F-theory at order α'^3

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> Based on work with: R. Minasian and T. Pugh arXiv:1506.06756

In this talk...

- I discuss quantum corrections (in g_s and α '). Need control over them to understand vacuum structure of string theory & its EFT.
- Here I focus on the Type IIB corner of the string landscape:
 - \exists efficient probe into the strong g_s regime \rightarrow F-theory; [Vafa `96]
 - All distinctive features of GUT models can be accommodated; [Donagi,Wijnholt; Beasley, Heckman, Vafa `08]
 - One has the Large Volume Scenario as promising paradigm of moduli stabilization. [Balasubramanian, Berglund, Conlon, Quevedo `05]
- Final target is the 4D, N=1 EFT of Type IIB on CY₃ with D7/O7:
 - Knowledge of quantum corrections is still very limited, especially for the effective Kähler potential.
 - They play a key role in the phenomenology of these vacua.

Context

• Most notable example is a certain $O(\alpha^{3})$ correction to the CY₃ volume:

$$K = -2\log\left(\mathcal{V}_3 - \frac{\zeta(3)\,\chi(CY_3)}{32\pi^3\,g_s^{3/2}}\right)$$

- Due to closed strings \Rightarrow N=2 sector; [Antoniadis, Ferrara, Minasian, Narain `97]
- Shown to survive orientifolding $N=2 \rightarrow N=1$. [Becker, Becker, Haack, Louis `02]
- <u>Main aim</u>: Study possible, genuinely N=1 modifications of CY₃ volume.
- Using F-theory is convenient, as:
 - It "geometrizes" all g_s effects of Type IIB, and it only needs to be corrected in $\alpha\space$;
 - It includes open string effects, through 7-brane backreaction.

F-theory

12D theory: Auxiliary T² fibered over the 10D string space

T² cplx-str = axio-dilaton $\tau = C_0 + ie^{-\phi}$

varies holomorphically with $SL(2, \mathbb{Z})$ transitions



 \nexists fundamental description (yet): \nexists (1,11) sugra.

Still, I2D are able to encode $SL(2,\mathbb{Z})$ -invariant I0D structures.

Two derivatives

- Let's warm up with $O(\alpha^{\prime 0})$.
 - Start with the I2D Ricci scalar, integrated on I0D subspace.

$$S_0^{(12)} \sim \frac{1}{l_s^8} \int R_{\rm sc}^{(12)} *_{10} 1 \quad \stackrel{\mathsf{T}^2}{\stackrel{\text{red.}}{\longrightarrow}} \quad S_0^{(10)} \sim \frac{1}{l_s^8} \int \left(R_{\rm sc}^{(10)} - 2P \cdot \bar{P} \right) *_{10} 1$$

KK-reduction using the T² metric $\frac{1}{\operatorname{Im} \tau} \left(\begin{array}{cc} 1 & \operatorname{Re} \tau \\ \operatorname{Re} \tau & |\tau|^2 \end{array} \right)$
 $P \text{ is the U(1)}_{\mathsf{R}} \text{ -covariant quantity } P = \frac{i}{2\operatorname{Im} \tau} \nabla \tau \text{ of charge 2.}$

- $U(I)_{R}$ -invariance + IOD diff-invariance \Rightarrow "I2D diff-invariance"
- Warning : No I2D lift of the I0D measure !

Eight derivatives

- At $O(\alpha^{3})$ we have non-trivial string dynamics.
 - Start with the "I2D" action:

$$S_3^{(12)} = \frac{1}{(2\pi)^7 \cdot 3 \cdot 2^{11} \cdot l_s^2} \int f_0(\tau, \bar{\tau}) \left[t_8 t_8 + \frac{1}{96} \epsilon_{12} \epsilon_{12} \right] (R^{(12)})^4 *_{10} 1$$

- ► f_0 holds all g_s corrections compatible with SUSY + SL(2)-invariance. For $(\operatorname{Im} \tau)^{-1} = g_s << 1$: $f_0(\tau, \bar{\tau}) \approx \underbrace{\frac{2\zeta(3)}{g_s^{3/2}}}_{\text{tree-level}} + \underbrace{\frac{2\pi}{3}g_s^{1/2}}_{1-\operatorname{loop}} + \underbrace{\mathcal{O}(e^{-1/g_s})}_{D(-1)}$
- Reduction on $T^2 \Rightarrow$ complicated sum of $U(I)_R$ -invariant couplings, written in terms of $R^{(10)}$, P and DP.
- Perfect match with known 4pt, tree-level result. [Policastro, Tsimpis `06 `08]
- Prediction of IOD couplings of $R \& \tau$ beyond 4pt.

4D N=I compactifications

- Reduce F-theory on smooth CY_4 , elliptically fibered over B_3 , with zero-section: $B_3 \rightarrow CY_4$.
 - Eight-derivative action yields correction to volume of base:

$$S_{0+3}^{(4)} = \frac{1}{2\pi\alpha'} \int \left(\mathcal{V}_b - \frac{1}{64\pi^3} \int_{B_3} f_0(\tau,\bar{\tau}) c_3(CY_4) |_{B_3} \right) R_{sc}^{(4)} *_4 1$$

- For constant $\tau \Rightarrow \text{old N=2 correction } (B_3 = CY_3);$
- N=I vacua: Correction is non-topological (τ varies over B₃);
- Weyl rescaling induces correction to 4D, N=1 Kähler potential.

Weak coupling

- Correction simplifies when going to weak string coupling.
 - Sen (`97): Restrict CY₄-complex structure s.t. $g_s \rightarrow 0$ is well-defined.
 - → Type IIB on $CY_3 \rightarrow B_3$ branched double cover.
 - \rightarrow O7-plane wrapping branch locus: In cohomology $D_{O7} = C_1(B_3)$.

$$\left| \tilde{\mathcal{V}}_3 = \mathcal{V}_3 - \frac{\zeta(3)}{32\pi^3 g_s^{3/2}} \left(\chi(\mathrm{CY}_3) + 2 \int_{\mathrm{CY}_3} D_{\mathrm{O7}}^3 \right) + \mathcal{O}(g_s^{-1/2}) \right|$$

- Topological correction at closed-string tree-level.
- Next-to-leading: tree-level of open + unorientable strings.
- Arises from graviton two-point function in orientifold backgrounds.
- Absent in toroidal models!

Caveats / Open questions

- Our I2D $O(\alpha'^3)$ action is only checked at 4pt!
 - Test it beyond 4pt : 5pt amplitude computation or M/F duality.
- Employ I2D logic to study the F_3 / H_3 -flux sector in I0D at $O(\alpha'^3)$.
- Correction to Weyl rescaling does not fully determine Kähler potential!
 - Derive corrections to kinetic terms of moduli. [Berg, Haack, Kang, Sjörs `14]
- Include corrections to the vacuum solution, such as warping effects.
 [Grimm, Pugh, Weissenbacher `14]; [Martucci `14]
- Our correction persists on singular CY₄! Are there more? [Junghans, Shiu`14]
- Study the 4D scalar potential. [Ciupke, Louis, Westphal `15]