

# ***Electron Polarimetry at JLEIC (Jefferson Lab Electron-Ion Collider)***

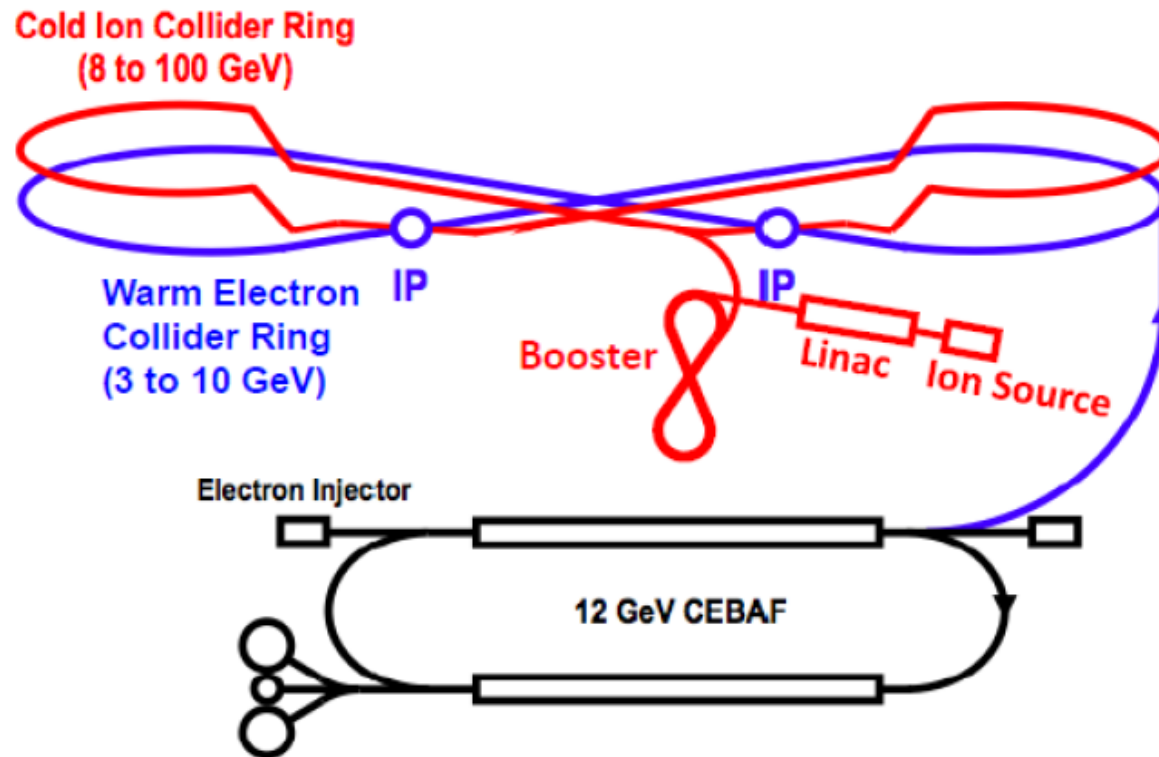
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Jefferson lab

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# Outline

- JLEIC electron beam
- Chicane for low- $Q^2$  tagging and Compton polarimetry
- Compton polarimetry - experience at JLab
- Polarimeter “baseline” design
- Future polarimeter R&D

