

Charged-Current Deep Inelastic Scattering at Third Order

Joshua Davies

*Department of Mathematical Sciences
University of Liverpool*

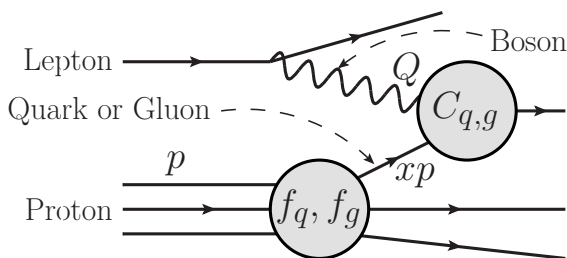


Collaborators: Andreas Vogt (UoL),
Sven Moch (University of Hamburg), Jos Vermaseren (nikhef)

Deep Inelastic Scattering 2016, DESY Hamburg

INTRODUCTION

Deep Inelastic Scattering – lepton scatters from a proton:



Boson: γ, H, Z^0 (Neutral Current) or W^\pm (Charged Current)

Cross-section $\sigma \sim \sum_a F_a = \sum_a [C_{a,q} \otimes f_q + C_{a,g} \otimes f_g]$.

$C_{a,j}$ – “Coefficient Function”
 f_j – “Parton Distribution Function”
 F_a – “Structure Function”

INTRODUCTION

Anomalous dimensions/Splitting functions – scale evolution of PDFs:

$$\frac{d}{d \ln Q^2} f_j = \sum_k \gamma_{jk} f_k$$

Series expansions in α_s :

$$C_{a,j} = c_{a,j}^{(0)} + \alpha_s c_{a,j}^{(1)} + \alpha_s^2 c_{a,j}^{(2)} + \alpha_s^3 c_{a,j}^{(3)} + \dots$$

$$\gamma_{jk} = \alpha_s \gamma_{jk}^{(0)} + \alpha_s^2 \gamma_{jk}^{(1)} + \alpha_s^3 \gamma_{jk}^{(2)} + \dots$$

KNOWN SO FAR

Natural to consider *combinations* $F_a^{W^++W^-}$ and $F_a^{W^+-W^-}$.

Known, to third-order (massless DIS):

- ▶ γ probe: $C_{2,q}, C_{2,g}, C_{L,q}, C_{L,g}$ [hep-ph/0504.242]
- ▶ H probe: $C_{\phi,q}, C_{\phi,g}$ [hep-ph/0912.0369]
- ▶ W^\pm probe (+): $C_{2,q}^+, C_{2,g}^+, C_{L,q}^+, C_{L,g}^+, C_{3,ns}^+$ [hep-ph/0812.4168]

Unknown/Incomplete:

- ▶ Z^0 probe: $C_{2,q}, C_{2,g}, C_{L,q}, C_{L,g}, C_{3,q}$
- ▶ W^\pm probe (-): $C_{2,ns}^-, C_{L,ns}^-, C_{3,ns}^-$

CHARGED-CURRENT DIS

Previous work at α_s^3 . Moch, Rogal:

[[hep-ph/0704.1740](https://arxiv.org/abs/hep-ph/0704.1740)]

- ▶ $N = 3, 5, \dots, 11$ of $C_{2,ns}^-$, $C_{L,ns}^-$.
- ▶ $N = 2, 4, \dots, 10$ of $C_{3,ns}^-$.

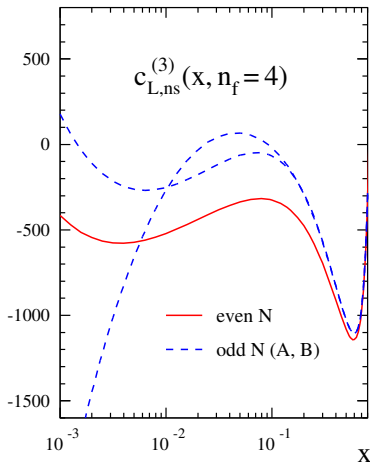
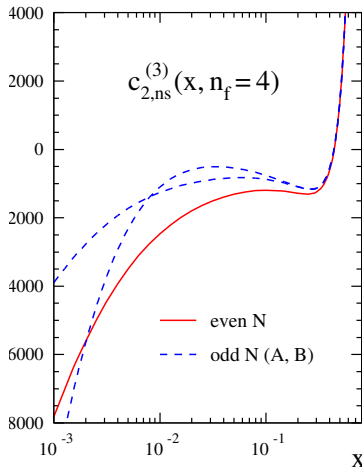
Used to produce approximations to full analytic result.

