



Pion structure from leading neutron electroproduction

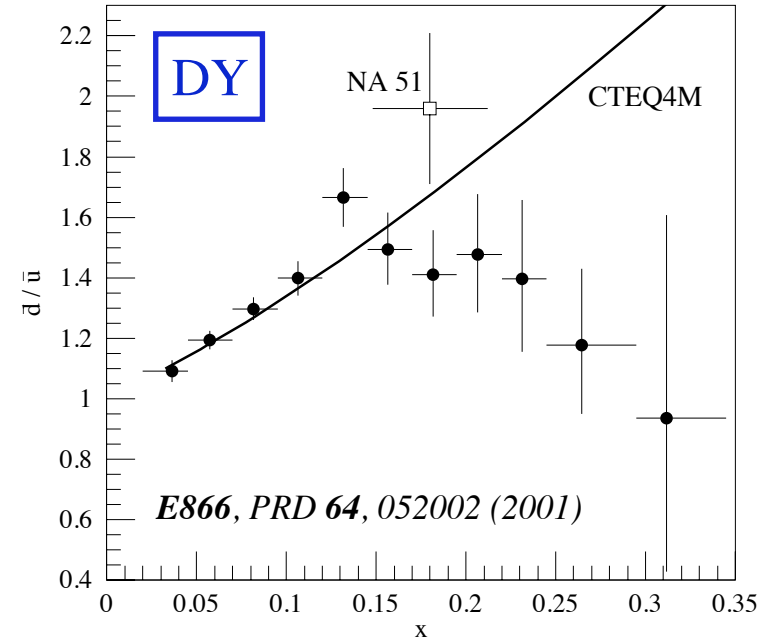
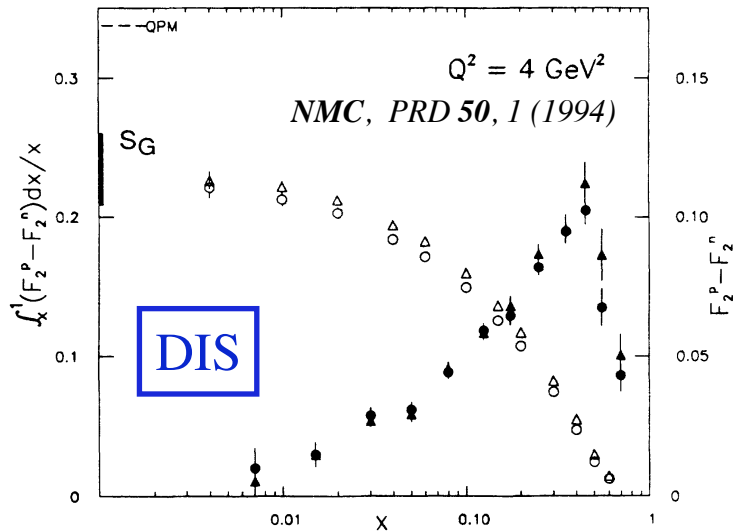
Wally Melnitchouk



*with Chueng Ji (NCSU), Josh McKinney (UNC),
Nobuo Sato (JLab), Tony Thomas (Adelaide)*

SU(2) flavor asymmetry

- One of the seminal discoveries of last 25 years has been the SU(2) flavor asymmetry in the proton sea, $\bar{d} \neq \bar{u}$



$$\int_0^1 \frac{dx}{x} (F_2^p - F_2^n) = \frac{1}{3} - \frac{2}{3} \int_0^1 dx (\bar{d} - \bar{u})$$

$$= 0.235(26)$$

violation of Gottfried sum rule

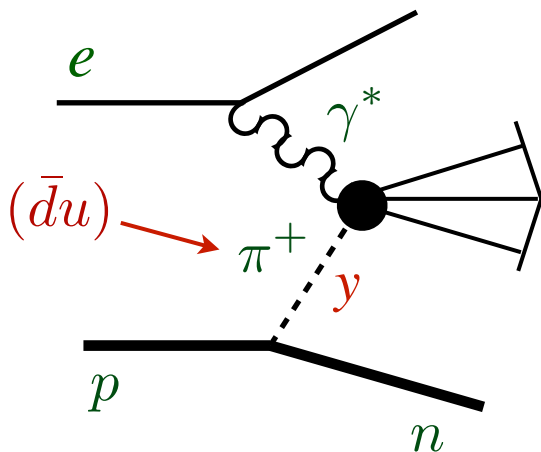
$$\frac{\sigma^{pd}}{2\sigma^{pp}} \approx \frac{1}{2} \left(1 + \frac{\bar{d}(x_t)}{\bar{u}(x_t)} \right) \quad x_b \gg x_t$$

$$\rightarrow \int_0^1 dx (\bar{d} - \bar{u}) = 0.118 \pm 0.012$$

SU(2) flavor asymmetry

- One of the seminal discoveries of last 25 years has been the SU(2) flavor asymmetry in the proton sea, $\bar{d} \neq \bar{u}$
 - paradigm shift – nucleon *not* simply 3 valence quarks + homogenous $\bar{q}q$ sea!
 - vital role played by nonperturbative dynamics, *e.g.* chiral symmetry breaking & nucleon's pion cloud
 - asymmetry actually predicted a decade earlier from “Sullivan” process

A.W. Thomas, PLB 126, 97 (1983)



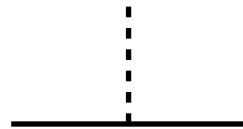
$$(\bar{d} - \bar{u})(x) = \int_x^1 \frac{dy}{y} f_{\pi^+ n}(y) \bar{q}_v^\pi(x/y)$$

pion momentum distribution in nucleon,
or $p \rightarrow \pi^+ n$ “splitting function”

Chiral effective theory

- Splitting function can be computed in chiral effective theory of QCD (*e.g.* chiral perturbation theory)
- At lowest order, effective (low-energy) πN Lagrangian

$$\mathcal{L}_{\pi N} = \frac{g_A}{2f_\pi} \bar{\psi}_N \gamma^\mu \gamma_5 \vec{\tau} \cdot \partial_\mu \vec{\pi} \psi_N - \frac{1}{(2f_\pi)^2} \bar{\psi}_N \gamma^\mu \vec{\tau} \cdot (\vec{\pi} \times \partial_\mu \vec{\pi}) \psi_N$$



$$g_A = 1.267$$

$$f_\pi = 93 \text{ MeV}$$

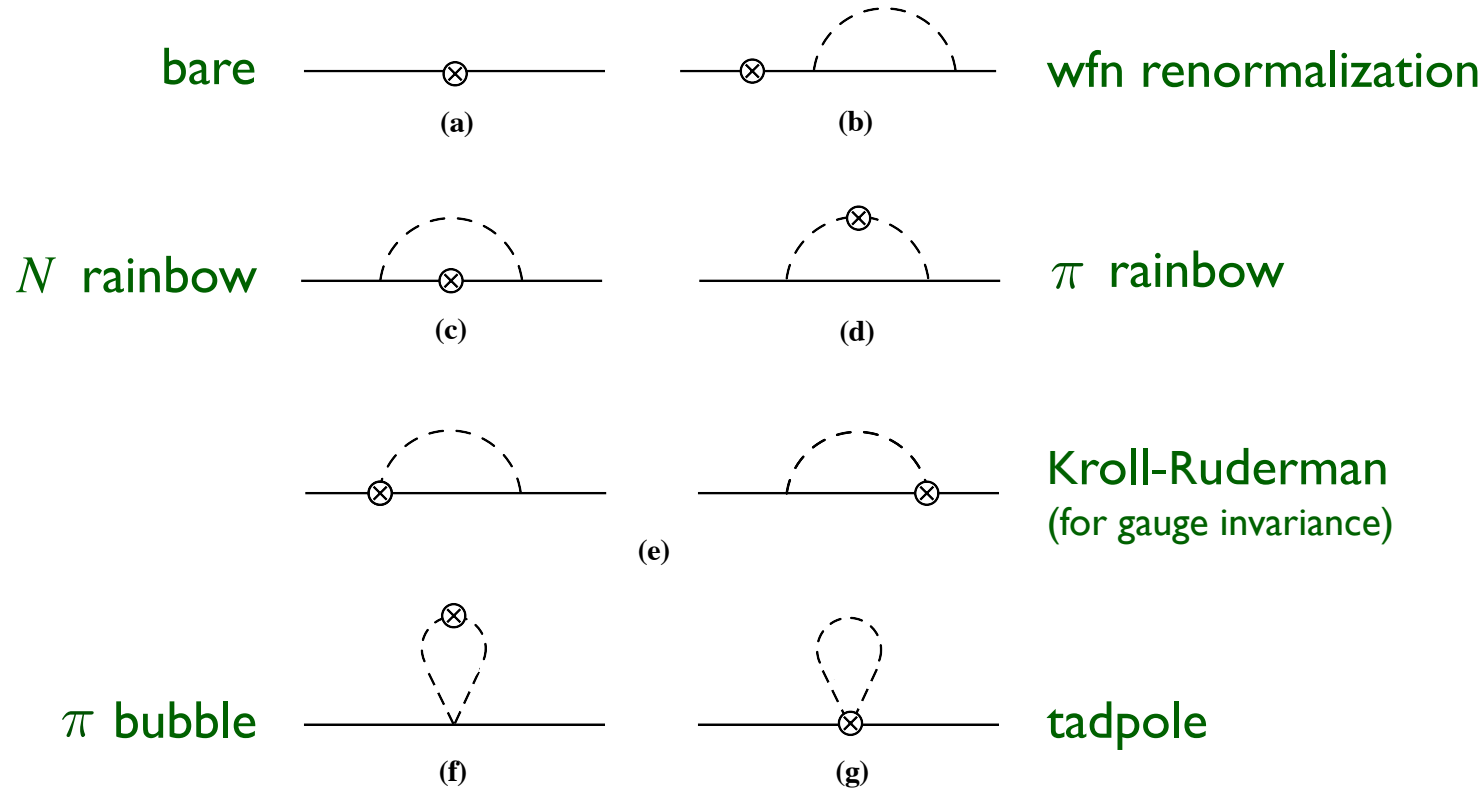
$$= -g_{\pi NN} \bar{\psi}_N i\gamma_5 \vec{\tau} \cdot \vec{\pi} \psi_N + \sigma NN \text{ term} + \text{higher order}$$