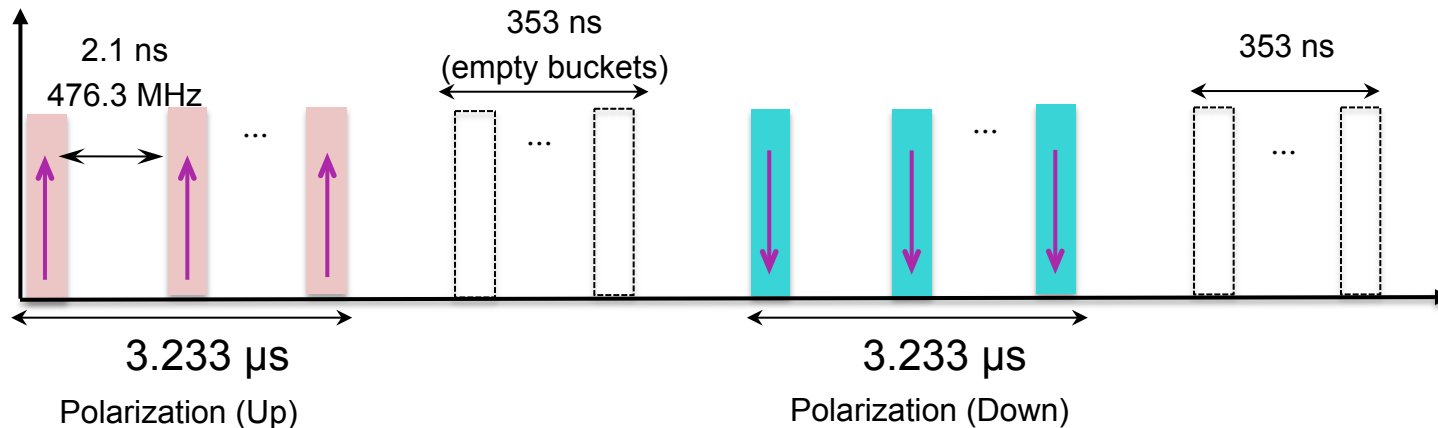
# JLEIC Electron Beam Structure and Polarization



- Storage ring: 476.3 MHz = 2.1 ns bunch structure
- 3 A at 5 GeV and 720 mA at 10 GeV
- 2 macrobunches with one polarization; each macrobunch = 3.2  $\mu$ s

# Electron Beam Time structure

bunch train & polarization pattern in the collider ring



Bunch spacing = 2.1 ns

Macrobunches with opposite polarization = 3.233  $\mu\text{s}$  long

1. Average polarization of beam in ring can be measured with single laser helicity
2. Polarization of each macrobunch can be determined independently by flipping laser helicity

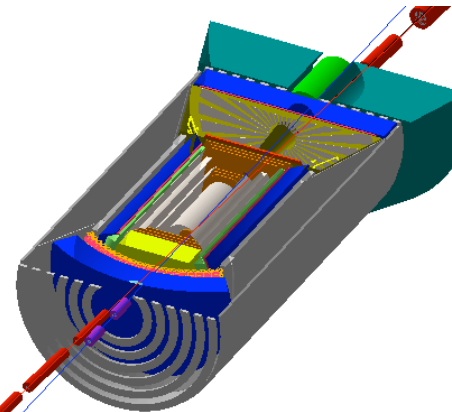
Note: revolution time = 7.17  $\mu\text{s}$ . Flipping laser helicity may require times of order 40-50  $\mu\text{s}$ , or longer

# Low $Q^2$ Tagger and Polarimetry

Low  $Q^2$ /nearly-real photon tagging:

→ Electrons scattered at very small angles (with small energy loss) not in the acceptance of main detector

→ Use of chicane downstream of IP allows detection of these electrons



Electron polarimetry:

