AMELIE

An Axion Modulation hELIoscope Experiment <u>idea</u>

J. Galan University of Zaragoza 11th Patras Workshop on Axions, WIMPs and WISPs

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Pionering helioscope searches were using stationary helioscopes



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The helioscope reference papers

VOLUME 51, NUMBER 16

PHYSICAL REVIEW LETTERS

Experimental Tests of the "Invisible" Axion

P. Sikivie Physics Department, University of Florida, Gainesville, Florida 32611 (Received 13 July 1983)

THIRD SERIES, VOLUME 39, NUMBER 8

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Design for a practical laboratory detector for solar axions

K. van Bibber Physics Department, Lawrence Livermore National Laboratory, University of California, Livermore, California 94550

> P. M. McIntyre Physics Department, Texas A&M University, College Station, Texas 77843

D. E. Morris Physics Division, Lawrence Berkeley Laboratory, University of California, Berkeley, California 94720

G. G. Raffelt Institute for Geophysics and Planetary Physics, Lawrence Livermore National Laboratory, University of California, Livermore, California 94550 and Astronomy Department, University of California, Berkeley, California 94720 (Received 19 September 1988)

1989 Axior ray Shielding MWPC Window Dispersion-Matching Gas (H, or He)

1983

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17 October 1983

15 APRIL 1989

Two helioscope techniques for axion detection



Absorbed or transmitted photon component



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A magnetic TPC for an helioscope proposal

The axion-photon <u>component transmitted</u> goes typically as L², thus a long pipe is usually the best suitable geometry.

However, for the axion-photon <u>component</u> <u>absorbed</u> is just proportional to L, thus (for high Γ) we can use **any volume geometry** (in principle).

$$P_{\gamma} = \frac{(B/2M)^2}{q^2 + \Gamma^2/4} \Gamma L$$

This technique is especially interesting if $\Gamma L >> 1$. High Z gases (for many reasons). A possible conceptual TPC-magnet design



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Daily modulation due to Earth rotation

The effective magnetic field for axion-photon conversion modulates. We must consider only the transversal component to the axion propagation direction that changes along the day.



Annual axion signal modulation (daily averaged)



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Full angular view sensitivity (4π)



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Effective conversion probability (Geometrical factor)

There is a geometrical factor to consider due to the different coherence lengths traversing the chamber.

5.5*10^-18 5.0*10^-18 Conversion probability 4.5*10^-18 10 mbar 4.0*10^-18 50 mbar 100 mbar 3.5*10^-18 For a vertical cylinder 3.0*10^-18 2.5*10^-18 5 10 15 20 0 Hour of the day

Conversion probability is reduced when ΓL is not >> 1. And this relation does not apply $(\Gamma L = 0.15)^2$

$$P_{\gamma} = \frac{(B/2M)^2}{q^2 + \Gamma^2/4} \Gamma L$$

Efficiency loss in Xenon

| 10 mbar | 21.5% |
|----------|-------|
| 50 mbar | 11.9% |
| 100 mbar | 6.9% |
| 500 mbar | 1.4% |
| 5 bar | 0.13% |

Small diurnal effect and negligible annual variations

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Conversion probability resonances



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State of the art in low background large volume TPCs



AMELIE Search sensitivity prospects

Run conditions (Xenon)

16 Pressure steps from 10mbar to 5bar 150 days per step ~ 8 years B = 5T

Background

AMELIE-PROTO : 1 cpd/m³/keV AMELIE : 0.1 cpd/m³/keV AMELIE-IAXO : 0.1 cpd/m³/keV

<u>Volume</u> AMELIE-PROTO : 21 dm³ AMELIE : 0.785 m³ AMELIE-IAXO : 45.23 m³



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Conclusions

- An independent technique for solar axion searches is foreseen.
- <u>Full angular sensitivity</u> (allows to observe the full solar disk + other exotic sources). Also <u>not high accuracy alignment</u> required.
- Improved <u>gas density stability and broader axion mass coverage</u> with a single setting. Leading to a simplified data taking process.
- Enables <u>searches to very low energies</u> (few ~100eV). Since there is no separation between detector and conversion volume.
- It would allow to proof KSVZ axions above 50meV-100meV.
- Rare event searches underground with TPCs are becoming very exciting, using a magnetic field will make these searches even more interesting.
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Backup Slides

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Production probability in a magnetic gas

$$i\partial_{z} \begin{pmatrix} A(z) \\ a(z) \end{pmatrix} = \begin{pmatrix} E_{a} - m_{\gamma}^{2}/2E_{a} - i\Gamma/2 & B' \\ B' & E_{a} - m_{a}^{2}/2E_{a} \end{pmatrix} \begin{pmatrix} A(z) \\ a(z) \end{pmatrix}$$

$$P_{a}(z) = |a(z)|^{2} = \frac{B'^{2}}{q^{2} + \Gamma^{2}/4} [1 + e^{-\Gamma z} - 2e^{-\Gamma z/2}\cos(qz)] \qquad P_{\gamma}(z) = |A(z)|^{2} = e^{-\Gamma z} \left\{ 1 - \frac{B'^{2}}{q^{2} + \Gamma^{2}/4} e^{\frac{\Gamma z}{2}\cos(qz) + \mathcal{O}(B'^{4})} \right\}$$

Photon in the dense plasma will be quickly absorbed. The coherence length will depend on the path-length of the photon.