Weinberg, *PRL* **18**, 88 (1967)

→ pseudoscalar interaction often used for simplicity – results generally different from pseudovector theory

Chiral effective theory

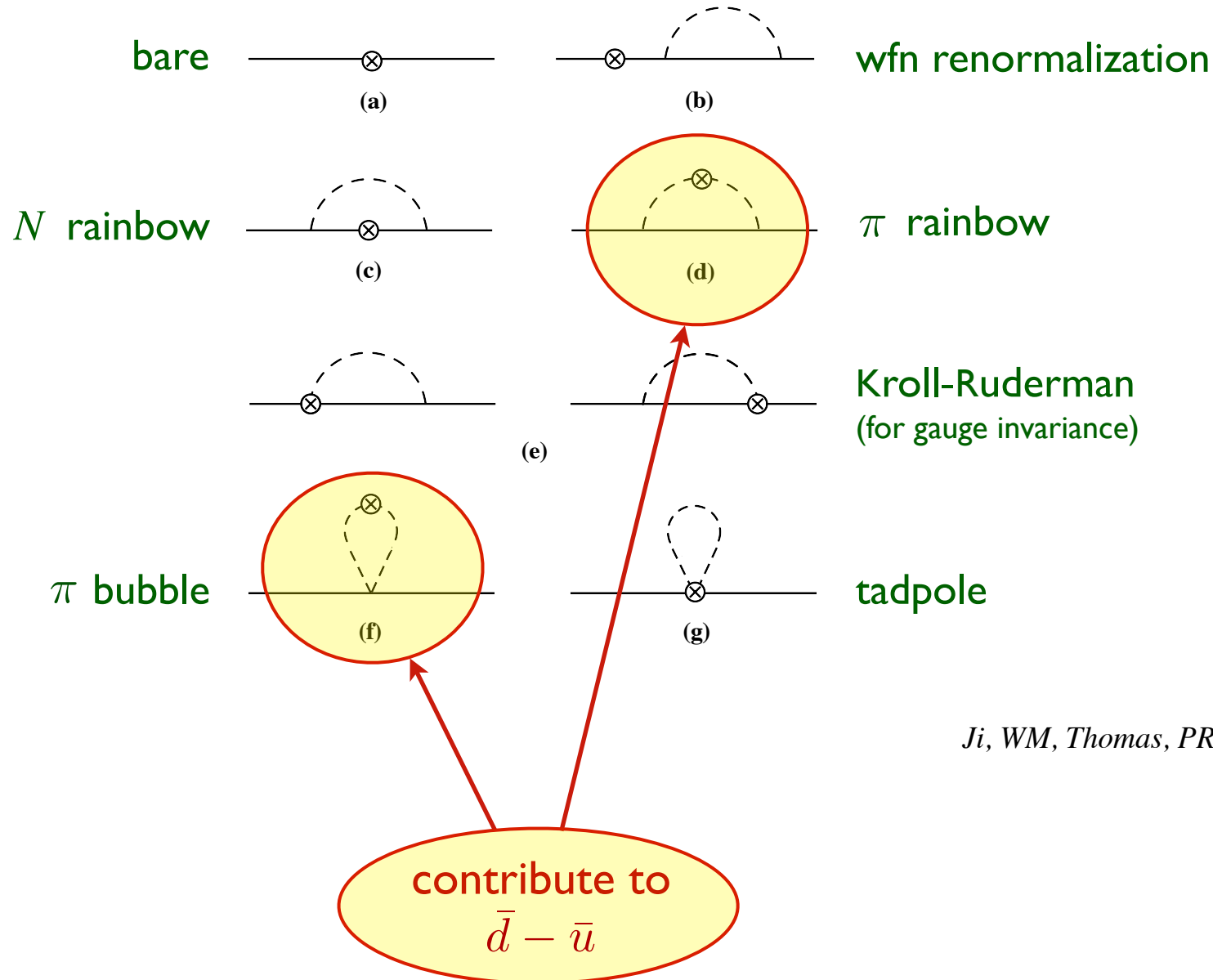
■ Coupling of e.m. current to nucleon dressed by pions



Ji, WM, Thomas, PRD 88, 076005 (2013)

Chiral effective theory

■ Coupling of e.m. current to nucleon dressed by pions



Ji, WM, Thomas, PRD 88, 076005 (2013)

Pion splitting functions

- Splitting function for pion rainbow diagram has on-shell and δ -function contributions

$$f_{\pi}(y) = f^{(\text{on})}(y) + f^{(\delta)}(y)$$

$$f^{(\text{on})}(y) = \frac{g_A^2 M^2}{(4\pi f_{\pi})^2} \int dk_{\perp}^2 \frac{y(k_{\perp}^2 + y^2 M^2)}{[k_{\perp}^2 + y^2 M^2 + (1-y)m_{\pi}^2]^2} \mathcal{F}^2$$

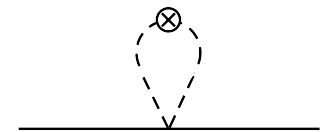
equivalent in
PV & PS theories

$$f^{(\delta)}(y) = \frac{g_A^2}{4(4\pi f_{\pi})^2} \int dk_{\perp}^2 \log \left(\frac{k_{\perp}^2 + m_{\pi}^2}{\mu^2} \right) \delta(y) \mathcal{F}^2$$

singular $y = 0$ term
only in PV theory

- Bubble diagram contributes only at $y = 0$ (hence $x = 0$)

$$f^{(\text{bub})}(y) = \frac{8}{g_A^2} f^{(\delta)}(y)$$



Salamu, Ji, WM, Wang
PRL **114**, 122001 (2015)