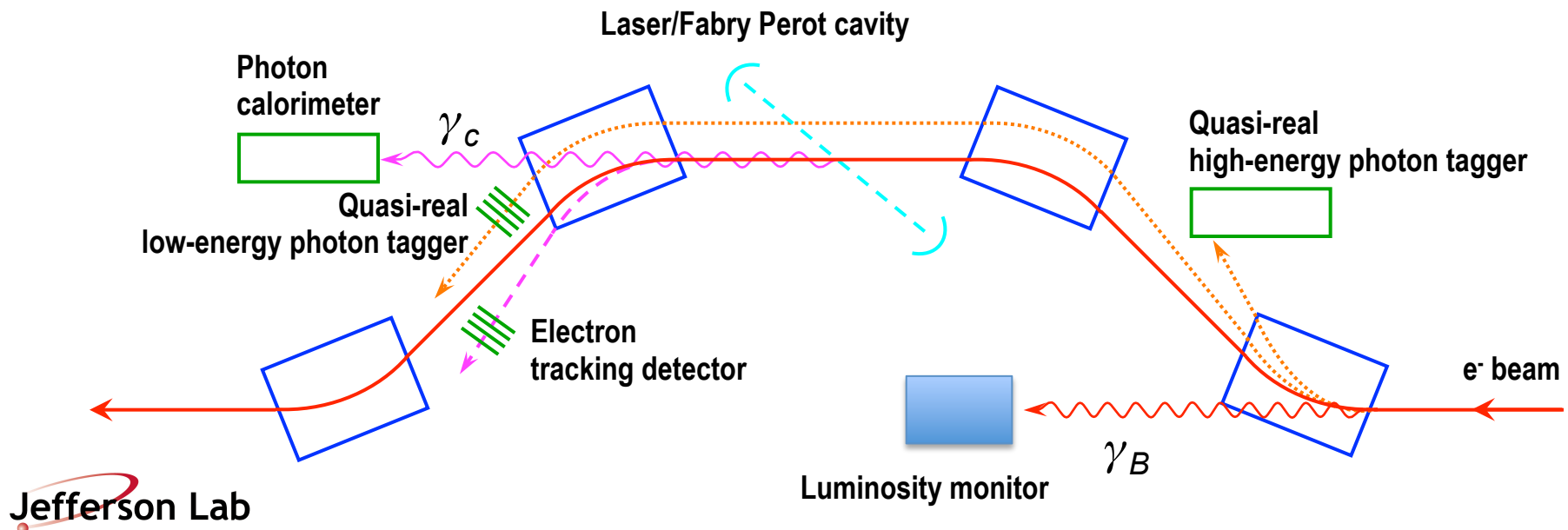
→ Unprecedented statistical precision available at JLEIC implies we may become dominated by systematic uncertainties

→ Luminosity measurement also potentially polarization dependent

→ Precision electron polarimetry crucial

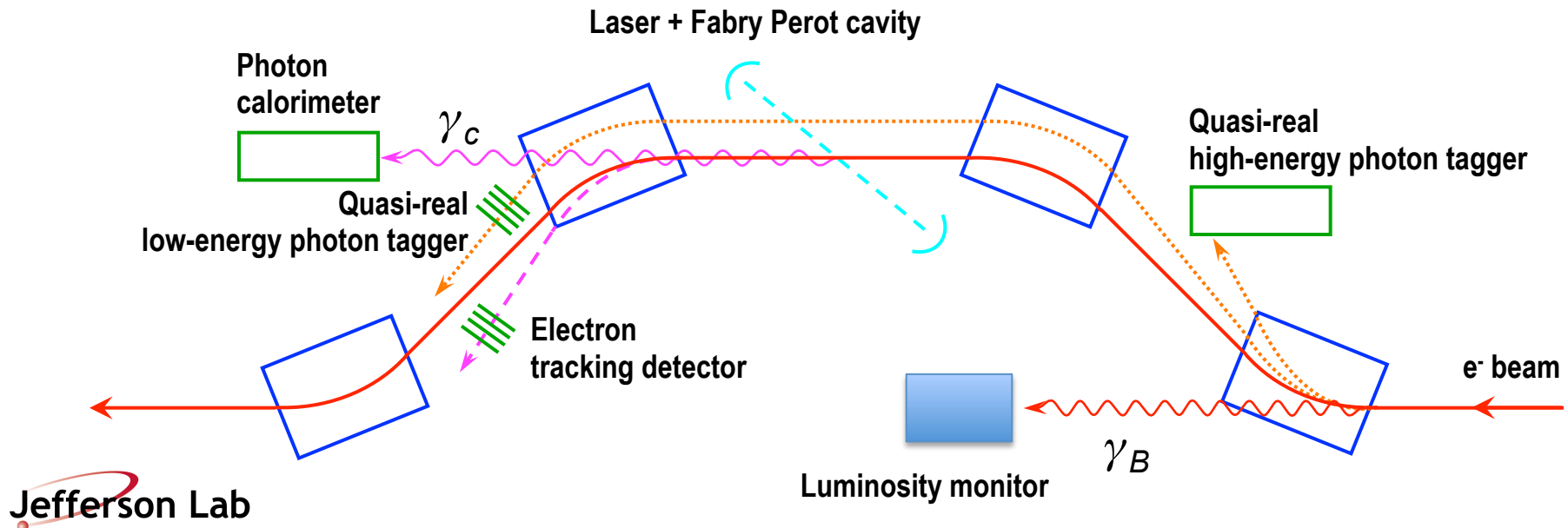
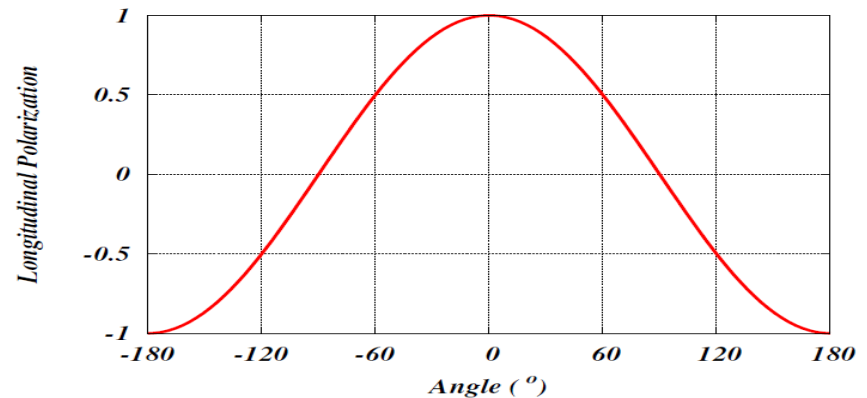
# Chicane for Low $Q^2$ tagger and Compton Polarimeter