Mean probability to convert a photon into axion

In practice we just integrate the probability of a photon to reach a distance z by the conversion probability at z

$$P_{\gamma a} = \int_0^\infty P_\gamma(z) P_a(z) dz \longrightarrow \qquad P_{\gamma a} = \frac{1}{6} \frac{B'^2}{q^2 + \Gamma^2/4}$$

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Solar axion flux integration

Considering only the transversal component we obtain the mean conversion probability per photon integrated to 4π

$$P_{\gamma a} = \frac{1}{6} \frac{B_{\perp}'^2}{q^2 + \Gamma^2/4} = \frac{\pi^2}{96} \frac{B_o'^2}{q^2 + \Gamma^2/4}$$

photon irradiance in thermal equilibrium

$$\frac{d\phi_a}{dE} = \frac{1}{4\pi d_{\odot}^2} \int_{V_{Sun}} P_{\gamma a} \, d^3 \mathbf{r} \, \frac{1}{(2\pi)^2} \frac{E^3}{e^{E/T} - 1} \qquad \begin{array}{l} \mbox{Flux in} \\ \mbox{cm}^{-2} \, \mathrm{s}^{\text{-1}} \, \mathrm{keV}^{\text{-1}} \end{array}$$

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Solar axion production in the Sun inner magnetic fields

S.Couvidat, S. Turck-Chieze, A. Kosovichev, APJ 599, 1434 (2003)

 B_0 Center^a Name (R_{\odot}) $(P_{\rm mag}/P_{\rm gas})_{\rm max}$ (T)104 0.224 asinic₁D₁..... 10 0.200 1.004 0.230 $.90 \times 10$ seisinie₁D₁₁..... 0.236Seismic B12. ₹x=10° 2.49 x 10 → Seismic₁B₁₃..... 1×10^{3} 0.236 2.80×10^{-4} 0.712 6.15 × 10⁻⁵ Seismic₁B₂..... 30 Seismic₁B₂₁..... 50 1.71×10^{-4} 1.34×10^{-4} Seismic₁B₃..... 2 0.96 3.02×10^{-4} Seismic₁B₃₁..... 3 ^a Radius at which P_{mag} is maximum. 9). The precision we have on the solar sound speed rules out

TABLE 4 Models with Magnetic Field Assuming scattering factor for H $f_1 = 1$ and a purely Hydrogen Sun composition

$$m_{\gamma} \simeq 28.77 \sqrt{\rho(r) [\text{g/cm}^3]} \,\text{eV}$$



9). The precision we have on the solar sound speed rules out a magnetic field with such a profile and an intensity at large as $B_0 = 10^4$ T. Actually, we can put an upper limit for a (toroidal) magnetic field in the radiative zone of about 3×10^3 T (seismic₁B₁₂ model). If $B_0 \le 10^3$ T, then c_s is not sensitive enough to P_{mag} , and we cannot draw any conclusion about J. Galan, \angle D-June

Solar axion production in the Sun inner magnetic fields



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Solar axion spectrum generated by the different magnetic regions



Differential axion flux production as a function of R



Spherical TPC Detection principle

Micromegas : Best results in terms of energy resolution and spatial resolution (few um)





Micromegas : Amplification given by (HV1-HV2)/d

Spherical TPC : Amplification given by HV and Rsensor. Definition of amplification gap and drift region relays on **strong field variation near the sensor**.

Low noise (few eV) Spherical geometry

Simplified read-out price for decreased spatial resolution

> Intense field. Enough to produce avalanche processes.



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A spherical TPC for particle physics (neutrons, neutrinos, DM, ...)

Existing spheres running in ground and underground laboratories (NEWS collaboration)

- Large volume
- Simplified read-out
- Gas pressure flexibility : from few mbar to several bar
- Good energy resolution
- Low energy threshold
- Field cage and vessel are one single entity.

Home made Ar-37 source: irradiating Ca-40 powder with fast neutrons 7x10⁶neutrons/s

Irradiation time 14 days. Ar-37 emits K(2.6 keV) and L(260 eV) X-rays (35 d decay time)



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A spherical TPC for rare event searches (SEDINE)

SEDINE detector installed at Underground Modane Laboratory

Volume : 100 dm³ Pressure : From few mbar to 10 bar Gas : He, Ne, Ar, ...