Pion splitting functions

■ Infrared behavior is model independent

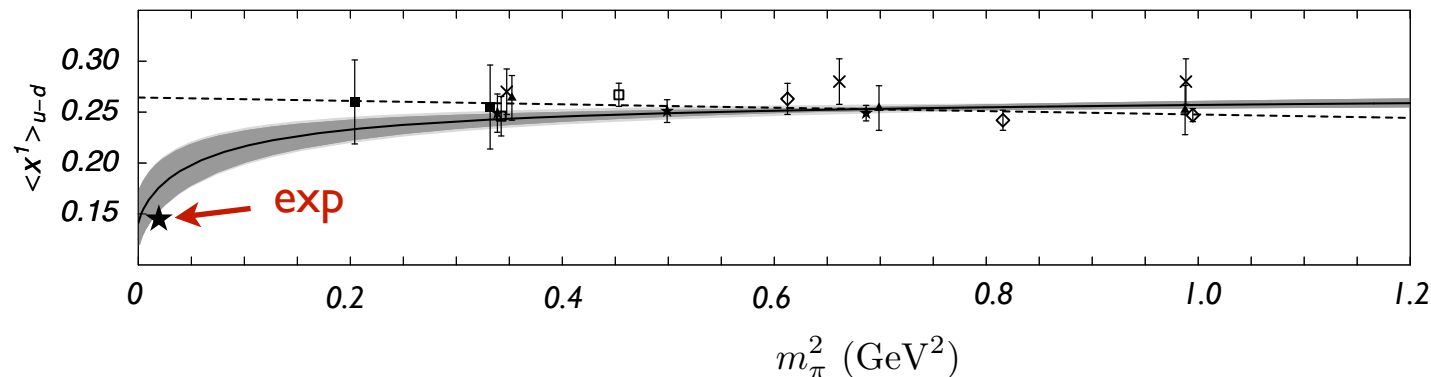
→ leading nonanalytic (LNA) structure of moments

$$\int_0^1 dx (\bar{d} - \bar{u}) = \frac{(3g_A^2 - 1)}{(4\pi f_\pi)^2} m_\pi^2 \log(m_\pi^2/\mu^2) + \text{analytic in } m_\pi^2$$

Thomas, WM, Steffens, PRL 85, 2892 (2000)

→ can only be generated by pion loops – nonzero π cloud contribution predicted by QCD!

→ vital *e.g.* for chiral extrapolation of lattice data



*Detmold et al.
PRL 87, 172001 (2001)*

Pion splitting functions

- Ultraviolet-divergent integrals for point-like particles
 - Finite size of nucleon provides natural scale to regularize integrals, but does not prescribe form of regularization
- freedom in choosing regularization prescription

$$\mathcal{F} = \Theta(\Lambda^2 - k_{\perp}^2)$$

k_{\perp} cutoff

$$\mathcal{F} = \left(\frac{\Lambda^2 - m_{\pi}^2}{\Lambda^2 - t} \right)$$

monopole in $t \equiv k^2 = -\frac{k_{\perp}^2 + y^2 M^2}{1 - y}$

$$\mathcal{F} = \exp \left[(t - m_{\pi}^2) / \Lambda^2 \right]$$

exponential in t

$$\mathcal{F} = \exp \left[(M^2 - s) / \Lambda^2 \right]$$

exponential in $s = \frac{k_{\perp}^2 + m_{\pi}^2}{y} + \frac{k_{\perp}^2 + M^2}{1 - y}$

$$\mathcal{F} = \left[1 - \frac{(t - m_{\pi}^2)^2}{(t - \Lambda^2)^2} \right]^{1/2}$$

Pauli-Villars

$$\mathcal{F} = y^{-\alpha_{\pi}(t)} \exp \left[(t - m_{\pi}^2) / \Lambda^2 \right]$$

Regge

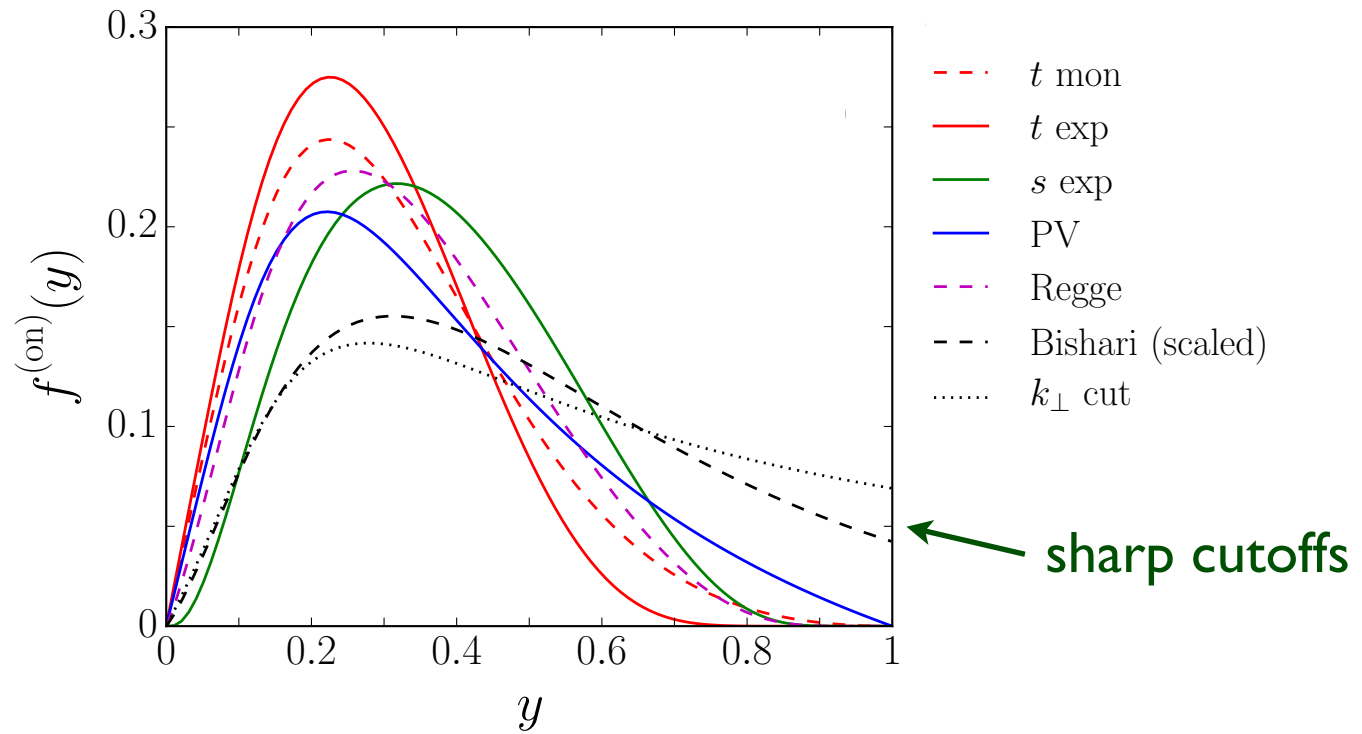
$$\mathcal{F} = y^{-\alpha_{\pi}(t)}$$

Bishari

Pion splitting functions

- Detailed shape of splitting function depends on regularization, but common general features

e.g. on-shell
function

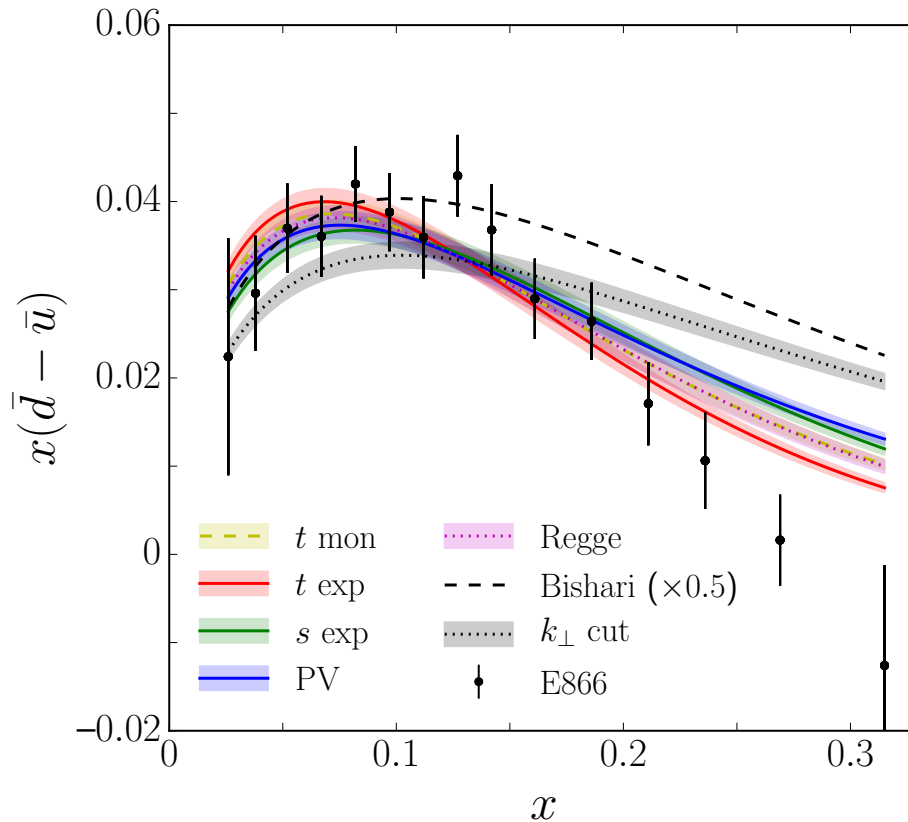


- At $x > 0$, only on-shell part contributes

$$(\bar{d} - \bar{u})(x) = [2f^{(\text{on})} \otimes \bar{q}_v^{\pi}](x)$$

Flavor asymmetry

- E866 $\bar{d} - \bar{u}$ data can be fitted with range of regulators



average pion “multiplicity”

$$\langle n \rangle_{\pi N} = 3 \int_0^1 dy f_N^{(\text{on})}(y) \\ \sim 0.25 - 0.3$$

