

Two-Loop integrals for precision Higgs boson phenomenology

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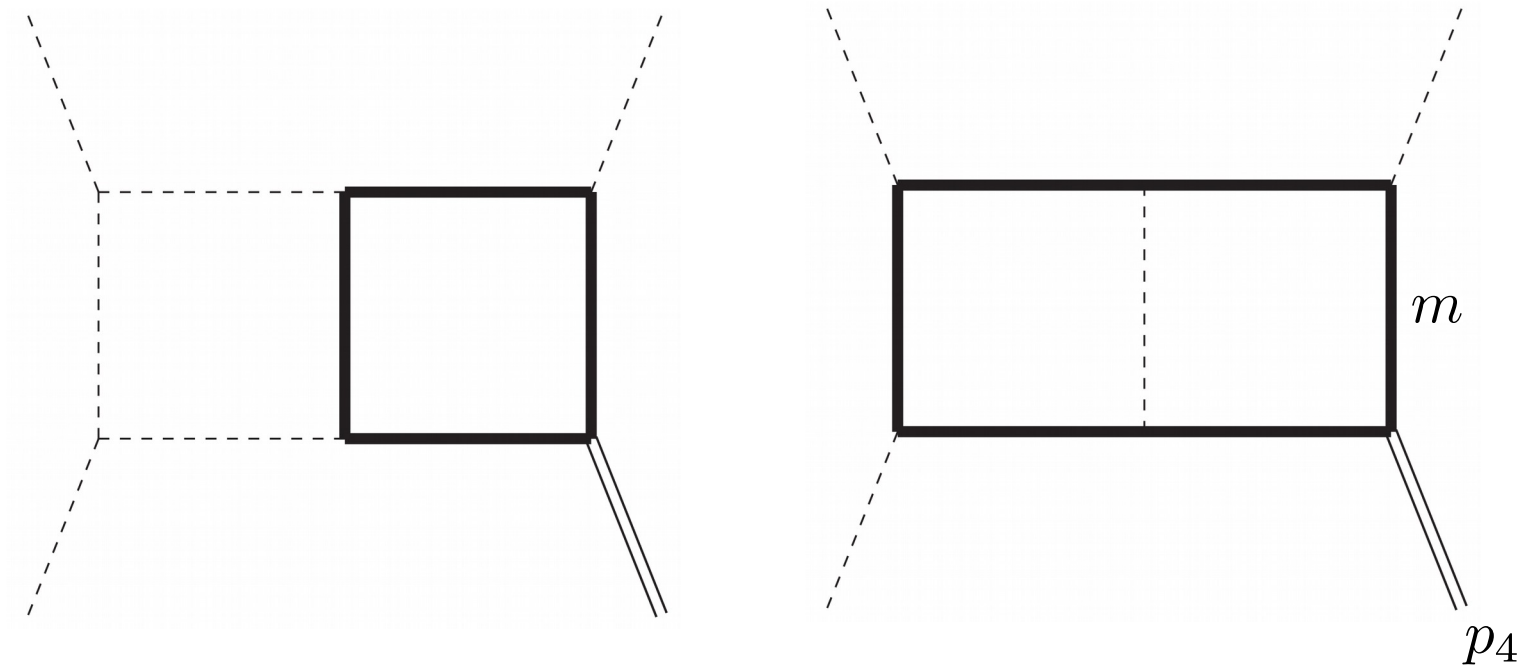
Based on work in collaboration with

R. Bonciani, V. Del Duca, H. Frellesvig, J.M. Henn, V. Smirnov

Loops and Legs

Leipzig, 26 April 2016

Integrals in this talk



$$(s, t, p_4^2, m^2)$$

Motivations

- NLO QCD corrections to Higgs decay width in a massive boson and a photon, with exact top quark mass
[Bonciani, Del Duca, Frellesvig, Henn, FM, Smirnov, 2015]
[Gehrmann et al, 2015]
- Higgs + jet at NLO, with exact top quark mass

