

Top quark mass effects in Higgs boson pair production up to NNLO

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Loops and Legs in Quantum Field Theory, Leipzig
April 26, 2016

[based on work done in collaboration with J. Grigo and M. Steinhauser; [arXiv:1508.00909](https://arxiv.org/abs/1508.00909)]



1 Motivation and introduction

2 Differential factorization

3 Soft-virtual approximation

4 Asymptotic expansion

5 Calculation

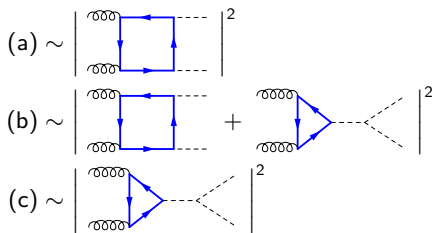
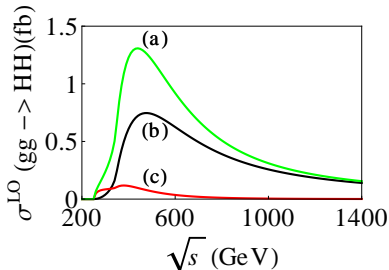
6 Results

Motivation

Higgs Potential in the Standard Model:

$$V(H) = \frac{1}{2}m_H^2 H^2 + \lambda v H^3 + \frac{1}{4}\lambda H^4, \quad \lambda^{\text{SM}} = \frac{m_H^2}{2v^2} \approx 0.13, \quad v: \text{Higgs vev.}$$

- Verify mechanism of spontaneous symmetry breaking in the SM
- Measure the Higgs self-coupling \Rightarrow sensitive process



Theory status

Prospects for the LHC @ 14 TeV:

- $b\bar{b}\gamma\gamma$ -channel, 600 fb^{-1} : $\lambda \neq 0$ [Baur, Plehn, Rainwater; '04]
- $b\bar{b}\gamma\gamma$ -, $b\bar{b}\tau^+\tau^-$ -channels: “promising”;
 $b\bar{b}W^+W^-$ -channel: “not promising” [Baglio, Djouadi, Gröber, Mühlleitner, Quevillon, Spira; '13]
- 600 fb^{-1} : $\lambda > 0$; 3000 fb^{-1} : $\lambda_{-20\%}^{+30\%}$ (ratio with Higgs cross section) [Goertz, Papaefstathiou, Yang, Zurita; '13]
- And many others, e.g.:
[Dolan, Englert, Spannowsky; '12] [Papaefstathiou, Yang, Zurita; '13]
[Barr, Dolan, Englert, Spannowsky; '13] [Barger, Everett, Jackson, Shaughnessy; '14]
[Englert, Krauss, Spannowsky, Thompson; '15] [...]
- Until now: Higgs pair production not observed in $b\bar{b}b\bar{b}$ - and $b\bar{b}\gamma\gamma$ -channels (as expected in the SM) [ATLAS; '15] [CMS; '15]

⇒ **Wait for HL-LHC**

Theory status

Known since long:

- LO result with exact M_t dependence [Glover, van der Bij; '88] [Plehn, Spira, Zerwas; '98]
- NLO result in $M_t \rightarrow \infty$ limit [Dawson, Dittmaier, Spira; '98]

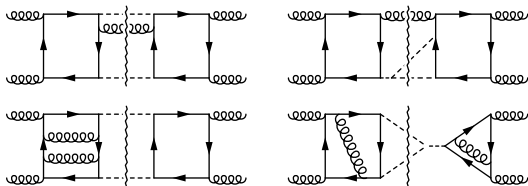
$$\sigma_H \approx 20^{\text{LO}} \text{ fb} + 20^{\text{NLO}, M_t \rightarrow \infty} \text{ fb} \quad \text{for } \sqrt{s_H} = 14 \text{ TeV}, \mu = 2m_H$$

More recently:

- NLO + NNLL ($M_t \rightarrow \infty$) \approx NLO +20% [Shao, Li, Li, Wang; '13]
- NNLO w/ or w/o soft-virtual approx. ($M_t \rightarrow \infty$) \approx NLO +20% [de Florian, Mazzitelli; '13]
- $\mathcal{O}(1/M_t^8)$ corrections at NLO \approx NLO +10% [Grigo, JH, Melnikov, Steinhauser; '13]
- NLO real exact in M_t , NLO virt. for $M_t \rightarrow \infty \approx$ NLO **-10%** [Maltoni, Vryonidou, Zaro; '14]
- Cross-check of virtual NNLO corr.; **NNLO matching coefficient** for $ggHH$ -coupling \approx NNLO +1% [Grigo, Melnikov, Steinhauser; '14]
- Improved $\mathcal{O}(1/M_t^{12})$ NLO, $\mathcal{O}(1/M_t^4)$ NNLO soft-virt. corrections (this talk) [Grigo, JH, Steinhauser; '15]
- NLO full M_t result (cf. previous talk by Matthias) [Borowka, Greiner, Heinrich, Jones, Kerner, Schlenk, Schubert, Zirke; '16]

Generalities

- Operate on full-theory diagrams at NLO and NNLO
- Virt. corr. in two independent calculations:
 - amplitude (differential; 2-/3-loop)
 - forward scattering (total; 4-/5-loop)
- Real corr.: only via forward scattering at NLO
- Perform expansion for $M_t \rightarrow \infty$; improve upon effective theory results for NLO [Dawson, Dittmaier, Spira; '98], NNLO [de Florian, Mazzitelli; '13] (not expected to be valid for $\sqrt{s} \gtrsim M_t$)
- Laporta reduction to master integrals for the “soft” subdiagrams
- Remaining “hard” massive tadpoles via MATAD
- Master integrals as series around $\sqrt{s} = 2m_H$



Differential factorization

Factorization of the LO result exact in M_t for:

Total cross section

$$\sigma^{(i)} = \Delta^{(i)} \sigma_{\text{exact}}^{(0)} = \frac{\sigma_{\text{exact}}^{(0)}}{\sigma_{\text{exp}}^{(0)}} \int_{4m_H^2}^s dQ^2 \frac{d\sigma_{\text{exp}}^{(i)}}{dQ^2}$$

$$\text{with } \Delta^{(i)} = \frac{\sigma_{\text{exp}}^{(i)}}{\sigma_{\text{exp}}^{(0)}}, \quad \sigma_{\text{exp}}^{(i)} = \sum_{n=0}^N c_n^{(i)} \rho^n, \quad \rho = \frac{m_H^2}{M_t^2}$$

Differential cross section

$$\sigma^{(i)} = \int_{4m_H^2}^s dQ^2 \frac{\left(\frac{d\sigma_{\text{exact}}^{(0)}}{dQ^2} \right)}{\left(\frac{d\sigma_{\text{exp}}^{(0)}}{dQ^2} \right)} \frac{d\sigma_{\text{exp}}^{(i)}}{dQ^2}$$

“Cure” the invalidity of the $M_t \rightarrow \infty$ expansion for the large- Q^2 region

- Virt. corr. via amplitude: **access to Q^2 -dependence $\sim \delta(s - Q^2)$**
- Real corr. via optical theorem (naively): **only total cross section**

⇒ Use the soft-virtual approximation [de Florian, Mazzitelli; '12]

Soft-virtual approximation

- Split σ into its contributions (works also for $d\sigma/dQ^2$):

$$\begin{aligned}\sigma &= \sigma^{\text{virt+ren}} + \sigma^{\text{real+split}} = \text{finite} \\ &= \underbrace{\Sigma_{\text{div}} + \Sigma_{\text{fin}}}_{=\Sigma_{\text{SV}}=\text{finite}} + \underbrace{\Sigma_{\text{soft}} + \Sigma_{\text{hard}}}_{=\Sigma_{\text{H}}=\text{finite}}\end{aligned}$$