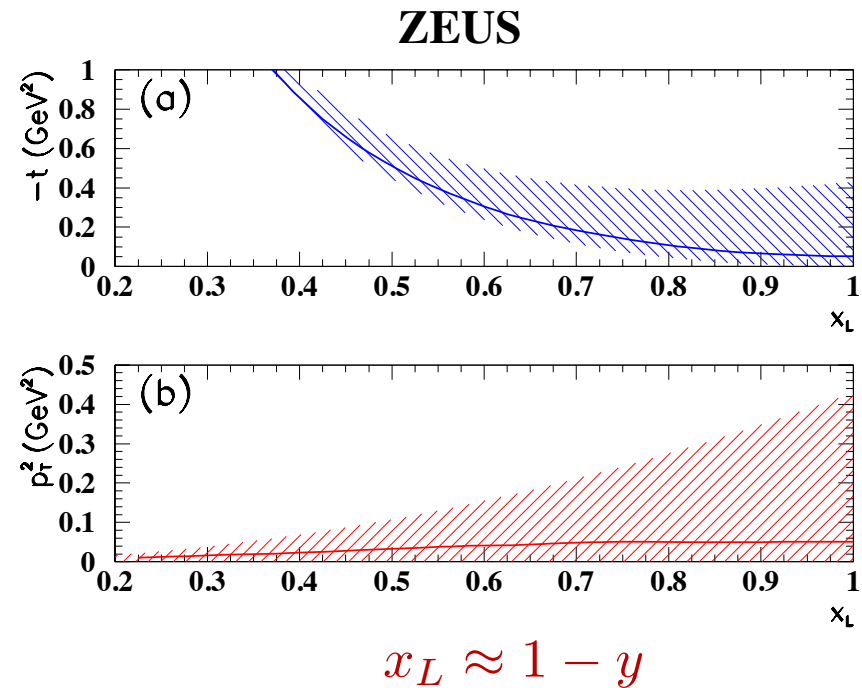
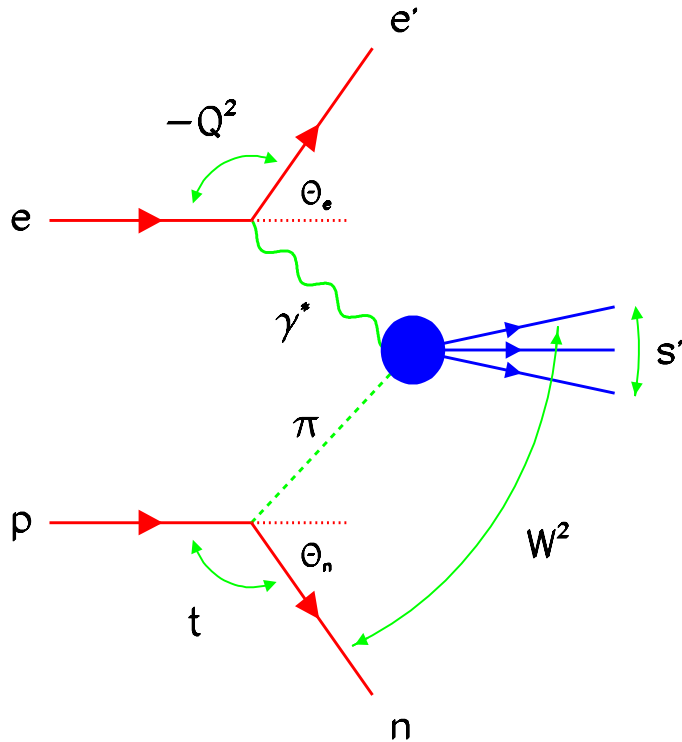
- with exception of k_{\perp} cutoff and Bishari models, all others give reasonable fits, $\chi^2 \lesssim 1.5$
- large- x asymmetry to be probed by FNAL *SeaQuest* expt.

Flavor asymmetry

- E866 $\bar{d} - \bar{u}$ data can be fitted with range of regulators
- Is pion cloud the only explanation for the asymmetry?
 - are there other data that can discriminate between different mechanisms?
 - semi-inclusive production of “leading neutrons” (LN) at HERA!

Leading neutron production at HERA

- ZEUS & H1 collaborations measured spectra of neutrons produced at very forward angles, $\theta_n < 0.8$ mrad



- can data be described within same framework as E866 flavor asymmetry?
- simultaneous fit never previously been performed

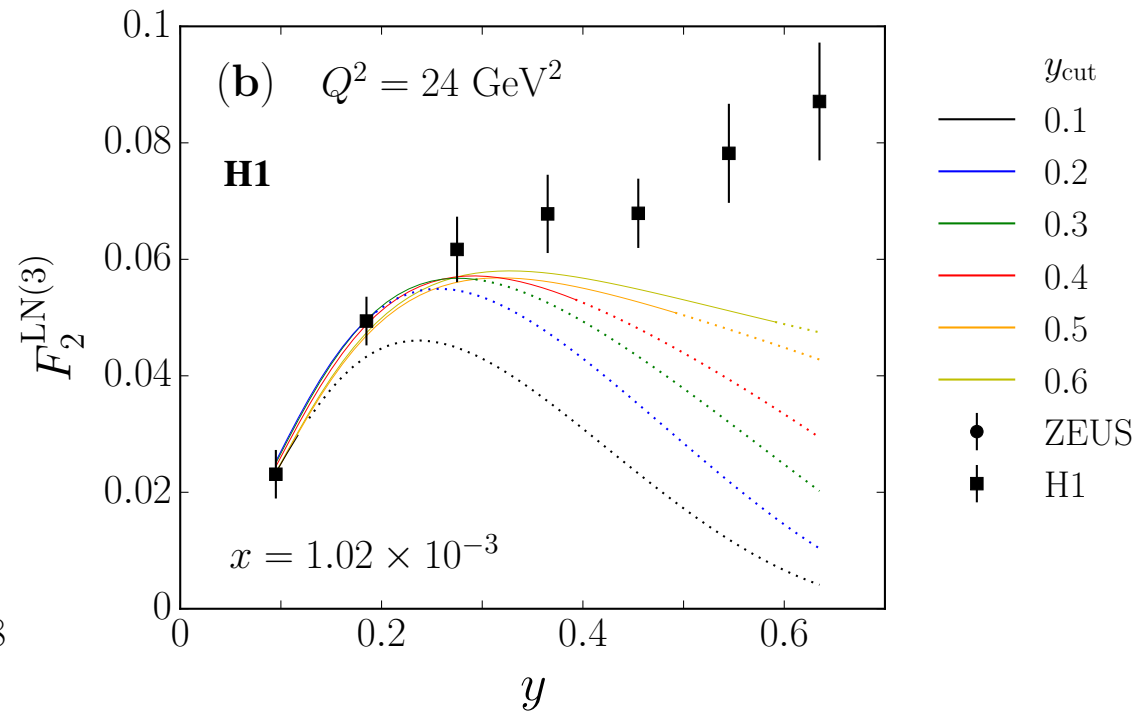
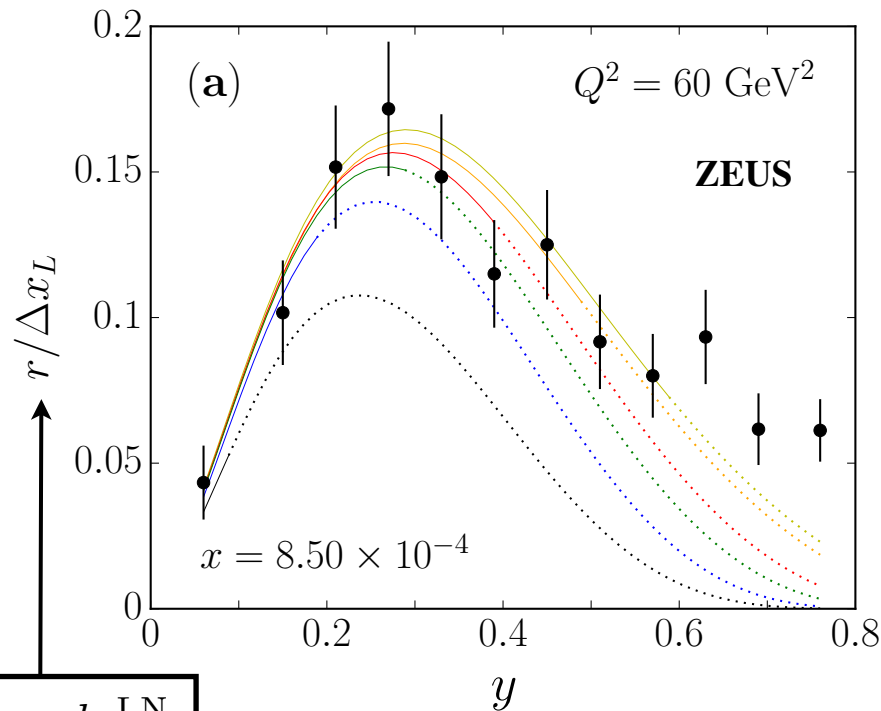
Leading neutron production at HERA

