

# On the Inverse Seesaw Mechanism and Axion Like Particles

Alex G. Dias

UFABC – Brazil

11<sup>th</sup> Patras Workshop on Axions, WIMPs and WISPs

Zaragoza, Spain, 25 June 2015



- 1 Axion Like Particles and the Inverse Seesaw Mechanism
- 2 A Model with  $\mathbb{Z}_N$  Discrete Gauge Symmetry
- 3 Conclusions

Based in a work done in collaboration with:

C. D. R. Carvajal, C. C. Nishi, B. L. Sánchez-Vega, JHEP 15

Prospective experiments, astrophysical limits and hinted phenomena on ALP.

### Prospective experiments, astrophysical limits and hinted phenomena on ALP.

- ALPS-II experiment [Bähre et al. 13]
- CAST and the helioscope IAXO [Armengaud et al 14]
- Haloscope [Asztalos et al 11]
- Observatories PIXIE and PRISM [Kogut et al. 11; André et al. 14; Mirizzi et al 09; Tashiro et al 13; Ejlli, Dolgov 13]
- ALP from massive stars, Supernova SN 1987A [A. Friedland et al. 13; Grifols et al 96; Brockway et al. 96; Payez et al. 14]
- Universe transparency for very energetic  $\gamma$ -ray, [Mirizzi et al 07, De angelis et al 07; Horns, Meyer 12; Rubtsov, Troitsky 14; Biteau, Williams 15;; Simet et al 07, Sanchez-Conde et al 09; Mirizzi, Montanino 09; Horns et al 14; Tavecchio et al 14 ]
- Soft X-ray excess through cosmic ALP background [Conlon, Marsh 13; Angus et al. 13; Kraljic 14]

ALP physics described by

$$\mathcal{L} = \frac{1}{2} \partial_\mu a \partial^\mu a - \frac{1}{2} m_a^2 a^2 - \frac{g_{a\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \mathcal{L}_{a\psi}(\partial_\mu a, \psi)$$

– This effective Lagrangian might result from a  $U(1)$  global symmetry broken spontaneously by a VEV of a scalar field hosting the ALP  $\sim a(x)$

$$\sigma(x) = \frac{1}{\sqrt{2}} [v_\sigma + \rho(x)] e^{i \frac{a(x)}{v_\sigma}}, \quad v_\sigma \gg 246 \text{ GeV.}$$

– ALP-photon coupling suppressed by  $\langle \sigma \rangle = v_\sigma$

$$g_{a\gamma} = \frac{\alpha C_{a\gamma}}{2\pi v_\sigma}, \quad \frac{C_{a\gamma}}{v_\sigma} = \frac{1}{f_a}$$

–  $m_a$  is due a small explicit breaking of  $U(1)$

ALP physics described by

$$\mathcal{L} = \frac{1}{2} \partial_\mu a \partial^\mu a - \frac{1}{2} m_a^2 a^2 - \frac{g_{a\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \mathcal{L}_{a\psi}(\partial_\mu a, \psi)$$

– This effective Lagrangian might result from a  $U(1)$  global symmetry broken spontaneously by a VEV of a scalar field hosting the ALP  $\sim a(x)$

$$\sigma(x) = \frac{1}{\sqrt{2}} [v_\sigma + \rho(x)] e^{i \frac{a(x)}{v_\sigma}}, \quad v_\sigma \gg 246 \text{ GeV}.$$

– ALP-photon coupling suppressed by  $\langle \sigma \rangle = v_\sigma$

$$g_{a\gamma} = \frac{\alpha C_{a\gamma}}{2\pi v_\sigma}, \quad \frac{C_{a\gamma}}{v_\sigma} = \frac{1}{f_a}$$

–  $m_a$  is due a small explicit breaking of  $U(1)$

But this breaking must be controlled in some way!

– ALP has to be an *ultra-light particle* to be relevant in certain astrophysical phenomena, and also to be discovered in the prospective experiments.

