

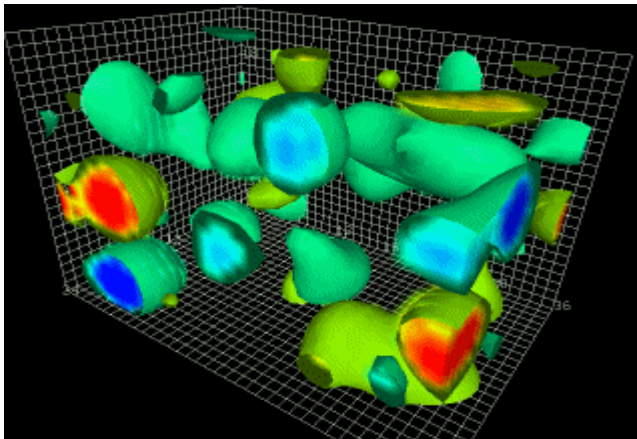
QCD Instantons at the LHC

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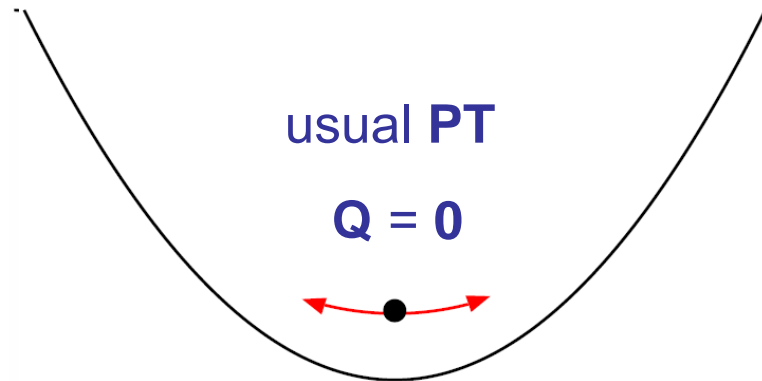
DESY

1. Introduction

- Huge effort to experimentally verify SM perturbation theory (**PT**) for hard processes..
- Yet 't Hooft '76: **hard** processes must exist that cannot be described by **PT**, despite $\alpha \ll 1$! Induced by **topological** fluctuations of gauge fields.
- Rich vacuum structure of **non-abelian** gauge theories (QCD) \leftarrow **Topology!**
Gauge fields carry an integer topological charge $\mathbf{Q} = 0, \pm 1, \dots \leftrightarrow$ winding number
Classical ground states energetically degenerate, but topologically distinct (fig)!



Potential energy:

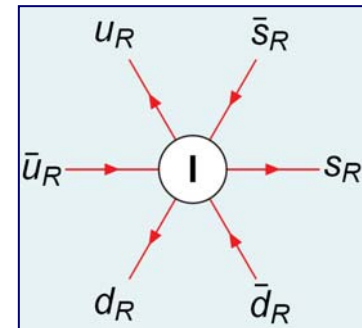


- (Anti) Instantons [Belavin *et al.* '75] \longleftrightarrow a **basic aspect** of the SM!
explicitly known tunneling transitions between vacua differing by $\mathbf{Q} = +1$ (-1),
“instantaneous” in time and space.

- ❑ **Theoretically:** QCD instantons known to play an important rôle in interface regime: partons \rightarrow hadrons. E.g. chiral symmetry breaking, spectroscopy.. [Diakonov '96]
Also: Instanton-driven **gluon saturation** at small x [F. Sch & Utermann '01-'04]
- ❑ **Yet... Direct experimental evidence for instantons still lacking!**
- ❑ **However:** [A. Ringwald & F. Sch. '94]

