

Search for SUSY in jets+MET final state (13 TeV)

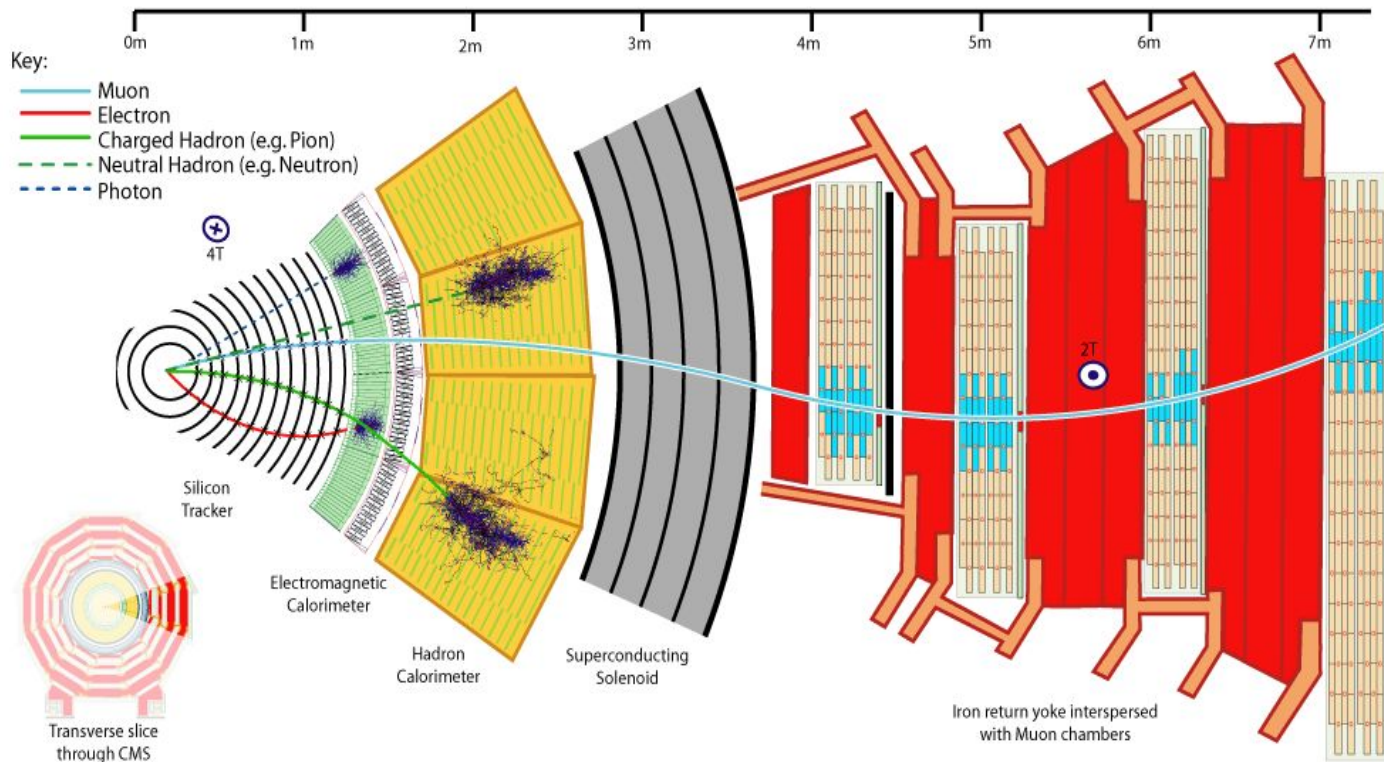
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Tata Inst. of Fundamental Research, India

DIS2016 ,12th April 2016, DESY Hamburg

(On behalf of CMS collaboration)

CMS detector



For jets +MET search :

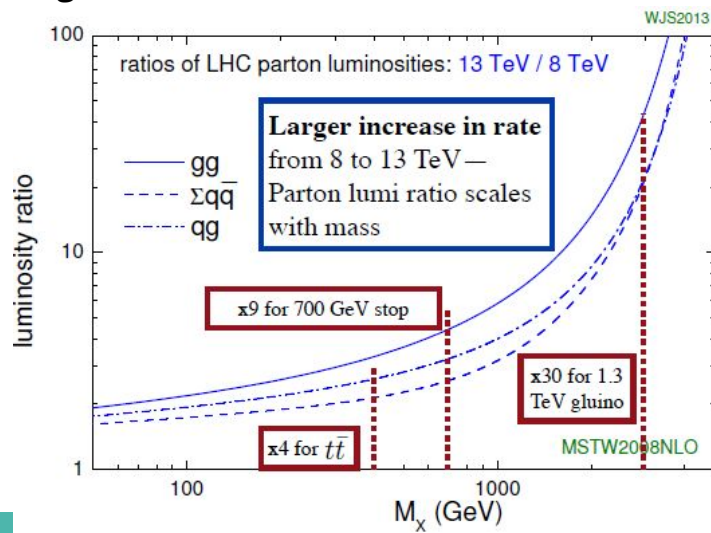
Jets : deposit major part of energy in calorimeter

MET : is calculated from all reconstructed particles in an event

.....
Tracker: helps to reduce pile-up from jets

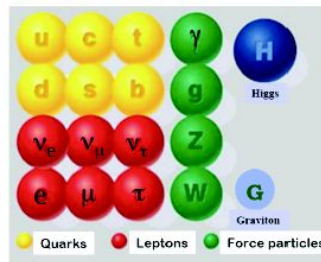
Introduction (8TeV to 13 TeV)

- 20 fb⁻¹ of 8TeV Data but no sign of SUSY (R parity Conserving)
- But we did discover something , a light higgs of mass ~ 125 GeV
- Naturalness puts some **SUSY particles at TeV scale accessible at LHC**
- 13 TeV parton luminosity gives a σ boost at high masses

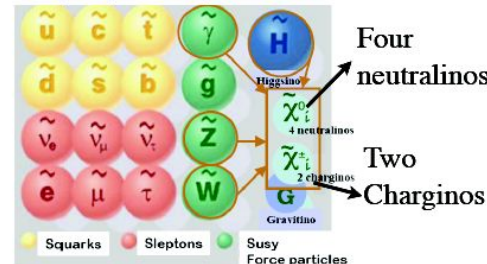


Supersymmetry : A super-partner of every SM particle differing by spin-half

Standard Model Particles



Supersymmetry Particles



SUSY is a broken symmetry : Expect new particles in ~TeV range !