- Universe transparency hint for very high energy  $\gamma$ -rays [Meyer et al 13]

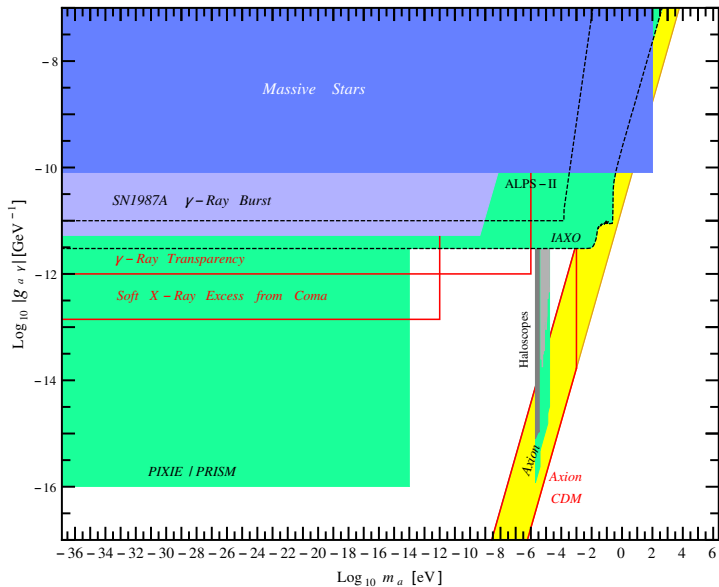
$$|g_{a\gamma}| \gtrsim 10^{-12} \text{ GeV}^{-1}; \quad m_a \lesssim 10^{-7} \text{ eV}.$$

- Soft X-ray excess hint from Coma cluster [Conlon et al 13]

$$|g_{a\gamma}| \gtrsim 10^{-13} \text{ GeV}^{-1} \sqrt{0.5/\Delta N_{\text{eff}}}; \quad m_a \lesssim 10^{-12} \text{ eV},$$



# Axion Like Particle: parameter space



AD, Machado, Nishi, Ringwald, Vaudrevange, arXiv:1403.5760

- Certain gravity induced nonrenormalizable operators need to be suppressed in order to have such ultra-light particles (no expectations that gravity respect continuous global symmetries).

$$V_{grav} = g \frac{\sigma^D}{M_{Pl}^{D-4}} + \text{H.c.} \rightarrow m_a \approx |g| D \frac{v_\sigma}{\sqrt{2}} \times \left[ \frac{v_\sigma}{\sqrt{2} M_{Pl}} \right]^{\frac{D}{2}-2}$$

- Certain gravity induced nonrenormalizable operators need to be suppressed in order to have such ultra-light particles (no expectations that gravity respect continuous global symmetries).

$$V_{grav} = g \frac{\sigma^D}{M_{Pl}^{D-4}} + \text{H.c.} \rightarrow m_a \approx |g| D \frac{v_\sigma}{\sqrt{2}} \times \left[ \frac{v_\sigma}{\sqrt{2} M_{Pl}} \right]^{\frac{D}{2}-2}$$

- In particular, the effects of gravity induced nonrenormalizable operators over the strong CP problem solution with the axion were observed many years ago [Ghigna et al.; Holman et al.; Kamionkowski et al. 92; Kallosh et al. 95] .

Let us imagine that the mass scales *beyond the Standard Model* are given by ratios

$$\text{masses} \sim \langle \sigma \rangle \times \left( \frac{\langle \sigma \rangle}{M_{Pl}} \right)^m, \quad M_{Pl} = 2.4 \times 10^{18} \text{ GeV}$$

$\sqrt{2}\langle \sigma \rangle = v_\sigma$  is within the ALP decay constant:  $10^9 \text{ GeV} \lesssim v_\sigma \lesssim 10^{14} \text{ GeV}$

– Mass scales beyond the Standard Model appear in the neutrinos physics.

# The Inverse Seesaw Mechanism

Neutrinos mass generation mechanism involving two energy scales beyond the Standard Model.

[Mohapatra, Valle 86]

# The Inverse Seesaw Mechanism

Neutrinos mass generation mechanism involving two energy scales beyond the Standard Model.

[Mohapatra, Valle 86]

– Texture

$$\begin{pmatrix} 0 & \mathbf{M}_D^T & 0 \\ \mathbf{M}_D & 0 & \mathbf{M}^T \\ 0 & \mathbf{M} & \mu \end{pmatrix}$$

–  $\mathbf{M}_D \sim$  electroweak scale,  $\mathbf{M} \sim 0.1\text{-}1$  TeV scale,  $\mu \sim$  keV scale (breaks lepton number).

