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The high energy radiation pattern from BFKLex

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Outline

- Short introduction to the iterative solution of the Balitsky-Fadin-Kuraev-Lipatov (BFKL) equation
- The collinear double logarithms in the BFKL gluon Green's function
- The actual implementation in BFKLex, a BFKL based Monte Carlo code in C++ (G.C, A. Sabio Vera)
- The anatomy of the gluonic ladder and new observables
- Conclusions and Outlook

Brief introduction (LL)

- BFKL formalism: in high energy scattering terms like $\alpha_s^n \log^n(s) \sim \alpha_s^n (y_A - y_B)^n$ need to be resummed.
- $$\sigma(Q_1, Q_2, Y) = \int d^2\vec{k}_A d^2\vec{k}_B \underbrace{\phi_A(Q_1, \vec{k}_A) \phi_B(Q_2, \vec{k}_B)}_{\text{PROCESS-DEPENDENT}} \underbrace{f(\vec{k}_A, \vec{k}_B, Y)}_{\text{UNIVERSAL}}$$
- $\phi_{A,B}(Q_{1,2}, \vec{k}_{A,B})$ are the process-dependent impact factors
- The gluon Green's function $f(\vec{k}_A, \vec{k}_B, Y)$ is universal and depends on the scales $\vec{k}_{A,B}$ and the energy $\sim e^{Y/2}$

Brief introduction (LL)

- The main goal is to have a way to calculate the gluon Green's function (GGF).
- The GGF is the solution to the BFKL equation. Use the iterative form:

$$f = e^{\omega(\vec{k}_A)Y} \left\{ \delta^{(2)}(\vec{k}_A - \vec{k}_B) + \sum_{n=1}^{\infty} \prod_{i=1}^n \frac{\alpha_s N_c}{\pi} \int d^2 \vec{k}_i \frac{\theta(k_i^2 - \lambda^2)}{\pi k_i^2} \right. \\ \left. \int_0^{y_i-1} dy_i e^{(\omega(\vec{k}_A + \sum_{l=1}^i \vec{k}_l) - \omega(\vec{k}_A + \sum_{l=1}^{i-1} \vec{k}_l))y_i} \delta^{(2)}\left(\vec{k}_A + \sum_{l=1}^n \vec{k}_l - \vec{k}_B\right) \right\}$$

$\omega(\vec{q}) = -\frac{\alpha_s N_c}{\pi} \log \frac{q^2}{\lambda^2}$ is the gluon Regge trajectory

From LO to NLO

- At NLL terms like $\alpha_s^{n-1} \log^n(s)$ need to be resummed
- The BFKL kernel changes, the real part becomes much more complicated. Dominant term:

$$\theta(k_i^2 - \lambda^2) \rightarrow \theta(k_i^2 - \lambda^2) - \underbrace{\frac{\bar{\alpha}_s}{4} \ln^2 \left(\frac{\vec{k}_A^2}{(\vec{k}_A + \vec{k}_i)^2} \right)}_{\text{NLL}}$$

- These collinear logarithms are an issue, they need to be resummed as well.

Dealing with the collinear logarithms

- Remove from the NLO BFKL kernel the double log:

$$-\frac{\bar{\alpha}_s^2}{4} \frac{1}{(\vec{q} - \vec{k})^2} \ln^2 \left(\frac{q^2}{k^2} \right)$$

- Replace it with:

$$\left\{ \left(\frac{q^2}{k^2} \right)^{-b \bar{\alpha}_s \frac{|k-q|}{k-q}} \sqrt{\frac{2(\bar{\alpha}_s + a \bar{\alpha}_s^2)}{\ln^2 \left(\frac{q^2}{k^2} \right)}} J_1 \left(\sqrt{2(\bar{\alpha}_s + a \bar{\alpha}_s^2) \ln^2 \left(\frac{q^2}{k^2} \right)} \right) \right. \\ \left. - \bar{\alpha}_s - a \bar{\alpha}_s^2 + b \bar{\alpha}_s^2 \frac{|k-q|}{k-q} \ln \left(\frac{q^2}{k^2} \right) \right\} \frac{1}{(\vec{q} - \vec{k})^2}$$

$$a = \frac{5}{12} \frac{\beta_0}{N_c} - \frac{13}{36} \frac{n_f}{N_c^3} - \frac{55}{36} \text{ and } b = -\frac{1}{8} \frac{\beta_0}{N_c} - \frac{1}{6} \frac{n_f}{N_c^3} - \frac{11}{12}$$

Implementation in the BFKLex

- We work at NLO: NLO kernel with the additional resummation of the double logs (NLO+Double Logs)
- In particular we use the scale invariant part of the NLO BFKL kernel, $\beta_0 = 0$ and $n_f = 0$
- We use the anti- k_t jet algorithm in the FastJet implementation (Cacciari, Salam, Soyez)
The jet radius is set to $R=0.7$

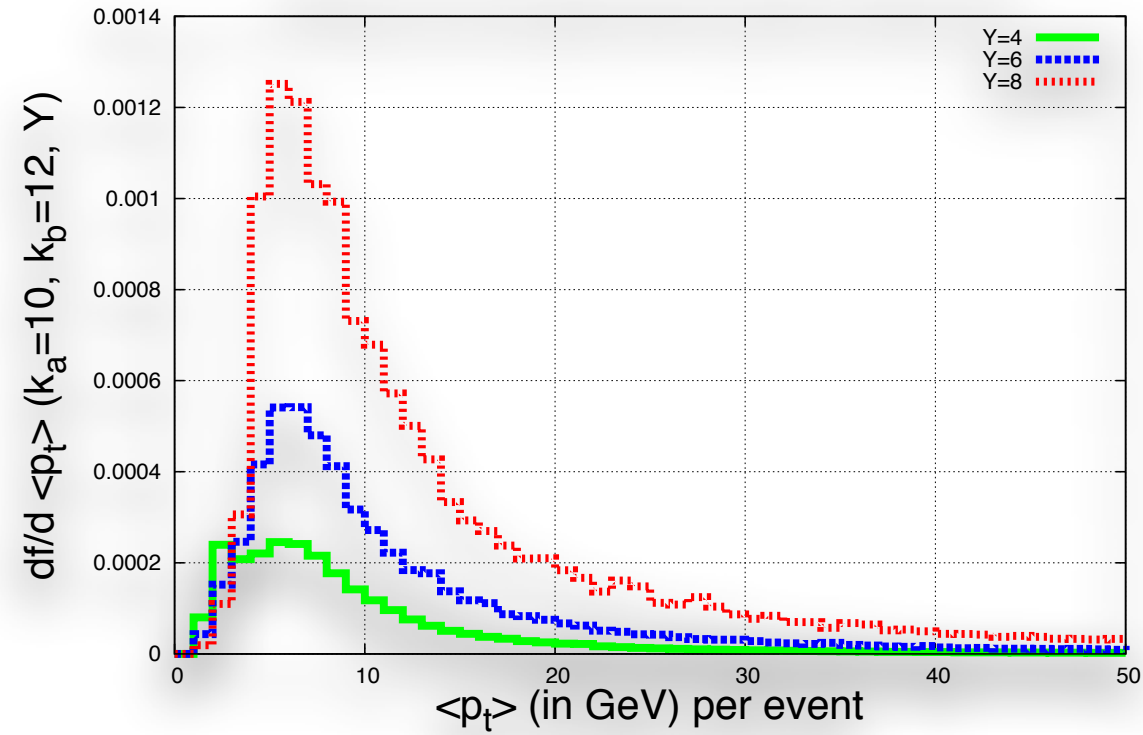
Three characteristics of the jet radiation pattern in the multi-Regge kinematics

