Searching for optimal conditions for exploration of double-parton scattering in four-jet production at the LHC

Rafał Maciuła

Institute of Nuclear Physics (PAN), Kraków, Poland

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Outline

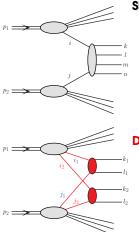
- \bullet Double-parton scattering mechanism: k_T -factorization
- 2 Theory vs. CMS four-jet data
- 3 Predictions for LHC Run2: optimal conditions for DPS
- Summary

Based on:

Maciuła, Szczurek, Phys. Lett. B 749 (2015) 57-62 (collinear approximation) Kutak, Maciuła, Serino, Szczurek, Hameren, paper in preparation (complementary studies in k_T -factorization)



Four-jet production: Mechanisms under consideration



Single-Parton Scattering (SPS $2 \rightarrow 4$)

- Kutak, Maciuła, Serino, Szczurek, Hameren, arXiv:1602.06814 (hep-ph) (accepted in JHEP)
- AVHLIB (A. van Hameren): https://bitbucket.org/hameren/avhlib
- High-Energy-Fact. (HEF): LO k_T -factorization (2 \rightarrow 4)
- first time: off-shell initial state partons
- more details: talk by Mirko Serino

Double-Parton Scattering (DPS $4 \rightarrow 4$)

- ullet Factorized Ansatz with experimental setup of $\sigma_{
 m eff}$
- LO k_T -factorization approach (2 \rightarrow 2 \otimes 2 \rightarrow 2)
- more precise studies of kinematical characteristics and correlation observables

<u>extension of</u> our previous studies based on LO collinear approach (ALPGEN):

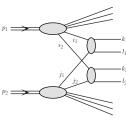
Maciuła, Szczurek, Phys. Lett. B 749 (2015) 57-62



Four-jet production in double-parton scattering (DPS)

Factorized ansatz (pocket-formula)

In a simple probabilistic picture:



process initiated by two simultaneous hard parton-parton scatterings in one proton-proton interaction ⇒

$$\underbrace{\overset{-}{-}_{l_1}}_{l_1} \sigma^{\text{DPS}}(pp \to 4 \text{jets} X) = \frac{C}{\sigma_{\text{eff}}} \cdot \sigma^{\text{SPS}}(pp \to \text{dijet} X_1) \cdot \sigma^{\text{SPS}}(pp \to \text{dijet} X_2)$$

two subprocesses are not correlated and do not interfere

analogy: frequently considered mechanisms of double charm, double gauge boson production and double Drell-Yan anihillation

valid also differentially:

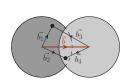
$$\frac{\textit{d}\sigma^{\textit{DPS}}(\textit{pp} \rightarrow 4 \text{jets X})}{\textit{d}\xi_1 \textit{d}\xi_2} = \sum\nolimits_{\stackrel{j_1,j_1,k_1,l_1}{l_2,j_2,k_2,l_2}} \frac{\textit{C}}{\sigma_{\textit{eff}}} \, \frac{\textit{d}\sigma(\textit{i}_1\textit{j}_1 \rightarrow \textit{k}_1\textit{l}_1)}{\textit{d}\xi_1} \, \frac{\textit{d}\sigma(\textit{i}_2\textit{j}_2 \rightarrow \textit{k}_2\textit{l}_2)}{\textit{d}\xi_2} \; ,$$

where
$$C = \left\{ \begin{array}{ll} \frac{1}{2} & \text{if } i_1 j_1 = i_2 j_2 \wedge k_1 l_1 = k_2 l_2 \\ 1 & \text{if } i_1 j_1 \neq i_2 l_2 \vee k_1 l_1 \neq k_2 l_2 \end{array} \right\}$$
 and $i, j, k, l = g, u, d, s, c, \bar{u}, \bar{d}, \bar{s}, \bar{c}.$

ullet combinatorial factors C include identity of the two subprocesses



Factorized ansatz and double-parton distributions (DPDFs)