- Σ_{div} universal for color-less final state [de Florian, Mazzitelli; '12]

- Compute $\sigma^{\text{virt+ren}}$ as ρ -expansion

- Solve $\sigma^{\text{virt+ren}} = \Sigma_{\text{div}} + \Sigma_{\text{fin}}$ for Σ_{fin}

- Σ_{div} and Σ_{soft} (soft coll. counterterms + soft real corr.)
 \sim exact σ^{LO} (include M_t effects)

$$Q^2 \frac{d\sigma}{dQ^2} = \sigma^{\text{LO}} z G(z) \quad \text{with} \quad z = \frac{Q^2}{s}, \quad G(z) = G_{\text{SV}}(z) + G_{\text{H}}(z)$$

$$\sigma_{(\text{SV})} = \int_{1-\delta}^1 dz \sigma^{\text{LO}}(zs) G_{(\text{SV})}(z) \quad \text{with} \quad \delta = 1 - \frac{4m_H^2}{s}$$

- $G_{\text{SV}}^{(i)}(z)$ constructed from $\sigma_{\text{fin}}^{(i)}$ and σ^{LO} only

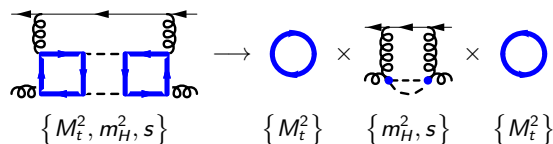
[de Florian, Mazzitelli; '12] [Grigo, JH, Steinhauser; '15]

Asymptotic expansion

- Expand at integrand level for all contributing regions (scalings of loop momenta)
≡ series expansion in analytic result
- Hierarchy: $M_t^2 \gg s, m_H^2 \Rightarrow$ series in $\rho = m_H^2/M_t^2$
- Effectively reduce number of loops and scales

- Here: regions correspond to subgraphs (in general more than one)
- Hard mass expansion: subgraphs must contain all heavy lines

Example: NLO real with one region

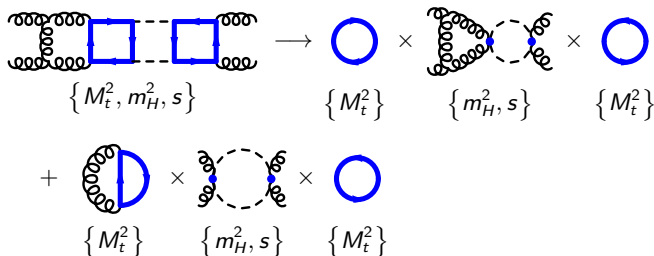


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Example: NLO virt. with two regions



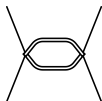
Calculation (toolchain)

Reduction

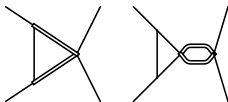
- 1 Generate Feynman diagrams
QGRAF [Nogueira; '93]
- 2 Select diagrams with specific cuts for $gg \rightarrow gg$ (17 millions)
filter [JH, Pak; (unpublished)]
- 3 Map diagrams to topologies (\leftarrow graph information)
exp [Harlander, Seidensticker, Steinhauser; '98]
- 4 Reduction to scalar integrals (\leftarrow generic topologies)
FORM [Kuipers, Ueda, Vermaseren, Vollinga; '13]
- 5 Reduction to master integrals (\leftarrow basic topologies)
 $gg \rightarrow gg$: rows [JH, Pak; (unpublished)]
 $gg \rightarrow hh$: FIRE [Smirnov]
- 6 Minimal basis of master integrals
TopoID [JH, Pak]

Calculation (via optical theorem)

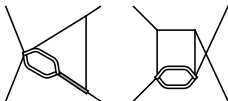
LO topology:



Virt. NLO topologies:



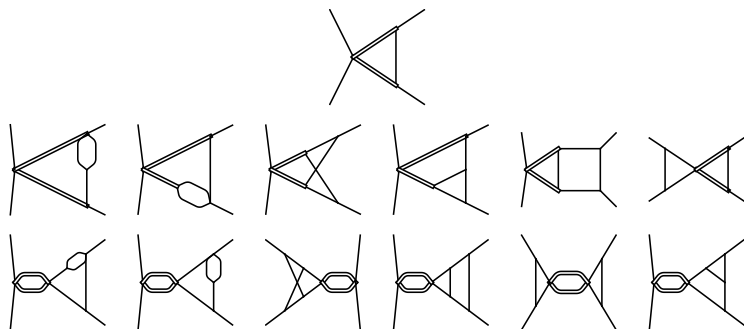
Real NLO topologies:



(shaded Higgs line cannot be cut)

Calculation (via optical theorem)

Virt. NNLO topologies:

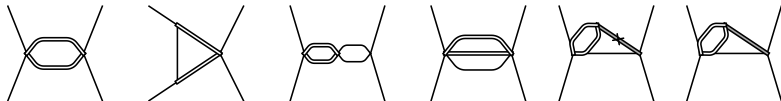


Note:

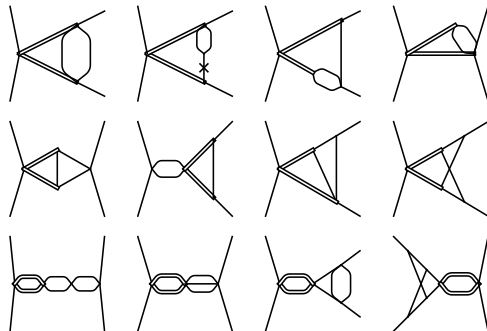
- Different regions in asymptotic expansion \Rightarrow different loop-orders
- Here: multiplied with 1- to 3-loop massive tadpoles

Calculation (via optical theorem)

Virt. and real LO-NLO master integrals:



Virt. NNLO master integrals (in addition):



Calculation (via optical theorem)

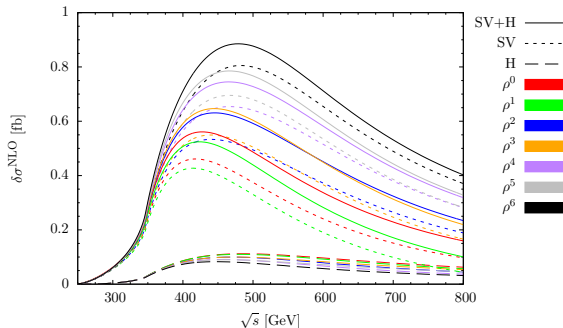
- NLO: 3 real and 1 virtual (+ 2-loop tadpoles) master integrals
- NNLO master integrals:
 - 3-loop: 1-loop phase space \times 2-loop form factors
 - 2-loop: massless sub-loops; same approach as one NLO
- Phase space integrals depend on $s = (q_1 + q_2)^2$ and m_H
- Derive 1-dimensional integral representation: e.g. for the simplest NLO case (cf. fourth diagram on previous slide)

$$I_1 = \mathcal{N} s^{1-2\epsilon} \delta^{5/2-3\epsilon} \int_0^1 \frac{d\mu}{\sqrt{1-\mu\delta}} (1-\mu)^{1/2-\epsilon} \mu^{1-2\epsilon}, \quad \delta = 1 - \frac{4m_H^2}{s}$$

- Expand up to $\mathcal{O}(\delta^{100})$ on NLO, $\mathcal{O}(\delta^{10})$ on NNLO
 - \Rightarrow very good convergence, small impact on numerics
 - \Rightarrow analytic results for partonic cross sections
- Note: All master integrals for the amplitude available in the literature