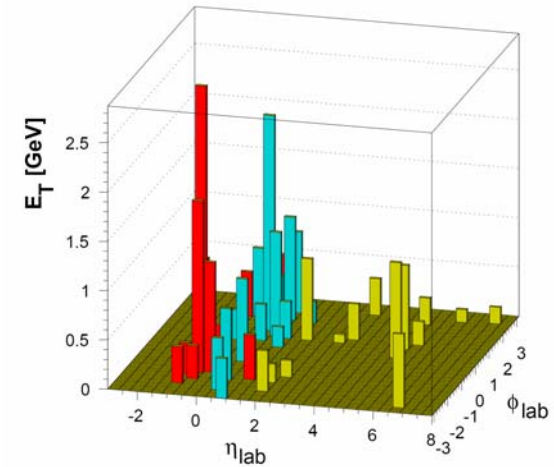
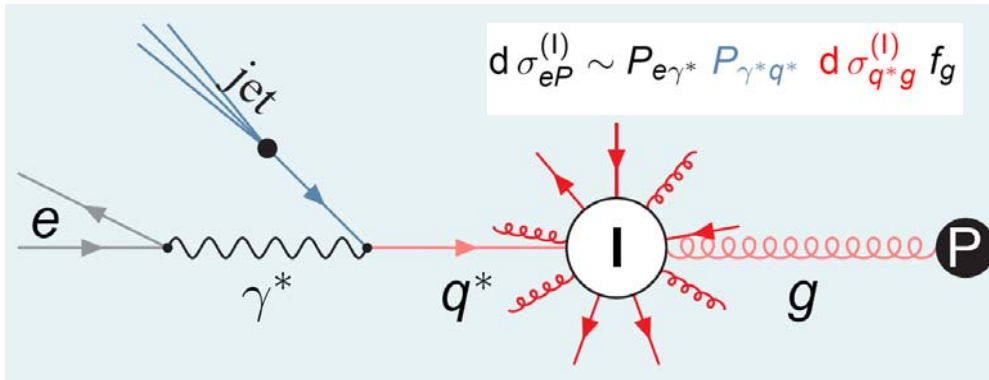
Characteristic **short**-distance manifestation of instantons may be exploited for experimental search!

- o QCD I 's induce **hard**, calculable, **chirality-violating processes**, forbidden in usual **PT** ! ['t Hooft '76]
- o Theoretical prediction of rate and characteristic event signature achieved in **DIS** (**strict** I perturbation theory) \rightarrow
rate in measurable range at HERA. [Ringwald & F. Sch. '94 - '01]
- o Two dedicated I - search experiments by **H1, ZEUS** **demonstrated** that required **exp. sensitivity is within reach!**



- ❑ Study in detail **discovery potential** for QCD instanton processes **@LHC!**
- ❑ Looking forward to data reanalysis in 2007 with $\sim 750 \text{ pb}^{-1}$ **@ HERA!**

1.1 Instantons at DESY



- ❑ A quick reminder:
- ❑ Extensive theoretical work
Ringwald & F. Sch. '94 - '01 & Moch (PhD '94-'97) & Utermann (PhD '00 – '03) & Petermann (PhD '04 – '07),
- ❑ Instanton Monte Carlo generator for HERA: **QCDINS** [Ringwald & F. Sch.] (fig. right)
- ❑ 2 **instanton search experiments** based on our work:

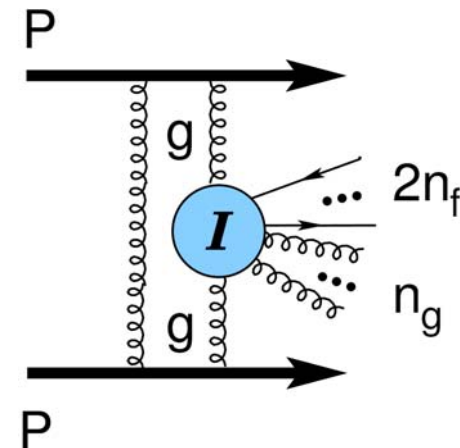
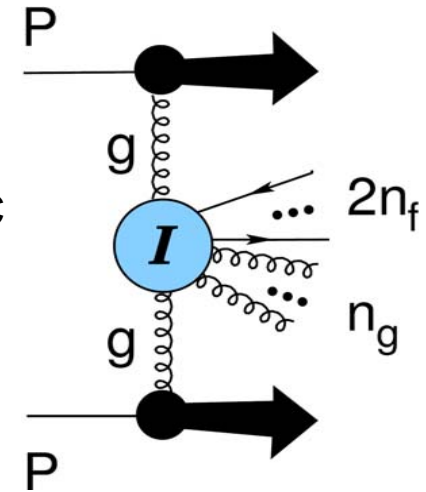
- **H1** $\left\{ \begin{array}{l} \int \mathcal{L} dt = 21 \text{ pb}^{-1}, \theta_{e^+} > 156^\circ, 0 < y < 0.6 \\ x > 10^{-3}, 10 \lesssim Q^2 \lesssim 100 \text{ GeV}^2 \end{array} \right.$
- **ZEUS** $\left\{ \begin{array}{l} \int \mathcal{L} dt = 38 \text{ pb}^{-1}, y > 0.05 \\ x > 10^{-3}, Q^2 \gtrsim 100 \text{ GeV}^2 \end{array} \right.$

[**H1 Coll.**, Eur.Phys. J. C 25 (2002) 495; **ZEUS Coll.**, Eur. Phys. J. C 34 (2004)]

- ❑ While **H1** saw significant excess over **MC**'s in accord with our predictions, large remaining uncertainties from normal **DIS** event generators remained challenging...