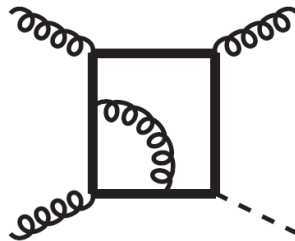
$$gg \rightarrow Hg$$

2 Loop QCD corrections with exact top mass

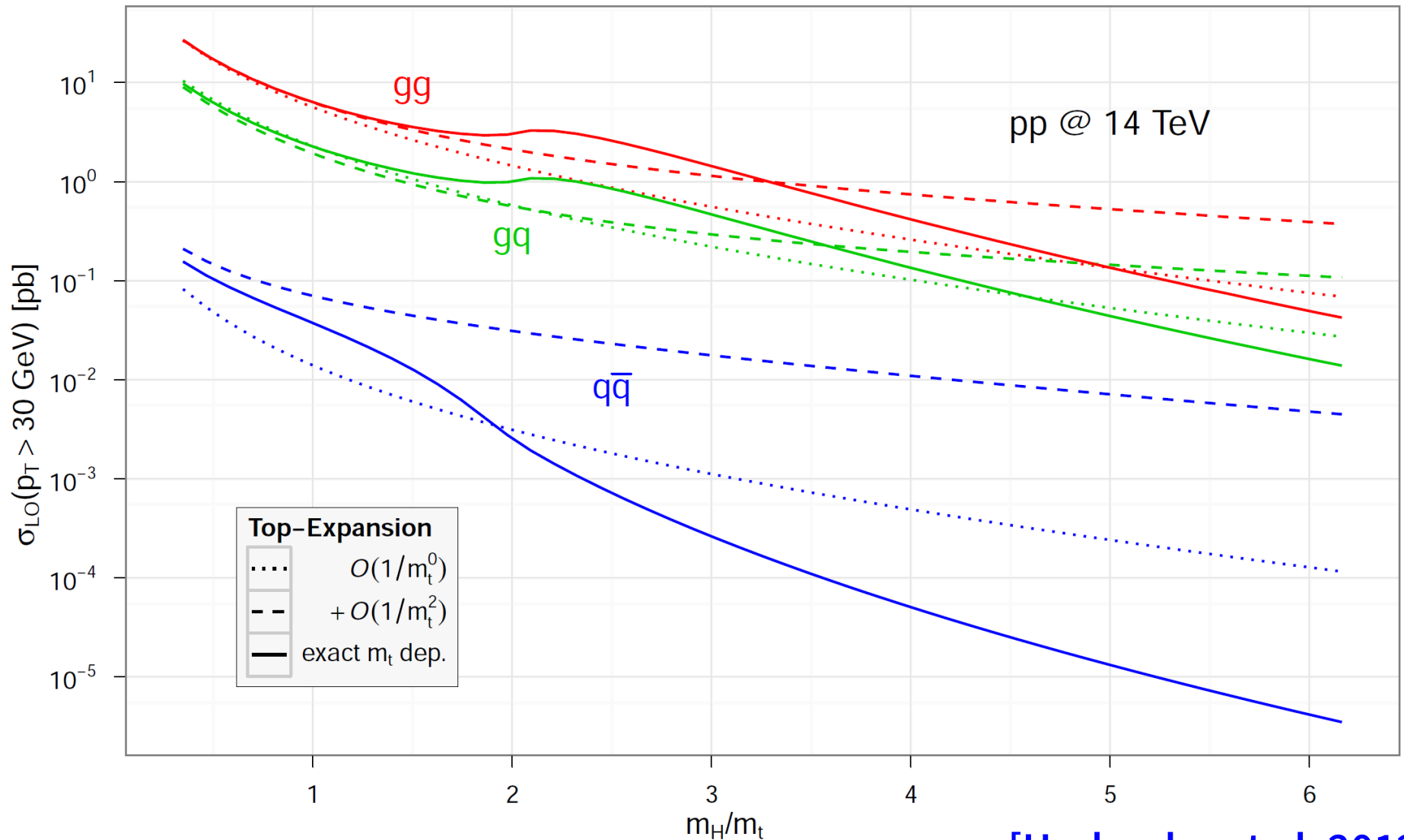
Theoretical uncertainties in this channel are among the largest errors of Higgs analysis ! [\[Heinemeyer et al,2013\]](#)

High transverse Higgs momentum (p_T) distribution has been shown to be very sensitive to new physics effects in many BSM models [\[Azatov et al,2013\]](#)[\[Grojean et al,2014\]](#)....

