### Measurement of the ridge correlations in pp and p+Pb collisions with the ATLAS detector at the LHC

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### Motivation

- Quark Gluon Plasma (QGP) produced and probed in heavy ion collisions
- Signatures of QGP: collective expansion (radial flow, elliptic flow, ...), jet quenching, etc.

P



#### Study plasma properties:

- particle azimuthal distribution  $\rightarrow$  using v<sub>n</sub> Fourier coefficients
- two/multi particle correlations

Singles: 
$$\frac{dN}{d\phi} \propto 1 + \sum_{n} 2v_n \cos[n(\phi - \Phi_n)]$$
  
Pairs:  $\frac{dN}{d\Delta\phi} \propto 1 + \sum_{n} 2v_n^a v_n^b \cos[n(\Delta\phi)]$ 

 $v_2$  – elliptic flow

### The ridge in Pb+Pb, p+Pb and pp collisions



- □ Two-particle correlations show long-range correlation structure along  $\Delta \eta$  at  $\Delta \Phi \approx 0 \rightarrow$  the near-side ridge
- What is the origin of ridge structures in small systems? Is it the same mechanism as in Pb+Pb? Is it single-particle azimuthal anisotropy?
- Detailed ridge studies are performed:
  - $\square$  p<sub>T</sub>,  $\eta$ , N<sub>ch</sub>, energy dependence, single-particle v<sub>n</sub>

### ATLAS detector and datasets

ATLAS pp 2.76 TeV data ATLAS p+Pb 5.02 TeV data  $L_{int} = 4 \text{ pb}^{-1}$  $L_{int} = 28 \text{ nb}^{-1}$ p: 4 TeV Pb: 1.57 TeV/N ATLAS pp 13 TeV data  $L_{int} = 14 \text{ nb}^{-1}$ Inner Detector (ID) | η | **< 2.5** Both analyses used charged particles measured in **Forward Calorimeter** Inner Detector (ID) with tracking (FCal) coverage  $|\eta| < 2.5$  and  $p_T > 0.4$  GeV  $3.2 < |\eta| < 4.9$ 

Forward Calorimeter (FCal) is used to define "event activity" of p+Pb collisions

# Event activity in **p+Pb** and **pp** collisions

N<sup>rec</sup>



- **Event activity** in p+Pb collisions is defined in two ways:
  - E<sub>T</sub><sup>Pb</sup>: Total transverse energy in the FCal on Pb-going side
  - N<sub>ch</sub><sup>rec</sup>: number of charged particles with p<sub>T</sub> > 0.4 GeV in ID
    - For pp collisions only N<sub>ch</sub><sup>rec</sup> intervals are used to characterize event activity
    - High multiplicity events in pp:
      - $\begin{array}{l} \textbf{2.76 TeV} \rightarrow \ \textbf{N}_{ch}^{rec} > 50 \\ \textbf{13 TeV} \rightarrow \ \textbf{N}_{ch}^{rec} > 60 \end{array}$

# Two-particle correlation (2PC) function



- S is constructed using all pairs of charged particles in the same event
- B obtained from "mixed events", measures distribution of uncorrelated pairs
- Both S and B are corrected for tracking inefficiencies detector acceptance effects largely cancel in the S/B ratio

### 2PC analysis – p+Pb collisions



### **2PC** analysis – **p+Pb** collisions



12/04/16

# **2PC** analysis – **pp** collisions

![](_page_8_Figure_1.jpeg)

12/04/16

# □ Y( $\Delta \phi$ ) is used to measure the strength of the long-range component N<sup>a</sup> - total number of trigger particles

$$Y(\Delta\phi) = \left(\frac{\int B(\Delta\phi)d\Delta\phi}{N^a \int d\Delta\phi}\right)C(\Delta\phi) - b_{ZYAM}$$

M b<sub>ZYAM</sub> – pedestal arising from uncorrelated pairs

Compared central and peripheral PTY for  $2 < |\Delta \eta| < 5$ 

Peripheral events show 'away-side' peak
→ dijets contribution

Central events feature the ridge (nearside peak) and excess on away-side

Excess on near-side and away-side independent on event activity, E<sub>T</sub><sup>Pb</sup>

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$$Y(\Delta\phi) = \left(\frac{\int B(\Delta\phi)d\Delta\phi}{N^a \int d\Delta\phi}\right)C(\Delta\phi) - b_{ZYAM}$$

![](_page_10_Figure_3.jpeg)

N<sup>a</sup> - total number of trigger particles b<sub>ZYAM</sub> – pedestal arising from uncorrelated pairs

Compared central and peripheral PTY for 2<|  $\Delta \eta$  | < 5

Peripheral events show 'away-side' peak
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$$Y(\Delta\phi) = \left(\frac{\int B(\Delta\phi)d\Delta\phi}{N^a \int d\Delta\phi}\right)C(\Delta\phi) - b_{ZYAM}$$

![](_page_11_Figure_3.jpeg)

N<sup>a</sup> - total number of trigger particles b<sub>ZYAM</sub> – pedestal arising from uncorrelated pairs

Compared central and peripheral PTY for 2<  $| \Delta \eta | < 5$ 

Peripheral events show 'away-side' peak  $\rightarrow$  dijets contribution

Central events feature the ridge (nearside peak) and excess on away-side

Excess on near-side and away-side independent on event activity,  $E_{T}^{Pb}$ 

$$Y^{sub}(\Delta\phi) = Y(\Delta\phi) - \alpha Y_{peri}^{corr}(\Delta\phi) \propto \left[1 + \sum_{n} 2\psi_{n,n} \cos(n\Delta\phi)\right]$$

![](_page_12_Figure_1.jpeg)

$$Y(\Delta\phi) = \left(\frac{\int B(\Delta\phi)d\Delta\phi}{N^a \int d\Delta\phi}\right)C(\Delta\phi) - b_{ZYAM}$$