■ Measured LN differential cross section (integrated over p_{\perp})

$$\frac{d^3\sigma^{\text{LN}}}{dx dQ^2 dy} \sim F_2^{\text{LN}(3)}(x, Q^2, y)$$

$$2f_N^{(\text{on})}(y) F_2^{\pi}(x/y, Q^2) \text{ for } \pi \text{ exchange}$$

e.g.

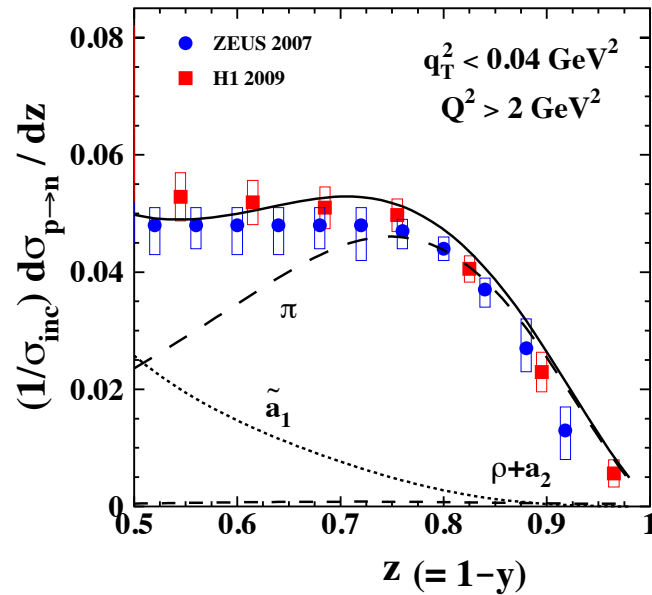


$$r = \frac{d\sigma^{\text{LN}}}{d\sigma^{\text{inc}}}$$

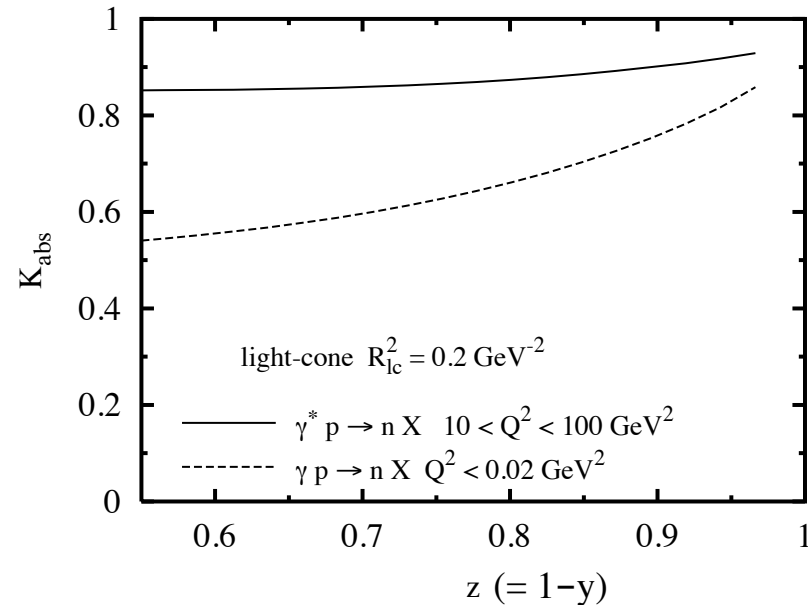
→ quality of fit depends on range of y fitted

Leading neutron production at HERA

- At large y , non-pionic mechanisms contribute (e.g. heavier mesons, absorption)



Kopeliovich et al., PRD 85, 114025 (2012)

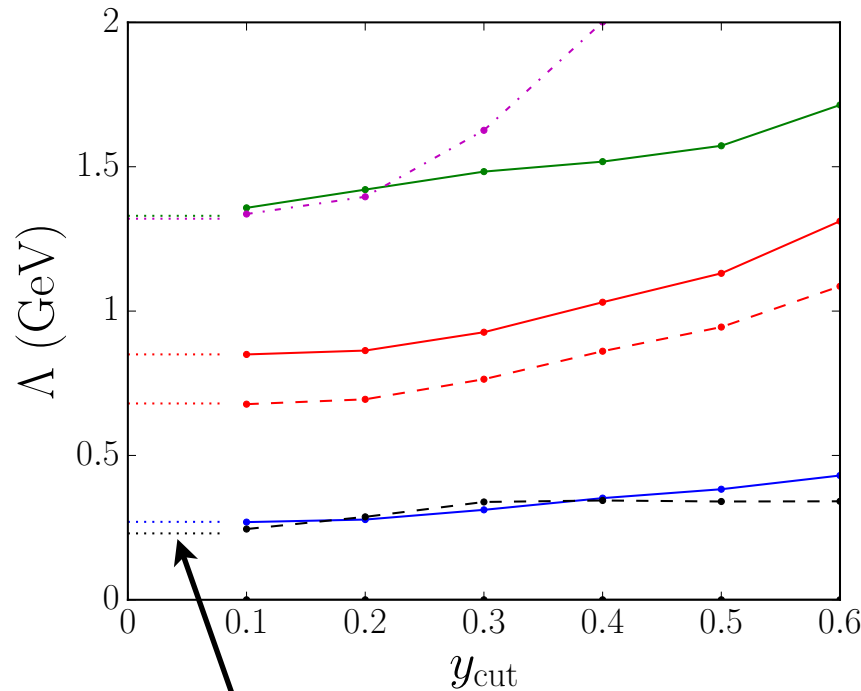


D'Alesio, Pirner, EPJA 7, 109 (2000)

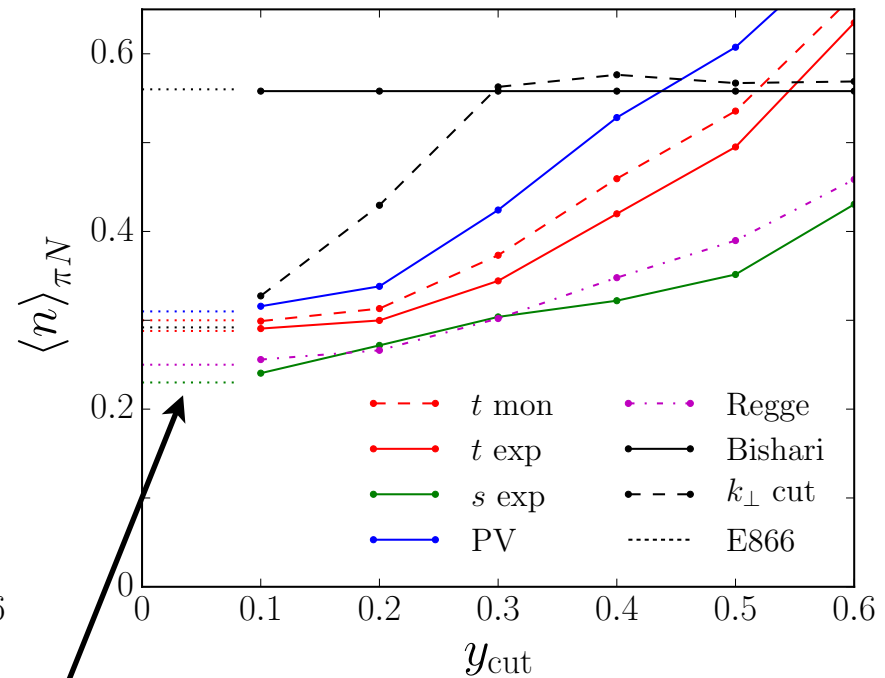
- To reduce model dependence, fit the value of y_{cut} up to which data can be described in terms of π exchange

