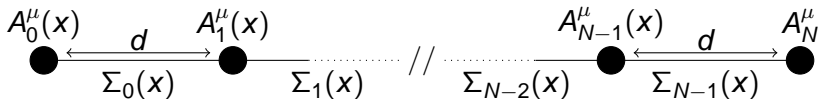
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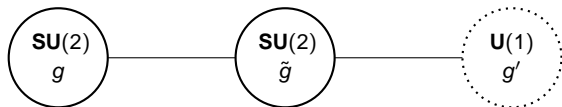
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Minimal Three-site higgsless model [Chivukula et al, 2006]

- three lattice sites
- Standard Model + one set of KK modes (no KK gluons or photon)

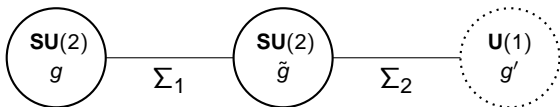


Model structure:



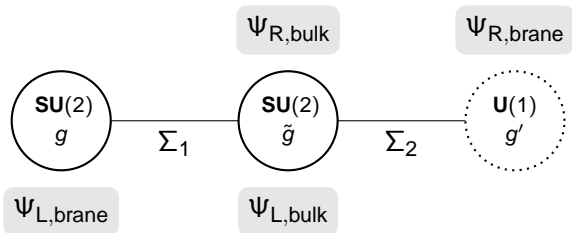
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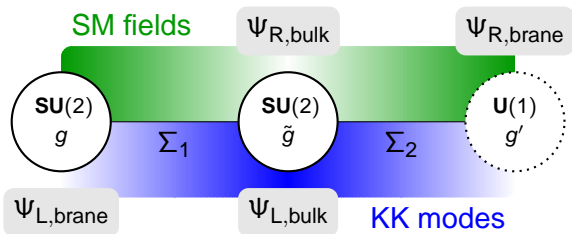
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Diagonalization:

- SM particles strongly localized at branes, KK partners in bulk
- couplings between SM fermions and KK bosons very small (exception: top, bottom)

Constraints from LEP & LEP2:

Tree Level: Chivukula, Simmons 2005 (hep-ph/050411)

Consistency with electroweak precision constraints



vanishing couplings between SM fermions and heavy W (W')

- “ideal” delocalization of the fermions along the lattice
- **no** W' in the s channel!

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One Loop: Abe, Matsuzaki, Tanabashi 2008 (arXiv:0807.2298)

Consistency with electroweak precision constraints



$W' \bar{f} f$ coupling small but **finite** ($\mathcal{O}(10^{-2})$)

“nonideal” delocalization: **both** Z' and W' may be visible in the s channel

Pros of the model:

- Breaks electroweak symmetry without invoking **any** scalar field
- Consistent with electroweak precision observables
- Maintains exact gauge invariance while chopping off all but the lightest Kaluza-Klein modes
- Very limited particle spectrum \leftrightarrow easy to implement into a Monte Carlo eventgenerator

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→ So:

No candidate for a **fundamental** model of nature, but might teach us something about the physics of scenarios without a Higgs that are not already forbidden by LEP data.

Free parameters in addition to standard model:

- W' mass: $380 \text{ GeV} \leq m_{W'} \leq 600 \text{ GeV}$
- $W'\bar{f}f$ coupling $g_{W'ff}(\epsilon_L)$ (ϵ_L : fermion delocalization parameter)

$$\frac{g_{W'ff}}{g} = \mathcal{O}(1\%)$$

(heavily constrained by precision data; g : isospin gauge coupling)