Splitting in soft-virtual and real contributions at NLO

- Partonic NLO correction; total factorization



Using $\mu = 2m_H$ (also in the following)

- Different behavior for higher orders in ρ expansion:

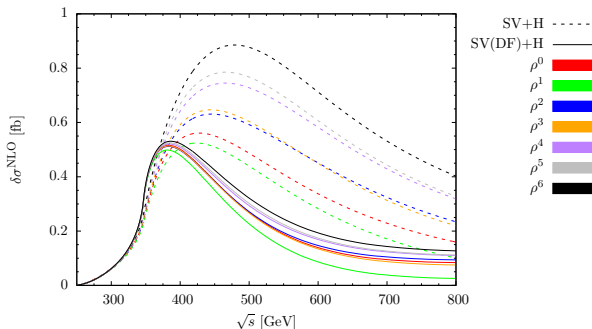
SV increasing

H decreasing (flat for $\sqrt{s} \gtrsim 400$ GeV)

⇒ SV numerically dominant

Total vs. differential factorization (DF) at NLO

- DF applied only to SV part;
H treated via total factorization (i.e. identical)



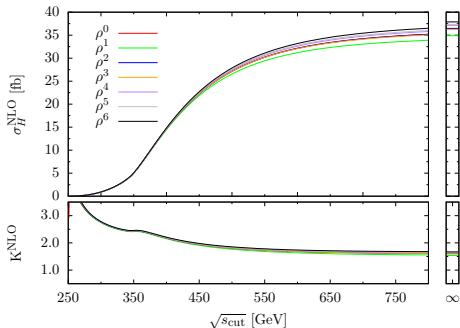
- Maxima of DF curves at lower \sqrt{s} ; smaller cross sections
- ⇒ Improvement of convergence:
difference of ρ^0 and corrections to ρ^6 (for $\sqrt{s} = 400$ GeV):
0.25 fb vs. 0.05 fb
- (Partonic K-factor: behavior at top quark pair threshold not washed out)

Hadronic cross section and K-factor

- Technical upper cut on \sqrt{s} (good proxy to Q^2):

$$\sigma_H(s_H, s_{\text{cut}}) = \int_{4m_H^2/s_H}^1 d\tau \left(\frac{d\mathcal{L}_{gg}}{d\tau} \right) (\tau) \sigma(\tau s_H) \theta(s_{\text{cut}} - \tau s_H)$$

- $\sqrt{s_{\text{cut}}} \rightarrow \infty$: total cross section for 14 TeV



- Spread of ρ -orders \Rightarrow $\pm 10\%$ uncertainty of EFT at NLO due to M_t

Revisiting NLO

Lessons from NLO for NNLO:

- SV approximation constructed for $z \rightarrow 1$;
 $G_{SV}(z)$ can be replaced by $f(z) G_{SV}(z)$ with $f(1) = 1$

Splitting into SV and H not unique

- Tune $f(z)$ at NLO such that $\sigma \approx \Sigma_{SV}$

$\Rightarrow f(z) = z$ accurate to 2%

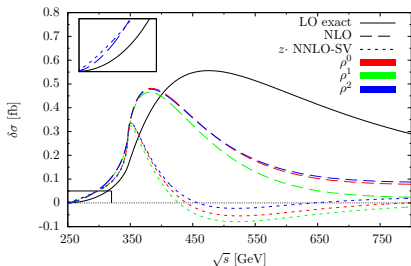
- Replace RGE logarithms ($\sqrt{s} \approx Q^2$ in the soft limit):

$\Rightarrow \log(\mu^2/s) \rightarrow \log(\mu^2/Q^2)$

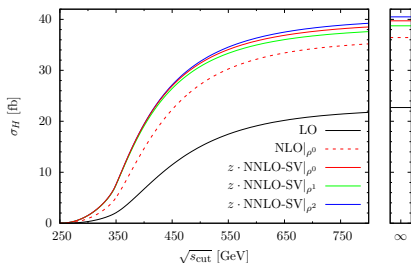
Discrepancy to [\[Maltoni, Vryonidou, Zaro; '14\]](#):