Leading neutron production at HERA

- Fit requires higher momentum pions with increasing y_{cut}



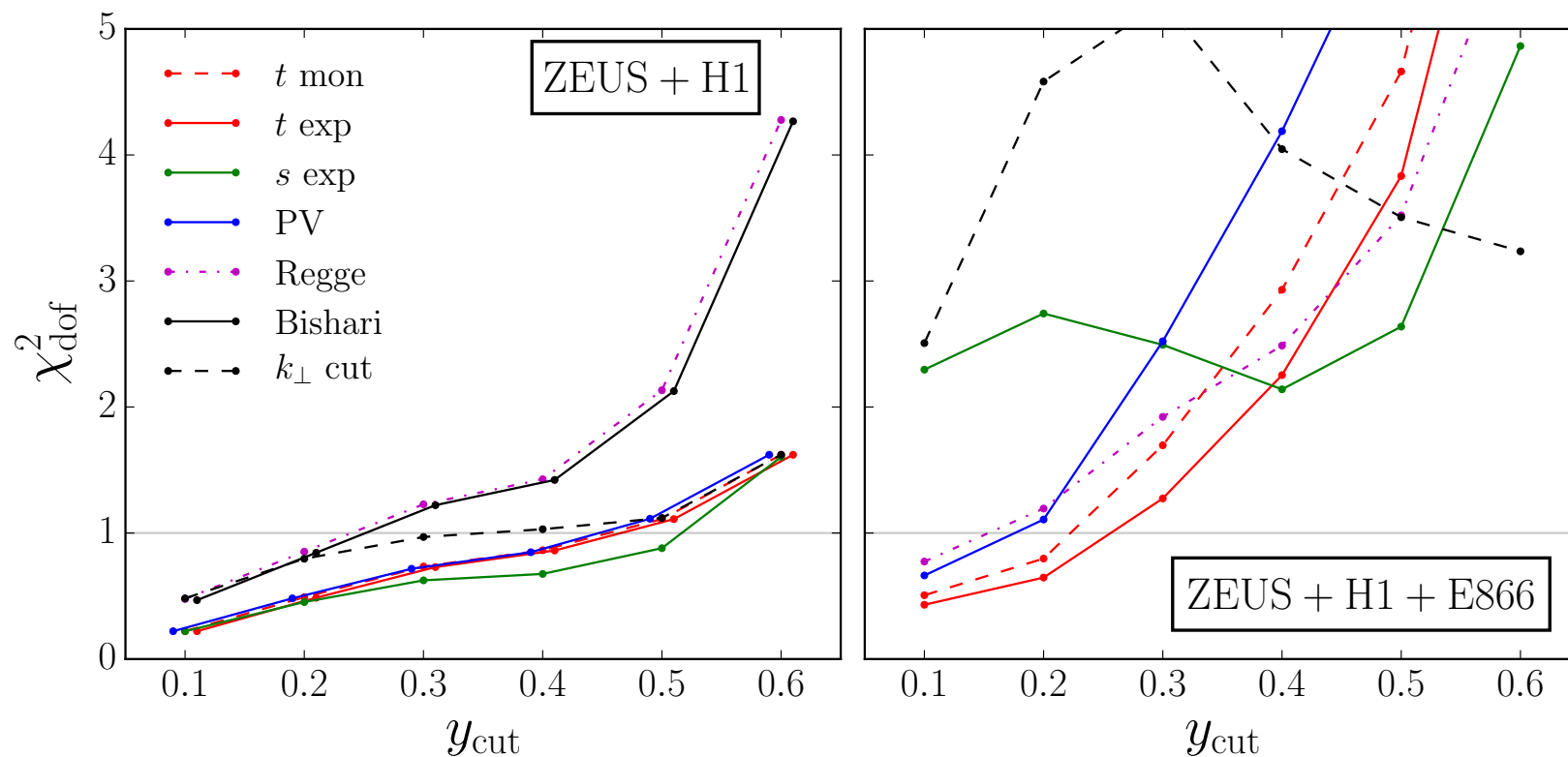
values from fit to E866 data only



→ larger values of y_{cut} more in conflict with E866 data

Leading neutron production at HERA

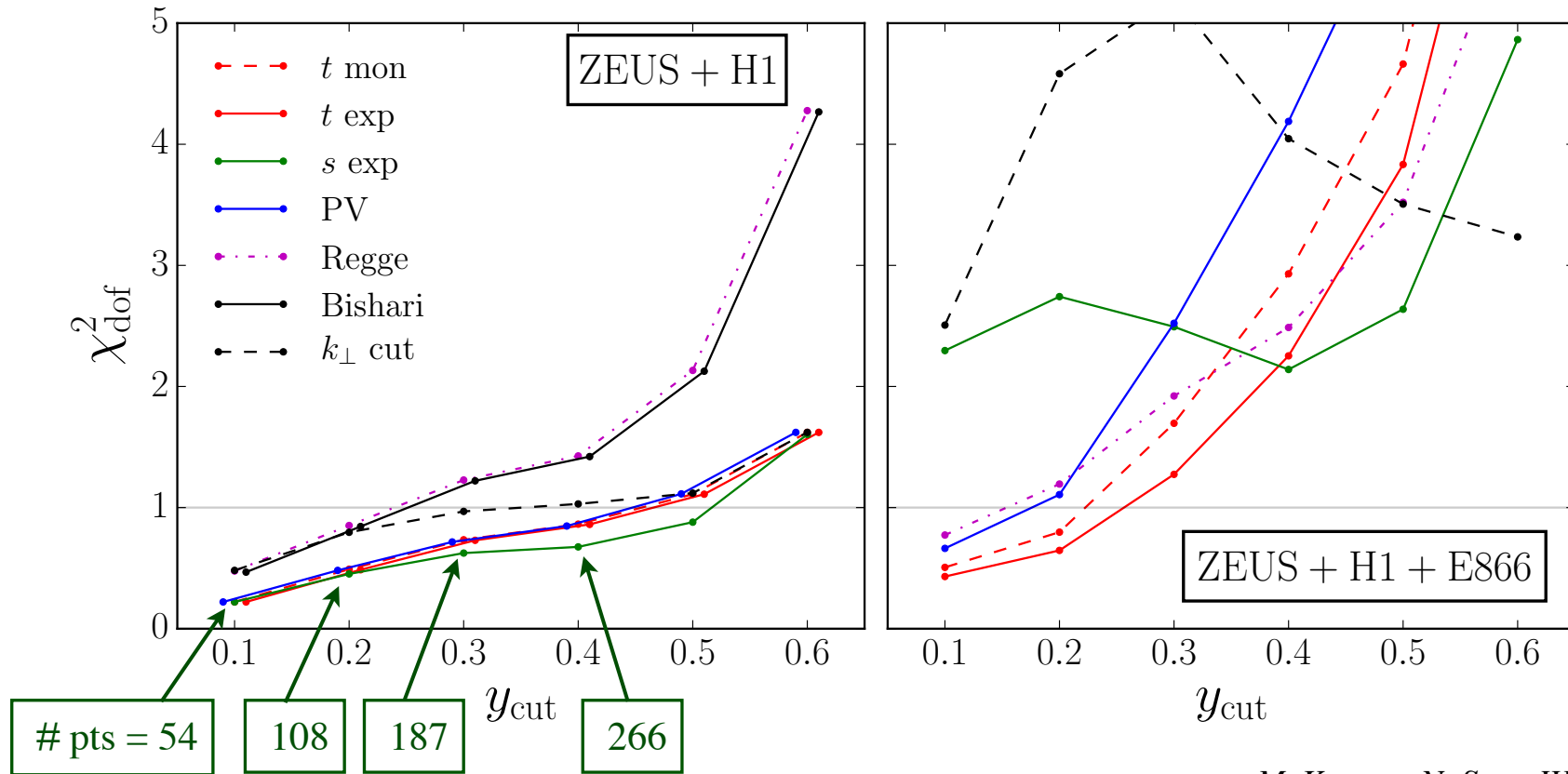
■ Combined fit to HERA LN and E866 Drell-Yan data