2. Setting the stage for the LHC

- ❑ “Gold-plated” I - induced events @LHC \leftrightarrow best compromise between rate and bg-freedom!
- ❑ Very promising new method for bg-suppression @HERA & @LHC [Barakbayev, Boos, Lohrmann, Petermann & F.Sch. '06] in prep.
- ❑ Three main aspects to select dominant I – subprocess,
 - 1) counting powers of α_s and α_{em} ,
 - o each external **non-perturbative gluon**: $\sigma \propto \frac{1}{\alpha_s}$
 - o each external **quark (zero mode)**: $\sigma \propto \mathcal{O}(1)$
 - 2) enhancements from **parton densities**
 - 3) subprocess dependent **power-growth** in $\hat{s} = E^2$, at lower energies for exclusive, exp-growth for inclusive processes!
- ❑ (1) & (2) \rightarrow $g + g$ initial state strongly **dominant** as in usual **PT**
- ❑ Focus on general topology as in fig. (top). Interesting variant with **2 rapidity gaps and central “fireball”** in fig. (bottom): “diffractive instanton” events...



2.1 Energy constraint

- The (inclusive) sum over multi-gauge boson final states is known to **exponentiate**,

$$\sigma'_{t \text{ Hoof}}^{(I)} \sim \exp \left[-\frac{4\pi}{\alpha} \right] \quad \Rightarrow \quad \sigma^{(I)} \sim \exp \left[-\frac{4\pi}{\alpha} F_{\text{hg}} \left(\frac{E}{m_{\text{sph}}} \right) \right]$$

with the **“holy grail”** function F_{hg} and $0 \leq F_{\text{hg}} \leq 1$, turning **‘t Hooft’s** tunneling factor (left) into a much weaker suppression (right), since $F_{\text{hg}} \downarrow$ for $E \uparrow$.

- The scale m_{sph} denotes the **“sphaleron mass” (barrier height)** for QFD and QCD:

$$m_{\text{sph}} \approx \frac{3\pi}{4} \frac{1}{\alpha \rho_{\text{eff}}} \approx \begin{cases} 4 \frac{m_W}{\alpha_W} \approx 10 \text{ TeV} & (m_{\text{Higgs}} \approx 115 \text{ GeV}) & \text{for QFD} \\ \mathcal{O}(1) \mathcal{Q} & (\text{virtuality } \mathcal{Q}) & \text{for QCD} \end{cases}$$

ρ_{eff} = **effective** instanton size

[Klinkhamer & Manton '84 (QFD); Ringwald & F. Sch '94 (QCD)]

- Resummation of gauge bosons (gluons) in final state via **“valley approximation”** & optical theorem. \rightarrow [Yung '88]

valley action $S_{\text{valley}}^{I\bar{I}}$ formally known for all $\frac{E}{m_{\text{sph}}} \leftrightarrow F_{\text{hg}}$ **“known”!**
 [Khoze & Ringwald '91, Verbaarschot '91]

Energy constraint

- Crucial: when does the valley approximation break down??

2 direct, independent evidences for requirement: $\frac{E}{m_{\text{sph}}} \leq \mathcal{O}(1)$

QCD: [F. Sch. & Utermann '02]; QFD: [Rubakov, Rebbi *et al.* '03]

- This energy constraint is implemented in our LHC analysis.

Since cross section increases with E until $E \approx m_{\text{sph}}$, always work effectively near the “sphaleron”!

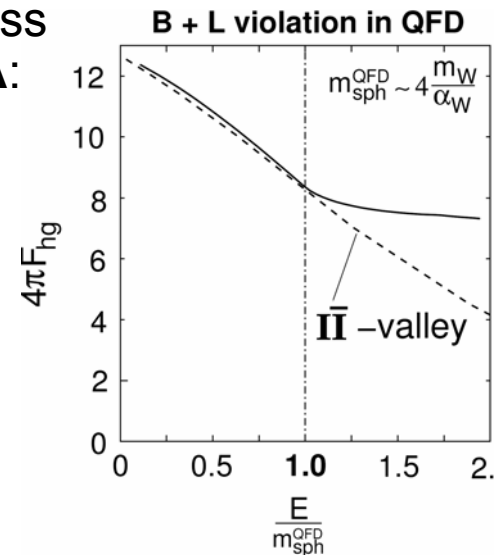
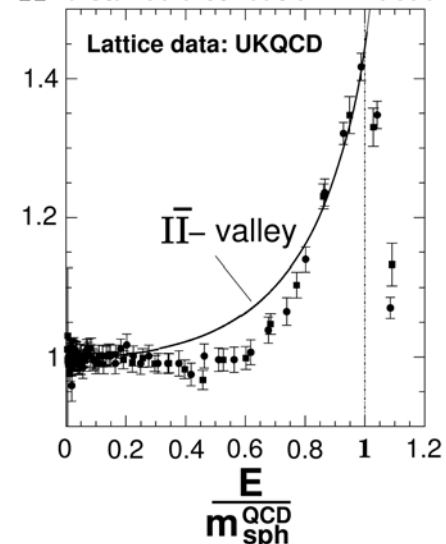
- Matches widely believed “square root rule” for instanton cross sections, which was also implemented for predictions @HERA:

$$E \text{ large} : \sigma^{(I)} \sim \sqrt{\sigma_{t \text{ Hooft}}^{(I)}} \sim \exp \left[-2\frac{\pi}{\alpha} \right], \text{ i.e. } F_{\text{hg}} \rightarrow \frac{1}{2}$$

- Note: observability of B+L violation in QFD at the LHC would require a far less conservative assumption, i.e.

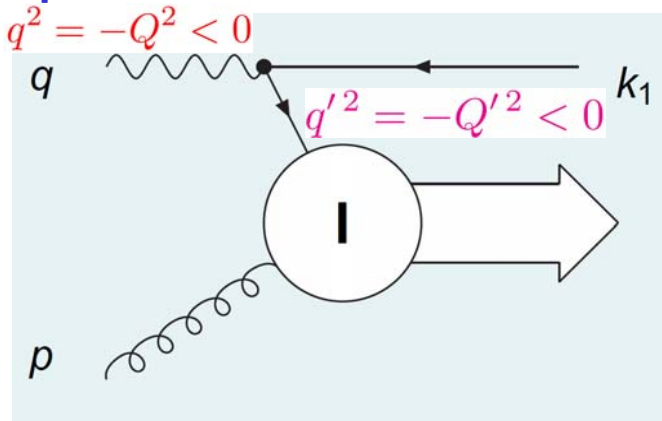
$$F_{\text{hg}} \xrightarrow{E \rightarrow \infty} 0$$

$\bar{\Pi}\bar{\Pi}$ -distance distribution in vacuum



2.2 From HERA to the LHC by crossing

I-process @HERA



1 HERA variables

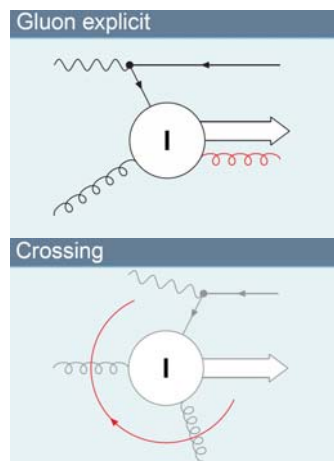
$$-Q^2 \equiv (q - k_1)^2 < 0$$

$$x' \equiv \frac{Q^2}{Q^2 + s'} \in [0, 1]$$

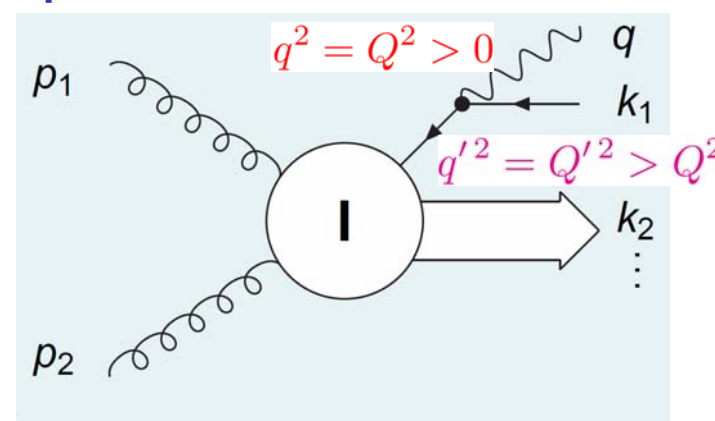
I - subprocess virtuality
spacelike **timelike**

$$Q^2 \equiv (q + k_1)^2 > 0$$

$$x' \equiv \frac{Q^2}{Q^2 - t'} \in [0, 1]$$



I-process @LHC



2 LHC variables

□ @HERA: initial state, virtual $\gamma^*(Q)$,

@LHC: consider $\gamma(q_T), \gamma^*(Q)_{\rightarrow l+l-}, W^\pm(m_W)$ in final state together with a q-jet.

□ @HERA, the I-subprocess virtuality Q'^2 **not** bounded by Q^2 , hence **need** Q'^2 reconstruction and a respective cut!