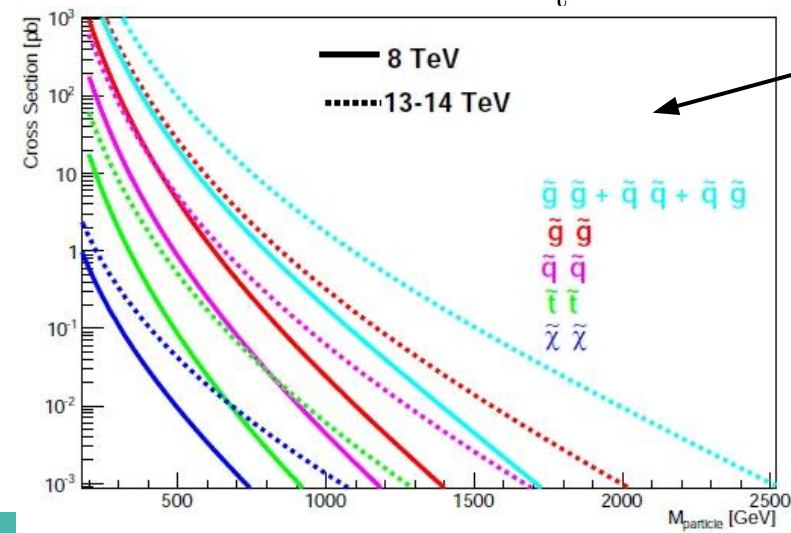
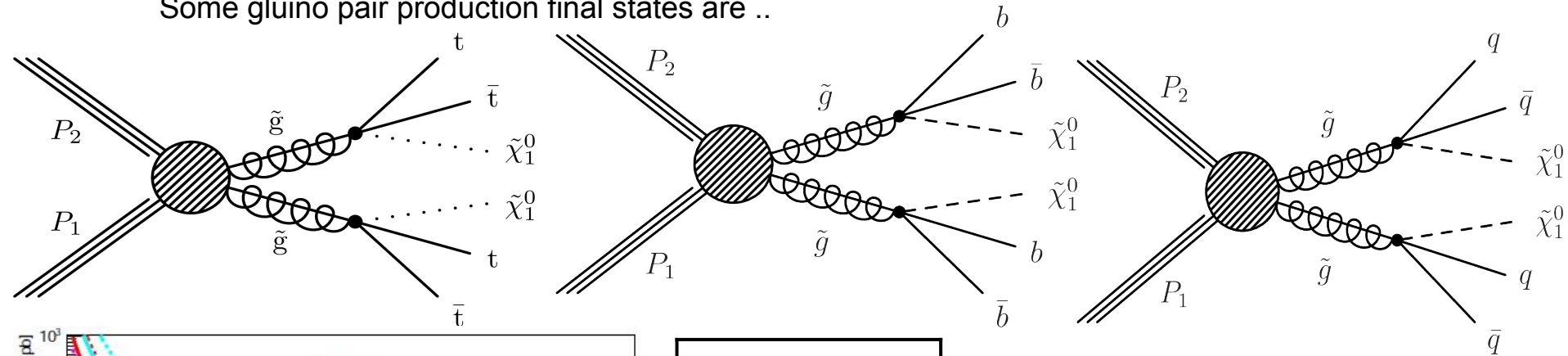
“Natural SUSY” requires :

- light stops, sbottoms (<1 TeV)
- not so heavy gluinos (1.5-2 TeV)
- light $\tilde{\chi}_1^0, \tilde{\chi}_1^\pm$ (few hundred GeV)

In R-Parity conserving models, lightest supersymmetry particle (LSP) is stable. Popularly lightest neutralino $\tilde{\chi}_1^0$ is an LSP.

SUSY scenarios

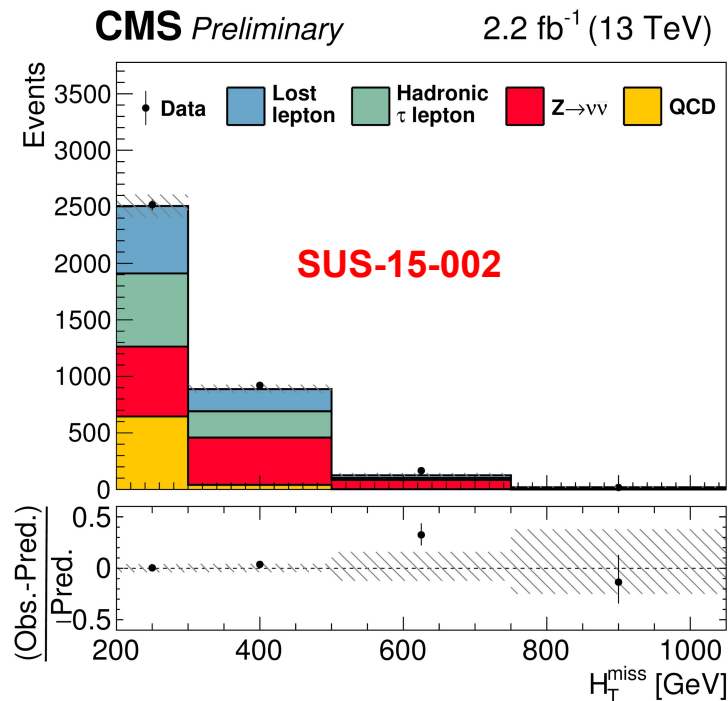
Some gluino pair production final states are ..



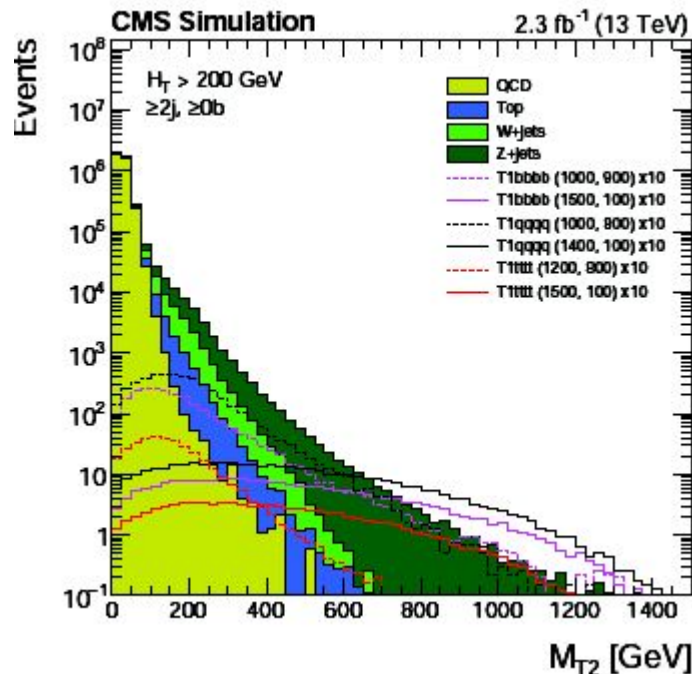
arXiv:1411.1427

- Inclusive searches designed to cover a broad range of scenarios for strongly produced SUSY
- Significant increase in gluino pair production cross section from 8 TeV to 13 TeV: ~50x for 1.5 gluinos.
- Expect to extend sensitivity of with a few fb⁻¹ of data.