- Introduce three quantities related to the jet activity along the ladder. These characterize uniquely the event (but not fully).

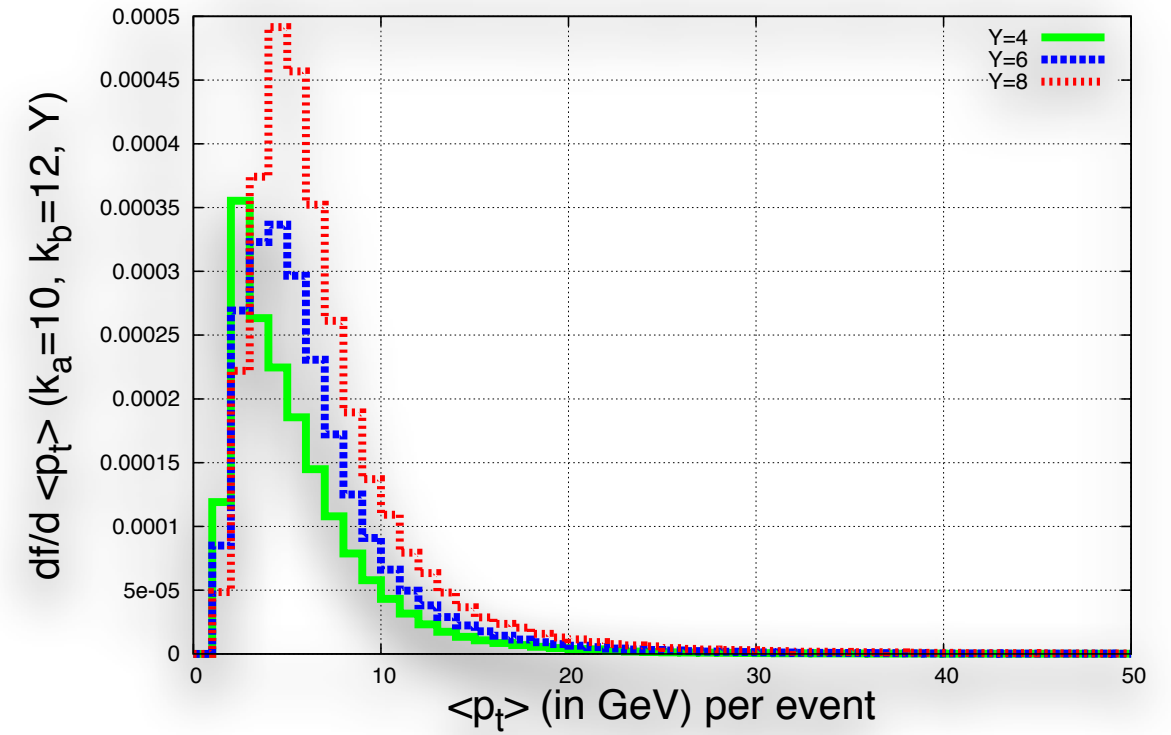
- average p_t $\langle p_t \rangle = \frac{1}{N} \sum_{i=1}^N |k_i|$
- average azimuthal angle $\langle \phi \rangle = \frac{1}{N} \sum_{i=1}^N \phi_i$
- rapidity ratio between subsequent jets $\langle \mathcal{R}_y \rangle = \frac{1}{N+1} \sum_{i=1}^{N+1} \frac{y_i}{y_{i-1}}$
 $y_0 = y_a, y_{N+1} = y_b = 0$ and $y_{i-1} > y_i$

Average p_t

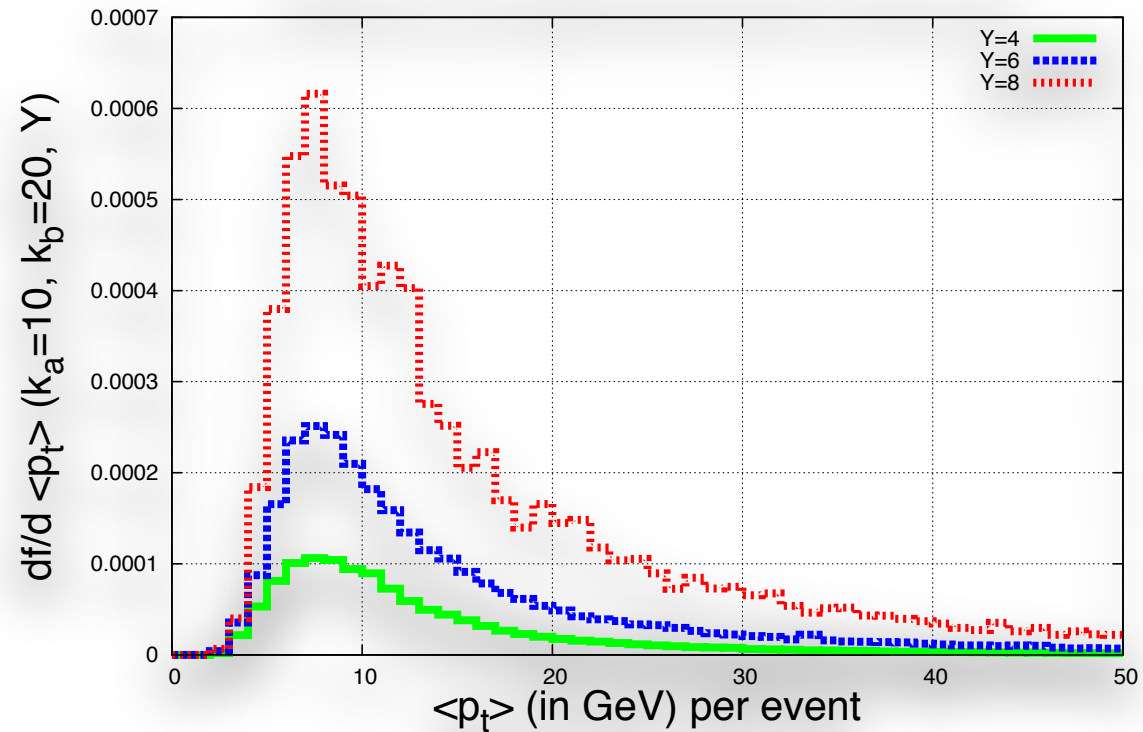
Distribution at LO in the average transverse momentum of emitted mini-jets



