

Non-thermal cosmic neutrino background

Andreas Trautner

based on arXiv:1509.00481
with: Mu-Chun Chen (UCI) and Michael Ratz (TUM).



DESY Theory WS

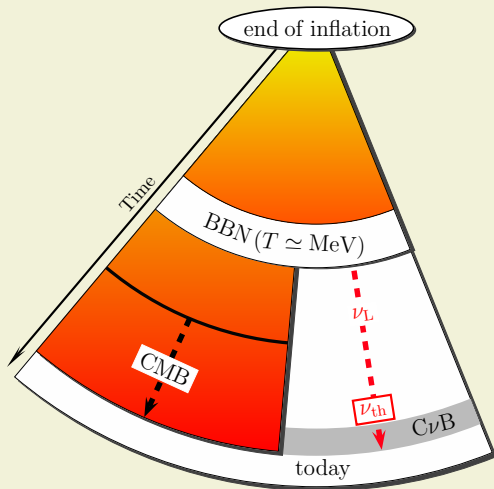
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Outline

- Review of CMB and thermal $C_{\nu B}$
- Idea of a non-thermal neutrino background
- Constraints from $N_{\text{eff}} \Rightarrow$ maximal relic density
- Detection prospects
- Possible origin: inflationary preheating
- Conclusion

Standard picture



- CMB

$$t_{\text{dec}} \sim 3.8 \cdot 10^5 \text{ a}$$

today:

$$T_\gamma = 2.73 \text{ K} \simeq 2.35 \cdot 10^{-4} \text{ eV}$$

$$n_\gamma = 410 \text{ cm}^{-3}$$

- standard CνB

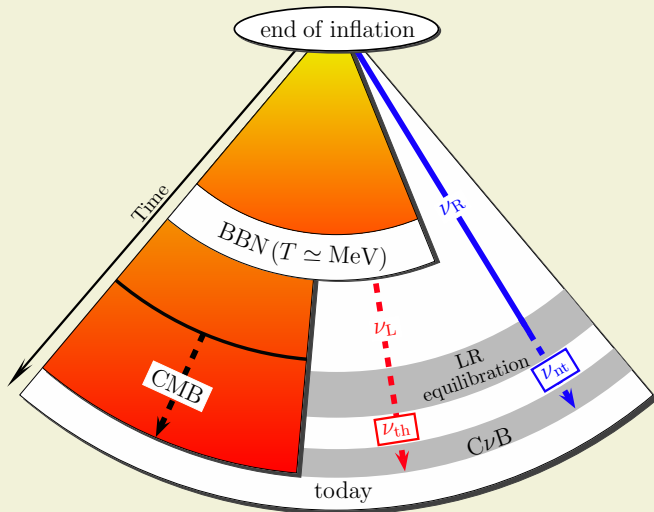
$$t_{\text{dec}} \sim 1 \text{ s } (T_{\text{dec}} \sim 1 \text{ MeV})$$

today:

$$T_{\nu_L} = T_\gamma \cdot (4/11)^{1/3} \sim 1.95 \text{ K}$$

$$n_{\nu_L} \sim 336 \text{ cm}^{-3}$$

Standard picture + non-thermal Dirac Neutrinos



Standard picture

Let's assume **Dirac** neutrinos

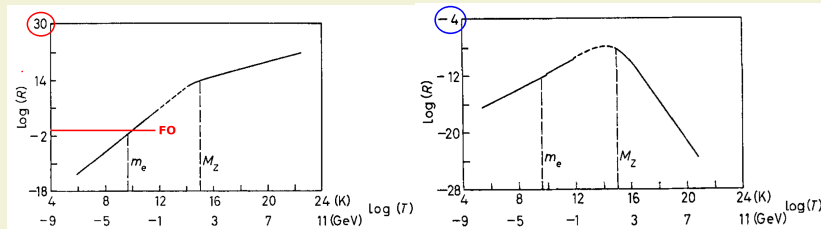
$$\mathcal{L}_\nu = Y_\nu^{ij} \begin{pmatrix} \bar{e}_L^i \\ \bar{\nu}_L^i \end{pmatrix} \cdot \tilde{H} \nu_R^j + \text{h.c.} .$$

$$\text{Planck} + \Lambda\text{CDM} \implies m_\nu \lesssim \mathcal{O}(0.1) \text{ eV} \iff \text{E.V.}(Y_\nu) \lesssim 10^{-12}$$

Note: Standard C ν B initially consists only of **left-chiral** ($\equiv$ **left-handed**) neutrinos. ν_R are too weakly coupled!