Search Variable $H_T^{\text{miss}}, M_{T2}$



$$H_T^{\text{miss}} = \left| -\sum \mathbf{p}_T^{\text{jets}} \right|$$



SUS-15-003

$$M_{T2}(m_X) = \min_{\vec{p}_T^{X(1)} + \vec{p}_T^{X(2)} = \vec{p}_T^{\text{miss}}} \left[\max \left(M_T^{(1)}, M_T^{(2)} \right) \right]$$

Search variable α_T and Razor

$$\alpha_T = \frac{E_T^{j_2}}{M_T}$$

$$M_T = \sqrt{\left(\sum_{i=1}^2 E_T^{j_i}\right)^2 - \left(\sum_{i=1}^2 p_x^{j_i}\right)^2 - \left(\sum_{i=1}^2 p_y^{j_i}\right)^2}$$

$E_T^{j_2}$ is the transverse energy of less energetic jet 2

M_T is transverse mass of dijet system

SUS-15-005

Razor Variables:
For dijet system

$$M_R \equiv \sqrt{(|\vec{p}_{j_1}| + |\vec{p}_{j_2}|)^2 - (p_z^{j_1} + p_z^{j_2})^2}$$

$$M_T^R \equiv \sqrt{\frac{E_T^{\text{miss}}(p_T^{j_1} + p_T^{j_2}) - \vec{E}_T^{\text{miss}} \cdot (\vec{p}_T^{j_1} + \vec{p}_T^{j_2})}{2}}$$

$$R \equiv \frac{M_T^R}{M_R}$$

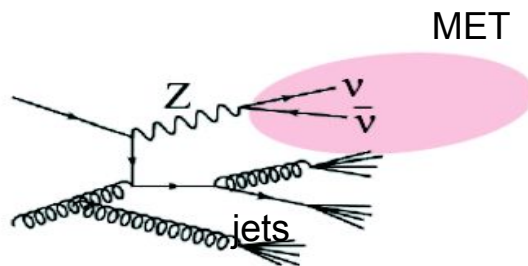
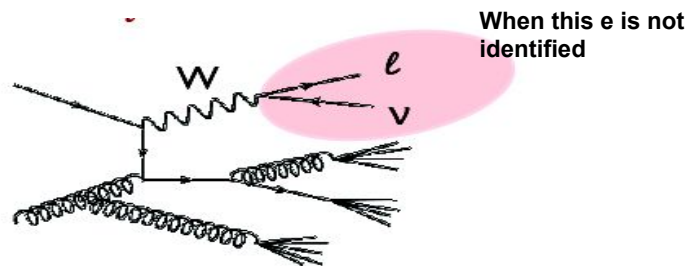
SUS-15-004

SM backgrounds

Dominant backgrounds are

- **Top/W +jets background**
 - Lepton is not identified
 - In $W(\tau\nu)$ +jets, when τ decays to hadronic final states
- **$Z(\nu\nu)$ +jets is an irreducible background**
 - This process has the same final state as the signal
- **QCD enters the search region because**
It can give rise to fake MET in the event through extreme tails of jet resolution
 - In QCD events there is no real MET
 - When jet energy is mismeasured, it gives rise to Fake MET

I will use MHT(**SUS-15-002**) analysis as an example to discuss about search region phase space and background estimation methods (. . .in next slides)

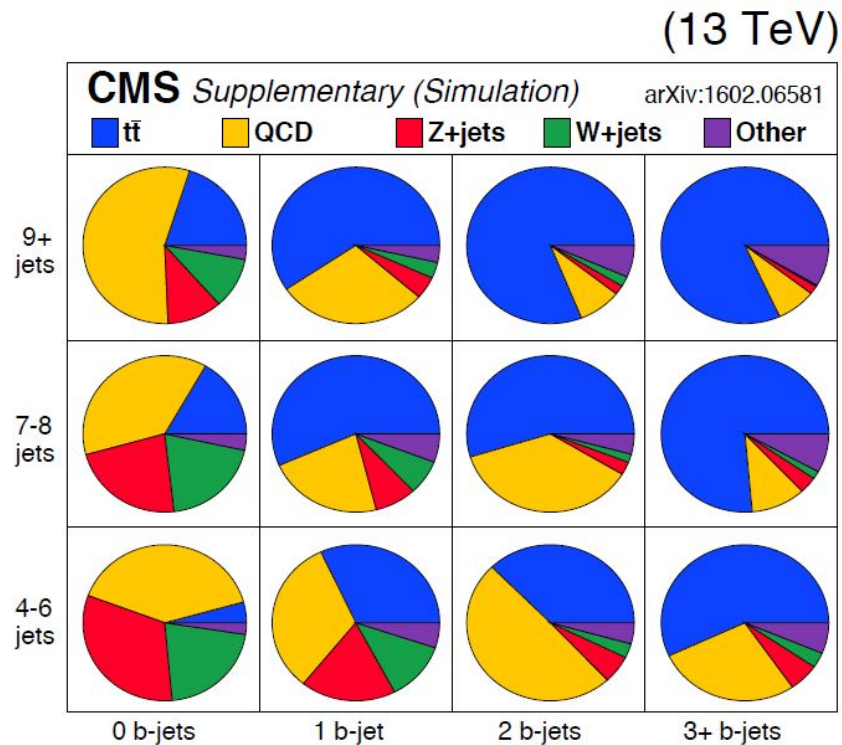
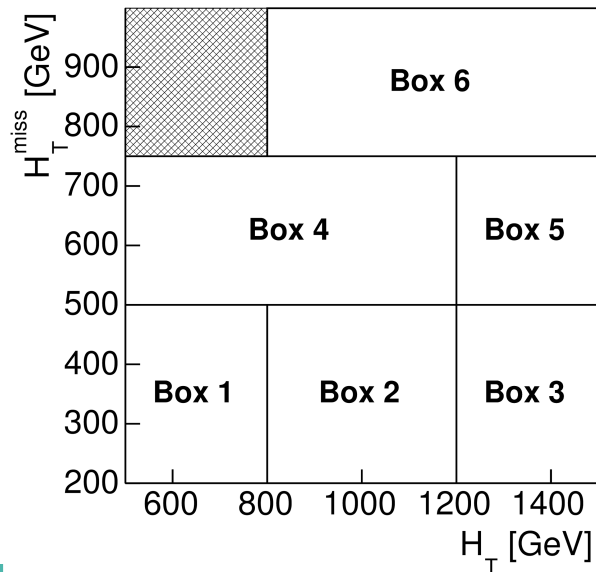