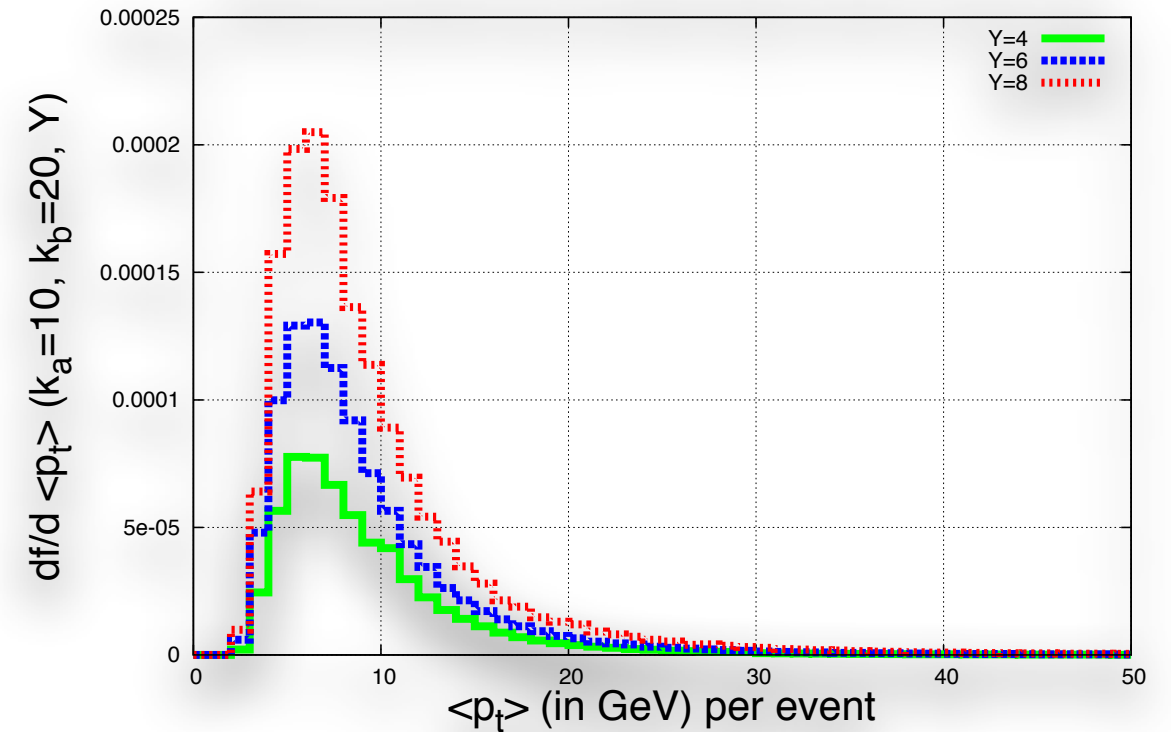
Distribution at NLO+Double Logs in the average transverse momentum of emitted mini-jets



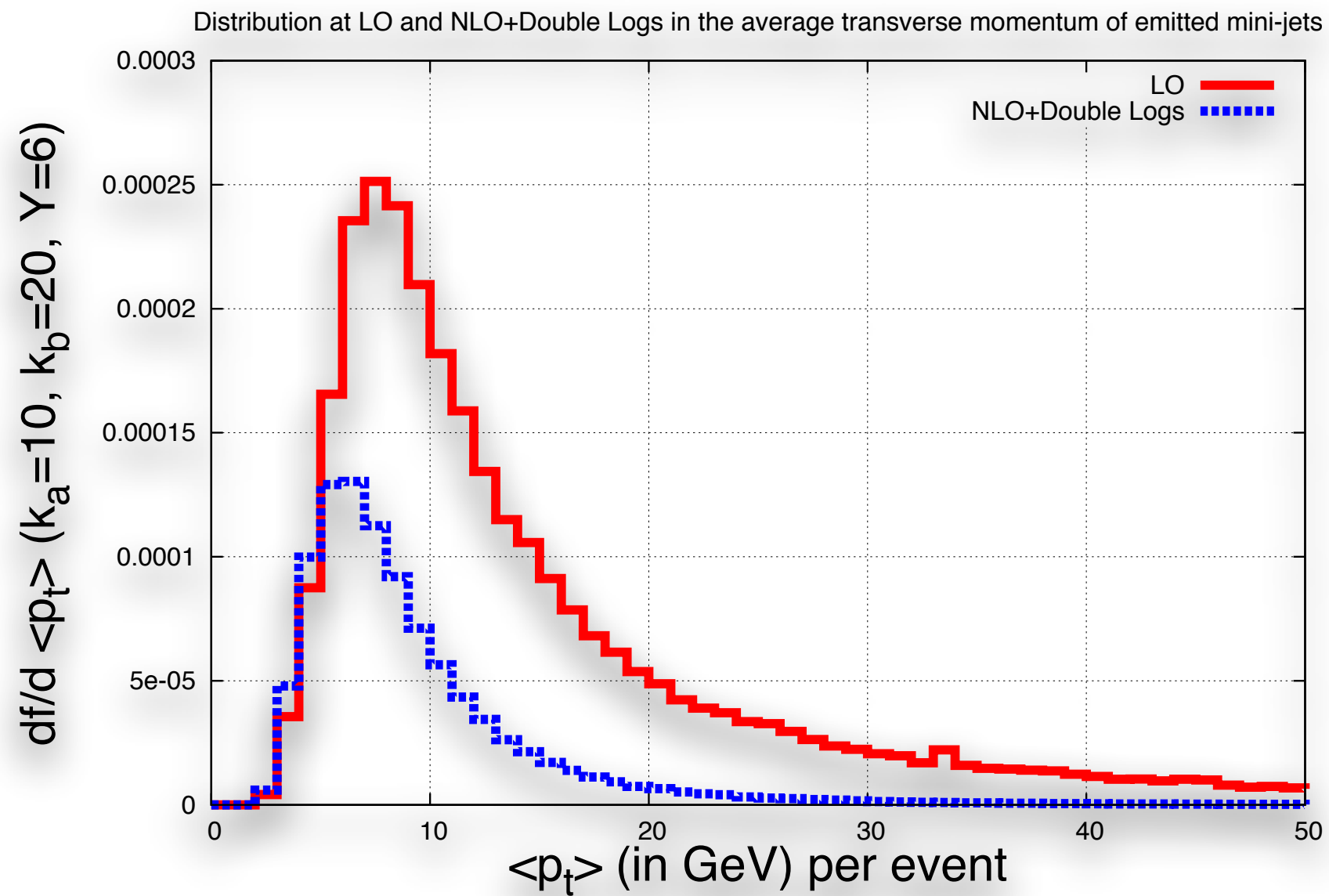
Distribution at LO in the average transverse momentum of emitted mini-jets