- Real corr.: treated exactly; Virt. corr.: EFT result
- Claim: **-10%** correction at NLO
- Difference due to treating M_t dependence for real and virt. corr. as consistent $1/M_t$ expansion or not

NNLO SV corrections



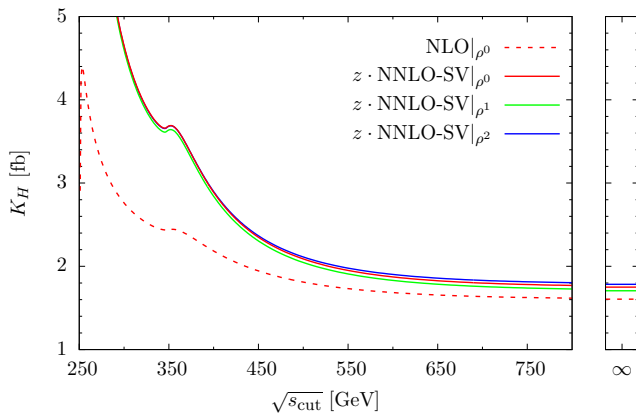
- EFT result plus ρ^- and ρ^2 -terms
- Peaks at smaller \sqrt{s}
- Same pattern of ρ -corrections



- Conv. up to $\sqrt{s_{\text{cut}}} \approx 400$ GeV
 - ρ^- and ρ^2 -corrections: $\pm 2.5\%$
- \Rightarrow M_t -uncertainty at NNLO: 5%

(NNLO corr. $\approx 20\%$)

NNLO SV K-factor



- Behavior at top threshold not washed out
- Strong raise close to threshold \leftarrow steeper NNLO correction

■ For total cross section: $K_H^{\text{NNLO}} \approx 1.7 - 1.8$

Conclusion and Outlook

- Real corrections at NLO via opt. theorem; employ SV approx. at NNLO
- Computed expansion to $\mathcal{O}(1/M_t^{12})$ at NLO, $\mathcal{O}(1/M_t^4)$ at NNLO
- Improvement of convergence by differential factorization
- Uncertainty of EFT results due to M_t :

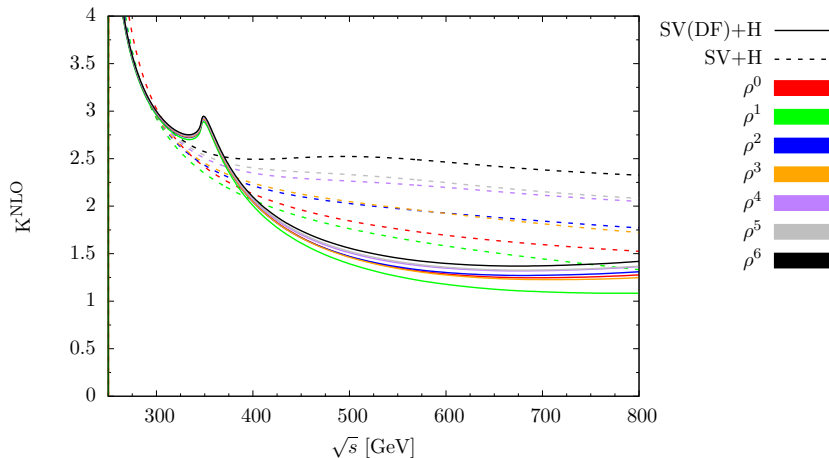
10% at NLO, 5% at NNLO

- Small- \sqrt{s} behavior as benchmark for exact calculations

- $1/M_t$ corrections for real contributions at NNLO?
- Adapt approximation procedure to match exact NLO calculation
⇒ Improved NNLO prediction