Known analytically at NLO only in the infinite top mass limit [\[Gehrmann et al,2012\]](#)



Finite top mass effects in H+jet



[Harlander et al, 2012]
[Frederix et al, 2016]

$$H \rightarrow Z\gamma$$

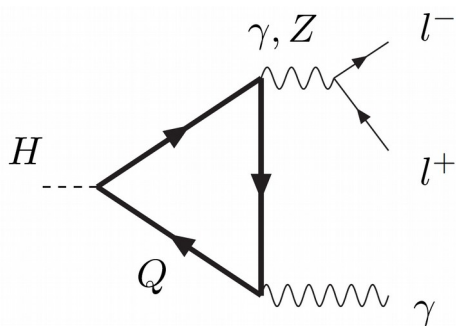
2 Loop QCD , EW light quark corrections with exact top mass

Master integrals are a subset of gg -> Hg

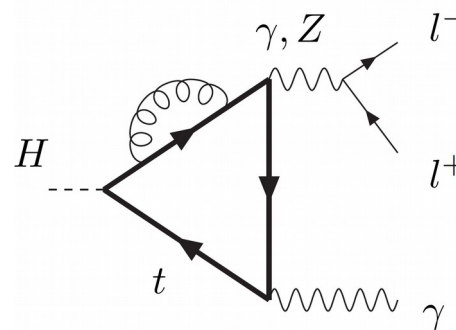
This channel is suitable for precision experiments

Lepton pair disentangled from the QCD background

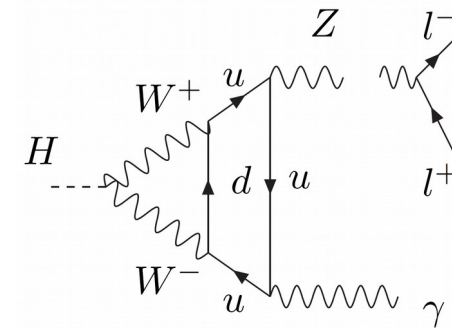
[CMS coll., 2013]



err ~ 10%



err ~ 5%



[Denner et al 2011]

Outline of the talk

- **Differential equations**
- **Canonical form**
- **Different representations of the solution**
 - Chen iterated integrals
 - (DE cannot be directly integrated in terms of GPL, even if not elliptic)
- **Analytic continuation**
 - Automatic analytic continuation

Differential Equations (notation)

- Compute the derivative of (n) Feynman integrals
[Kotikov, 1991] [Remiddi, 1997] [Gehrmann, Remiddi, 2000]

$$\partial_x f(x, \epsilon) = A(x, \epsilon) f(x, \epsilon)$$

- Each integral satisfies an n-th order DE
- Equivalent to a 1-st order $n \times n$ system of DE

Canonical form

- Direct solution is not feasible
- How do we decouple it ?
- An effective approach is decoupling the different eps orders of the solution [[Henn, 2013](#)]

$$\partial_x f(x)_i = \epsilon A(x) f(x)_{i-1}$$

- Solution in terms of Chen iterated integrals [[Chen, 1977](#)]
(path encodes boundaries)

$$f_i(x) = \epsilon \int_C dA(x) f_{i-1}$$

Canonical form - algorithms

- Uniform weight Integrals satisfy canonical DE
(\mathbb{Q} -linear combinations of transcendental functions of same weight)
- Integrals with constant leading singularities should have uniform weight
(Maximal cut)
[\[Arkani-Hamed, Bourjaily, Cachazo, Trnka, 2010\]](#) [\[Arkani-Hamed et al, 2012\]](#)
- Univariate problems in Fuchsian form (algorithmic)
[\[Lee, 2014\]](#)
- Linear DE (algorithmic) [\[Argeri et al, 2014\]](#)
$$\partial_x f(x) = (A_0(x) + \epsilon A_1(x))f(x)$$
- Block diagonal linear DE [\[Gehrmann et al, 2014\]](#)

Chen integrals, polylogarithms

- N fold integrals

$$f_i(x) = \epsilon \int_C dA(x) f_{i-1}$$

- Solve the DE in terms of Goncharov polylogarithms: numeric evaluation [Vollinga et al., 2005], Hopf algebra structure [Goncharov, 2005] (**only if** DE are in Fuchsian form)

$$\partial_x f(x)_i = \epsilon \sum_k \frac{a_k}{x - c_k} f(x)_{i-1} \quad f_i = \epsilon \sum_k a_k \int_C d \log(x - c_k) f_{i-1}$$

constant matrices
independent of x

- Are we always this lucky ? No (of course)