Distribution at NLO+Double Logs in the average transverse momentum of emitted mini-jets

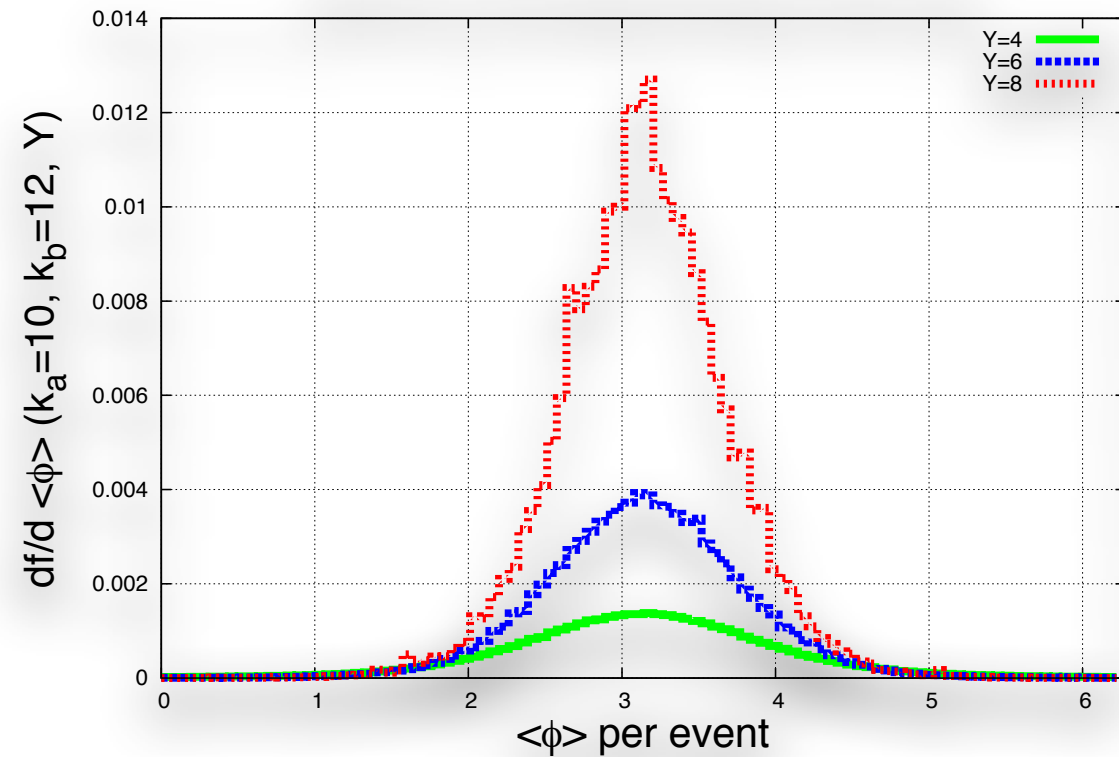


Average p_t

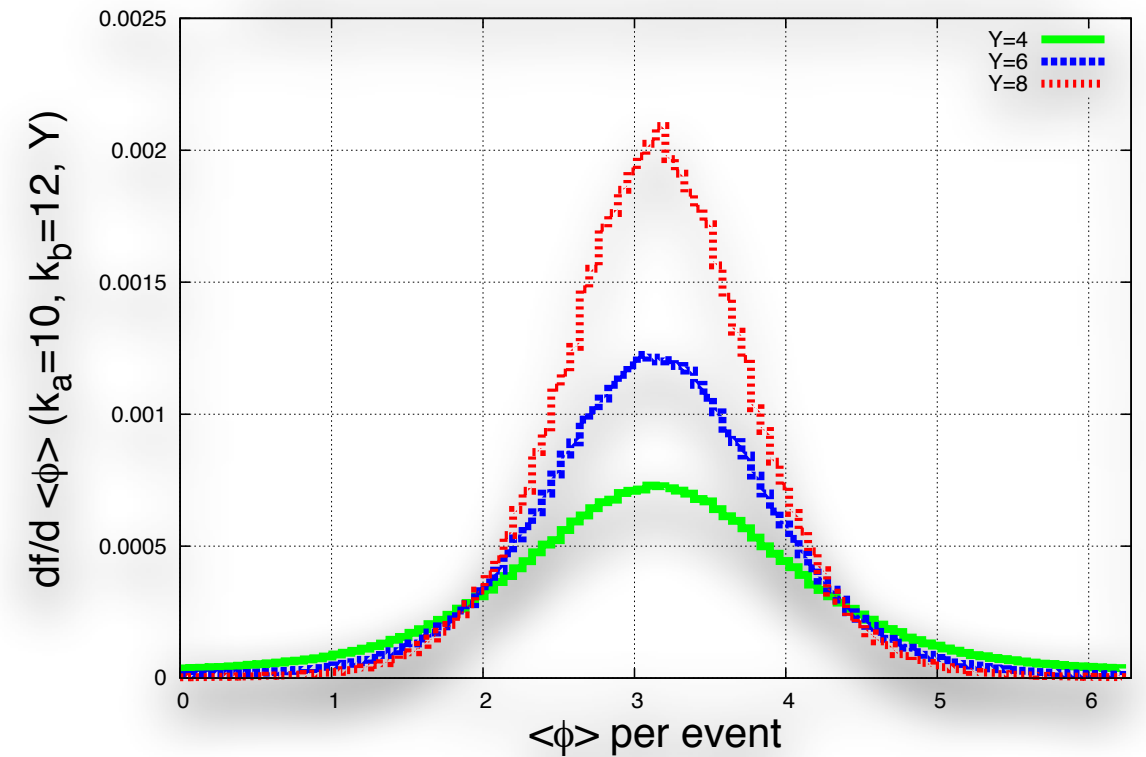


Average azimuthal angle