(We compute in Mellin space. N th Mellin Moment:

$$F_a(N, Q^2) = \int_0^1 x^{N-1} F_a(x, Q^2) dx.$$

For $F_a^{W^+ - W^-}$: F_2, F_L live on odd moments, F_3 on even moments.)

CHARGED-CURRENT DIS



even N : $W^+ + W^-$, odd N : $W^+ - W^-$ [[hep-ph/0704.1740](https://arxiv.org/abs/hep-ph/0704.1740)]

Approximations unconstrained below $x \approx 10^{-2}$.

COMPUTATION

Diagram preparation:

- ▶ QGRAF + FORM: produce diagram database.

Mellin moment computation:

- ▶ MINCER (FORM): verify Moch, Rogal's moments and compute more.

Analytic computation:

- ▶ All- N FORM code: used in MVV's calculation, directly produces analytic results in N .

THE OPTICAL THEOREM

Relates structure functions to the *forward scattering amplitude*:

$$\left| \begin{array}{c} \text{wavy line} \\ \bullet \\ \nearrow \text{line} \rightarrow \text{line} \end{array} \right|^2 \sim \text{Im} \begin{array}{c} \text{wavy line} \\ \bullet \\ \nearrow \text{line} \rightarrow \text{line} \bullet \\ \searrow \text{line} \\ \text{wavy line} \end{array}$$

Number of forward scattering “diagrams”:

- ▶ F_2, F_L : 209
- ▶ F_3 : 197

Big improvement on diagram database of Moch, Rogal (1300)
→ faster moment computation.

Check: gauge parameter in propagators, ensure cancellation.

RENORMALIZATION

Forward scattering amplitudes T_a will contain poles in Dim.-Reg. parameter ε (where $D = 4 - 2\varepsilon$), due to loop integrals (UV) and collinear momenta.

F_a are physically measurable – renormalize!

- ▶ UV poles: coupling constant renormalization
- ▶ Collinear poles: mass factorization

We can write:

$$F_a = T_a \left(\varepsilon^{-1}, \varepsilon \right) \hat{f} = C_a \left(\varepsilon \right) Z \left(\varepsilon^{-1} \right) \hat{f} = C_a \left(\varepsilon \right) f.$$

Absorb poles into a redefinition of the non-perturbative parton distribution function: $f = Z \cdot \hat{f}$.

RENORMALIZATION

Z is a function of the (already known) anomalous dimensions γ_{ns} since

$$\frac{dZ}{d \ln Q^2} = -\gamma_{ns} Z$$

gives

$$Z = 1 + \alpha_s \frac{1}{\epsilon} \gamma_{ns}^{(0)} + \alpha_s^2 \left[\frac{1}{2\epsilon^2} \left(\gamma_{ns}^{(0)} - \beta_0 \right) \gamma_{ns}^{(0)} + \frac{1}{2\epsilon} \gamma_{ns}^{(1)} \right] + \mathcal{O}(\alpha_s^3).$$

Setting $\gamma_{ns}^{(i)}$ order-by-order removes remaining poles from forward scattering amplitude, fixes coefficient function.

$$T_a(\epsilon^{-1}, \epsilon) = C_a(\epsilon) Z(\epsilon^{-1})$$

Reproducing $\gamma_{ns}^{(i)}$ is a strong consistency check of the calculation.

MOMENT CALCULATION

Verify that the new diagram database yields the same moments. Compute with MINCER.

Extend from first 5 to first 14 moments: (non trivial!)

