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Measurement of photon production cross sections with the ATLAS detector

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On behalf of the ATLAS Collaboration

• Outline

- \rightarrow Physics with photons
- \rightarrow Photons with the ATLAS detector
- \rightarrow Inclusive photon production at 8 TeV
- \rightarrow Inclusive photon production at 13 TeV
- \rightarrow Summary



Photon production in pp collisions at LHC

- \bullet Photon production in pp collisions allows
- \rightarrow tests of perturbative QCD
- \rightarrow to extract information on the proton PDFs
- Possibilities to study inclusive production of photons or in association with jets
- Prompt photons represent a cleaner probe of the hard interaction than jet production
- Of special interest as the 2nd largest source of background to $H o \gamma\gamma$



Other sources of photons

- Quarks and gluons are sources of photons
- \rightarrow Quarks and gluons hadronize mostly into pions and, by isospin symmetry, 1/3 are π^0 's, which decay into two photons $\Rightarrow \gamma$'s are produced copiously inside jets!
- \rightarrow Quarks have electric charge and radiate photons
 - \Rightarrow fragmentation function $D_{q/g}^{\gamma}(z,\mu_f)$

 \Rightarrow Distinct feature: these photons are inside jets, i.e. not isolated!





- It is essential to require the photon to be isolated. It is achieved by requiring $E_T^{iso} \equiv \sum_i E_T^i < E_T^{max}$ with the sum over the particles (except the photon!) inside
 a cone of radius R = 0.4 centered on the photon in the $\eta \phi$ plane
 - ightarrow For the analyses presented here: $E_T^{
 m max} = 4.2 \cdot 10^{-3} \cdot E_T^{\gamma} + 4.8~{
 m GeV}$
- The isolation requirement suppresses the contribution of photons inside jets: π^0 (as well as other neutral mesons) decays and the fragmentation contribution

NLO QCD calculations for inclusive photon production

$$egin{split} \sigma_{pp o \gamma+\mathrm{X}} &= \sum_{i,j,a} \int_{0}^{1} dx_{1} \; f_{i/p}(x_{1},\mu_{F}^{2}) \int_{0}^{1} dx_{2} \; f_{j/p}(x_{2},\mu_{F}^{2}) \; \hat{\sigma}_{ij o \gamma a} + \ &\sum_{i,j,a,b} \int_{z_{min}}^{1} dz \; D_{a}^{\gamma}(z,\mu_{f}^{2}) \int_{0}^{1} dx_{1} \; f_{i/p}(x_{1},\mu_{F}^{2}) \int_{0}^{1} dx_{2} \; f_{j/p}(x_{2},\mu_{F}^{2}) \; \hat{\sigma}_{ij o ab} \end{split}$$

- The calculations include NLO corrections for both direct-photon and fragmentation contributions; <u>beware</u> the components are <u>not</u> distinguishable beyond LO
- Isolation at "parton level": E_T^{iso} calculated with the (few) final-state partons
- Using the JetPhox program (S. Catani, M. Fontannaz, J. Ph. Guillet and E. Pilon) with

- \rightarrow Corrections for hadronisation and underlying event needed
- Theoretical uncertainties:
- \rightarrow terms beyond NLO; varying μ_R, μ_F, μ_f by factors 2 and 1/2 (singly or simultaneously)
- \rightarrow PDF-induced uncertainties; estimated using set of PDF eigenvectors
- \rightarrow uncertainty on α_s ; estimated using PDFs in which different values of α_s are assumed
- \rightarrow uncertainty on non-perturbative correction; estimated with different MCs

Photons with the ATLAS detector

Photon reconstruction in the ATLAS LAr Calorimeter

• Layout of the ATLAS electromagnetic calorimeter (Lead-liquid Argon)

- ightarrow barrel section, $|\eta| < 1.475$
- ightarrow two end-cap sections, $1.375 < |\eta| < 3.2$
- \rightarrow three longitudinal layers
- $\begin{array}{l} \mbox{ First layer: high granularity in } \eta \\ \mbox{direction, width 0.003-0.006 (except for} \\ 1.4 < |\eta| < 1.5 \mbox{ and } |\eta| > 2.4) \end{array}$

- Second layer: collects most of the energy, granularity 0.025×0.025 in $\eta \times \phi$

- Third layer: used to correct for leakage
- Cluster of EM cells without matching track:
- \rightarrow "unconverted" photon candidate
- Cluster of EM cells matched to pairs of tracks (from reconstructed conversion vertices in the inner detector) or matched to a single track consistent with originating from a photon conversion
- \rightarrow "converted" photon candidate





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Photon identification in the ATLAS LAr Calorimeter

• To discriminate signal vs background: shape variables from the lateral and longitudinal energy profiles of the shower in the calorimeters; "loose" and "tight" identification criteria.