# The Inverse Seesaw Mechanism

Neutrinos mass generation mechanism involving two energy scales beyond the Standard Model.

[Mohapatra, Valle 86]

– Texture

$$\begin{pmatrix} 0 & \mathbf{M}_D^T & 0 \\ \mathbf{M}_D & 0 & \mathbf{M}^T \\ 0 & \mathbf{M} & \mu \end{pmatrix}$$

–  $\mathbf{M}_D \sim$  electroweak scale,  $\mathbf{M} \sim 0.1\text{-}1$  TeV scale,  $\mu \sim$  keV scale (breaks lepton number).

$\Rightarrow m_{\nu} \sim$  sub-eV,  $m_{\nu\text{heavy}} \sim$  TeV.

Neutrinos mass generation mechanism involving two energy scales beyond the Standard Model.

[Mohapatra, Valle 86]

– Texture

$$\begin{pmatrix} 0 & \mathbf{M}_D^T & 0 \\ \mathbf{M}_D & 0 & \mathbf{M}^T \\ 0 & \mathbf{M} & \mu \end{pmatrix}$$

–  $\mathbf{M}_D \sim$  electroweak scale,  $\mathbf{M} \sim 0.1\text{-}1$  TeV scale,  $\mu \sim$  keV scale (breaks lepton number).

$\Rightarrow m_\nu \sim$  sub-eV,  $m_{\nu\text{heavy}} \sim$  TeV.

$$\mathcal{L} \supset \overline{N}_R \mathbf{M}_D \nu_L + \overline{S}_R \mathbf{M} N_R^c + \frac{1}{2} \overline{S}_R \mu S_R^c + \text{h.c.}$$



# The Inverse Seesaw Mechanism

Neutrinos mass generation mechanism involving two energy scales beyond the Standard Model.

[Mohapatra, Valle 86]

– Texture

$$\begin{pmatrix} 0 & \mathbf{M}_D^T & 0 \\ \mathbf{M}_D & 0 & \mathbf{M}^T \\ 0 & \mathbf{M} & \mu \end{pmatrix}$$

–  $\mathbf{M}_D \sim$  electroweak scale,  $\mathbf{M} \sim 0.1\text{-}1$  TeV scale,  $\mu \sim$  keV scale (breaks lepton number).

$\Rightarrow m_\nu \sim$  sub-eV,  $m_{\nu\text{heavy}} \sim$  TeV.

$$\mathcal{L} \supset \overline{N}_R \mathbf{M}_D \nu_L + \overline{S}_R \mathbf{M} N_R^c + \frac{1}{2} \overline{S}_R \mu S_R^c + \text{h.c.}$$

–  $N_{iR}, S_{iR}$  are six neutrino field singlets in addition to the SM one  $\nu_L$ .

# The Inverse Seesaw Mechanism

Neutrinos mass generation mechanism involving two energy scales beyond the Standard Model.

[Mohapatra, Valle 86]

– Texture

$$\begin{pmatrix} 0 & \mathbf{M}_D^T & 0 \\ \mathbf{M}_D & 0 & \mathbf{M}^T \\ 0 & \mathbf{M} & \mu \end{pmatrix}$$

–  $\mathbf{M}_D \sim$  electroweak scale,  $\mathbf{M} \sim 0.1\text{-}1$  TeV scale,  $\mu \sim$  keV scale (breaks lepton number).

$\Rightarrow m_{\nu} \sim$  sub-eV,  $m_{\nu\text{heavy}} \sim$  TeV.

$$\mathcal{L} \supset \overline{N}_R \mathbf{M}_D \nu_L + \overline{S}_R \mathbf{M} N_R^c + \frac{1}{2} \overline{S}_R \mu S_R^c + \text{h.c.}$$

- $N_{iR}$ ,  $S_{iR}$  are six neutrino field singlets in addition to the SM one  $\nu_L$ .
- It can leave a trace at the TeV scale.

# The Inverse Seesaw Mechanism

Neutrinos mass generation mechanism involving two energy scales beyond the Standard Model.

[Mohapatra, Valle 86]

– Texture

$$\begin{pmatrix} 0 & \mathbf{M}_D^T & 0 \\ \mathbf{M}_D & 0 & \mathbf{M}^T \\ 0 & \mathbf{M} & \mu \end{pmatrix}$$

–  $\mathbf{M}_D \sim$  electroweak scale,  $\mathbf{M} \sim 0.1\text{-}1$  TeV scale,  $\mu \sim$  keV scale (breaks lepton number).