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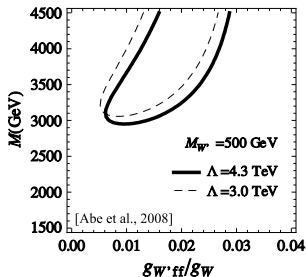
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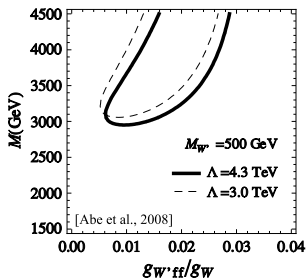
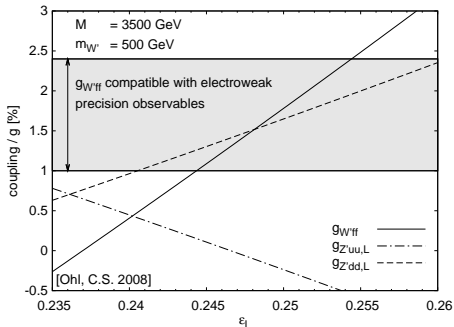
Particle spectrum:

- standard model particles
- heavy gauge bosons W', Z' : quasidegenerate $\frac{m_{W'}}{m_{Z'}} = 1 + \mathcal{O}(1\%)$,
relative width $\frac{\Gamma}{m} = \mathcal{O}(1\%)$
- heavy fermions: masses around M , $\frac{\Gamma}{m} = \mathcal{O}(10\%)$

- Bounds from electroweak precision observables on $\frac{g_{W'ff}}{g'}$ for $m_{W'} = 500$ GeV, [Abe et al., arXiv:0807.2298]
- allowed region varies (increases) with (lower) $m_{W'}$



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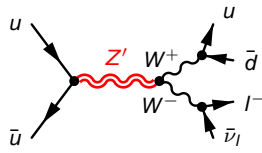


- influence of ϵ_L on left-handed $Z'f\bar{f}$ coupling
- changes to $g_{Z'uu,L}$ and $g_{Z'dd,L}$ tend to compensate each other in pp collisions
- **no** change to right-handed couplings which are of the same order of magnitude

Interesting processes with heavy vectors in s channel:

higher	→ total cross section (signal)	→ lower
$pp \rightarrow jjl\nu_l$	$pp \rightarrow jjl\nu_l$	$pp \rightarrow jjll$

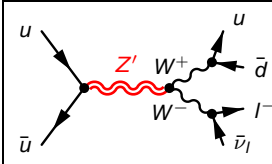
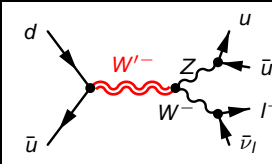
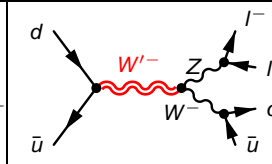
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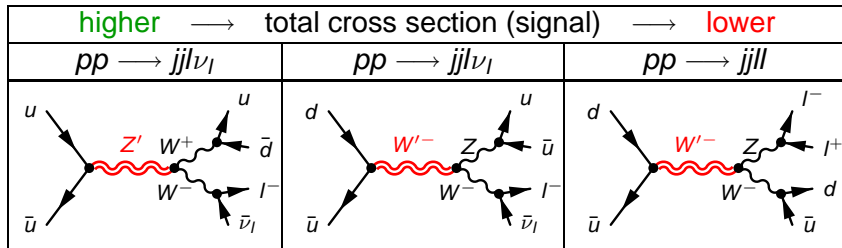
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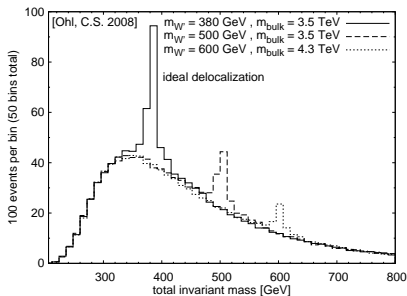


Monte Carlo simulation with WHIZARD / O'Mega:

- Full parton level simulations of $pp \rightarrow jjl\nu_l$ and $pp \rightarrow jjll$
→ includes irreducible and most reducible backgrounds
- Missing neutrino momentum: reconstructed from transverse momentum conservation and W, ν mass shell conditions
- Strict cuts for background reduction:
 $p_T \geq 50$ GeV for all particles and also for $p_{T,miss}$