McKenney, N. Sato, WM, C. Ji
PRD **93**, 700205 (2016)

Leading neutron production at HERA

■ Combined fit to HERA LN and E866 Drell-Yan data

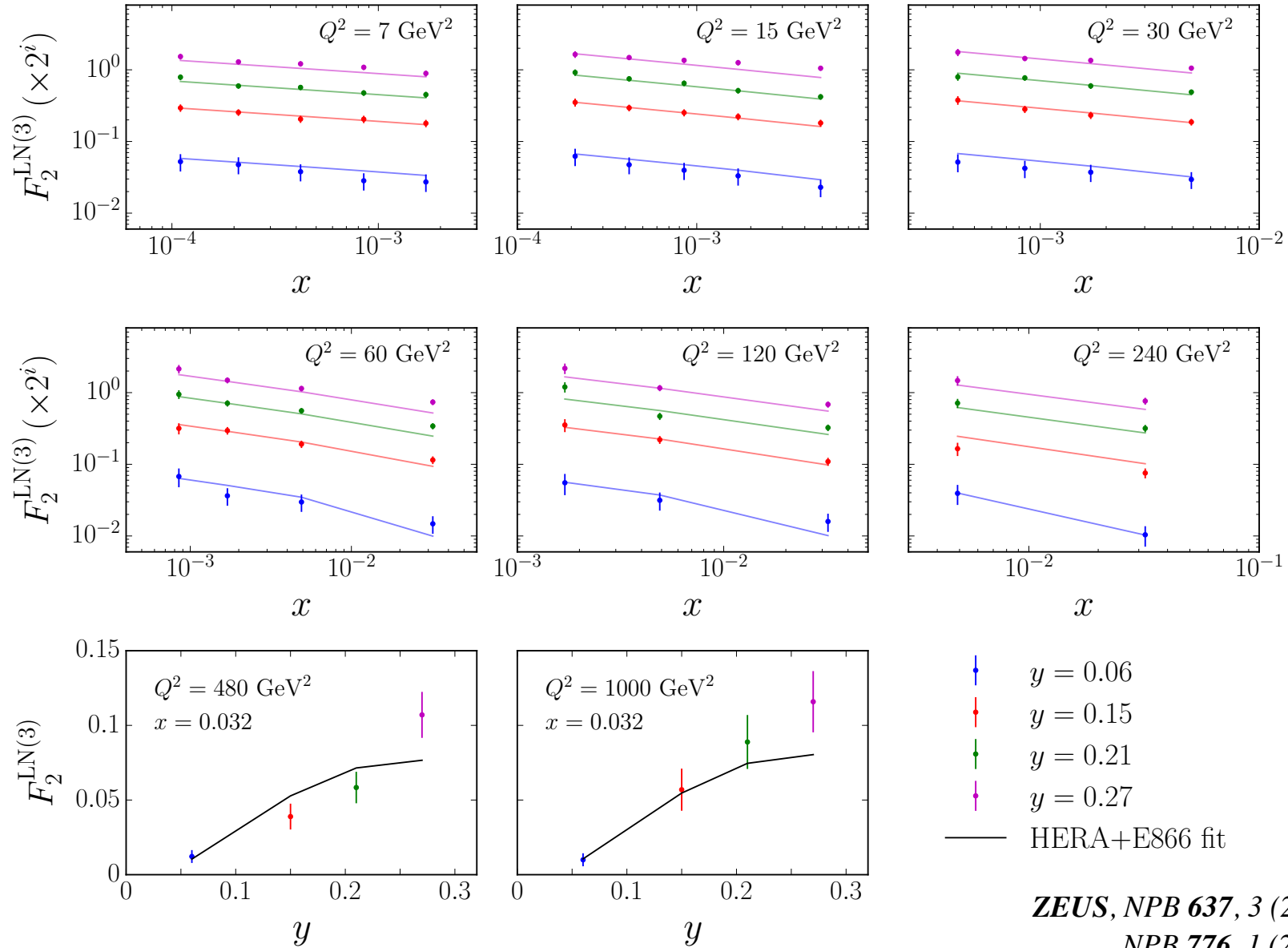


McKenney, N. Sato, WM, C. Ji
PRD **93**, 700205 (2016)

→ best fits for largest number of points afforded by t -dependent exponential (and t monopole) regulators

Leading neutron production at HERA

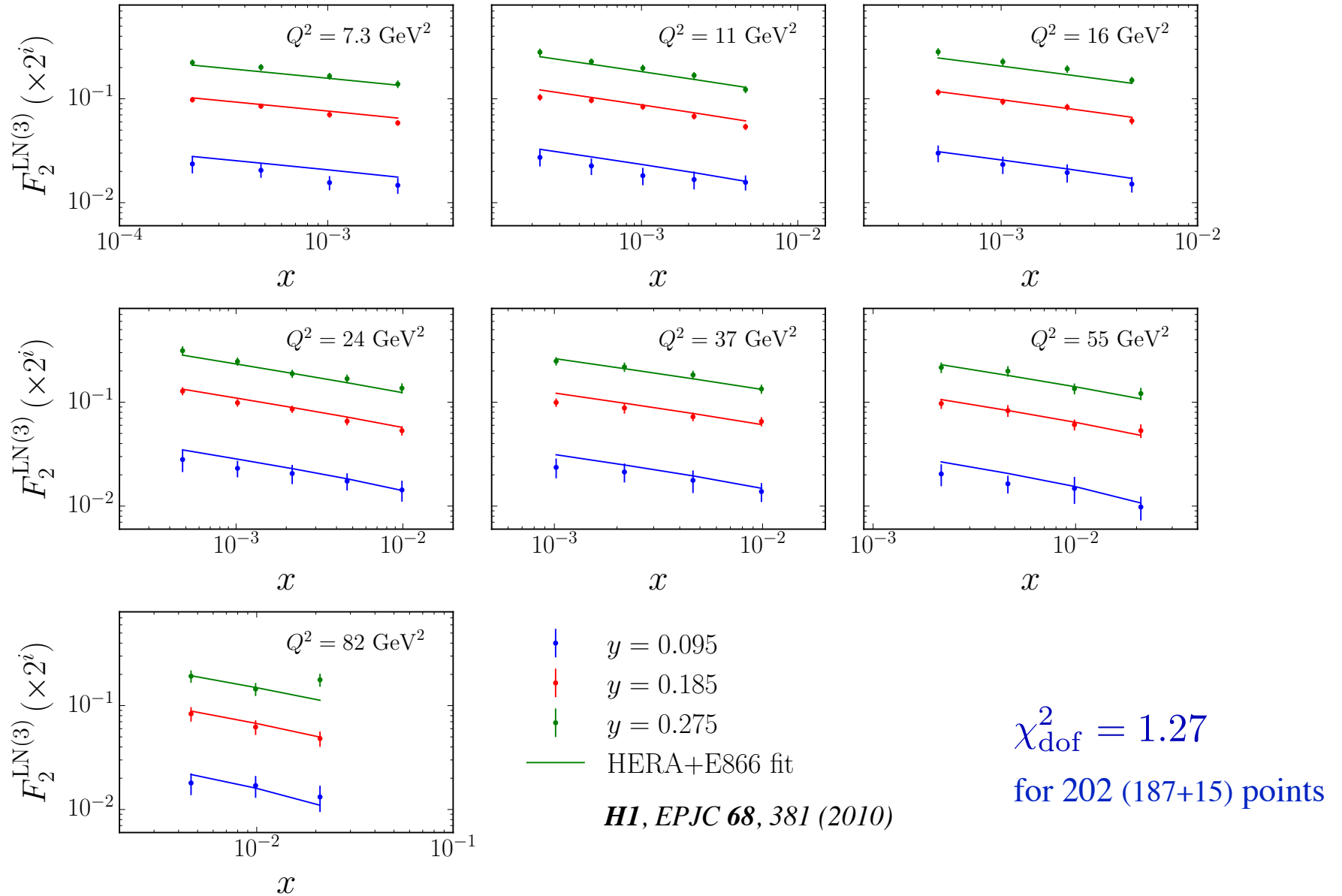
- Fit to ZEUS LN spectra for $y_{\text{cut}} = 0.3$ (t -dependent exponential)



