

# The EW Sudakov approximation in SHERPA

Jennifer Thompson

II. Physikalisches Institut, Universität Göttingen

*[jennifer.thompson@physik.uni-goettingen.de](mailto:jennifer.thompson@physik.uni-goettingen.de)*

April 13, 2016

DIS16 DESY

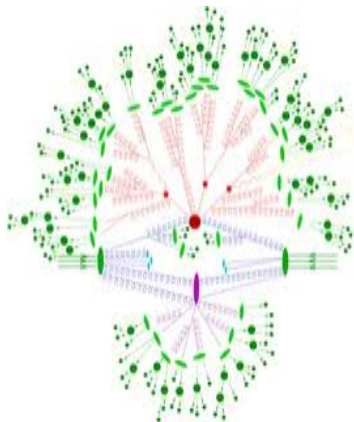


# Overview

- 1 Motivation
- 2 Next steps in Monte Carlo accuracy
- 3 Sudakov Logarithms
- 4 Full NLO EW Corrections
- 5 Conclusions

# Why a Monte Carlo?

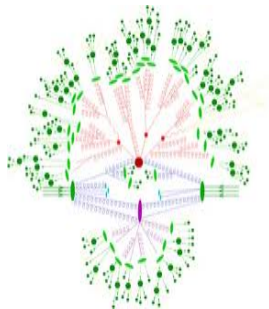
- 1 Analytical calculations deal with low numbers of final state particles.
- 2 Experiments have high multiplicity final state with jets and missing energy.
- 3 In order to test theoretical predictions we need a way to compare the 2.
- 4 Monte Carlo simulations provide this bridge between theory and experiment.



# Precision Calculations in Monte Carlo Simulation

There are several components to a Monte Carlo simulation

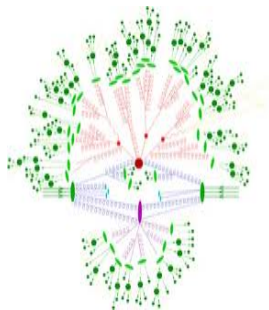
- Hard Process - Perturbative.
- Underlying Event.
- Parton shower - All order resummation.
- Decays of unstable particles.
- Hadronisation.
- Decays of unstable hadrons.



# Precision Calculations in Monte Carlo Simulation

There are several components to a Monte Carlo simulation

- **Hard Process - Perturbative.**
- Underlying Event.
- Parton shower - All order resummation.
- Decays of unstable particles.
- Hadronisation.
- Decays of unstable hadrons.



# Numerical NLO Calculation

## Numerical NLO

In a numerical implementation, divergences need to be dealt with at the integrand phase.

$$\sigma_{NLO} = \int (B + V) d\Phi_B + \int R d\Phi_R$$

The integrals over  $V$  and  $R$  separately contain divergences, which will cancel after the integration has been performed.

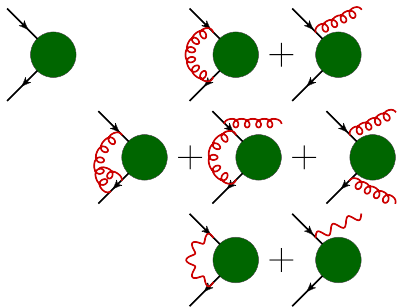
# Subtraction Method

$$\sigma_{\text{NLO}} = \int (\text{B} + \text{V} + \int \text{S} d\Phi_1) d\Phi_{\text{B}} + \int (\text{R} - \text{S}) d\Phi_{\text{R}}$$

- Introduce a term that:
  - ① Exactly matches the divergent structure of the real term.
  - ② Does not introduce any additional divergences.
  - ③ Can be analytically integrated over the 1 parton subspace.
- Subtract this term from the real contribution.
- Integrate over the 1 parton sub space and add this in to the virtual term.