This process has exactly same final states as the signal

SUS-15-002(H_T^{miss}) search bins

- 6 bins in HT/MHT
- 3 bins in N_{jet}
- 4 bins in $N_{\text{b-jet}}$

Total number of search bins = $6 \times 3 \times 4 = 72$

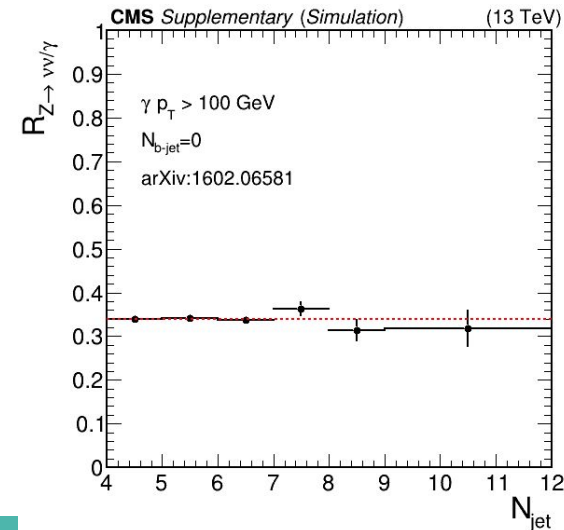
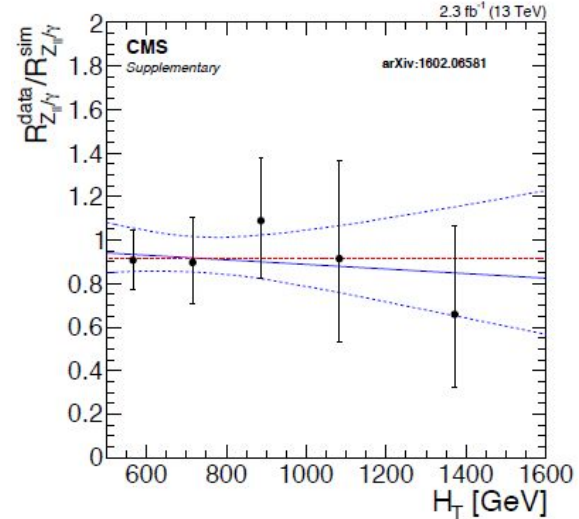


Z to neutrinos background

- A straightforward method is to use Z ($\rightarrow e+e-$ or $\mu+\mu-$) +Jets events as the topology of events is identical if one ignores leptons in the event
- But this process have few events in the control region
- γ + Jets control region has more events but also more systematic uncertainties
- A combination of both control region makes bkg estimation more effective

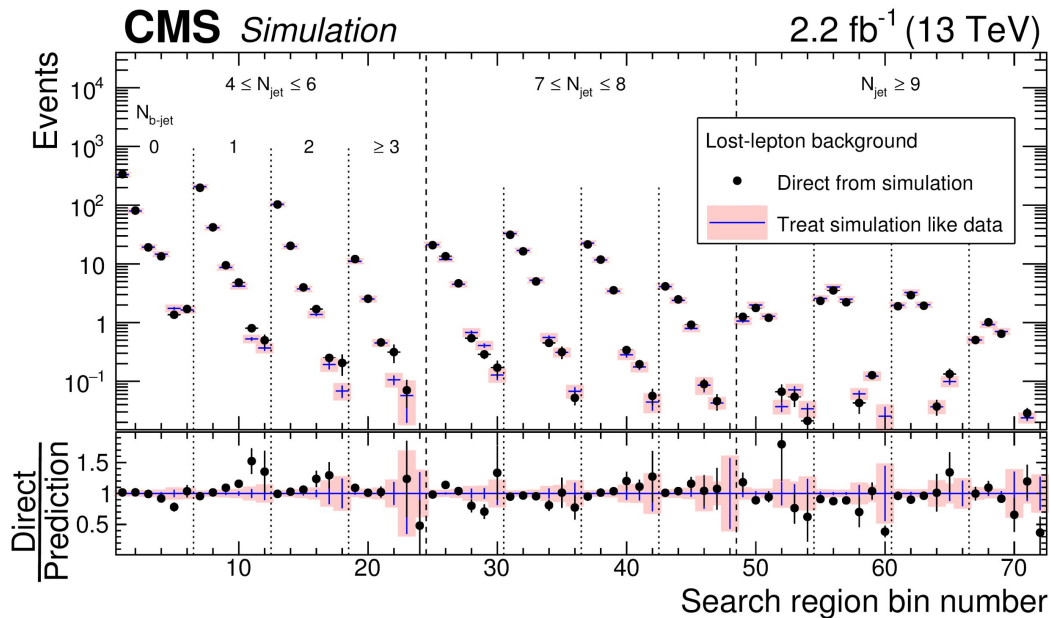
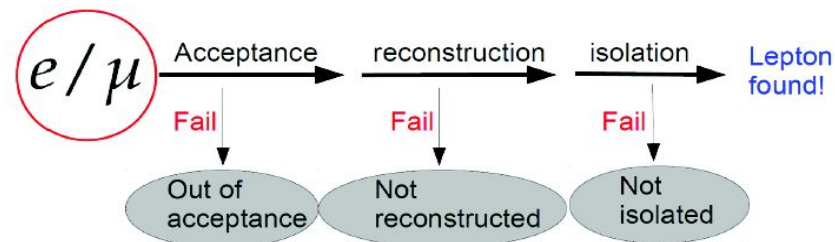
$$\boxed{N_{\gamma+\text{Jets}}} \times \boxed{\text{Correct for } \gamma \text{ identification, QCD contamination}} \times \boxed{\frac{\sigma(\text{Z+Jets})}{\sigma(\gamma+\text{Jets})}} = \boxed{N_{\text{Z}(\nu\nu)+\text{Jets}}}$$