- At JLEIC, Compton can share chicane with low  $Q^2$  tagger
- Laser-electron collisions in middle of chicane assures no spin rotation relative to IP
- No interference with electron detectors needed for low  $Q^2$  tagger



# Chicane for Low $Q^2$ tagger and Compton Polarimeter

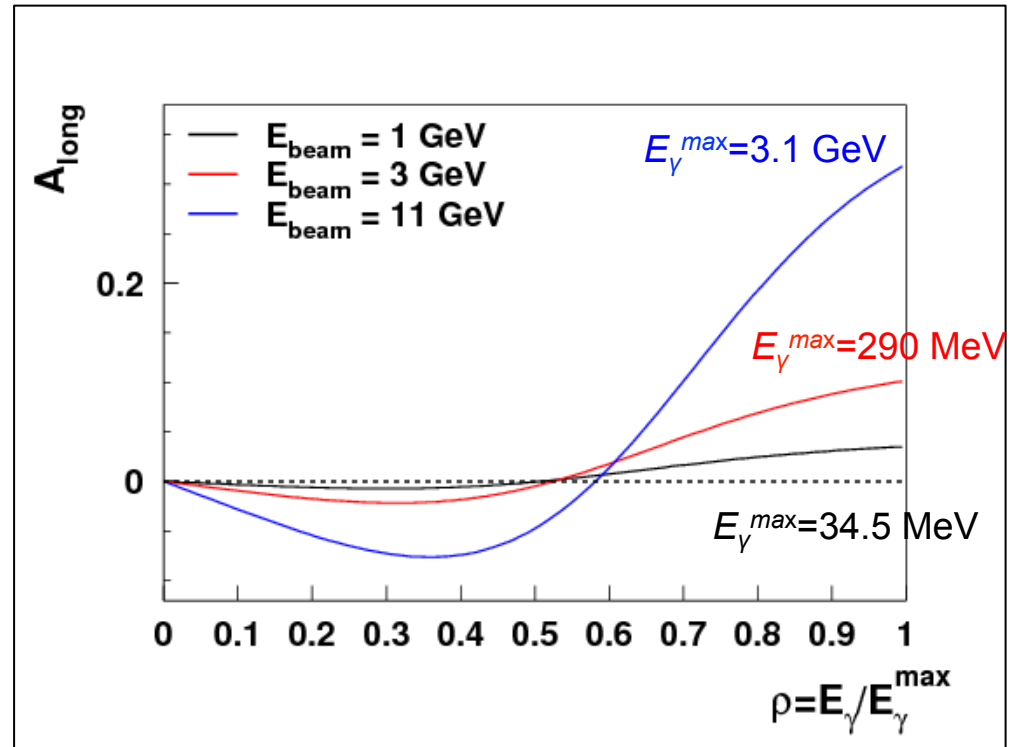
- At collision with laser, electron beam has same polarization direction as at IP due to net zero bend
- Use of spin rotators allows us to perform “spin dance” to verify longitudinal polarization at IP