$\Rightarrow m_{\nu} \sim$  sub-eV,  $m_{\nu\text{heavy}} \sim$  TeV.

$$\mathcal{L} \supset \overline{N}_R \mathbf{M}_D \nu_L + \overline{S}_R \mathbf{M} N_R^c + \frac{1}{2} \overline{S}_R \mu S_R^c + \text{h.c.}$$

–  $N_{iR}, S_{iR}$  are six neutrino field singlets in addition to the SM one  $\nu_L$ .

– It can leave a trace at the TeV scale.

– There is a symmetry leading to the required texture and scales from  $\langle \sigma \rangle \times \left( \frac{\langle \sigma \rangle}{M_{Pl}} \right)^m$

A Lagrangian involving the neutrinos fields,  $N_R$ ,  $S_R$ , and a charged heavy fermion  $E_R$  is found to be

A Lagrangian involving the neutrinos fields,  $N_R$ ,  $S_R$ , and a charged heavy fermion  $E_R$  is found to be

$$\begin{aligned} \mathcal{L} \supset & y_{ij} \overline{N_{iR}} \tilde{H}^\dagger L_j + \eta_{ij} \frac{\sigma^2}{M_{\text{Pl}}} \overline{S_{iR}} N_{jR}^c + \zeta_{ij} \frac{\sigma^{*3}}{2M_{\text{Pl}}^2} \overline{S_{iR}} S_{jR}^c \\ & + \kappa_i \frac{\sigma}{M_{\text{Pl}}} \overline{L}_i H E_R + \kappa \frac{\sigma^2}{M_{\text{Pl}}} \overline{E}_L E_R + \text{H.c.} \end{aligned}$$

$y_{ij}$ ,  $\eta_{ij}$ ,  $\zeta_{ii}$  of order  $\mathcal{O} \sim 0.1 - 1$ ,  $H$ ,  $L$  the Higgs lepton SM doublets.

A Lagrangian involving the neutrinos fields,  $N_R$ ,  $S_R$ , and a charged heavy fermion  $E_R$  is found to be

$$\begin{aligned} \mathcal{L} \supset & y_{ij} \overline{N_{iR}} \tilde{H}^\dagger L_j + \eta_{ij} \frac{\sigma^2}{M_{\text{Pl}}} \overline{S_{iR}} N_{jR}^c + \zeta_{ij} \frac{\sigma^{*3}}{2M_{\text{Pl}}^2} \overline{S_{iR}} S_{jR}^c \\ & + \kappa_i \frac{\sigma}{M_{\text{Pl}}} \overline{L}_i H E_R + \kappa \frac{\sigma^2}{M_{\text{Pl}}} \overline{E}_L E_R + \text{H.c.} \end{aligned}$$

$y_{ij}$ ,  $\eta_{ij}$ ,  $\zeta_{ii}$  of order  $\mathcal{O} \sim 0.1 - 1$ ,  $H$ ,  $L$  the Higgs lepton SM doublets.

**Spontaneous Symmetry Breaking:**  $\langle H \rangle = v/\sqrt{2}$ ,  $\langle \sigma \rangle = v_\sigma/\sqrt{2}$

A Lagrangian involving the neutrinos fields,  $N_R$ ,  $S_R$ , and a charged heavy fermion  $E_R$  is found to be

$$\mathcal{L} \supset y_{ij} \overline{N_{iR}} \tilde{H}^\dagger L_j + \eta_{ij} \frac{\sigma^2}{M_{Pl}} \overline{S_{iR}} N_{jR}^c + \zeta_{ij} \frac{\sigma^{*3}}{2M_{Pl}^2} \overline{S_{iR}} S_{jR}^c \\ + \kappa_j \frac{\sigma}{M_{Pl}} \overline{L}_j H E_R + \kappa \frac{\sigma^2}{M_{Pl}} \overline{E}_L E_R + \text{H.c.}$$

$y_{ij}$ ,  $\eta_{ij}$ ,  $\zeta_{ii}$  of order  $\mathcal{O} \sim 0.1 - 1$ ,  $H$ ,  $L$  the Higgs lepton SM doublets.