## N(N)LO QCD and EW

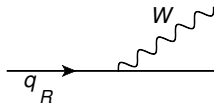
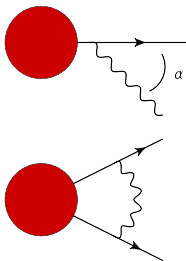
- Current Monte Carlo simulations are NLO QCD with a parton shower.
- Extending to NNLO QCD  $\rightarrow$  precision sensitive to EW corrections.  
( $\mathcal{O}(\alpha_s^2) \sim \mathcal{O}(\alpha)$ )





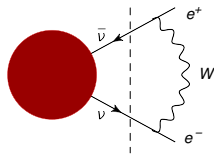
# Considerations in NLO EW

Masses of weak bosons regulate divergences



Weak interactions depend on helicity

Weak boson exchange can alter the external flavours



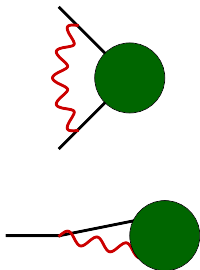
interferences with different Born MEs.

## Sudakov Approximation

- The virtual contribution yields logarithms  
→  $\log\left(\frac{s}{M_V^2}\right)$   
 $M_V$  is boson mass,  $\sqrt{s}$  is COM energy.
- Divergent in limit of vanishing mass.  
→ therefore dominant in the high energy regime
- Outside of this regime, a full NLO EW calculation is necessary.
- Non-logarithmic contributions always at  $\sim 1\%$  level, so EW Sudakov approximation is limited by this.

## Contributing Diagrams - Mass Singularities

Relevant diagrams as shown in J. Math. Phys. 3, 650 (1962)



- Produced from the exchange of bosons.
- Bosons can become soft and/or collinear, giving logarithms.
- This leads to divergences for vanishing mass.
- Real contribution for  $V = W^\pm, Z$  emission distinguishable.
- We are left with large logarithms of  $\log s/M_V^2$ .
- Call these Sudakov logarithms

# One-Loop Approximation

hep-ph/0010201v3

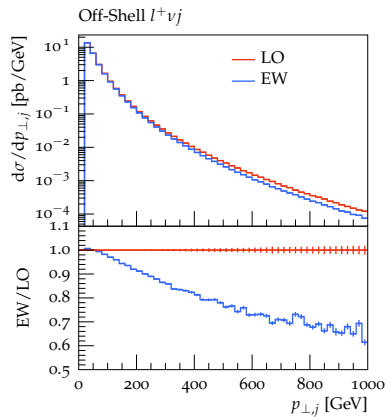
- My implementation follows A. Denner and S. Pozzorini.
- Analytical calculations for one and two loops for some processes available.
- First approximation is the one-loop calculation.

## One-loop

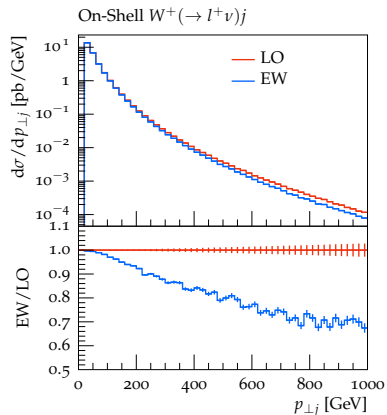
$$\frac{\alpha}{4\pi} \left[ A \log^2 \left( \frac{(p_i + p_j)^2}{M^2} \right) + B \log \left( \frac{(p_i + p_j)^2}{M^2} \right) \right]$$

$p_i$ ,  $p_j$  are the momenta for external legs  $i$  and  $j$ .  $M$  is the mass of the EW bosons.

# Results - $pp \rightarrow Wj$ at 14 TeV LHC



Off-shell production



On-shell production

## Sudakov Logarithmic Approximation vs. Full NLO

- Sudakov approximation is only valid in the high energy limit.
- Sudakov approximation is only valid for well-separated EW particles.
- Sudakov approximations do not capture full NLO behaviour.
- Sudakov approximation is far less computationally intensive.
- Sudakov approximations are simple to include on top of NLO QCD calculations.

→ Need implementations of both

## NLO EW with SHERPA

arXiv:1511.08692, arXiv:1505.05704, arXiv:1412.5157

- There is currently an implementation (not public) of full NLO EW computations with SHERPA+OpenLoops
- Currently papers on  $V$ +jets
- Current work to also interface to Recola
- Recola generates NLO MEs and will be a loop provider in the interface.
- Aim to automate EW corrections in the same way as NLO QCD corrections.

## Conclusions

- EW Corrections are of a similar size to NNLO QCD Corrections.
- EW Sudakov logarithms are a good approximation in the high-energy limit
- These have been implemented within SHERPA
- More accurate predictions require NLO EW + NLO QCD corrections.