@LHC: $Q'^2 > Q^2$ kinematically! **Timelike virtuality** Q'^2 enforced in I-subprocess, by requiring a final state vector boson (fig)

$$\gamma(q_T), \gamma^*(Q)_{\rightarrow l+l-}, W^\pm(m_W) + \mathbf{1 \text{ jet}}$$

2.2 Event signature

- In I - rest system: **“fireball”** decaying **isotropically** to

$$n_f \cdot (q + \bar{q}) + \mathcal{O}\left(\frac{1}{\alpha_S}\right) \cdot g = \mathcal{O}(20) \text{ partons} \sim \mathcal{O}(60 - 80) \text{ hadrons}$$

- $u, \bar{u}, d, \bar{d}, s, \bar{s}$ **flavor “democracy”** in **each** event!
strangeness $\Rightarrow K' s, \Lambda' s$

- Lego plot $(\eta, \varphi, E_T) \Rightarrow$ **“I-band”**

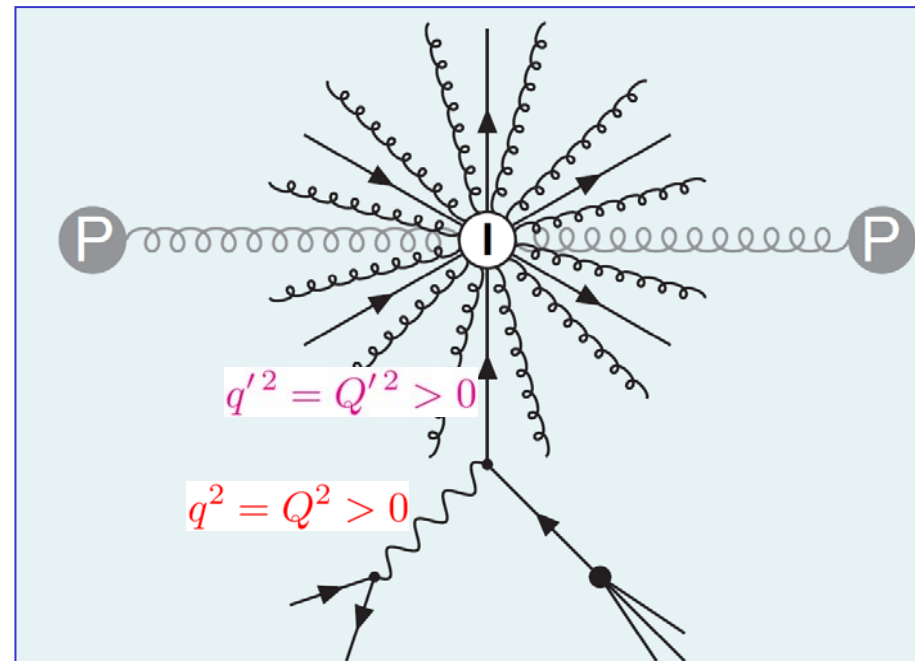
Isotropy \Rightarrow small width $\Delta\eta = \pm 1$ in (pseudo-) rapidity η and isotropy in azimuth φ , large total E_T

- Every exclusive I-process grows initially with energy like a high power:

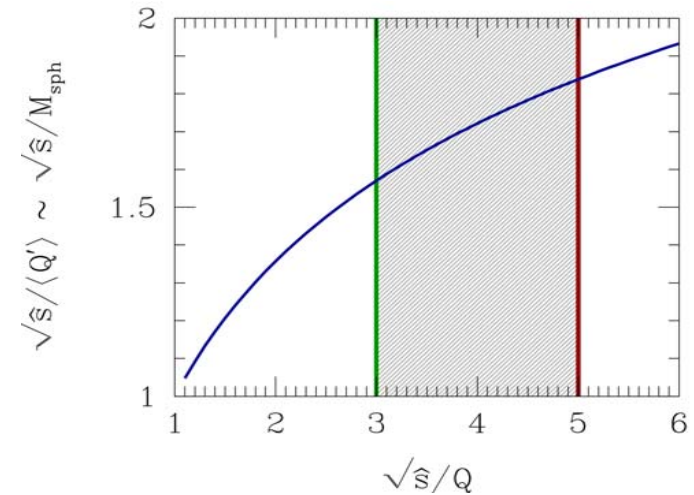
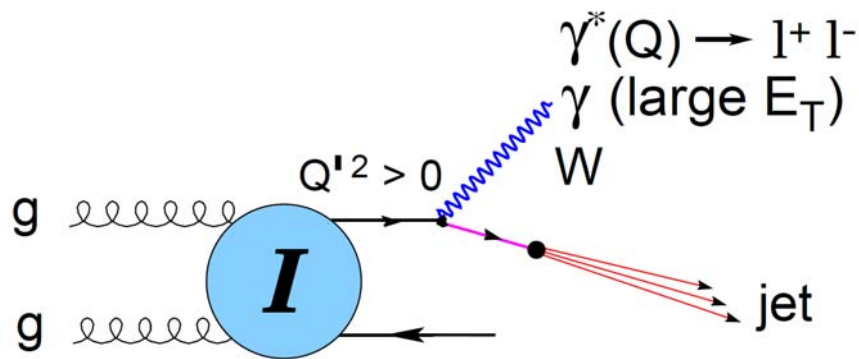
$$\hat{\sigma} \sim \hat{S}^{3n_f + 2n_g - 5}$$