Estimated Z($\nu\nu$)+jets events in data



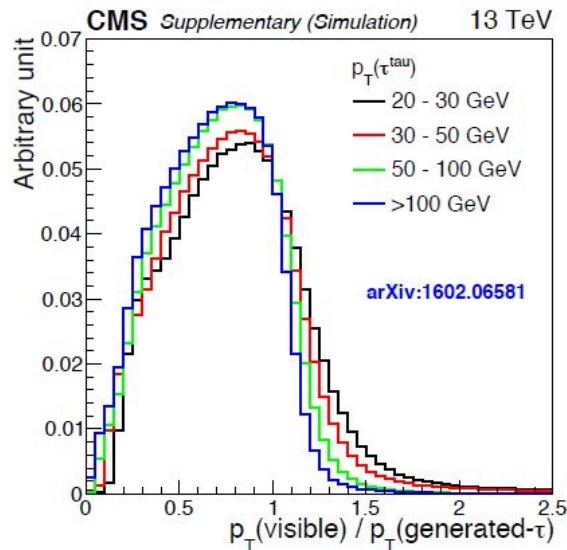
W/Top +jets bkg (lost lepton), data driven

- Enters the signal region when the lepton fails the lepton veto
- Use MC information to know the probability(ϵ_{eff}) of happening this
- Take single μ/e control sample of from data
- Use lepton ϵ_{eff} to trace back no of leptons that failed the

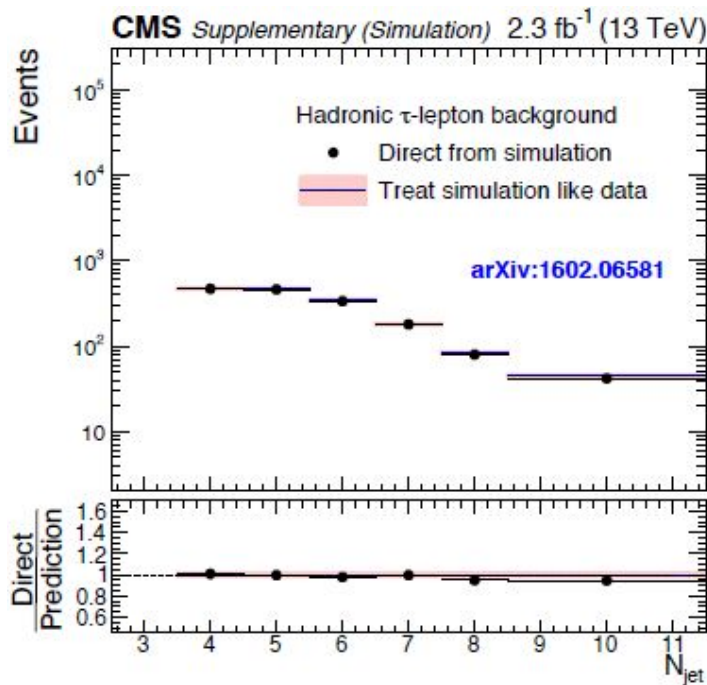


W+jets/top - hadronic tau background

- Results from hadronic decay of tau



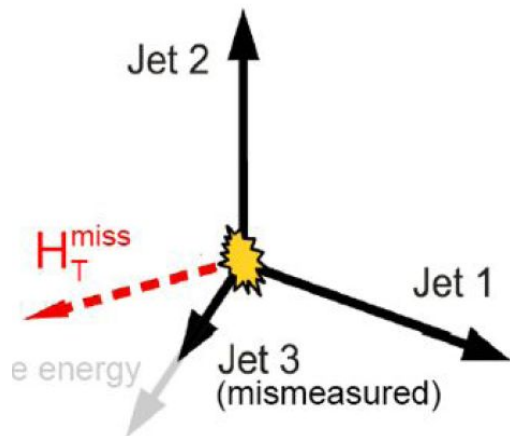
- Get single muon control sample from data
- Smear with hadronic tau template
- Gives the estimated number of had τ events



QCD multijet background

- QCD events do not have real missing energy but jet energy can be mismeasured
 - This results fake MET region
- This fake MET tends to be aligned with jet
- We reject 90% of background events with minimal effect on signal using cuts on low $\Delta\phi$

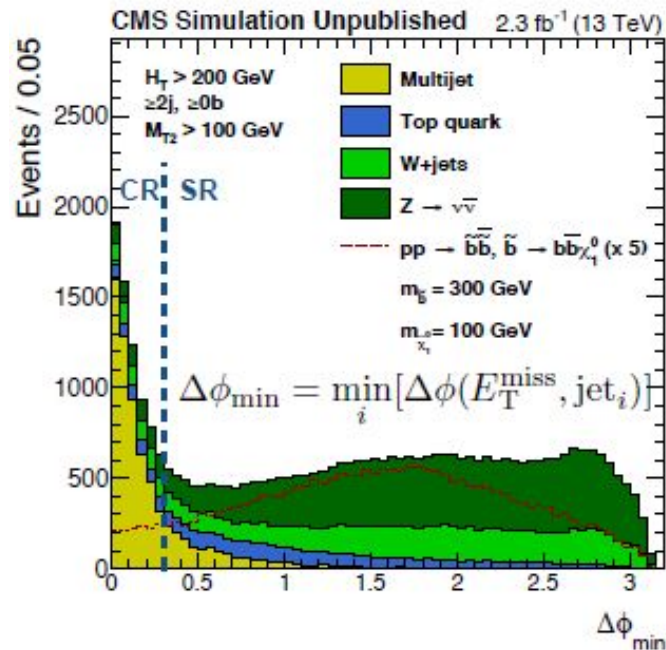
$\Delta\phi(H_T^{\text{miss}}, \text{ith jet}) = \text{Angle between } H_T^{\text{miss}} \text{ and ith jet}$