DPDF - emission of parton *i* with assumption that second parton *j* is also emitted:

$$\Gamma_{i,j}(b, x_1, x_2; \mu_1^2, \mu_2^2) = F_i(x_1, \mu_1^2) F_j(x_2, \mu_2^2) F(b; x_1, x_2, \mu_1^2, \mu_2^2)$$

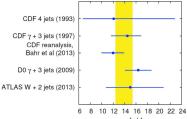
correlations between two partons
 C. Flensburg et al., JHEP 06, 066 (2011)

in general:

$$\frac{\sigma_{\text{eff}}(x_1, x_2, x_1', x_2', \mu_1^2, \mu_2^2)}{\sigma_{\text{eff}}(x_1, x_2, x_1', x_2', \mu_1^2, \mu_2^2)} = \left(\int d^2b \ F(b; x_1, x_2, \mu_1^2, \mu_2^2) \ F(b; x_1', x_2', \mu_1^2, \mu_2^2) \right)^{-1}$$

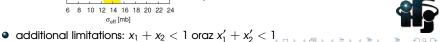
Factorized ansatz:

- DPDF in multiplicative form: $F_{ij}(b; x_1, x_2, \mu_1^2, \mu_2^2) = F_i(x_1, \mu_1^2) F_j(x_2, \mu_2^2) F(b)$
- $\sigma_{\rm eff} = \left[\int d^2b \left(F(b)\right)^2\right]^{-1}$, F(b) energy and process independent

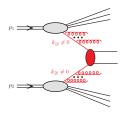


phenomenology: $\sigma_{\rm eff} \Rightarrow$ nonperturbative quantity with a dimension of cross section, connected with transverse size of proton $\sigma_{\rm eff} \approx 15 \ {\rm mb} \ (p_{\perp} - {\rm independent})$

a detailed analysis of $\sigma_{\rm eff}$: Seymour, Siódmok, JHEP 10, 113 (2013)



SPS dijet production: k_t -factorization (semihard) approach



$$k_t$$
-factorization $\longrightarrow \kappa_{1,t}, \kappa_{2,t} \neq 0$

Collins-Ellis, Nucl. Phys. B360 (1991) 3;

Catani-Ciafaloni-Hautmann, Nucl. Phys. B366 (1991) 135; Ball-Ellis, JHEP 05 (2001) 053

⇒ efficient approach for jet-jet or QQ correlations

multi-differential cross section

$$\frac{d\sigma}{dy_{1}dy_{2}d^{2}p_{1,t}d^{2}p_{2,t}} = \sum_{i,j} \int \frac{d^{2}\kappa_{1,t}}{\pi} \frac{d^{2}\kappa_{2,t}}{\pi} \frac{1}{16\pi^{2}(x_{1}x_{2}s)^{2}} \frac{|\mathcal{M}_{i^{*}j^{*}\rightarrow kl}|^{2}}{|\mathcal{M}_{i^{*}j^{*}\rightarrow kl}|^{2}}$$

$$\times \delta^{2}(\vec{\kappa}_{1,t} + \vec{\kappa}_{2,t} - \vec{p}_{1,t} - \vec{p}_{2,t}) \mathcal{F}_{i}(x_{1}, \kappa_{1,t}^{2}) \mathcal{F}_{j}(x_{2}, \kappa_{2,t}^{2})$$

- $\mathcal{F}_i(x_1, \kappa_{1,t}^2)$, $\mathcal{F}_j(x_2, \kappa_{2,t}^2)$ unintegrated (k_t -dependent) PDFs
- LO off-shell $|\mathcal{M}_{i^*j^* \to kl}|^2 \Rightarrow$ calculated numerically in AVHLIB analytical form: Nefedov, Saleev, Shipilova, Phys. Rev. D87, 094030 (2013) Quasi Multi Regge Kinematics (QMRK) with effective BFKL NLL vertices

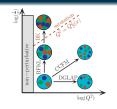
$$\sqrt{s} \gg p_T$$
. $M \gg \Lambda_{OCD}$ and $x \ll 1$