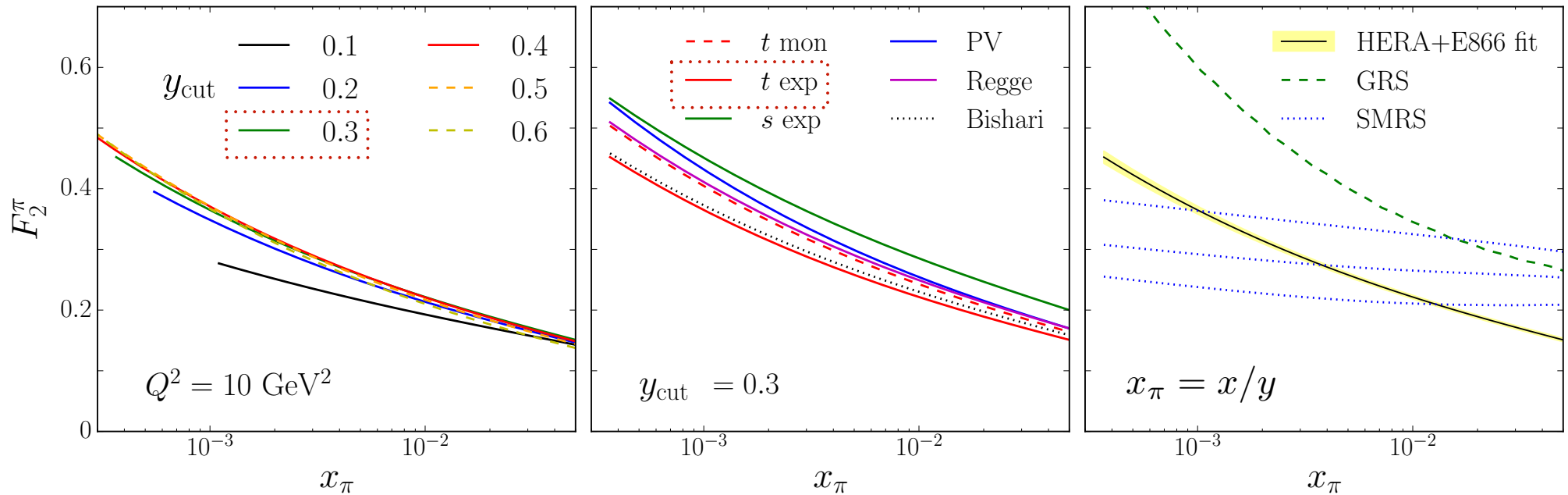
ZEUS, NPB 637, 3 (2002)
NPB 776, 1 (2007)

Leading neutron production at HERA

- Fit to H1 LN spectra for $y_{\text{cut}} = 0.3$ (t -dependent exponential)



Extracted pion structure function



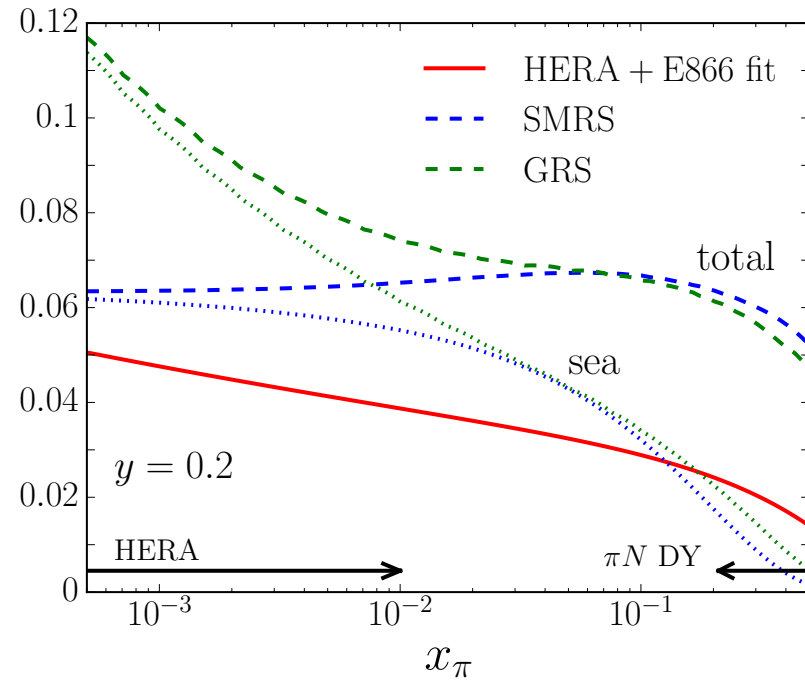
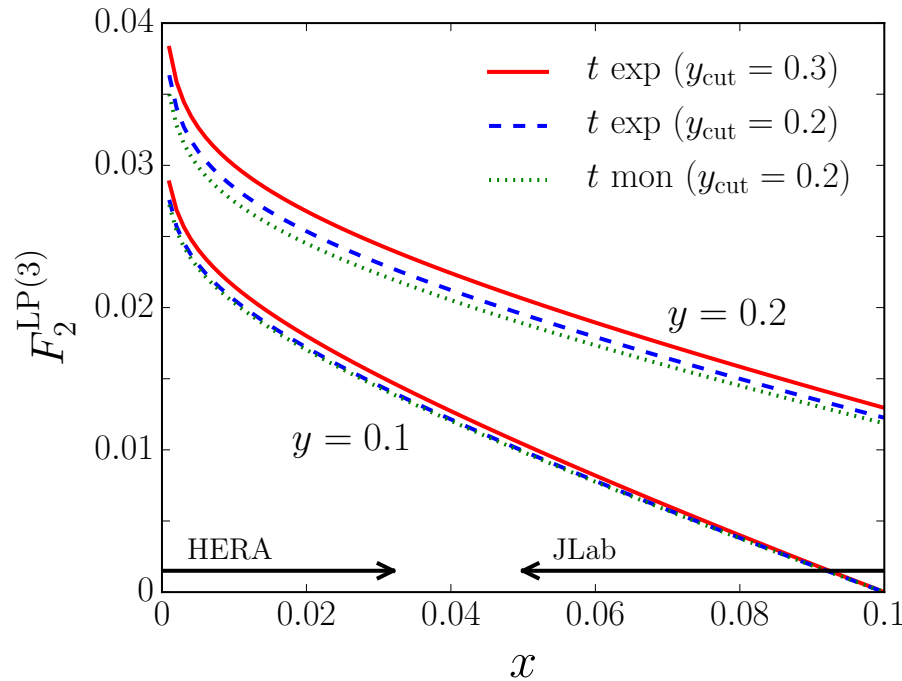
$$F_2^\pi = N x_\pi^a (1 - x_\pi)^b, \quad a = a_0 + a_1 \eta$$

$$\eta \sim \log(\log Q^2)$$

McKenney, N. Sato, WM, C. Ji
PRD **93**, 700205 (2016)

- stable values of F_2^π at $4 \times 10^{-4} \lesssim x_\pi \lesssim 0.03$ from combined fit
- shape similar to GRS fit to πN Drell-Yan data (for $x_\pi \gtrsim 0.2$), but smaller magnitude

Predictions at TDIS kinematics



$$F_2^{\text{LP}(3)} = f_{\pi-p}(y) F_2^\pi(x_\pi, Q^2)$$

McKenney, N. Sato, WM, C. Ji
PRD **93**, 700205 (2016)

→ JLab TDIS (“tagged” DIS, $ed \rightarrow eppX$) experiment
can fill gap in x_π coverage between HERA
and πN Drell-Yan kinematics

Outlook

- Combined analysis can be extended by including also πN Drell-Yan data
→ constrain large- x_π region ($x_\pi \gtrsim 0.2$)
- Generalize parametrization by fitting individual pion valence and sea quark PDFs, rather than F_2^π
- Ultimate goal will be to use all data sensitive to pion structure (including TDIS, EIC?) to constrain pion PDFs over full range $10^{-4} \lesssim x_\pi \lesssim 1$