Low $\Delta\phi$ = QCD control region

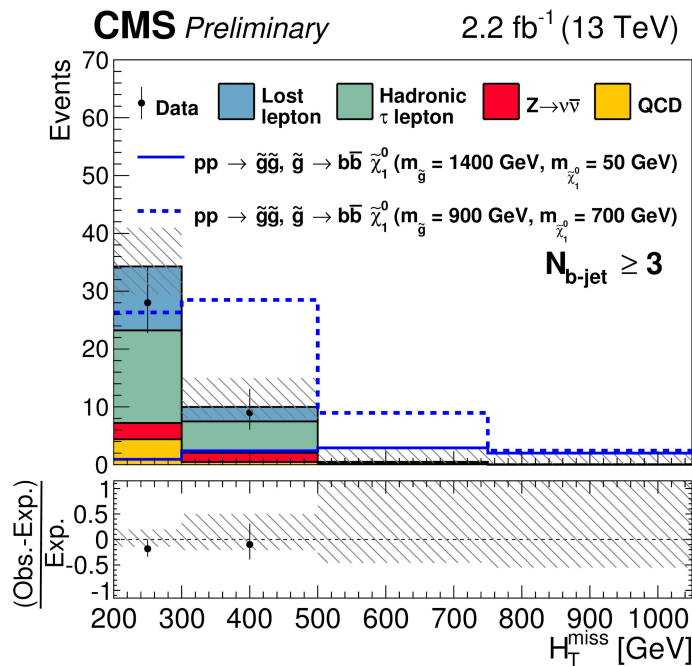
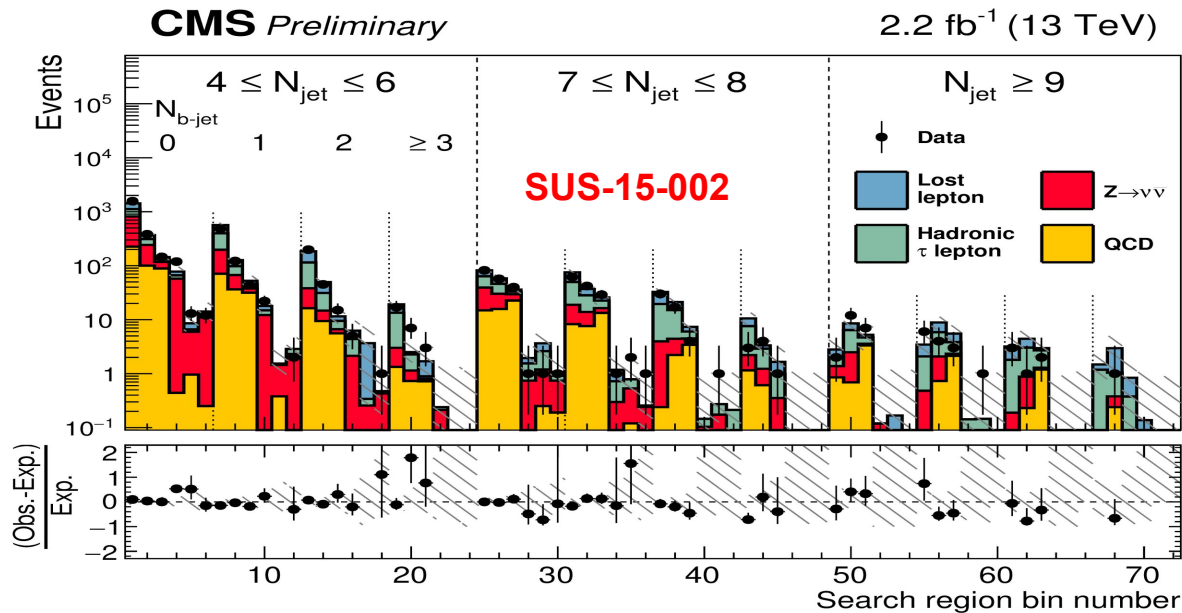
High $\Delta\phi$ = signal region