Parton-Reggeization Approach (k_T -factorization with Reggeized initial partons): an effective way to take into account amount part of radiative corrections at high energy Regge kinematics

 some part of higher-order corrections may be effectively included depending on UPDF model ⇒ possible emission of extra (hard) gluons



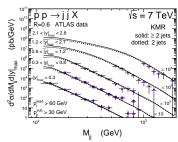
Unintegrated parton distribution functions (UPDFs)

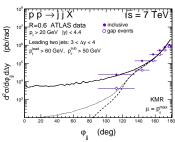


Different evolution equations (or their combinations):

- Kwieciński, Jung (CCFM, wide range of x)
- Kimber-Martin-Ryskin (DGLAP-BFKL, wide range of x)
- Kwieciński-Martin-Staśto (BFKL-DGLAP, small x-values)
- Kutak-Staśto (BK, saturation, only small x-values)

Lessons from inclusive dijet production at the LHC:





• KMR UPDFs work well for jet-jet correlation observables in dijet production



DPS in the framework of k_T -factorization

DPS production of four-jet system within k_T -factorization approach, assuming factorization of the DPS model:

$$\frac{d\sigma^{DPS}(pp\to 4\mathrm{jets}X)}{dy_1dy_2d^2p_{1,t}d^2p_{2,t}dy_3dy_4d^2p_{3,t}d^2p_{4,t}}=\frac{C}{\sigma_{\mathrm{eff}}}\cdot\frac{d\sigma^{SPS}(pp\to j\,j\,X_1)}{dy_1dy_2d^2p_{1,t}d^2p_{2,t}}\cdot\frac{d\sigma^{SPS}(pp\to j\,j\,X_2)}{dy_3dy_4d^2p_{3,t}d^2p_{4,t}}$$
 Each step of DPS (each individual scattering):

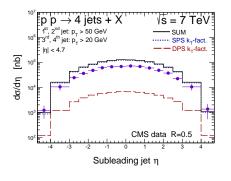
$$\frac{d\sigma^{SPS}(pp \to j j X)}{dy_{1}dy_{2}d^{2}p_{1,t}d^{2}p_{2,t}} = \frac{1}{16\pi^{2}\hat{s}^{2}} \int \frac{d^{2}k_{1t}}{\pi} \frac{d^{2}k_{2t}}{\pi} \overline{|\mathcal{M}_{l^{n}k^{*} \to j j}|^{2}} \\ \times \delta^{2} \left(\vec{k}_{1t} + \vec{k}_{2t} - \vec{p}_{1t} - \vec{p}_{2t}\right) \mathcal{F}(x_{1}, k_{1t}^{2}, \mu^{2}) \mathcal{F}(x_{2}, k_{2t}^{2}, \mu^{2})$$

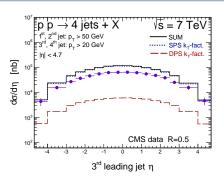
- 9 channels from the 2 \rightarrow 2 SPS \Rightarrow 45 channels for the 4 \rightarrow 4 DPS but only 14 contribute to \geq 95% of the cross section (see Mirko's talk)
- KMR UPDFs from CT10 NLO collinear PDFs
- ullet $n_{F}=4$ flavour scheme, running $a_{S}@NLO$ from MSTW08 package
- scales: $\mu=\mu_{\it R}=\mu_{\it F}={1\over 2}\sum_i p_{\it T}^i$ (sum over final state particles)
- ullet all details the same for 2 ightarrow 4 SPS and 4 ightarrow 4 DPS calculations