Higgs -> Zeta Gamma

- D-log (trivial algebraic functions)

$$\sqrt{x(x-4)}$$

- Landau variable makes Fuchsian form manifest (Goncharov Polylogarithms)

$$x(y) = -\frac{(1-y)^2}{y}$$

Higgs + jet

- 6 independent square roots:

$$\begin{aligned} & \sqrt{s(s - 4m^2)}, \quad \sqrt{m^4 (p_4^2 - s)^2 + 2m^2 st (p_4^2 - s) + st^2 (s - 4m^2)}, \\ & \sqrt{t(t - 4m^2)}, \quad \sqrt{st (4m^2 p_4^2 - 4m^2 (s + t) + st)}, \\ & \sqrt{p_4^2 (p_4^2 - 4m^2)}, \quad \sqrt{(p_4^2 - s) (-4m^2 + p_4^2 - s)}. \end{aligned}$$

- “No change of variables removes them” !
- No Fuchsian form
- **We cannot integrate DE in terms of Goncharov polylogarithms**

From DE to Symbols

- Thanks to the very peculiar form of DE we are still able to find a solution in terms of polylogarithms

DE in d-log form

$$df(x)_i = \epsilon \sum_k a_k d \log(\beta_k(x)) f(x)_{i-1}$$

constant matrices

logarithms (alphabet)

“Chen integral solution = Symbol” !

$$f_i = \epsilon \sum_k a_k \int_C d \log(\beta_k(x)) f_{i-1}$$

$$S(f_i) = \epsilon \sum_k a_k \beta_k(x) \otimes f_{i-1}$$

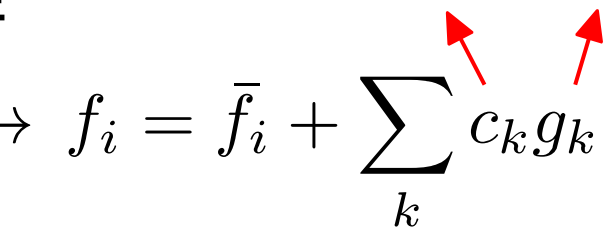
Higgs + jet (polylogarithms)

- We can use algorithms to find the corresponding polylogarithmic functions [[Goncharov et al. 1998-2010](#)][[Duhr et al 2011](#)]


$$S(f_i) \rightarrow \bar{f}_i$$

- We are still missing terms in the kernel of the symbol

- We can extract the information from DE free coefficients lower w. func.

$$\partial_x f(x)_i = \epsilon dA(x) f(x)_{i-1} \rightarrow f_i = \bar{f}_i + \sum_k c_k g_k$$



- Pure constants are fixed imposing boundary conditions

$$f_i = \bar{f}_i + \sum_k c_k g_k + b$$


Higgs + jet (polylogarithms)

- The alphabet of H+jet is large (~ 50 letters)
- Finding the function matching a given symbol amounts to constructing a general ansatz (polylogarithms arguments)

[Duhr et al 2011]

$$\text{Li}_n(a), \text{Li}_{22}(a, b) \quad a, b = \prod_i \beta_i^{m_i} \quad (m_i \in \mathbb{N})$$


DE logarithms arguments (alphabet)

$$1 - a, 1 - b = \prod_i \beta_i^{m'_i}, \quad 1 - ab = \prod_i \beta_i^{m''_i}$$

arguments factorize over the alphabet

- The combinatorics is very large, solving for the coefficients of the ansatz can be challenging!

Higgs + jet (PL+Chen integrals)

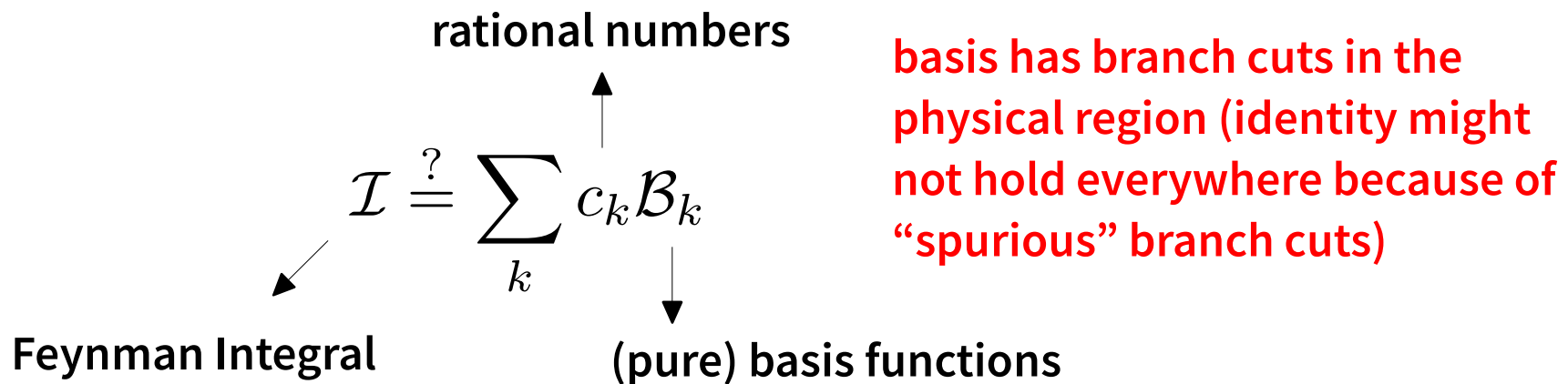
- Up to weight 2 solution in terms of polylogarithms
- Weight 3 and 4 are 1 fold integrals (Chen integrals), they can be computed numerically in many frameworks (e.g. Mathematica, Cquad)