• <u>"Loose" identification criteria</u>:

 $\rightarrow \text{leakage } R_{had} = E_T^{had} / E_T \text{ (1st layer hadronic calorimeter)} \\ \rightarrow R_\eta = E_{3\times7}^{S2} / E_{7\times7}^{S2}; S2 = \text{second layer of EM calorimeter} \\ \rightarrow \text{RMS width of the shower in } \eta \text{ direction in } S2$

• "Tight" identification criteria:

 \rightarrow the requirements applied in "Loose" are tightened $\rightarrow R_{\phi} = E_{3 \times 3}^{S2} / E_{3 \times 7}^{S2}$ and shower shapes in the first layer (to discriminate single-photon showers from overlapping nearby showers, such as $\pi^0 \rightarrow \gamma\gamma$)

 \rightarrow e.g. asymmetry between the 1st and 2nd maxima in the energy profile along $\eta~(S1)$











• Data-driven measurements of photon identification efficiency for converted and unconverted photons (radiative Z decays, extrapolation from e^{\pm} and matrix method) compared to estimations based on Monte Carlo simulations

Photon isolation in ATLAS

• $E_T^{iso}(R=0.4)$ computed using clusters of calorimeter cells (EM and HAD) in a cone R=0.4, excluding the contribution from the photon

• Subtraction of the leakage of the photon energy outside that region (few %)

• The underlying event and pileup (overlapping pp interactions in the same/neighbouring bunch crossings) contribute to E_T^{iso} ! Subtracted on event-by-event basis using the jet-area method of M. Cacciari et al

• After isolation requirement, residual background still expected



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Background subtraction

- Residual background still expected even after the tight identification and isolation requirements
- A data-driven method necessary to avoid relying on detailed simulations of the background processes
- The two-dimensional sideband method:
- \rightarrow photon identification γ_{ID} vs E_T^{iso} plane



ullet It is assumed that for background events there is no correlation between γ_{ID} and E_T^{iso}

$$rac{N_A^{bkg}}{N_B^{bkg}} = rac{N_C^{bkg}}{N_D^{bkg}} \qquad \Rightarrow R_{bkg} \equiv rac{N_A^{bkg} \cdot N_D^{bkg}}{N_B^{bkg} \cdot N_C^{bkg}} = 1$$

and the effects of the small signal contaminations can be accounted for by using

$$\frac{N_A - N_A^{sig}}{N_B - \epsilon_B N_A^{sig}} = \frac{N_C - \epsilon_C N_A^{sig}}{N_D - \epsilon_D N_A^{sig}} \quad \text{to extract the signal yield } N_A^{sig}$$

the leakage fractions ($\epsilon_K \equiv N_K^{sig}/N_A^{sig}$, K = B, C, D) are estimated using MC samples of signal \Rightarrow purity rises from 60% ($E_T^{\gamma} \sim 25$ GeV) to 100% ($E_T^{\gamma} \sim 300$ GeV)

Inclusive photon production at 8 TeV

Inclusive isolated-photon production in pp collisions at $\sqrt{s}=8~{ m TeV}$



• Significant improvement in experimental uncertainties over the previous measurements

• Good description (in log scale) of the data by NLO QCD calculations using JetPhox





• The uncertainty on the photon energy scale* (about 1% except in the region $1.56 < |\eta^{\gamma}| < 1.81$) is dominant at high E_T^{γ} * (ATLAS Collaboration, Eur. Phys. J. C74 (2014) 3071)

- The uncertainty on the correlation in the background $(\pm 10\%)$ dominates at low E_T^γ , but negligible at high E_T^γ
- ullet The uncertainty on the admixture of direct and fragmentation photons increases at low E_T^γ