**Spontaneous Symmetry Breaking:**  $\langle H \rangle = v/\sqrt{2}$ ,  $\langle \sigma \rangle = v_\sigma/\sqrt{2}$

The value  $v_\sigma = 3 \times 10^{10}$  GeV gives the natural scales for the inverse seesaw mechanism!

$$\mathbf{M} = \frac{\eta v_\sigma^2}{2M_{Pl}} = \eta \times 187 \text{ GeV}, \quad \mu = \frac{\zeta v_\sigma^3}{2^{\frac{3}{2}} M_{Pl}^2} = \zeta \times 1.6 \text{ keV}$$

- Construction based on a specific discrete gauge symmetry [Krauss, Wilczek 89; Coleman et al 92]
- Discrete symmetry derived from conditions to cancel the anomalies:  
[Ibáñez, Ross 91; Banks, Dine 92]

$$\mathbb{Z}_{13} \equiv e^{i2\pi \frac{k}{13}}, \quad k = 0, 1, \dots, 12.$$

It allows for having an ultra-light ALP, and the required operators for the inverse seesaw mechanism.

$$\text{SM} \otimes \mathcal{U}(1) \otimes \dots$$

$$\Downarrow$$

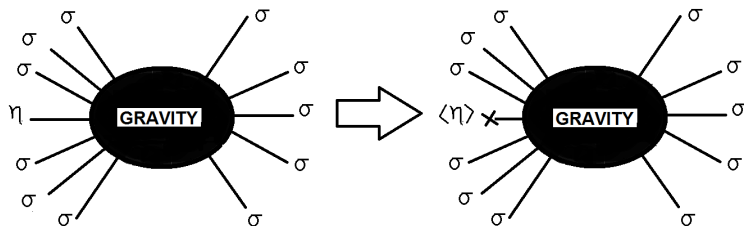
$$\text{SM} \otimes \mathbb{Z}_{13} \otimes \dots$$

$\mathcal{U}(1)$  respected by gravity  $\rightarrow$   $\mathbb{Z}_N$  respected by gravity as well.



## Inverse Seesaw Mechanism and ALP: required symmetry

Such “*discrete gauge symmetries*” might be respected by gravitational interactions. Local symmetries can masquerade as  $Z_N$  to an observer having only low energy probes [Krauss, Wilczek 89; Coleman et al. 92].



$$\mathcal{O}_{N+1} \equiv \frac{\eta(x) \sigma(x)^N}{M_{Pl}^{N-3}} \quad \Rightarrow \quad \mathcal{O}_N \equiv \frac{\langle \eta \rangle}{M_{Pl}^{N-3}} \sigma(x)^N$$

Being  $\mathcal{O}_{N+1}$  invariant under  $\mathcal{U}(1)$ ,  $\mathcal{O}_N$ , and all other interactions below the scale  $\langle \eta \rangle$ , will be invariant by  $Z_N$ .

- Fields charges under  $\mathbb{Z}_{13}$  symmetry:  $\psi_k \rightarrow e^{i2\pi\frac{Z_k}{13}}\psi_k$

$\mathbb{Z}_{13}$	$q_{iL}$	$d_{iR}$	$u_{iR}$	$H$	$L_i$	$l_{iR}$	$N_{iR}$	$S_{iR}$	$E_L$	$E_R$	$\sigma$
$Z_k$	2	2	2	0	7	7	7	-3	9	5	2

- $\mathbb{Z}_{13}$  charges in terms barion and an extension of lepton number:

$$Z_k = (6\mathcal{B} + 7\mathcal{L})_k$$

- Fields charges under  $\mathbb{Z}_{13}$  symmetry:  $\psi_k \rightarrow e^{i2\pi \frac{Z_k}{13}} \psi_k$

$\mathbb{Z}_{13}$	$q_{iL}$	$d_{iR}$	$u_{iR}$	$H$	$L_i$	$l_{iR}$	$N_{iR}$	$S_{iR}$	$E_L$	$E_R$	$\sigma$
$Z_k$	2	2	2	0	7	7	7	-3	9	5	2