Use high/low ratio to go from control region to signal region

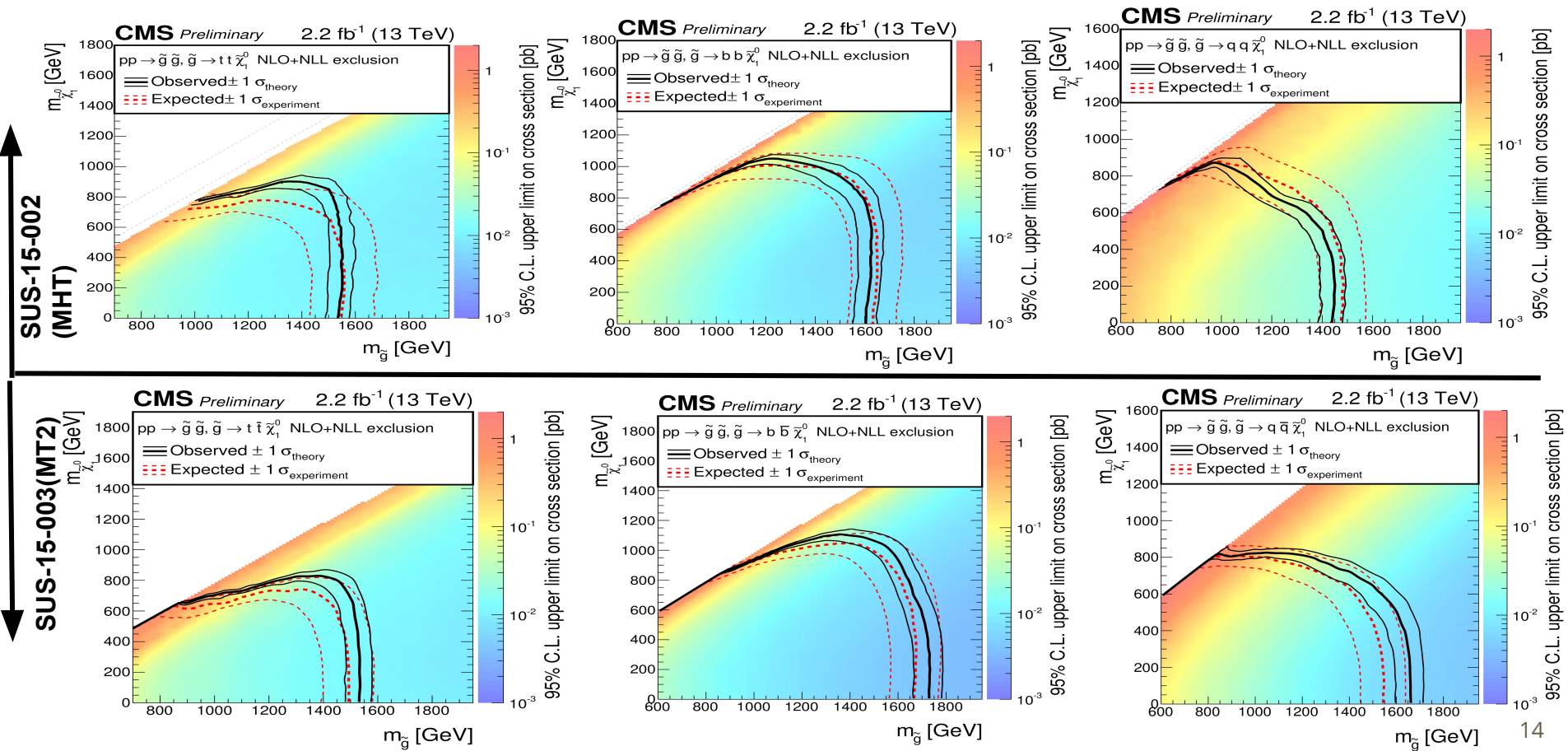


Data vs SM backgrounds

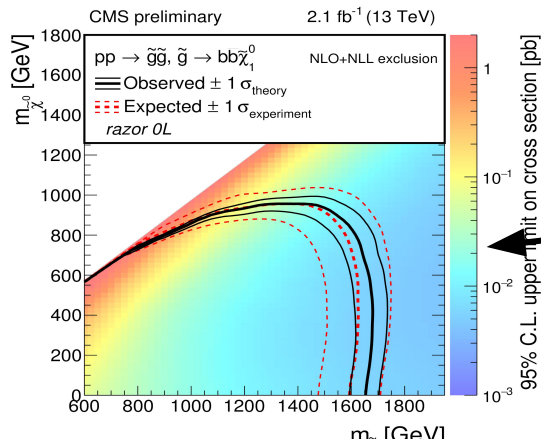
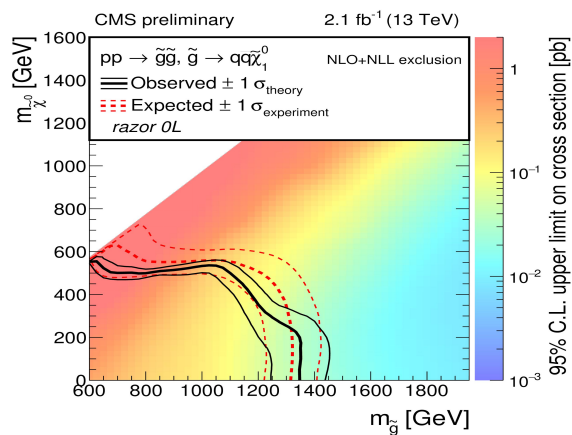
- No broad excess consistent with signal expectation
- Data agrees with SM background expectation



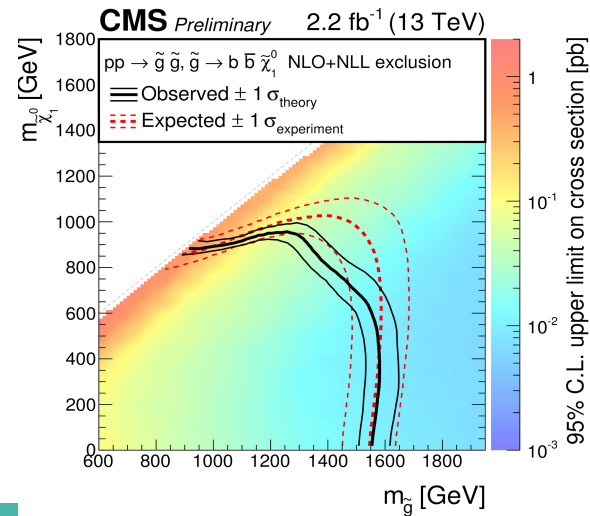
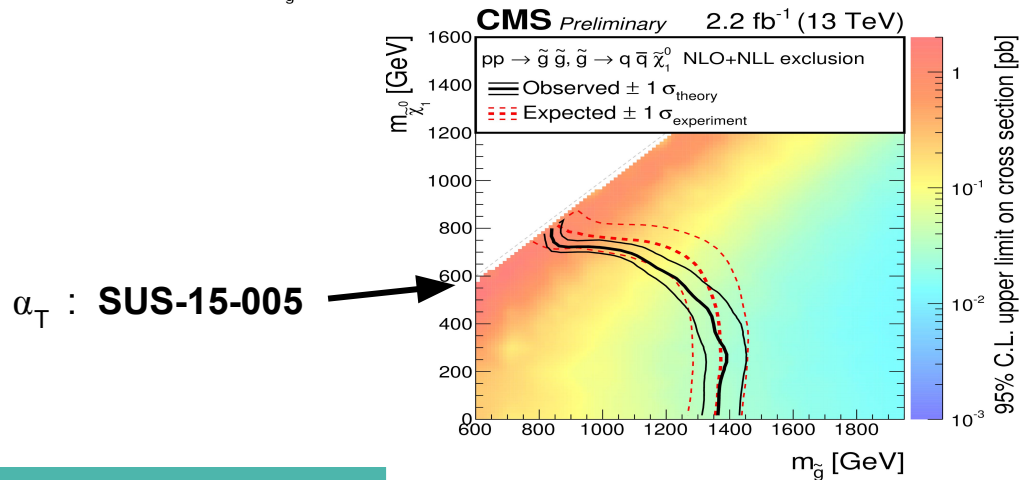
Interpretation: Gluino pair production