N<sup>a</sup> - total number of trigger particles

ZYAM assumes that number of correlated pairs is zero at minimum

□ The minimum of  $Y(\Delta \phi)$  may be modulated due to long-range correlations

Only near-side long range correlation can be extracted

- $\rightarrow$  away-side Y( $\Delta \phi$ ) dominated by dijets
- □ No information about full  $\Delta \phi$  dependence → new template fitting method

# Template fitting method - **pp**

![](_page_13_Figure_1.jpeg)

### Single-particle anisotropies v<sub>2</sub> in **pp & p+Pb**

□ If ridge in pp and p+Pb collisions also results from modulation of the single particle  $\phi$ -distribution, measured v<sub>2.2</sub> should factorize:

$$v_{2,2}(p_T^a, p_T^b) = v_2(p_T^a)v_2(p_T^b) \qquad \qquad v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^a, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^b, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^a) = v_{2,2}(p_T^b, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^b) = v_{2,2}(p_T^b, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^b) = v_{2,2}(p_T^b, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^b) = v_{2,2}(p_T^b, p_T^b) / \sqrt{v_{2,2}(p_T^b, p_T^b)} \\ \sim v_2(p_T^b) = v_$$

v<sub>2,2</sub> depends on both, a and b, particles but v<sub>2</sub> (p<sub>T</sub><sup>a</sup>) should be independent of reference p<sub>T</sub><sup>b</sup>

□ Factorization works fine for both, pp & p+Pb, systems → global anisotropy

### $v_n$ ( $p_T$ ) and event activity dependence in **p+Pb**

![](_page_15_Figure_1.jpeg)

# $v_2$ integrated over $p_T$ in **pp** and **p+Pb**

![](_page_16_Figure_1.jpeg)

- In pp system v<sub>2</sub> shows very weak dependence on energy and track multiplicity
- In p+Pb multiplicity dependence is stronger
  - $\rightarrow$  v<sub>2</sub> grows with N<sub>ch</sub><sup>rec</sup>

### $N_{ch}^{rec}$ and $p_T$ dependence of $v_2$ in **pp** and **p+Pb**

![](_page_17_Figure_1.jpeg)

 $\Box$  v<sub>2</sub> in pp system behaves similar as v<sub>2</sub> in Pb+Pb and p+Pb:

- □ Increase with p<sub>T</sub> up to ~3 GeV
- Decreases at higher p<sub>T</sub>
- v<sub>2</sub> (p<sub>T</sub>) shows very weak dependence on energy and track multiplicity (in pp)

### Summary

- Ridge was observed in high-multiplicity pp collisions at 2.76 and 13 TeV as well as in p+Pb collisions at 5.02 TeV
- □ ZYAM and template fitting methods were applied to extract long range correlations  $v_{2,2}$  modulated by  $cos(2\Delta \phi)$  in *p*+*Pb* and *pp* systems respectively
- $\Box$  v<sub>2.2</sub> exhibits factorization into single particle v<sub>2</sub>
- □ Elliptic flow, v<sub>2</sub>, was measured for both systems → higher harmonics measured in p+Pb
- □ The shape of v<sub>2</sub>(p<sub>T</sub>) distribution in pp and p+Pb is similar to the shape of v<sub>2</sub>(p<sub>T</sub>) distribution in Pb+Pb

→ This suggests that the ridge in pp, p+Pb and Pb+Pb collisions arises from similar dynamics ?

### **Backup slides**

# High Multiplicity Triggers in **pp** collisions

- Dedicated triggers (HMT) implemented using ATLAS L1 & HLT systems to increase high multiplicity events
- The structures in the distributions result from the different HMT trigger thresholds

□ High multiplicity events: 2.76 TeV  $\rightarrow$  N<sub>ch</sub><sup>rec</sup> > 30 13 TeV  $\rightarrow$  N<sub>ch</sub><sup>rec</sup> > 60

![](_page_20_Figure_4.jpeg)

### High multiplicity triggers in **p+Pb**

![](_page_21_Figure_1.jpeg)

- Six pairs of HMT thresholds
- Enhancement from individual HMT (top)
  - Reweighted by MB +HMT distribution (bottom)
  - Reweighting procedure comprise probability for a given event to be accepted by MB or HMT, prescale and trigger effiviency

# Event activity in **p+Pb** collisions

![](_page_22_Figure_1.jpeg)

Event activity in p+Pb collisions is defined in two ways:

 $\Box$  N<sub>ch</sub><sup>rec</sup>: number of charged particles with p<sub>T</sub> > 0.4 GeV in ID

E<sub>T<sup>Pb</sup>: Total transverse energy on the FCal on Pb-going side</sub>

 $\Box$  E<sub>T</sub><sup>Pb</sup> distribution is divided into 12 intervals

For pp collisions only N<sub>ch</sub><sup>rec</sup> intervals are used to characterize event activity

• Jet & dijets peaks are estimated from the peripheral collisions and subtracted from Y(  $\Delta \phi$  )

![](_page_23_Figure_2.jpeg)

### v<sub>n</sub>(p<sub>T</sub>) scaling between the **p+Pb** and **Pb+Pb**

![](_page_24_Figure_1.jpeg)

Compare  $v_n(p_T)_{p+Pb}$ with  $v_n(p_T/1.25)_{Pb+Pb}$ (Teaney et al. arXiv: 1312.6770[nucl-th])

- The shape after scaling is similar in both systems
- v<sub>2</sub> values differ only by a scale factor between the two systems
- This suggest similar origins of long range correlations

### **2PC** funcition and $Y(\Delta \phi)$ in p+Pb

![](_page_25_Figure_1.jpeg)