# Compton Polarimetry

Compton polarimetry ideal method for electron polarimetry at JLEIC

- Photon “target” very thin – no impact on electron beam
- High precision accessible – sub-1% precision has been achieved



Beam polarization extracted via double-spin asymmetry:

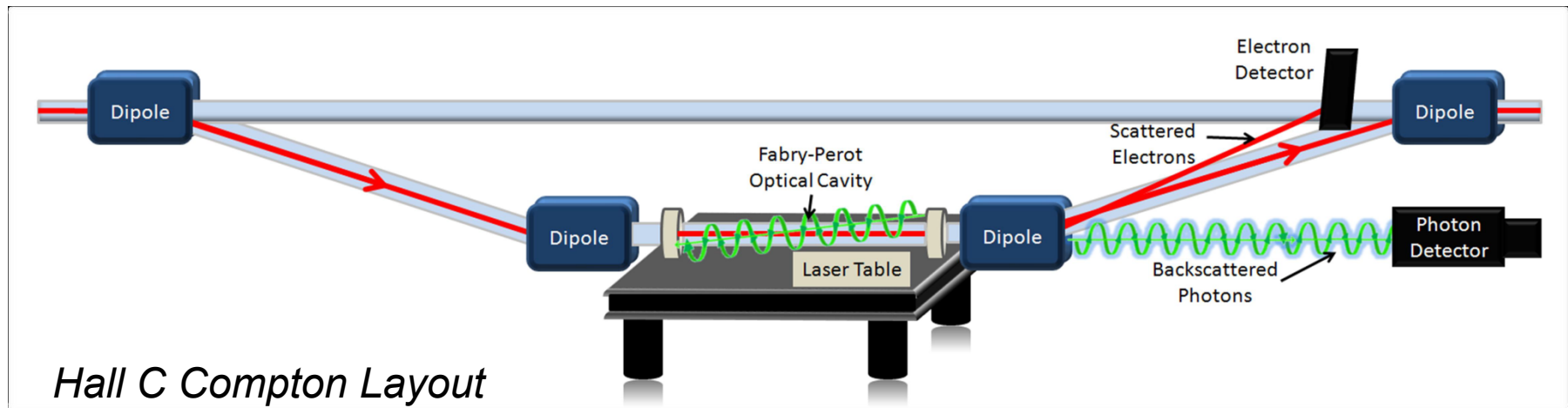
$$A_{\text{meas}} = P_{\text{laser}} P_{\text{beam}} A_{\text{th}} = \frac{\sigma^{\uparrow\downarrow} - \sigma^{\uparrow\uparrow}}{\sigma^{\uparrow\downarrow} + \sigma^{\uparrow\uparrow}}$$

Laser+electron spins parallel

Laser+electron spins anti-parallel



# Compton Polarimetry – Experience at JLab



JLab has built two similar Compton polarimeters in Halls A and C  
→ Both have achieved ~1% electron beam polarization measurements