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RaZOR : SUS-15-004



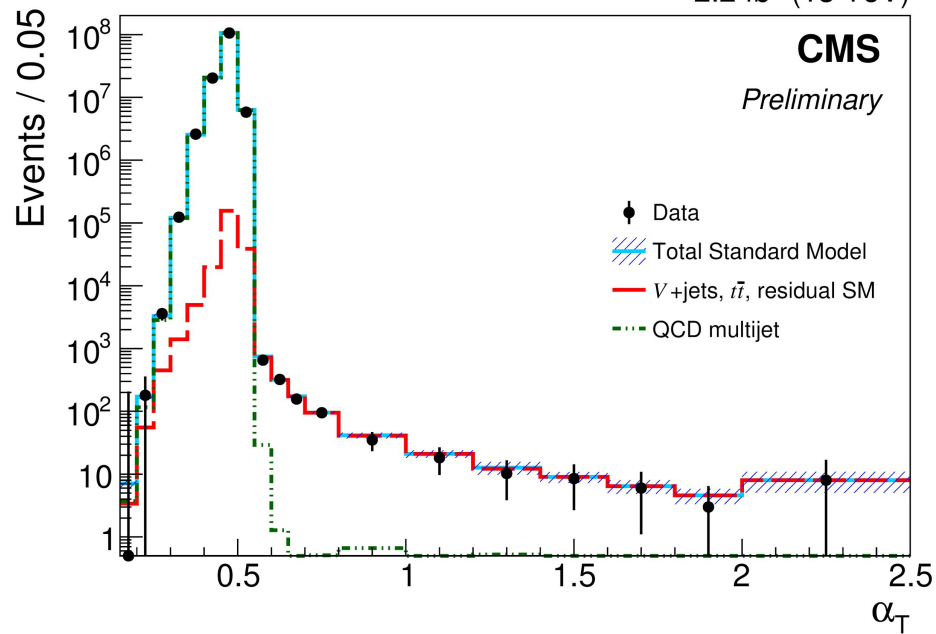
Conclusion

- **Results from a variety of analyses to search for supersymmetric partners in 2.3 fb^{-1} of 13 TeV data in CMS experiment**
- **The gluinos of mass up to 1.6 TeV excluded for low LSP masses**
- **The exclusion limits are improved as much as 150-200 GeV compared to the 8 TeV limits**

Back-up slides

α_T

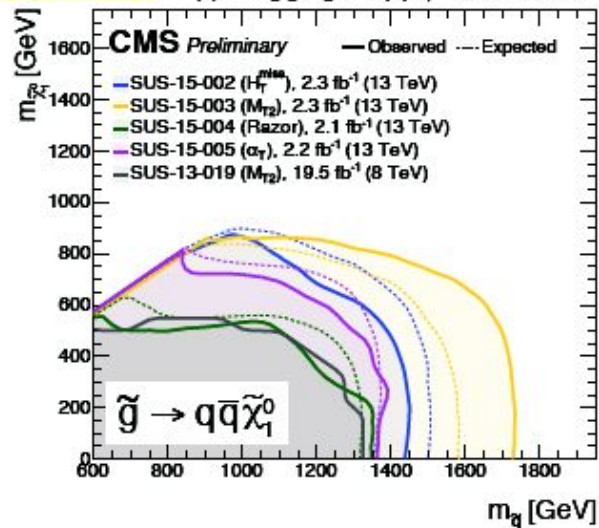
2.2 fb⁻¹ (13 TeV)



Interpretation : gluino pair production

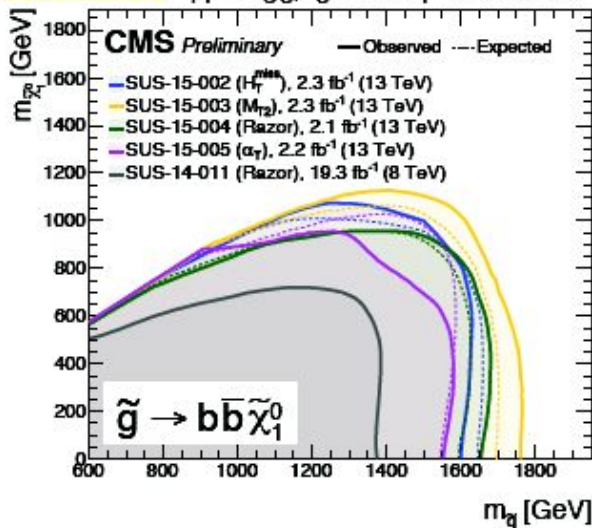
Updated

$pp \rightarrow \tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$ Moriond 2016



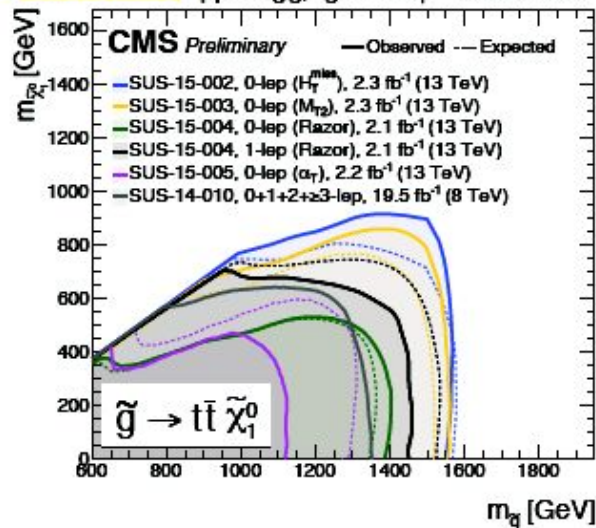
Updated

$pp \rightarrow \tilde{g}\tilde{g}, \tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$ Moriond 2016



Updated

$pp \rightarrow \tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ Moriond 2016



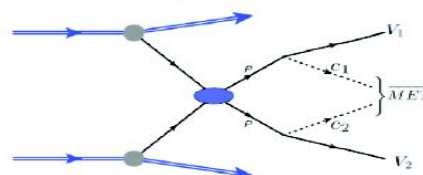
All hadronic search variables and SM backgrounds

$H_T = \sum_{i=\text{jets}} |p_T^{\text{jets}}|$ = scalar sum of p_T of all jets

N_{jet} = Number of jets in the event

$N_{\text{b-jet}}$ = Number of b-jets in the event

Analysis designed search variables

<p>All jets \rightarrow</p> $H_T^{\text{miss}} = -\sum p_T^{\text{jets}} $ <p>Missing Transverse Hadronic energy</p> <p>SUS-15-002</p>	<p>For dijet events</p> $\alpha_T = \frac{E_T^{j2}}{M_T}$ <p>E_T^{j2} is the transverse energy of less energetic jet 2</p> <p>M_T is transverse mass of dijet system</p> <p>SUS-15-005</p>
<p> $M_{T2}(m_X) = \min_{\vec{p}_T^{X(1)} + \vec{p}_T^{X(2)} = \vec{p}_T^{\text{miss}}} \left[\max(M_T^{(1)}, M_T^{(2)}) \right]$ </p>  <p>SUS-15-003</p>	<p> $M_R \equiv \sqrt{(\vec{p}_{j1} + \vec{p}_{j2})^2 - (p_z^{j1} + p_z^{j2})^2}$ </p> <p> $M_T^R \equiv \sqrt{\frac{E_T^{\text{miss}}(p_T^{j1} + p_T^{j2}) - \vec{E}_T^{\text{miss}} \cdot (\vec{p}_T^{j1} + \vec{p}_T^{j2})}{2}}$ </p> <p> $R \equiv \frac{M_T^R}{M_R}$ </p> <p>SUS-15-004</p>