CMS four-jets: SPS + DPS in the k_T -factorization

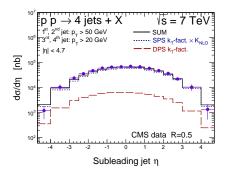


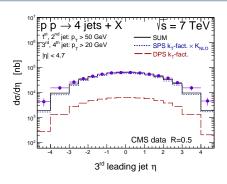


- the SPS component above the data ⇒ the same problem with the ALPGEN code (LO collinear approach) ⇒ exact SPS NLO calculations needed?
- first full SPS NLO (collinear) four-jets: Z. Bern et al., Phys. Rev. Lett. 109, 042001 (2012) NLO corrections \Rightarrow damping of the cross section \Rightarrow $K_{NLO} \approx 0.5$
- SPS 2 \rightarrow 2: $K_{NLO} \approx 1.1 1.2 \Rightarrow$ much less important for DPS
- much better description of exp. data for harder p_T cuts (see Mirko's talk)



CMS four-jets: SPS + DPS in the k_T -factorization

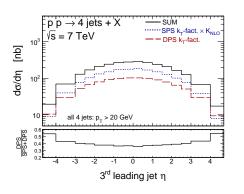


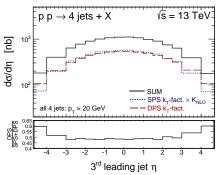


- Now: quite good description of the CMS data
- 3rd leading jet (softer): forward/backward region slightly underestimated
- very small DPS contribution \Rightarrow unsupportive CMS cuts: 1st, 2nd jet $p_7 > 50$ GeV
- **DPS favoured**: small p_T region (see Mirko's talk)



DPS effects in four-jet sample: lowering p_T cuts



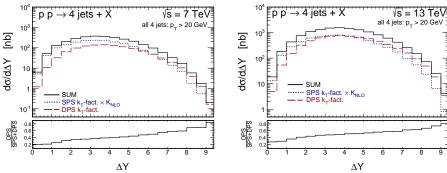


- all 4 jets: $20 < p_T < 100 GeV$
- 13 TeV: DPS contribution $\geq 50\%$
- DPS favoured: forward/backward rapidity region



DPS effects in four-jet sample: large rapidity distance

Rapidity difference between jets most remote in rapidity



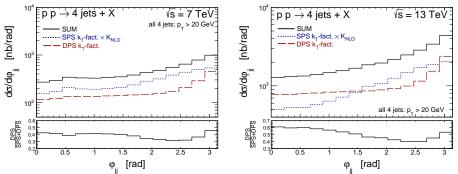
• 13 TeV:

 $\Delta Y > 6 \Rightarrow$ four-jet sample dominated by DPS



DPS effects in four-jet sample: large rapidity distance

Azimuthal angle between jets most remote in rapidity



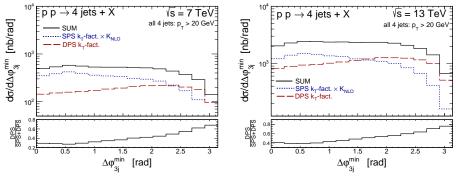
• 13 TeV:

 $\varphi_{ij} < \frac{\pi}{2} \Rightarrow$ four-jet sample dominated by DPS



DPS effects in four-jet sample: special angular correlation

Minimum azimuthal separation between any three jets



- variable proposed by ATLAS analysis: JHEP 12, 105 (2015) (see talk by Melissa Ridel)
- distinguishes events with pairs of nearby jets (which have large φ_{3j}^{min}) from the recoil of three jets against one (leading to small φ_{3i}^{min} values)
- 13 TeV:

 $\Delta arphi_{3j}^{min} > rac{\pi}{2} \Rightarrow$ four-jet sample dominated by DPS



Conclusions

A recipe for DPS dominated four-jet sample at $\sqrt{s}=13$ TeV:

- ① crucial: lower transverse momentum cuts for all 4 jets: $p_T > 20$ GeV
 - asymmetric configuration also acceptable: leading jet $p_T > 35$ GeV; 2^{nd} , 3^{rd} , 4^{th} jet $p_T > 20$ GeV however any further increasing of the p_T cuts leads to significant damping of the DPS contribution
- ② concentrate on large jet-jet rapidity separations: $\Delta Y > 6$
- lacksquare useful angular jet-jet correlations: $arphi_{jj} < rac{\pi}{2}$, $\Delta arphi_{3j}^{min} > rac{\pi}{2}$

Thank you for your attention!