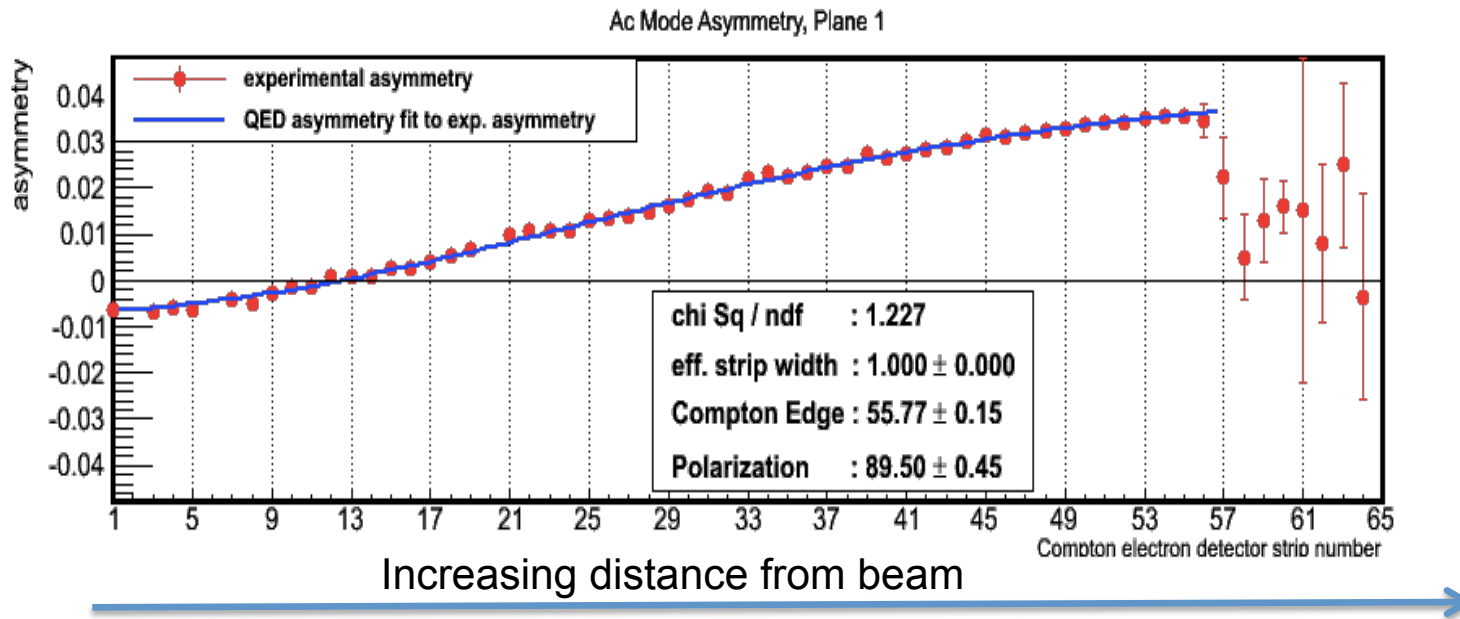
Important design considerations:

1. Dipole chicane allows simultaneous measurement of scattered electrons and backscattered photons
2. Electron-laser collision at center of chicane assures no difference in electron spin direction relative to beam before/after chicane
3. Continuous electron beam might require high power CW laser system due to background issues

# Precision Compton Polarimetry

- Precision goal for electron beam polarization is  $dP/P = 1\%$
- Sub-1% polarimetry has been achieved at:
  - SLD  $\rightarrow 0.52\%$  at 45.6 GeV (electron detection)
  - JLab Hall A  $\rightarrow 1\text{-}3$  GeV (electron and photon detection)
  - JLab Hall C  $\rightarrow 1$  GeV (electron detection)
- Sub-1% precision has only been achieved via photon detection using threshold-less, “integrating” technique
  - Large synchrotron backgrounds at JLEIC may make this impossible
- For now, the JLEIC Compton design emphasizes detection of the Compton scattered electron

# Compton Electron Detector



Hall C @ JLab: Diamond microstrips used for electron detector

Analysis employs a 2 parameter fit (polarization and Compton edge) to the differential spectrum

→ This has yielded good results → strip width (resolution) is important

→ Zero-crossing must be in acceptance to constrain the fit well

Dominant systematics related to the interplay between trigger and strip efficiency

# Laser and Backgrounds

- Choice of system depends on backgrounds in Compton polarimeter
- Main sources of background
  - Bremsstrahlung from residual gas in beampipe
  - Synchrotron radiation
  - Beam halo interacting with detector and/or apertures in beamline
- Two potential choices for laser system
  - Single pass, CW or pulsed laser → 10s of Watts easily achievable
  - High gain Fabry-Perot cavity → kW laser powers