- $\mathbb{Z}_{13}$  charges in terms barion and an extension of lepton number:

$$Z_k = (6\mathcal{B} + 7\mathcal{L})_k$$

- Global quasi-exact accidental symmetry:  $U(1)_X \supset U(1)_{\mathcal{L}}$

$U(1)_X$	$q_{iL}$	$d_{iR}$	$u_{iR}$	$H$	$L_i$	$l_{iR}$	$N_{iR}$	$S_{iR}$	$E_L$	$E_R$	$\sigma$
$X$	0	0	0	0	$\frac{7}{2}$	$\frac{7}{2}$	$\frac{7}{2}$	$-\frac{3}{2}$	$\frac{9}{2}$	$\frac{5}{2}$	1

- Fields charges under  $\mathbb{Z}_{13}$  symmetry:  $\psi_k \rightarrow e^{i2\pi \frac{Z_k}{13}} \psi_k$

$\mathbb{Z}_{13}$	$q_{iL}$	$d_{iR}$	$u_{iR}$	$H$	$L_i$	$l_{iR}$	$N_{iR}$	$S_{iR}$	$E_L$	$E_R$	$\sigma$
$Z_k$	2	2	2	0	7	7	7	-3	9	5	2

- $\mathbb{Z}_{13}$  charges in terms barion and an extension of lepton number:

$$Z_k = (6\mathcal{B} + 7\mathcal{L})_k$$

- Global quasi-exact accidental symmetry:  $U(1)_X \supset U(1)_{\mathcal{L}}$

$U(1)_X$	$q_{iL}$	$d_{iR}$	$u_{iR}$	$H$	$L_i$	$l_{iR}$	$N_{iR}$	$S_{iR}$	$E_L$	$E_R$	$\sigma$
$X$	0	0	0	0	$\frac{7}{2}$	$\frac{7}{2}$	$\frac{7}{2}$	$-\frac{3}{2}$	$\frac{9}{2}$	$\frac{5}{2}$	1

- $E_{L,R}$  chiral under  $U(1)_X$ , allowing for the ALP-photon coupling.

- ALP mass and effective coupling with electromagnetic field

$$V(\sigma) = g \frac{\sigma^{13}}{M_{Pl}^9} + \text{H.c.} \Rightarrow m_a \simeq 1.6 \times 10^{-16} |g|^{\frac{1}{2}} \left( \frac{v_\sigma}{3 \times 10^{10} \text{ GeV}} \right)^{5.5} \text{ eV}$$

– ALP mass and effective coupling with electromagnetic field

$$V(\sigma) = g \frac{\sigma^{13}}{M_{Pl}^9} + \text{H.c.} \Rightarrow m_a \simeq 1.6 \times 10^{-16} |g|^{\frac{1}{2}} \left( \frac{v_\sigma}{3 \times 10^{10} \text{ GeV}} \right)^{5.5} \text{ eV}$$

$$g_{a\gamma} \approx 1.5 \times 10^{-13} \left( \frac{3 \times 10^{10} \text{ GeV}}{v_\sigma} \right) \text{ GeV}^{-1}.$$

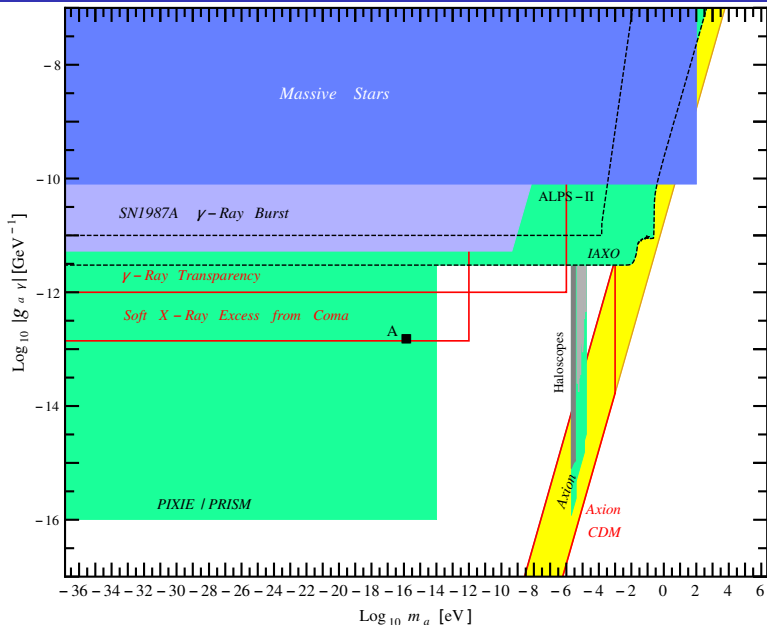
- ALP mass and effective coupling with electromagnetic field