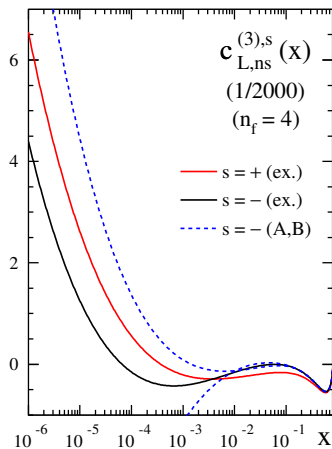
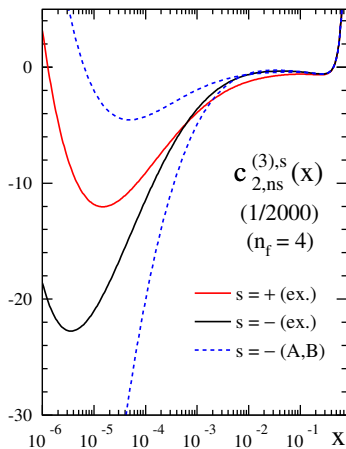
- ▶ $N = 3, 5, \dots, 29$ of $C_{2,ns}^-, C_{L,ns}^-$.
- ▶ $N = 2, 4, \dots, 28$ of $C_{3,ns}^-$.

Agree with Moch, Rogal's moments ✓

(Possibility to reconstruct analytic result from moments in the style of MVV's recent calculation of polarized anomalous dimension: [\[hep-ph/1409.5131\]](https://arxiv.org/abs/hep-ph/1409.5131)
→ unsuccessful).

Analytic expressions computed with all- N code reproduce these moments ✓

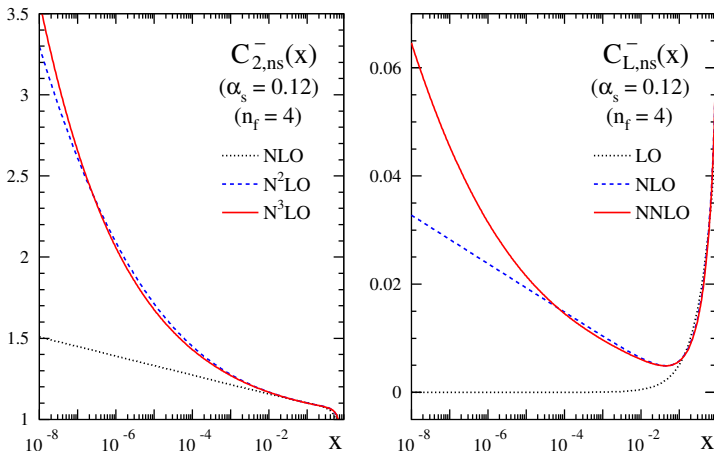
COMPARISON WITH APPROXIMATIONS



$s = +: W^+ + W^-$, $s = -: W^+ - W^-$.

Approximation error $> 5\%$ below $x \approx 0.1$.

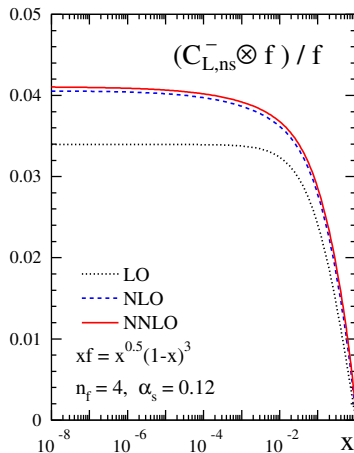
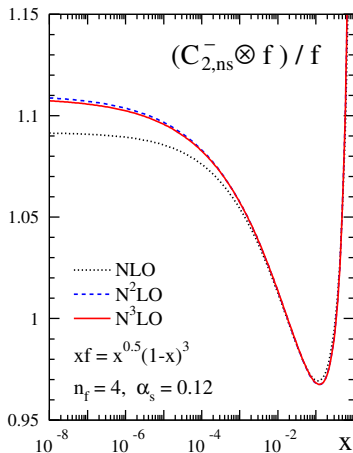
PERTURBATIVE STABILITY