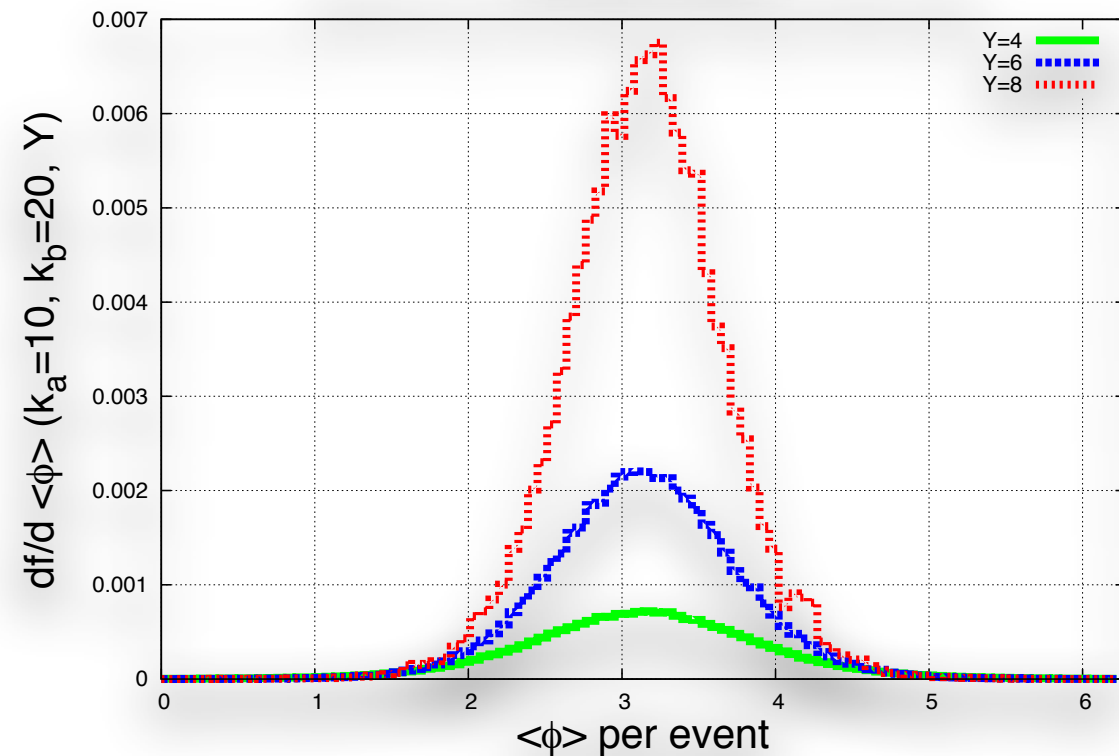
Distribution at LO in the average azimuthal angle of emitted mini-jets



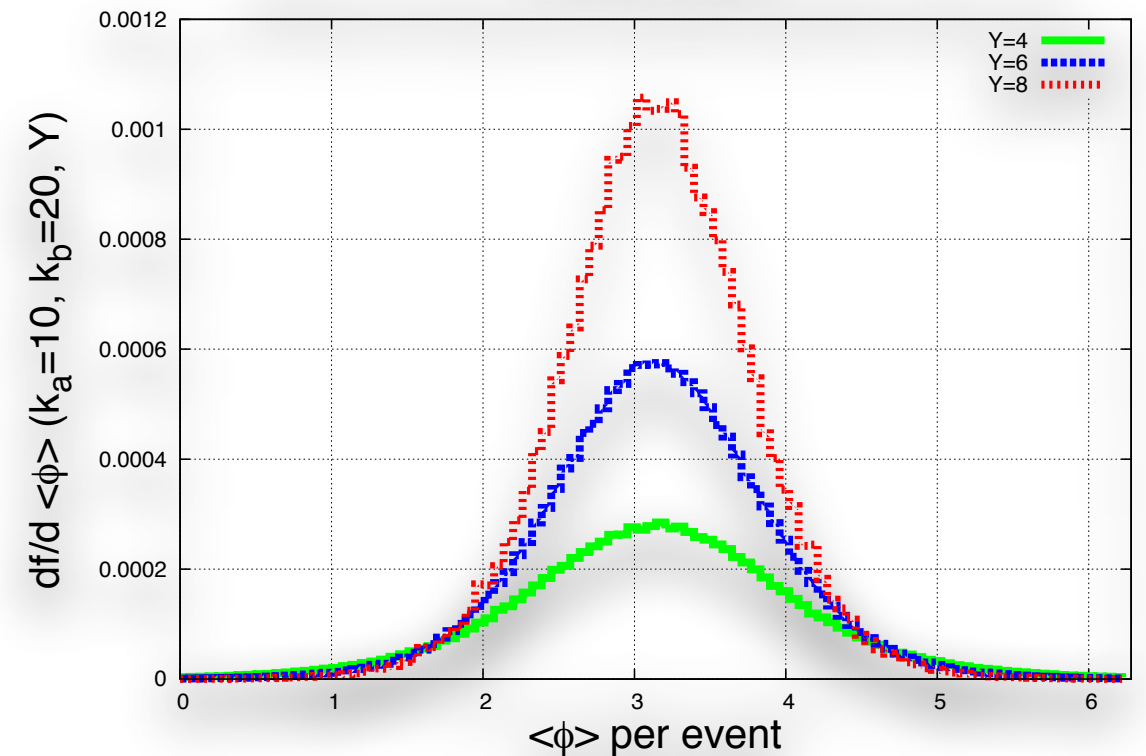
Distribution at NLO+Double Logs in the average azimuthal angle of emitted mini-jets