(like a contact term)

“fireball” + vector boson + 1 jet



3. The simplest I-induced LHC process



- Like for **HERA**, idealized simplest I-process @LHC, $n_f = 1 \oplus$ **no final-state gluons** is very instructive (fig left)! Being **exactly calculable** in I-perturbation theory, many of its features remain valid after **final state gluon resummation**. Focus on such aspects.

- Fig. right illustrates that **sphaleron mass** constraint may be characterized in terms of vector boson virtuality:

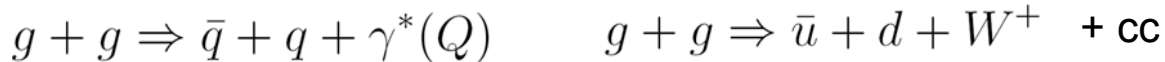
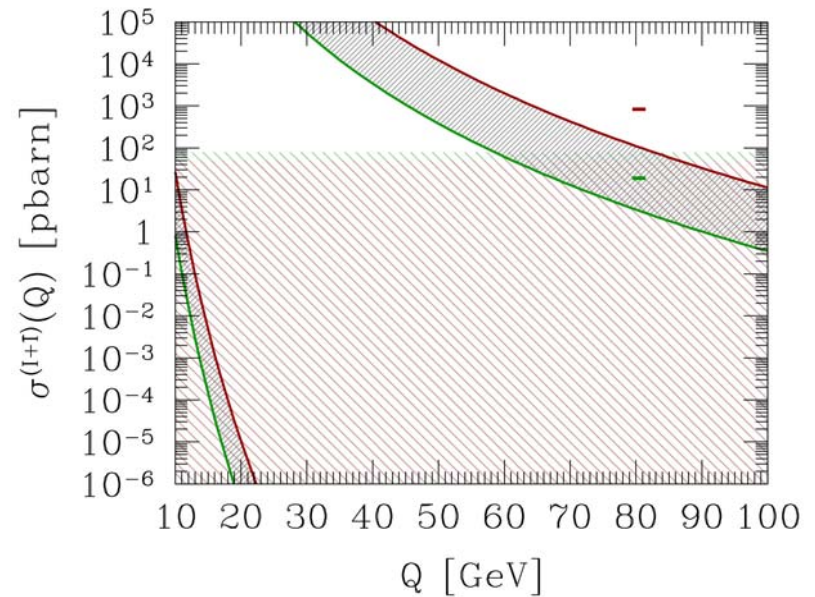
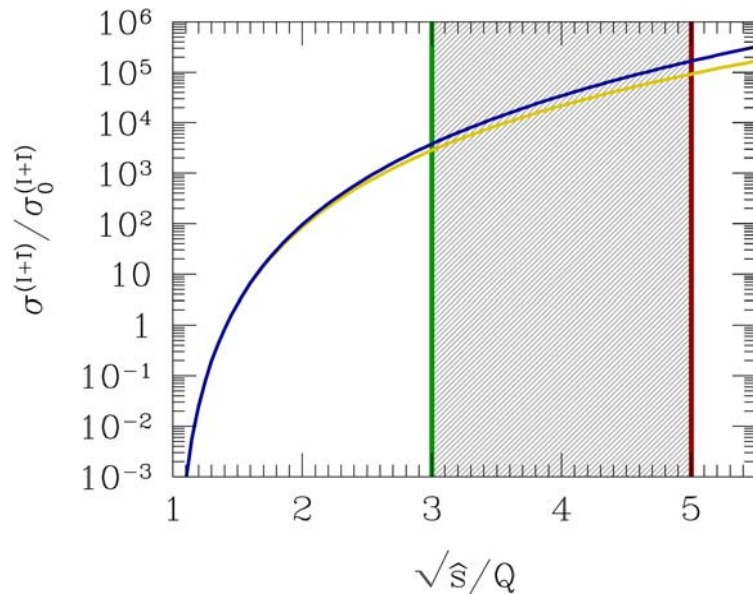
$$3 \lesssim \frac{\sqrt{\hat{s}}}{Q} \lesssim 5 \Leftrightarrow \frac{\sqrt{\hat{s}}}{m_{\text{sph}}} = \mathcal{O}(1) \text{ more/less conservatively}$$

resulting from (known) form of Vqq' vertex in I-background!

