# INTEGRATING IN THE HIGGS PORTAL TO FERMION DARK MATTER

based on JHEP 1509 (2015) 015 with Ayres Freitas and Jure Zupan

Susanne Westhoff



#### Universität Heidelberg

Workshop "Physics at the LHC and beyond" — Sep 29 - Oct 2, 2015 — DESY Hamburg

### HIGGS-PORTAL DARK MATTER AT THE LHC



### FERMION DARK MATTER



$$\mathcal{L}_{\text{eff}} = \frac{g_S}{\Lambda} (\bar{\chi}\chi) (H^{\dagger}H) + i \frac{g_P}{\Lambda} (\bar{\chi}\gamma_5\chi) (H^{\dagger}H)$$

Higgs portal interaction is not renormalizable.
 → UV completion includes mediator.

### Can we see Higgs-portal mediators at colliders?



[image credit: CERN]

## UV COMPLETIONS OF THE HIGGS PORTAL



Mediators with mass up to a few 100 GeV: Higgs portal is ,,open'' at the LHC.

### EXAMPLE: SINGLET-DOUBLET MODEL

Dark fermions mix through Yukawa interaction:

$$\mathcal{L} = -m_D \overline{\psi}_D \psi_D - m_S \overline{\psi}_S \psi_S - (y \overline{\psi}_D H \psi_S + \text{ h.c.})$$

$$\begin{array}{l} \langle H \rangle = v/\sqrt{2} \\ & \longrightarrow \end{array} \end{array} \qquad \left[ \begin{array}{c} \chi_l^0 = \cos \theta \psi_D^0 - \sin \theta \psi_S \\ \chi_h^0 = \sin \theta \psi_D^0 + \cos \theta \psi_S \end{array} \right]$$

Mixing controls coupling to Higgs and gauge bosons:



 $Z_2$  symmetry ensures DM stability.

### OBSERVING PARTICLE DARK MATTER

#### annihilation

relic density and indirect detection

scattering direct detection



production colliders



Spin-independent DM-nucleus scattering:

![](_page_6_Picture_2.jpeg)

$$\sigma_{0} = \frac{\mu_{A}^{2}}{\pi} |Zf_{p} + (A - Z)f_{n}|^{2}$$

 $\sqrt{k^2} \simeq 10 - 50 \,\mathrm{MeV} \ll M_{\mathrm{EW}}$ 

### Currently strongest bound on weak-scale DM scattering: LUX experiment: $\sigma_0(m_\chi \approx 100\,{ m GeV}) \lesssim 10^{-45}{ m cm}^2$ [LUX coll., arXiv:1310.8214]

![](_page_6_Picture_6.jpeg)

[picture: lux.brown.edu]

### DARK FERMION-NUCLEON SCATTERING

![](_page_7_Figure_1.jpeg)

![](_page_7_Figure_2.jpeg)

#### Effective interactions:

$$f_{p,n} \sim \frac{g_{\chi}^Z g_q^Z}{M_Z^2}$$

Dirac singlet:

Majorana singlet:  $g_{\chi}^{Z} = -\frac{g}{2c_{W}}\cos^{2}\theta$ 

$$g_{\chi}^Z = 0$$

$$f_{p,n} \sim \frac{g_{\chi}^h}{M_h^2} \frac{m_q}{v}$$

$$g_{\chi}^{h} = \frac{y}{\sqrt{2}}\sin(2\theta)$$

$$g_{\chi}^{h} = \frac{y}{2}\sin(2\theta')$$

LUX bound  $\rightarrow$  DM must be singlet-like,  $\theta \approx \pi/2$ .

### RELIC DENSITY

Dirac dark matter annihilation:  $\chi \bar{\chi} \to Z \to q \bar{q}, \ell^+ \ell^-$ 

Majorana dark matter:

![](_page_8_Figure_3.jpeg)

Observed relic density:  $\Omega_{\chi}h^2=0.1199\pm0.0022$  [Planck coll., arXiv:1502.01589]

LUX bound strongly constrains annihilation rate.

Co-annihilation  $\chi_l^0 \chi^+$ ,  $\chi_l^0 \chi_m^0$  prevents over-abundance.

Exception: Higgs-resonance region for Majorana DM.

## DARK FERMION SEARCHES AT COLLIDERS

Relic density and direct detection: Small mass splittings  $m_m^0 - m_l^0, m^+ - m_l^0 \longrightarrow$  soft decay products.

Hard jet helps to trigger on soft-lepton events:

![](_page_9_Picture_3.jpeg)

[Schwaller, Zurita, arXiv:1312.7350, et al.]

Cross section too small for mono-jet searches at the LHC.

### SUMMARY DARK DIRAC FERMIONS

![](_page_10_Figure_1.jpeg)

Need high-energy collider to test this model conclusively.

### SUMMARY DARK MAJORANA FERMIONS

![](_page_11_Figure_1.jpeg)

Future direct detection experiments and/or high-energy collider can test this model.

![](_page_12_Figure_0.jpeg)

- Mediators can be searched for in signatures with soft leptons at the LHC.
- Future lepton and high-energy hadron colliders are needed to test such models conclusively.
- **Direct detection** experiments provide complementary information.