Inclusive isolated-photon cross sections vs NLO QCD

ATLAS Coll., STDM-2014-09



• Comparison to NLO QCD calculation using the JetPhox program

- \rightarrow a similar trend is observed at low E_T^{γ} in all $|\eta^{\gamma}|$ regions, the NLO QCD predictions underestimate the data by $\approx 20\%$
- \rightarrow the theoretical uncertainty (12-20%) prevents a more precise test of the SM predictions
- Halving the measured uncertainties compared to previous measurements
 - \Rightarrow useful constraint on proton PDFs once included in a global fit

Inclusive isolated-photon cross sections vs improved NLO QCD



• Comparison to <u>improved</u> NLO QCD calculations using the <u>PeTeR</u> program: resummation of QCD threshold logarithms at NNNLL and large electroweak Sudakov logarithms

- \rightarrow improved description of the data: PeTeR vs JetPhox
- \rightarrow reduction of the theoretical uncertainty: $\sim 20\%$ smaller than in NLO QCD (JetPhox)

Inclusive photon production at 13 TeV

Inclusive isolated-photon production in pp collisions at $\sqrt{s}=13~{ m TeV}$



- First study of inclusive isolated-photon production in pp collisions at 13 TeV using $\mathcal{L} = 6.4 \text{ pb}^{-1}$ and covering the ranges $125 < E_T^{\gamma} < 350 \text{ GeV}$ and $|\eta^{\gamma}| < 2.37$ (except for $1.37 < |\eta^{\gamma}| < 1.56$)
- ullet Background subtraction using a data-driven technique based on the $E_T^{
 m iso}$ distribution
- Good description of the shape of the measured distributions in E_T^{γ} and $|\eta^{\gamma}|$ by the predictions of Sherpa 2.1 in the new regime opened by the LHC in 2015

Summary



• Measurement of isolated photon production in pp collisions at $\sqrt{s} = 8$ TeV in the range $25 < E_T^{\gamma} < 1500$ GeV with improved experimental uncertainties

- \Rightarrow NLO QCD supplemented with resummation+electroweak corrections describe the data
- \Rightarrow theoretical \gg experimental uncertainties, room for improvement with better calculations
- \Rightarrow experimental information to constrain further the proton parton densities
- \bullet First look at inclusive isolated photon production in pp collisions at $\sqrt{s}=13~{\rm TeV}$



Corrections for non-perturbative effects; photon isolation

 The measurements are corrected for detector effects to the "particle" level

 → to isolated photons, where E^{iso}_T
 is calculated using all the final-state φ
 particles and the jet-area method is <u>also</u>
 <u>applied</u>
 This is performed using MC simulations



η

• Corrections for non-perturbative effects (hadronisation and underlying event)

$$C_{NP} = rac{\sigma_{\gamma+\mathrm{X}}(\mathrm{MC, particle} - \mathrm{level}, \mathrm{UE})}{\sigma_{\gamma+\mathrm{X}}(\mathrm{MC, parton} - \mathrm{level, no \, UE})}$$

 \rightarrow Less dependence on the modelling of the final state by having used the jet-area method to subtract the "extra" transverse energy contribution to E_T^{iso}

 \Rightarrow The resulting $C_{\rm NP}$ is found to be consistent with 1 for the measurements presented here

Impact of inclusive isolated photon measurements at LHC on PDFs



Analysis by D. d'Enterria and J. Rojo (Nucl. Phys. B860 (2012) 311)
Study of the impact on the gluon density of existing isolated-photon measurements from a variety of experiments, from √s = 200 GeV up to 7 TeV
→ those at LHC are the more constraining datasets
→ reduction of gluon uncertainty up to 20%
→ localised in the range x ≈ 0.002 to 0.05
⇒ improved predictions for low mass Higgs production in gluon fusion, PDF-induced uncertainty decreased by 20%



The ATLAS detector



• Inner detector (ID): tracking and particle identification in $|\eta| < 2.5$ • Calorimeters: electromagnetic (LAr) \rightarrow barrel $|\eta| < 1.475$, endcap $1.375 < |\eta| < 3.2$, forward $3.1 < |\eta| < 4.9$; hadronic (scintillator/steel, LAr/Cu, LAr/W) \rightarrow barrel $|\eta| < 0.7$ extended barrel $0.8 < |\eta| < 1.7$, endcap $1.5 < |\eta| < 3.2$ and forward $3.1 < |\eta| < 4.9$

Electron and photon calibration in ATLAS

ATLAS Collaboration, Eur. Phys. J. C74 (2014) 3071



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