- While rigorous calculation being still in progress, consider “poor man’s” gluon resummation in terms of dominant effect:

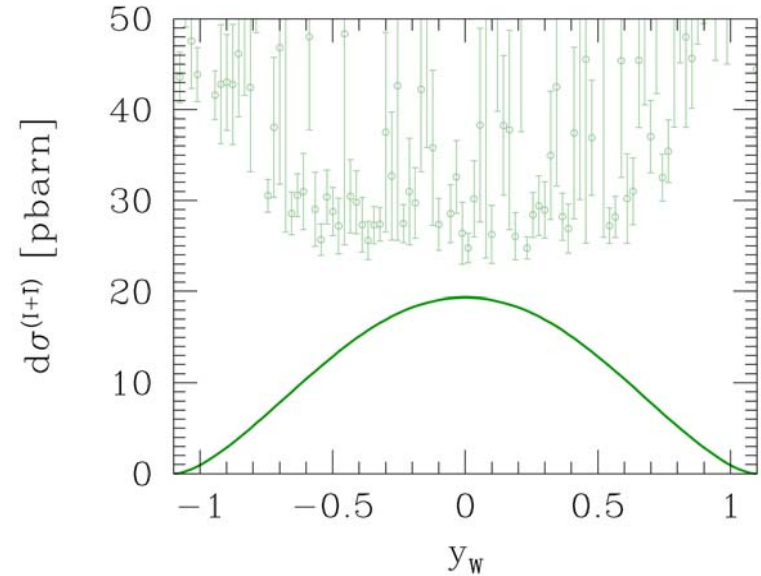
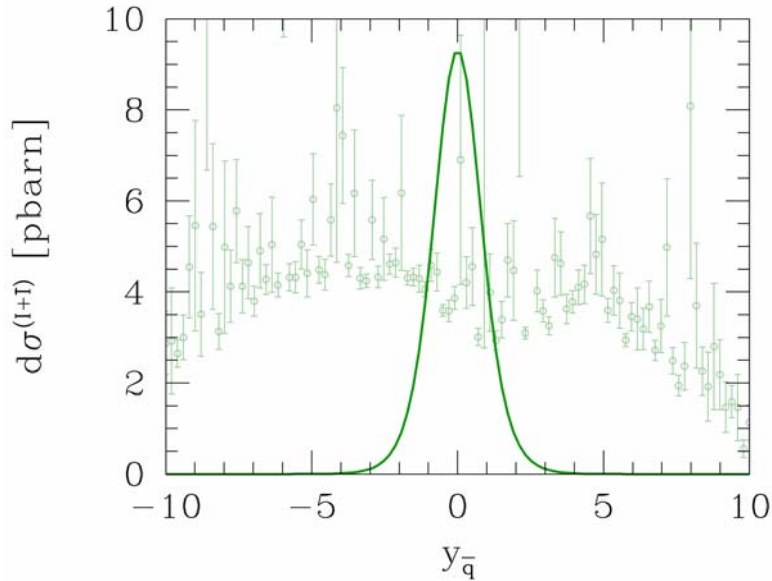
$$F_{\text{hg}} \rightarrow \frac{1}{2}$$

3.1 Results



- ❑ Display a set of characteristic results for **simplest processes** and enhancements via “**poor man’s**” **gluon resummation** that are expected to survive in fully realistic case.
- ❑ Fig. left: relative energy dependence and our two “sphaleron limits” (green, red lines)
- ❑ Fig. right: huge enhancement via “**poor man’s**” **gluon resummation**. W final state not too small!
shaded region: normal, perturbative background, calculated with **Comphep**