$$f_3 = \sum_k \beta_k \int d \log(b_k) f_2$$

$$f_4 = \sum_{k'} \alpha'_{k'} \int d \log(a'_{k'}) \sum_k \beta_k \int d \log(b_k) f_2$$

Analytic continuation

- We compute the integrals in the euclidean region ($s_{ij} < 0$)
- The physical region is reached by analytic continuation (Feynman prescription)
- The analytic continuation is in general complicated because of **spurious branch cuts** of the functional basis (symbol blinds to cuts of individual functions)
[Goncharov et al., 2010]



- We can solve the differential equations directly in the physical region (effectively circumventing the difficult task of analytic continuation!)

Solving in the physical region

- Physical Boundary conditions
- Analytic basis functions

$$\mathcal{I} = \sum_k c_k \mathcal{A}_k$$

analytic basis in the
physical region
(no spurious branch cuts)

- Identity holds in the entire physical region

Boundary conditions

- It is natural to expand wrt the (large) top mass
- When the external invariants approach 0, the expansion is exact (boundary point $s_{ij}=0$)
- Large mass expansion

For a given scalar graph depending on a set of momenta q and small masses m , the expansion in the large mass M limit is given by [\[V. Smirnov\]](#)

$$F_{\Gamma}(q, M, m, \epsilon) \stackrel{M \rightarrow \infty}{\sim} \sum_{\gamma} F_{\Gamma/\gamma}(q, m, \epsilon) T_{q^{\gamma}, m^{\gamma}} F_{\gamma}(q, M, m, \epsilon)$$

- The boundary point $s_{ij}=0$ is both physical and euclidean !

Analytic functions

- Basis (freedom to rotate to a preferred basis)


$$\text{Li}_n(x) = -G(0_{n-1}, 1, x), \quad \text{Li}_{2,2}(x, y) = G\left(0, \frac{1}{y}, 0, \frac{1}{xy}, 1\right)$$

- Choose polylogarithms with no branch cuts in the physical region (choose proper arguments)
- $\text{Li}_n(x)$, $x < 1$
- $\text{Li}_{2,2}(x, y)$, $y < 1$ and $xy < 1$
- Caveat, in principle functions with these requirements may not be a basis [[Bonciani, Del Duca, Frellesvig, Henn, FM, Smirnov, 2015](#)]

Analytic functions

- Analytic continuation of weight 3 and 4 is elementary
- Resulting Chen integrals requires just the continuation of logarithms and algebraic functions

$$f_4 = \sum_{k'} \alpha'_{k'} \int d \log(a'_{k'}) \sum_k \beta_k \int d \log(b_k) f_2$$



analytic in the
entire physical
region (by
construction)

Summary and outlook

- **New results for integrals of $gg \rightarrow Hg$ (only 4 integrals missing)**
- The integrals can be computed using a basis of UT integrals
- The solution is represented in terms of a minimal polylogs basis ($w=1, w=2$), and 1 fold integrals ($w=3, w=4$)

- The numeric evaluation is still quite slow, $O(\min)/(physical\ point)$, implementing it in c (Cquad) will significantly improve the evaluation speed
- Solution in terms of polylogarithms up to w_4 (combinatorics!)
- Integrate the DE in terms of GPL. Introducing a (non-physical) parameter such that the DE wrt it are in Fuchsian form. [[Papadopoulos 2014](#)]