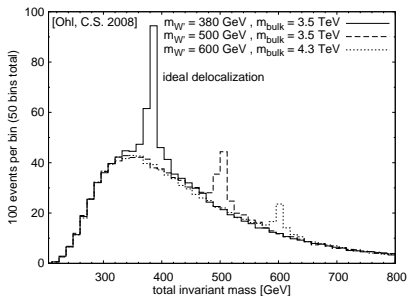
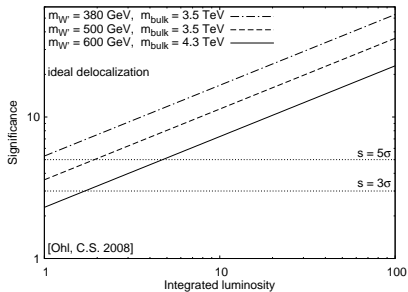
Simulation results for $pp \rightarrow jjl\nu_l$ (intermediary Z'):

- Simulations for integrated luminosity $\int dt \mathcal{L} = 100 \text{ fb}^{-1}$
- Signal drops with higher $m_{W'}$
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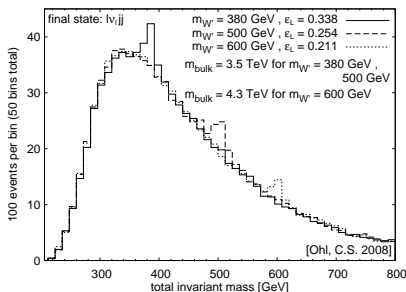
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- Signal significance vs. integrated luminosity
- \rightarrow Z' should be visible in the first $10 - 20 \text{ fb}^{-1}$

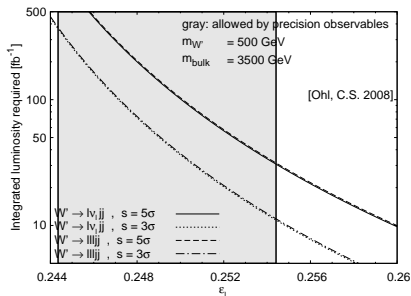
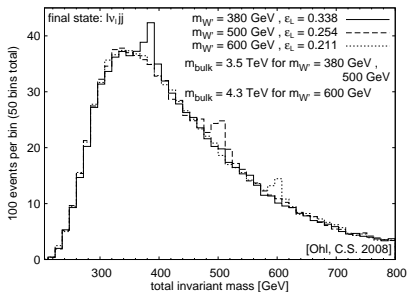
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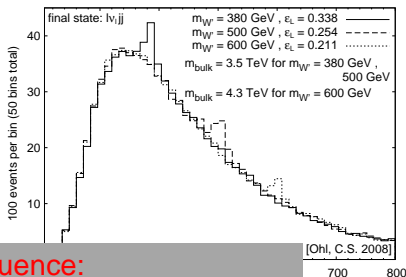
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- Int. luminosity required for 5σ (3σ) discovery vs. ϵ_L
- Performance of $jjl\nu_l$ and $jjll$:
 - ▶ similar for $m_{W'} = 500 \text{ GeV}$
 - ▶ $jjll$ slightly better for other values of $m_{W'}$

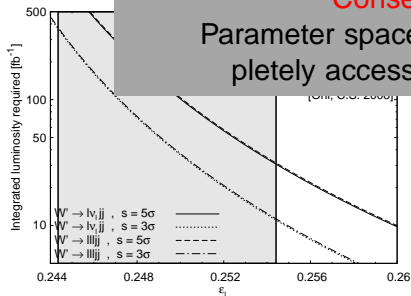
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Consequence:

Parameter space may not be completely accessible at the LHC!



required for 5σ

(3σ) discovery vs. ϵ_L

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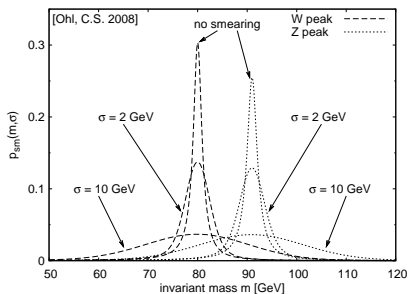
Separation of W' and Z' contributions to $pp \rightarrow jjl\nu_l$