$$V(\sigma) = g \frac{\sigma^{13}}{M_{Pl}^9} + \text{H.c.} \Rightarrow m_a \simeq 1.6 \times 10^{-16} |g|^{\frac{1}{2}} \left( \frac{v_\sigma}{3 \times 10^{10} \text{ GeV}} \right)^{5.5} \text{ eV}$$

$$g_{a\gamma} \approx 1.5 \times 10^{-13} \left( \frac{3 \times 10^{10} \text{ GeV}}{v_\sigma} \right) \text{ GeV}^{-1}.$$

- These values can provide explanation to the Soft X-ray excess from Coma cluster (being inside the region prospected by PIXIE/PRISM).

# A Model with one ALP and $\mathbb{Z}_{13}$ symmetry: benchmark point





- Heavy lepton mass controlled by the discrete symmetry

$$\mathcal{L} \supset \kappa_i \frac{\sigma}{M_{\text{Pl}}} \bar{L}_i H E_R + \kappa \frac{\sigma^2}{M_{\text{Pl}}} \bar{E}_L E_R + \text{H.c.}$$

$$\Downarrow$$

$$M_E = \kappa_E \frac{v_\sigma^2}{2M_{\text{Pl}}} \approx \kappa_E \times 187 \text{ GeV}$$

- Heavy lepton mass controlled by the discrete symmetry

$$\mathcal{L} \supset \kappa_i \frac{\sigma}{M_{\text{Pl}}} \bar{L}_i H E_R + \kappa \frac{\sigma^2}{M_{\text{Pl}}} \bar{E}_L E_R + \text{H.c.}$$

$$\Downarrow$$

$$M_E = \kappa_E \frac{v_\sigma^2}{2M_{\text{Pl}}} \approx \kappa_E \times 187 \text{ GeV}$$

- Lifetime

$$\tau_E \simeq \frac{32\pi}{3} \frac{M_{\text{Pl}}^2}{\kappa_i^2 v_\sigma^2} \frac{m_E^3}{(m_E^2 - m_h^2)^2} \times 6.5822 \times 10^{-25} \text{ s}$$

$$\leq 10^{-9} \text{ s} \quad \text{[CMS 2013]}$$

$$m_h = 125 \text{ GeV}, \quad m_E \gtrsim 250 \text{ GeV}$$

- Decay channel:  $E \rightarrow h + e^-$

- ① Appropriated texture and mass scales for the inverse seesaw mechanism triggered by  $v_\sigma$ ;

- 1 Appropriated texture and mass scales for the inverse seesaw mechanism triggered by  $\nu_\sigma$ ;
- 2 Ultra-light ALP which could explain the soft X-ray excess from Coma cluster, and well inside the PIXIE/PRISM prospective region;

- 1 Appropriated texture and mass scales for the inverse seesaw mechanism triggered by  $v_\sigma$ ;
- 2 Ultra-light ALP which could explain the soft X-ray excess from Coma cluster, and well inside the PIXIE/PRISM prospective region;
- 3 Heavy lepton  $E$  with mass near electroweak scale;

- 1 Appropriated texture and mass scales for the inverse seesaw mechanism triggered by  $v_\sigma$ ;
- 2 Ultra-light ALP which could explain the soft X-ray excess from Coma cluster, and well inside the PIXIE/PRISM prospective region;
- 3 Heavy lepton  $E$  with mass near electroweak scale;
- 4 Interesting restrictions for ALP parameters may be obtained using a scheme with  $\mathbb{Z}_N$ , or any other sort of, symmetries producing appropriate textures for neutrinos mass matrix/mixing.

- 1 Appropriated texture and mass scales for the inverse seesaw mechanism triggered by  $v_\sigma$ ;
- 2 Ultra-light ALP which could explain the soft X-ray excess from Coma cluster, and well inside the PIXIE/PRISM prospective region;
- 3 Heavy lepton  $E$  with mass near electroweak scale;
- 4 Interesting restrictions for ALP parameters may be obtained using a scheme with  $\mathbb{Z}_N$ , or any other sort of, symmetries producing appropriate textures for neutrinos mass matrix/mixing.

Thank you!