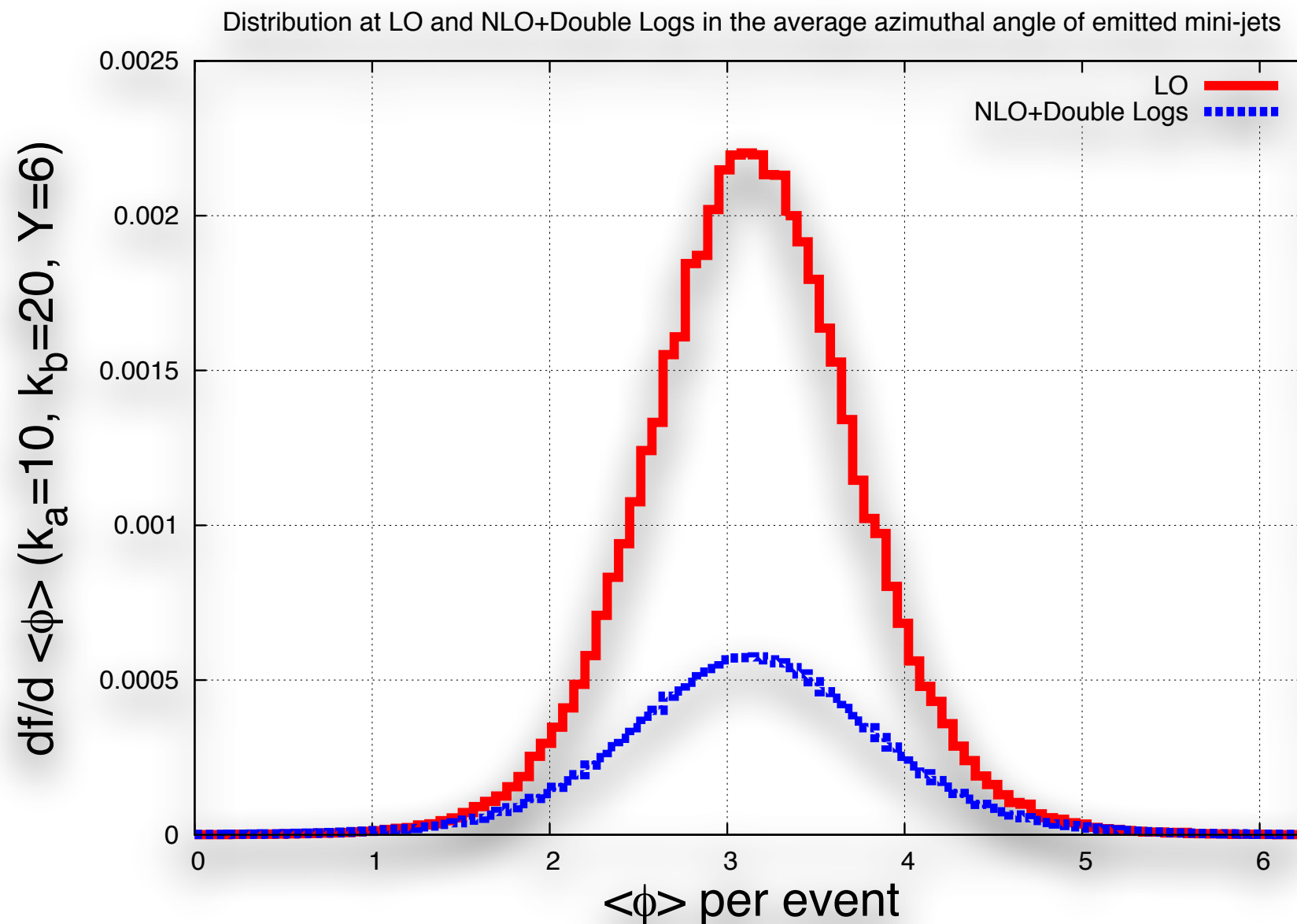
Distribution at LO in the average azimuthal angle of emitted mini-jets



Distribution at NLO+Double Logs in the average azimuthal angle of emitted mini-jets



Average azimuthal angle



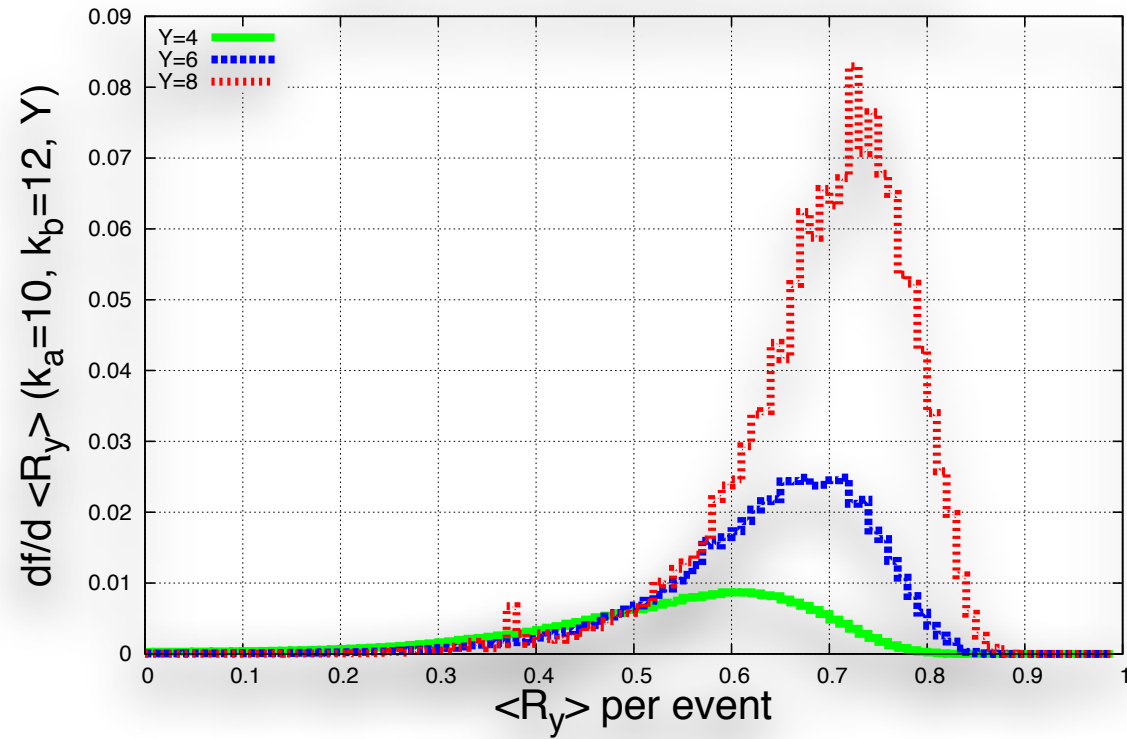
Average rapidity ratio for a typical multi-Regge kinematics event

- In a typical multi-Regge kinematics event we expect the total rapidity interval Y to be populated by N emissions equally spaced in rapidity intervals of length Δ such that $Y = (N+1) \Delta$. Then, the i -th emission is at $(N+1-i) \Delta$ and