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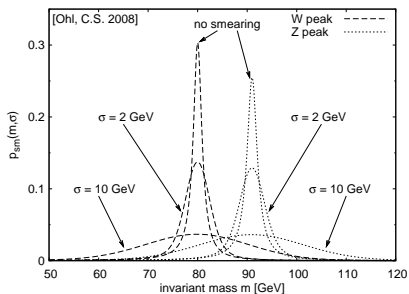
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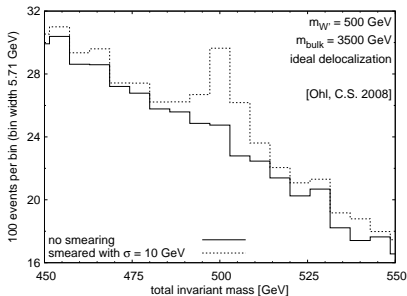
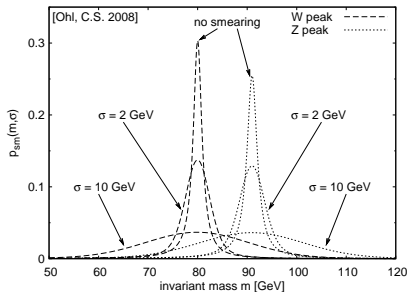
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- Separation of $jj = W$ and $jj = Z$ via invariant mass cuts
- Measurement error smears the sharp Breit-Wigner resonances of the W and Z
- Possible misidentification due to finite experimental resolution
- Consequence: misidentified Z' events can fake a W' !



Possible solution:

- Smearred Breit-Wigner can be integrated for given cuts
- Gives propability matrix T for W/Z identification:

$$\begin{pmatrix} N_{W,\text{meas.}} \\ N_{Z,\text{meas.}} \end{pmatrix} = \begin{pmatrix} P_{WW} & P_{WZ} \\ P_{ZW} & P_{ZZ} \end{pmatrix} \begin{pmatrix} N_W \\ N_Z \end{pmatrix} = T \begin{pmatrix} N_W \\ N_Z \end{pmatrix}$$

- Invert to reconstruct “true” event counts

$$\begin{pmatrix} N_W \\ N_Z \end{pmatrix} = T^{-1} \begin{pmatrix} N_{W,\text{meas.}} \\ N_{Z,\text{meas.}} \end{pmatrix}$$

- Tried it: **seems to work in our example!**

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Problems:

- Requires separation of signal \leftrightarrow assumptions about background shape
- Statistical error increased by reconstruction \rightarrow greatly reduced significance

Conclusions:

- Three Site Model is an interesting example of higgsless BSM modelbuilding
- Couplings to W'/Z' are severely constrained by electroweak precision observables \rightarrow measurement crucial
- s Channel measurement seems be possible, but not the whole parameter space may be accessible
- Quasi-degeneracy of W'/Z' makes separation difficult in some channels \leftrightarrow tricks like the T -matrix might help
- More in-depth study required including detector effects \rightarrow in progress [T. Trefzger, F. Bach, Würzburg]

Thank you for your attention!

Reconstructing the neutrino momentum:

constraints on p_ν :

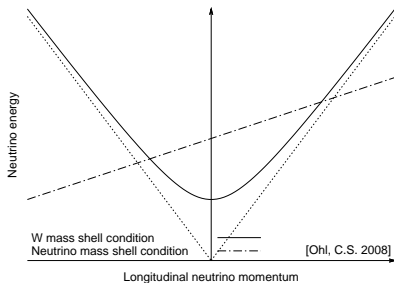
- ν mass shell condition:

$$p_\nu^2 = 0$$

- W mass shell condition:

$$(p_l + p_\nu)^2 = m_W^2$$

- transverse momentum conservation

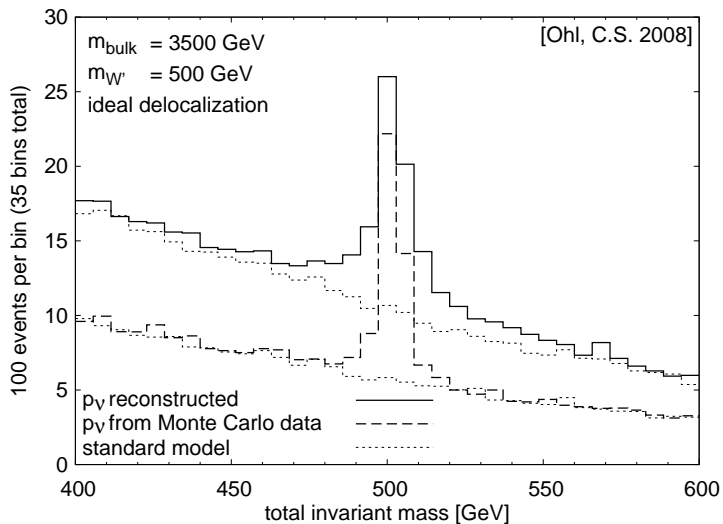


→ neutrino momentum can be solved for

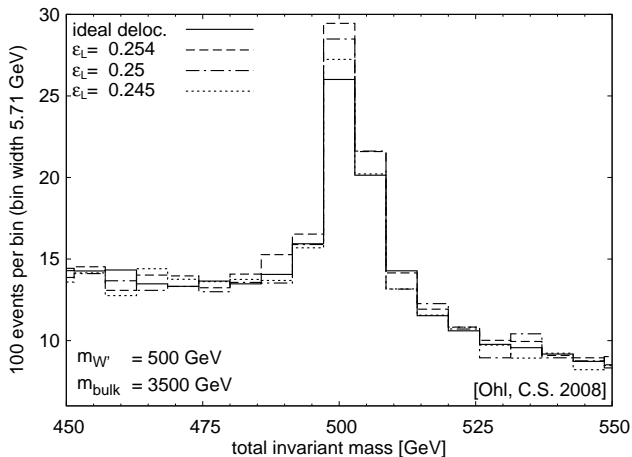
Problem: Two solutions

- Only **one** is correct
- **None** is preferred on kinematical grounds (in the general case)
- One (**our**) solution to this: count both into the histograms → more background!

Example: Z' production:

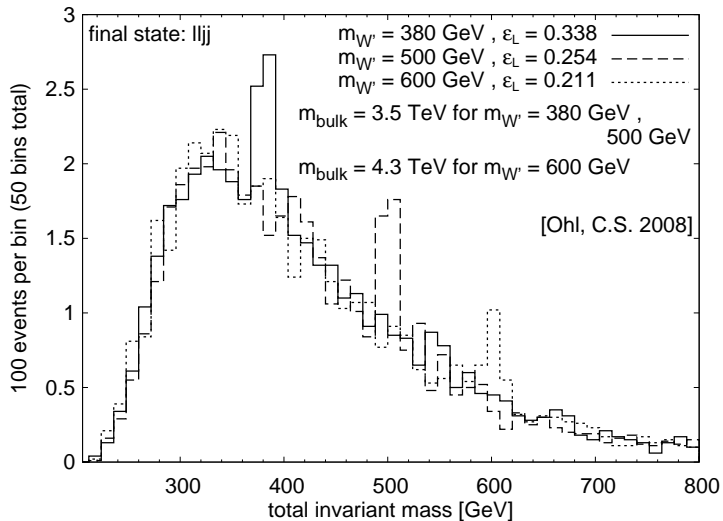


Effect of changing ϵ_L on the Z' production:



→ makes no big difference in the whole allowed parameter space

Invariant mass distributions for W' production in $pp \rightarrow jjll$:



Cuts:

- p_T cuts on all particles, including $p_{T,\text{miss}}$

$$p_T \geq 50 \text{ GeV}$$

- angular cuts on all visible particles

$$-0.95 \leq \cos \theta \leq 0.95$$

- small x cut on ingoing partons

$$x \geq 1.4 \times 10^{-3}$$

- mass shell cuts on fermion pairs

$$75 \text{ GeV} \leq m_W \leq 85 \text{ GeV} \quad , \quad 86 \text{ GeV} \leq m_Z \leq 96 \text{ GeV}$$