ν_L

ν_R



($R := \Gamma/H$)

[Antonelli, Fargion, Konoplich '81]

Non-thermal background

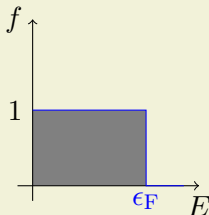
$\Gamma_{\nu_R} \ll H$ implies:

- **No thermal production** of ν_R , but also
- **No thermalization** of existing abundance of ν_R !

↪ Assume that there **is** a **non-thermal** abundance of ν_R ...

Easiest thing to do:

Fill ν_R states from the bottom up $\iff$ degenerate Fermi gas.



$$n_{\nu_R} = \frac{g}{6\pi^2} \varepsilon_F^3$$

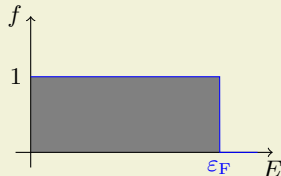
$$\rho_{\nu_R} = \frac{g}{8\pi^2} \varepsilon_F^4$$

(ultrarelativistic approx.)

($g = 2$ for a Weyl fermion)

Evolution of the non-thermal background

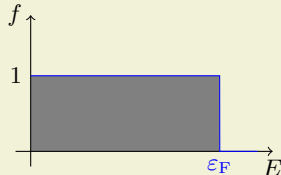
$$n_{\nu_R}(T) = n_{\nu_R}(T_{RH}) \cdot \left(\frac{R(T_{RH})}{R(T)} \right)^3$$



Evolution of the non-thermal background

$$n_{\nu_R}(T) = \frac{g \xi^3}{6 \pi^2} \frac{g_{*S}(T)}{g_{*S}(T_{RH})} T^3$$

$$\xi := \varepsilon_F / T_{RH}$$

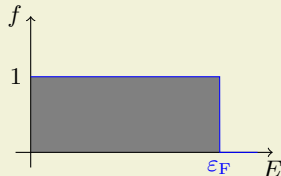


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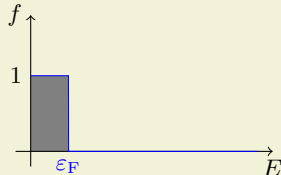
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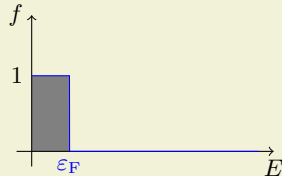


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$$\frac{n_{\nu_R}(T_\gamma)}{n_\gamma} = \frac{g \xi^3}{12 \zeta(3)} \frac{g_{*S}(T_\gamma)}{g_{*S}(T_{RH})}$$

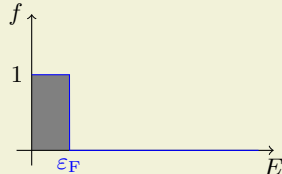
$$\Delta N_{\text{eff}}^{(\nu_R)} = \frac{8}{7} \frac{30}{8 \pi^4} \frac{g \xi^4}{2} \left(\frac{g_{*S}(T_{BBN})}{g_{*S}(T_{RH})} \right)^{4/3}$$

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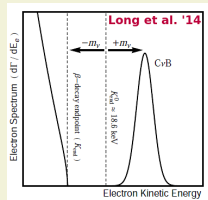
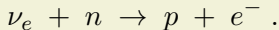
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$$\Delta N_{\text{eff}} \lesssim 0.7 \implies n_{\nu_R}(T_\gamma) \lesssim 0.53 n_\gamma \approx 217 \text{ cm}^{-3}$$

Planck: $\Delta N_{\text{eff}} = 0.2 \pm 0.5$ (95%CL) [Ade et al.'15]

Detection of the non-thermal background

PTOLEMY collaboration: proposal to measure C ν B. [Betts et al.'13]
 Capture of relic neutrinos via inverse β -decay



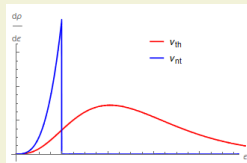
- relic ν 's are **non-relativistic**

⇒ Chiralities mix via the mass term; $\nu_L \rightarrow \nu_{th}$, $\nu_R \rightarrow \nu_{nt}$
 (Helicity, not chirality is conserved)

- Thermal** and **non-thermal** neutrinos **indistinguishable** by experiment with resolution $\mathcal{O}(0.1)$ eV.

$$\langle E_{\nu_{nt}} \rangle \lesssim \langle E_{\nu_{th}} \rangle \ll m_\nu \lesssim \mathcal{O}(0.1) \text{ eV} .$$

⇒ **Non-thermal** neutrinos give irreducible contribution to C ν B measurement.



Detection of the non-thermal background

- Proposal: [Long, Lunardini, Sabancilar '14]
Distinguish Dirac vs. Majorana ν 's by different $C\nu B$ count rate

$$\Gamma^M = 2\Gamma^D \approx 8\text{ yr}^{-1}.$$

- If there are **non-thermal** neutrinos, the count rate Γ^D could be enhanced by up to $\sim 64\%$.