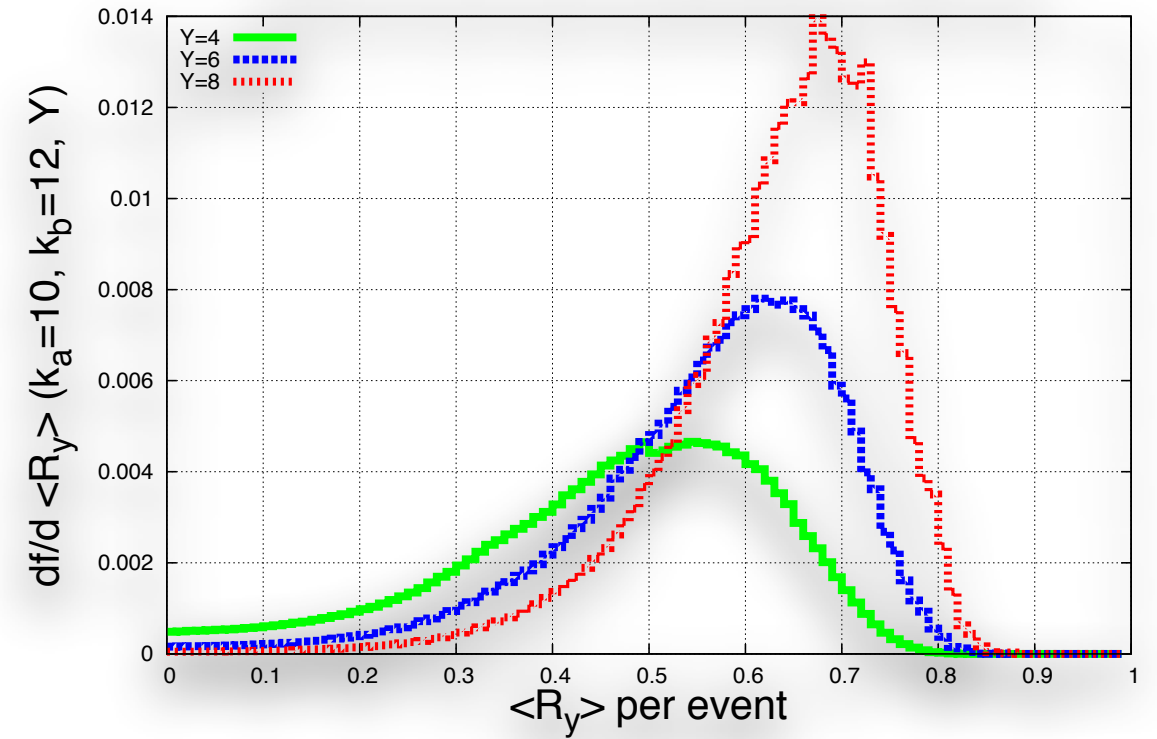
$$\begin{aligned}\langle \mathcal{R}_y \rangle_{\text{MRK}} &= \frac{1}{N+1} \sum_{i=1}^N \frac{i}{i+1} = \frac{N+1 - \psi(N+2) + \psi(1)}{N+1} \\ &= 1 + \frac{\Delta}{Y} \left(\psi(1) - \psi \left(1 + \frac{Y}{\Delta} \right) \right) \\ &\simeq 1 + \frac{\Delta}{Y} \left(\psi(1) + \ln \frac{\Delta}{Y} \right) + \dots\end{aligned}$$

Average rapidity ratio

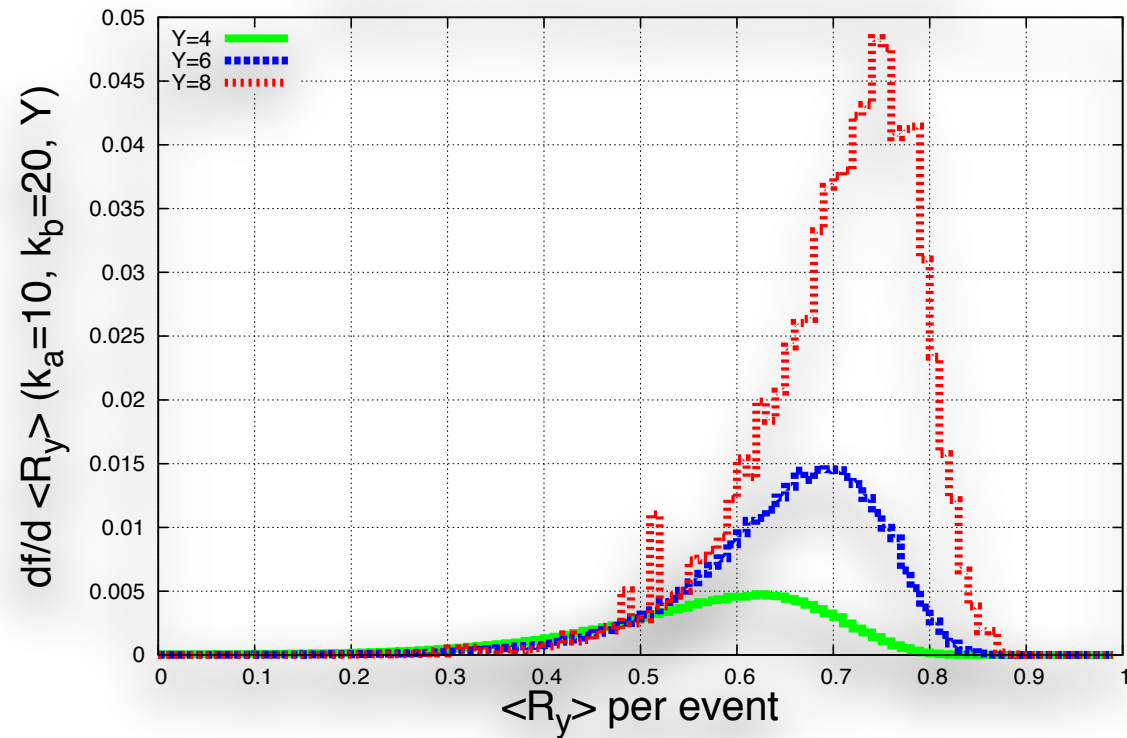
Distribution at LO in the average rapidity ratio of emitted mini-jets