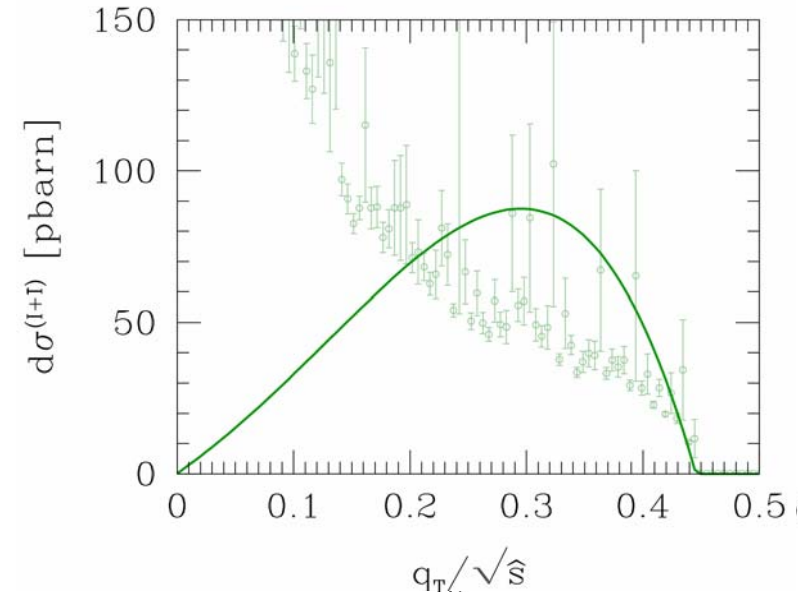
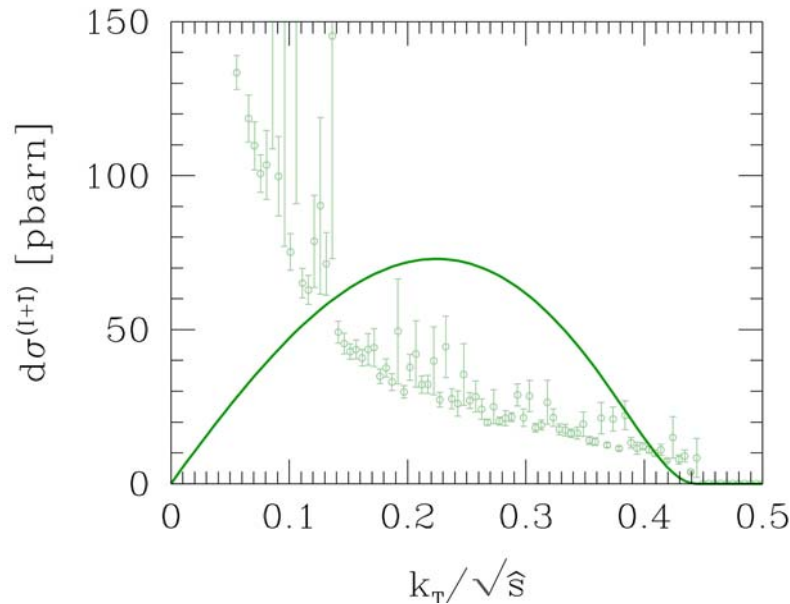
Results



$$g + g \Rightarrow \bar{u} + d + W^+ + \text{cc.}$$

- ❑ **Rapidity distributions** for final state quark-let and W, first for conservative “benchmark”: $\sqrt{\hat{s}} = 3 m_W$, and using “poor man’s” gluon resummation.
- ❑ Normal SM background calculated with **Comphep**. “Errors” are Monte Carlo errors.
- ❑ Note: widths of rapidity profiles are uniformly ± 1 unit, due to isotropy!
- ❑ Next: $\sqrt{\hat{s}} = 5 m_W$: **signal / background improves a lot!**

Results



$$g + g \Rightarrow \bar{u} + d + W^+ + \text{cc.}$$

- ❑ **Transverse momentum distributions** for final state quark-jet and W, first for conservative “benchmark”: $\sqrt{\hat{s}} = 3 m_W$, and using “poor man’s” gluon resummation.
- ❑ Normal SM background calculated with **Comphep**. “Errors” are Monte Carlo errors.
- ❑ Again for $\sqrt{\hat{s}} = 5 m_W$: **signal / background improves a lot!**

4. Conclusions & Outlook

- ❑ Search for instanton processes concerns a **basic non-perturbative aspect of QCD!**
- ❑ Calculations for the “simplest I-induced LHC process”: **completed & promising.**
- ❑ Explicit gluon resummation via the “valley method” and (saddle point) integration over the 9 (!) collective instanton coordinates is difficult (**timelike** virtuality!) and still in progress.
- ❑ Yet, the presented “poor man’s” gluon resummation estimates give rise to **optimism!**
- ❑ **Higher “fireballness”** of I-events is expected compared to **HERA**, due to increased phase space **@LHC!**
- ❑ Important relations of I-predictions **@LHC** to those **@HERA** due to **crossing.**

Outlook:

- ❑ Explore new method of background suppression for LHC
[**Barakbayev, Boos, Lohrmann, Petermann & F.Sch. '06**]
- ❑ With **Tancredi Carli/CERN** & friends: LHC observables...QCDINS@LHC (C++)...