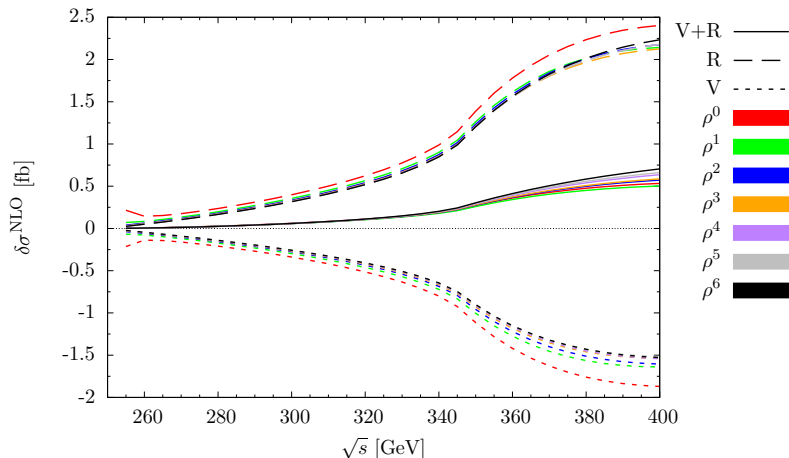
Thank you for your attention!

Partonic NLO K-factor



Note: Behavior around top quark pair threshold not washed out

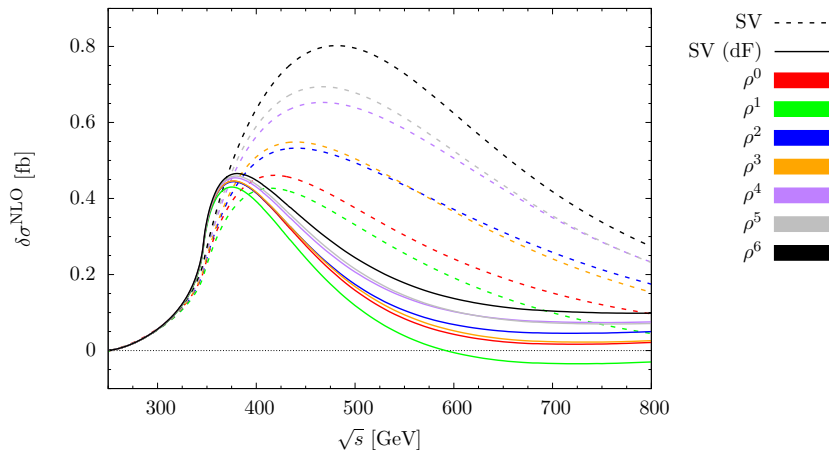
Splitting into real and virtual corrections at NLO



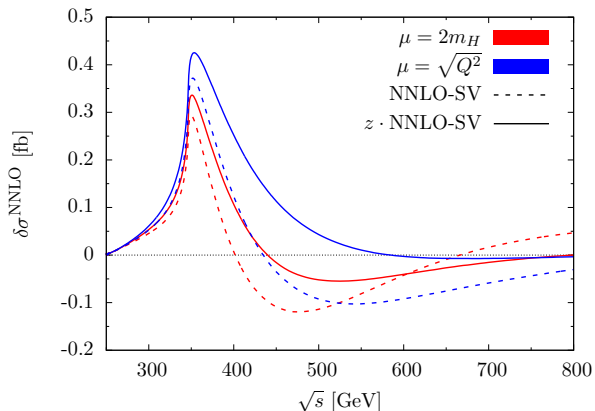
Note:

- R and V separately divergent; only finite contributions shown;
- Dominant positive shift from virtual corrections
- Different correction patterns (R negative, V positive)

Total vs. differential factorization without hard contributions at NLO



Partonic NNLO cross section for different scales choices and $f(z)$



Note: $f(z) = z$ (better proxy) leads to higher values