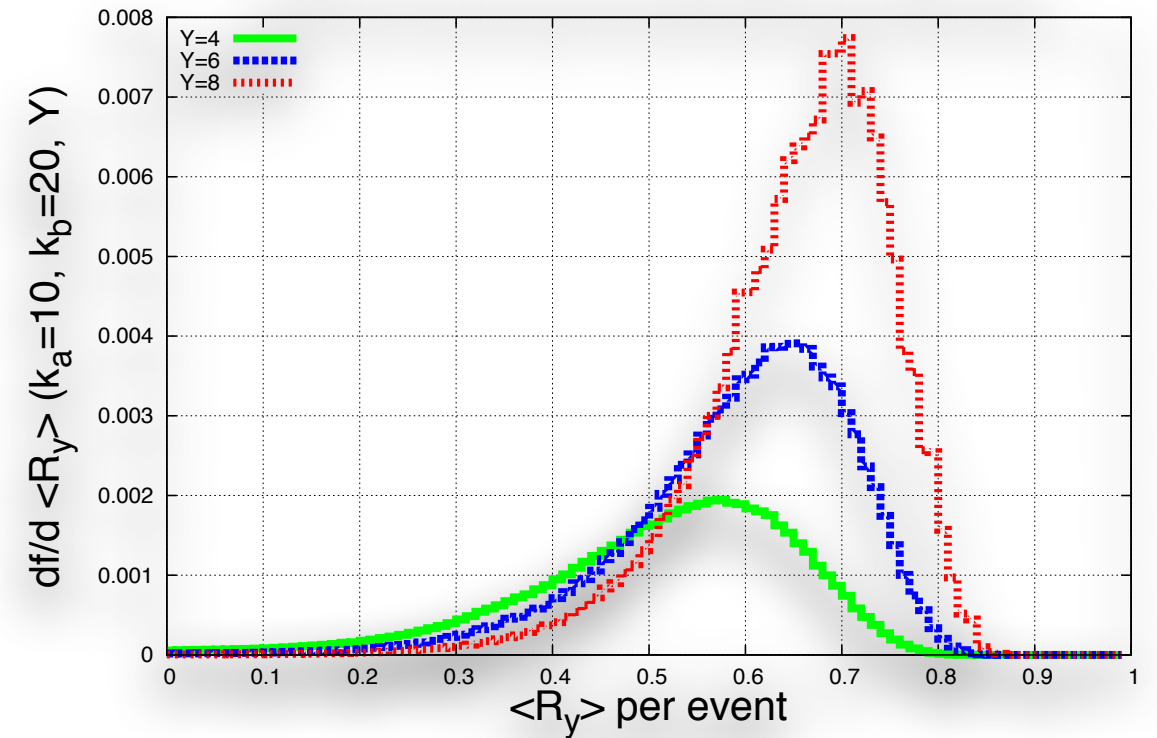
Distribution at NLO+Double Logs in the average rapidity ratio of emitted mini-jets



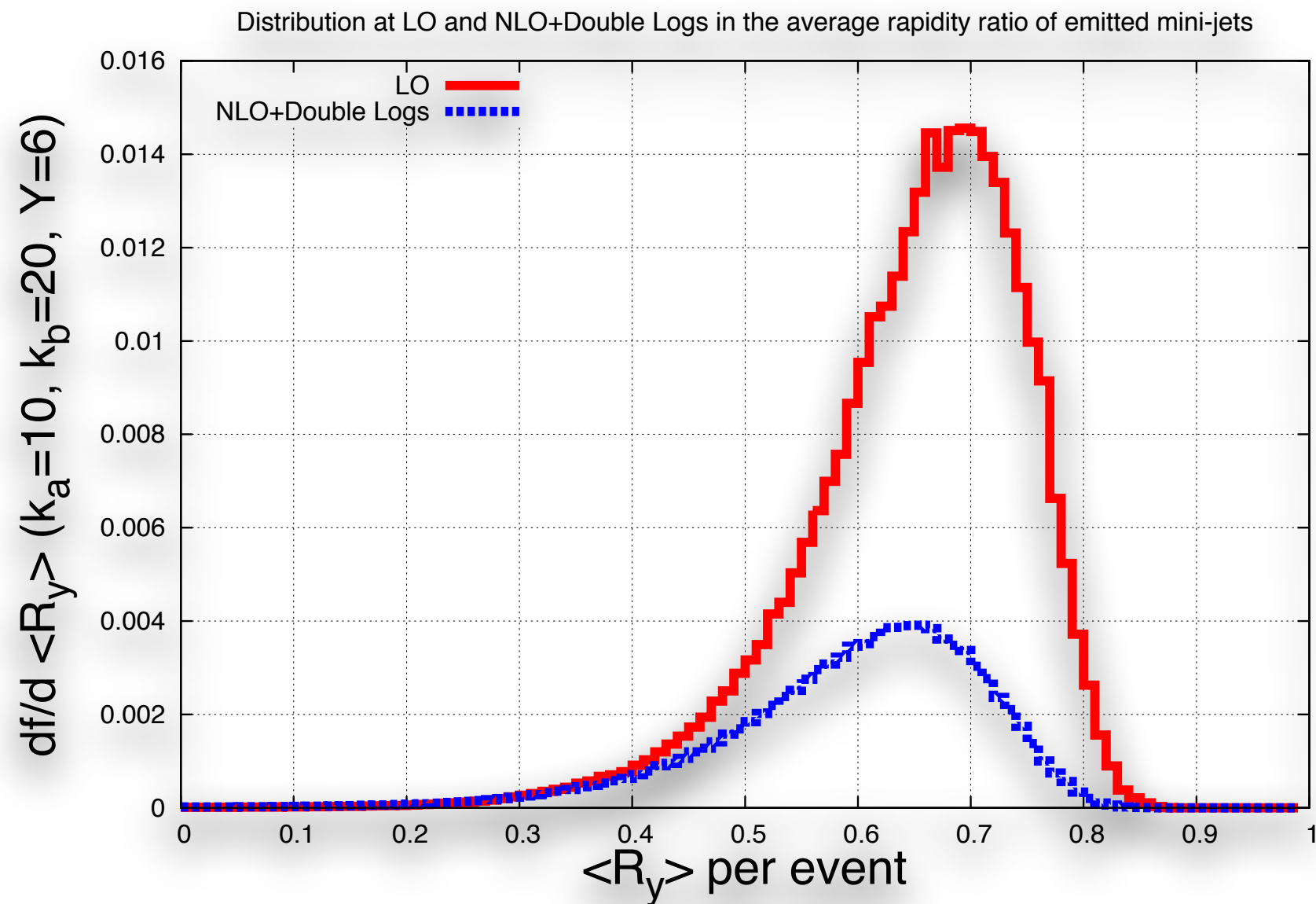
Distribution at LO in the average rapidity ratio of emitted mini-jets



Distribution at NLO+Double Logs in the average rapidity ratio of emitted mini-jets



Average rapidity ratio



Conclusions

- BFKLex is a MC tool with the NLO BFKL collinearly improved kernel implemented
- One can apply all sorts of kinematical cuts to pin down situation where the multi-Regge kinematics should be dominant
- Newly proposed observables can be studied in a fully differential way and new BFKL probes should emerge