$\Rightarrow$ Distinction could be inconclusive with low statistics.

Other effects:

- Relic neutrino clustering $\Rightarrow \mathcal{O}(1)$ factor. [Ringwald, Wong '04]
- SM extensions with Dirac neutrinos: ν_R could also be **thermally** coupled.
eg. [Anchordoqui, Goldberg, Steigman '12],
[Solaguren-Beascoa, Gonzalez-Garcia '12]
Maximal enhancement in this case: $\sim 28\%$. [Zhang&Zhou '15]

Production mechanism

Question: Is there a (well-motivated) mechanism to generate a **non-thermal** spectrum of fermions?

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✓ YES: “fermionic preheating”. [Greene&Kofman '98]
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Ingredients:

- + Inflation (chaotic, hybrid,...), e.g. $V(\phi) \sim m_\phi^2 \phi^2$
- + Inflaton coupling to fermion field $\lambda \phi \bar{\Psi} \Psi$.
- ↪ (non-perturbative) “parametric resonance” effect

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Theory of fermionic preheating

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... In an expanding universe, parametric excitation of fermions is stochastic. Created fermions very quickly, within tens of inflaton oscillations, fill up a sphere of radius $\approx q^{1/4}$ in momentum space. ...

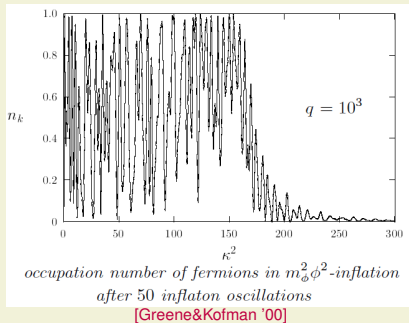
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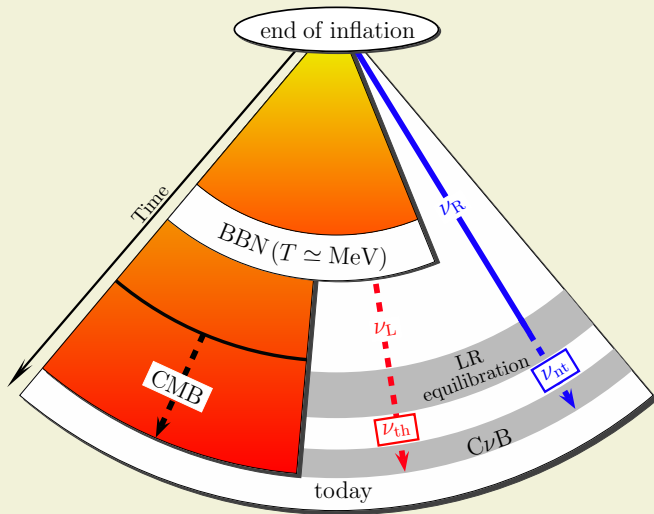
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$$q := \lambda^2 \phi_0^2 / m_\phi^2$$

$$\varepsilon_F \sim q^{1/4} m_\phi$$

Summary



- **non-thermal** CνB: $t_{\text{creation}} \sim t_{\text{infl.}}$; Today: $n_{\nu_{nt}} \lesssim 217 \text{ cm}^{-3}$.

Thank You!

Backup slides

Evolution of Spectra

Relic density

$$\frac{n_{\nu_{\text{nt}}}}{n_{\gamma}} \approx 1.2 \frac{g_{*S}(T_{\gamma})}{g_{*S}(T_{\text{BBN}})} \eta^{1/4} g^{1/4} \left(\Delta N_{\text{eff}}^{(\nu_{\text{R}})} \right)^{3/4}.$$

η : filling factor of Fermi-gas

Production mechanism

In our case we need a coupling

$$\mathcal{L} \supset \lambda \phi \overline{\nu_R^c} \nu_R + \text{h.c.} .$$

Neglecting spacial expansion, but for $\sim$ realistic ξ :

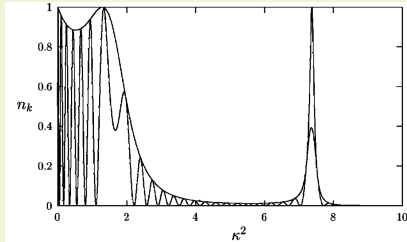


FIG. 2. The occupation number of fermions in $m_\phi^2 \phi^2$ theory as a function of κ^2 after 10 inflaton oscillations for resonance parameter $q = 10$.

[Greene&Kofman '00]

Reheating of the SM: via perturbative decay of ϕ , or $\phi^2 H^2$ coupling and the “scalar” parametric resonance.

[Brandenberger&Traschen '90]

[Kofman, Linde, Starobinsky '94]

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