N^3 LO correction to $C_{2,ns}^-$: $< 3\%$ for $x \in (6 \times 10^{-8}, 0.75)$

N^2 LO correction to $C_{L,ns}^-$: $< 5\%$ for $x \in (0.007, 0.9)$

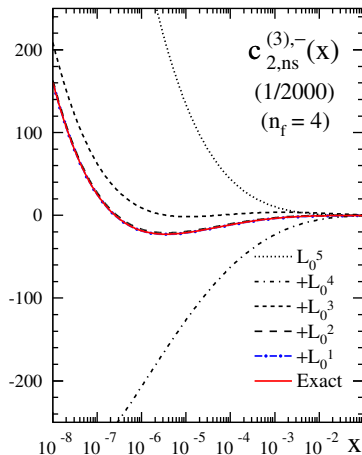
PERTURBATIVE STABILITY



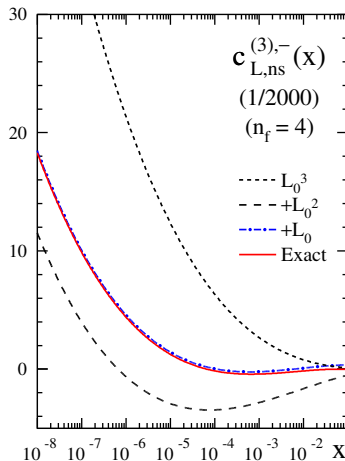
Converge better:

N^3 LO correction to $C_{2,ns}^-$: $< 1\%$ for $x \in (10^{-8}, 0.8)$

N^2 LO correction to $C_{L,ns}^-$: $< 3\%$ for $x \in (10^{-8}, 0.12)$

SMALL- x BEHAVIOUR

Error < 3% for $x < 3 \times 10^{-4}$.



Error < 3% for $x < 9 \times 10^{-10}$.

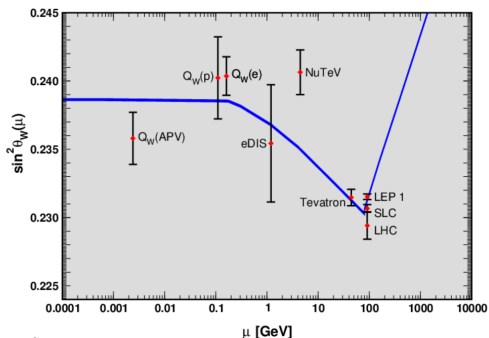
Small- x logarithms ($L_0 = \ln(x)$) are a poor approximation!

PASCHOS-WOLFENSTEIN RELATION

Ratio between NC and CC cross-sections:

$$R^- = \frac{\sigma(\nu P \rightarrow \nu X) - \sigma(\bar{\nu} P \rightarrow \bar{\nu} X)}{\sigma(\nu P \rightarrow \mu^- X) - \sigma(\bar{\nu} P \rightarrow \mu^+ X)}$$

Determines $\sin^2 \theta_W$. NuTeV anomaly?



PASCHOS-WOLFENSTEIN RELATION

Can now provide exact α_s^3 correction.

($N = 2$ of $\delta C_{2,ns}^{(3)}$, $\delta C_{L,ns}^{(3)}$ where $\delta C_{2,ns}^{(3)} = C_{2,ns}^{(3)W^++W^-} - C_{2,ns}^{(3)W^+-W^-}$)

Exact results correct Moch, Rogal's approximations to $\delta C_{2,ns}$ and $\delta C_{L,ns}$ by just 0.5% and 0.05%.

→ Don't expect much change in PW relation.

$$R^- = \left(\frac{1}{2} - \sin^2 \theta_W \right) + \frac{u^- - d^- + c^- - s^-}{u^- + d^-} \left\{ 1 - \frac{7}{3} \sin^2 \theta_W \right. \\ \left. + \left(\frac{1}{2} - \sin^2 \theta_W \right) \frac{8\alpha_s}{9\pi} \left[1 + 1.689\alpha_s + 3.661\alpha_s^2 \right] \right\}$$

Including the **third-order** result:

16% correction of [], 1% correction of { } over α_s^2 contribution.

CONCLUSION

Charged-Current DIS in $F_i^{W^+ - W^-}$ combination computed analytically at third order.

Exact expressions differ greatly from previous approximations at small x .

Coefficient functions largely stable in α_s expansion over physically relevant x range.

Paper to come soon:

- ▶ Full discussion of all results (inc. polarized $\Delta C_{g1}^{(3)}$)
- ▶ 0.1% accurate parametrizations for numerics
- ▶ FORM files for analytic results
- ▶ Fortran and C++ (GiNaC) files for numerics