# Compton Laser Options

- Single pass laser
  - Advantages: Able to rapidly flip helicity ( $\sim 10 \mu\text{s}$ ), relatively simple alignment
  - Disadvantages: Relatively low power  $\rightarrow$  backgrounds may become problem, requires small crossing angle  $\rightarrow$  interaction point stretched out – more care needed for good beam overlap
- Fabry-Perot cavity
  - Advantages: High power  $\rightarrow$  improved signal:noise
  - Disadvantages: Potential source of background (beam halo), technically complex, beam must be aligned to laser (or movable laser table needed)

# Laser and Backgrounds

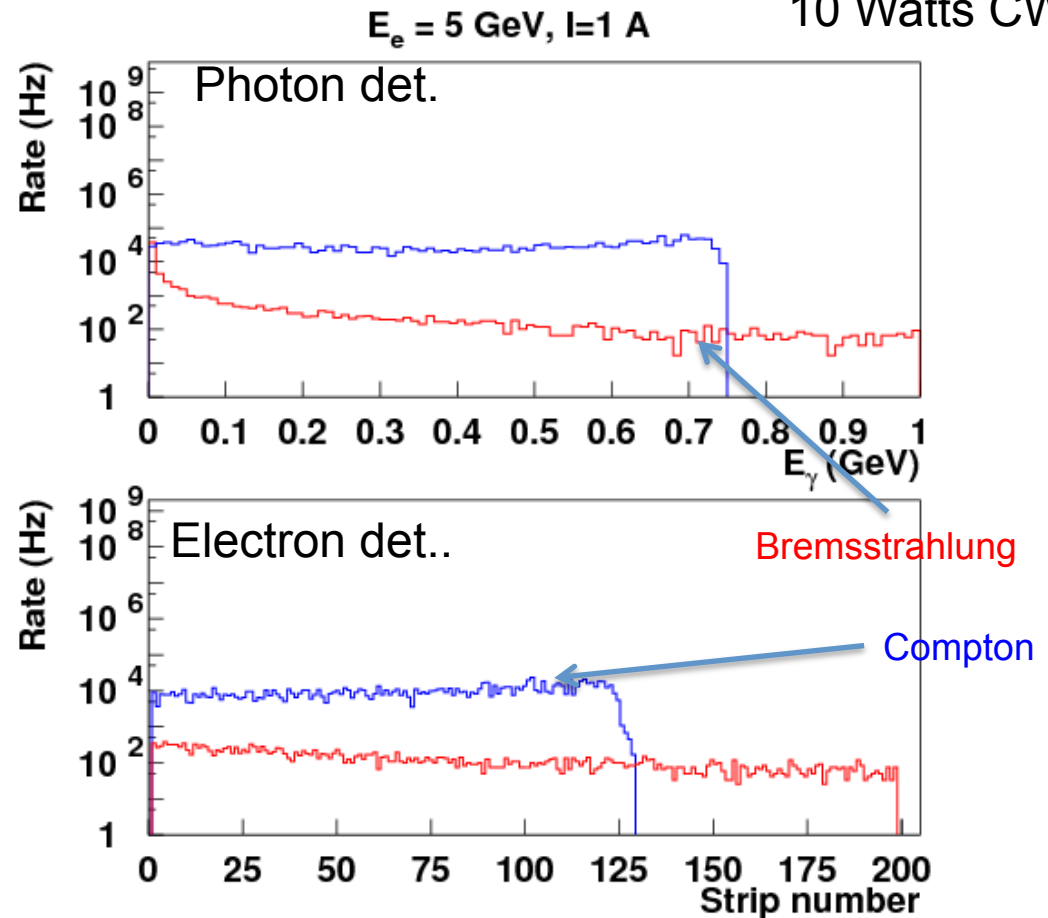
Historically, Compton polarimeters have been able to suppress backgrounds by matching laser pulse structure to beam

→ With modern, CW machines, there is little to be gained in this manner

Conventional CW lasers may be useable if backgrounds can be controlled

Rates and backgrounds: MEIC

Green laser  
10 Watts CW



# Laser and Backgrounds

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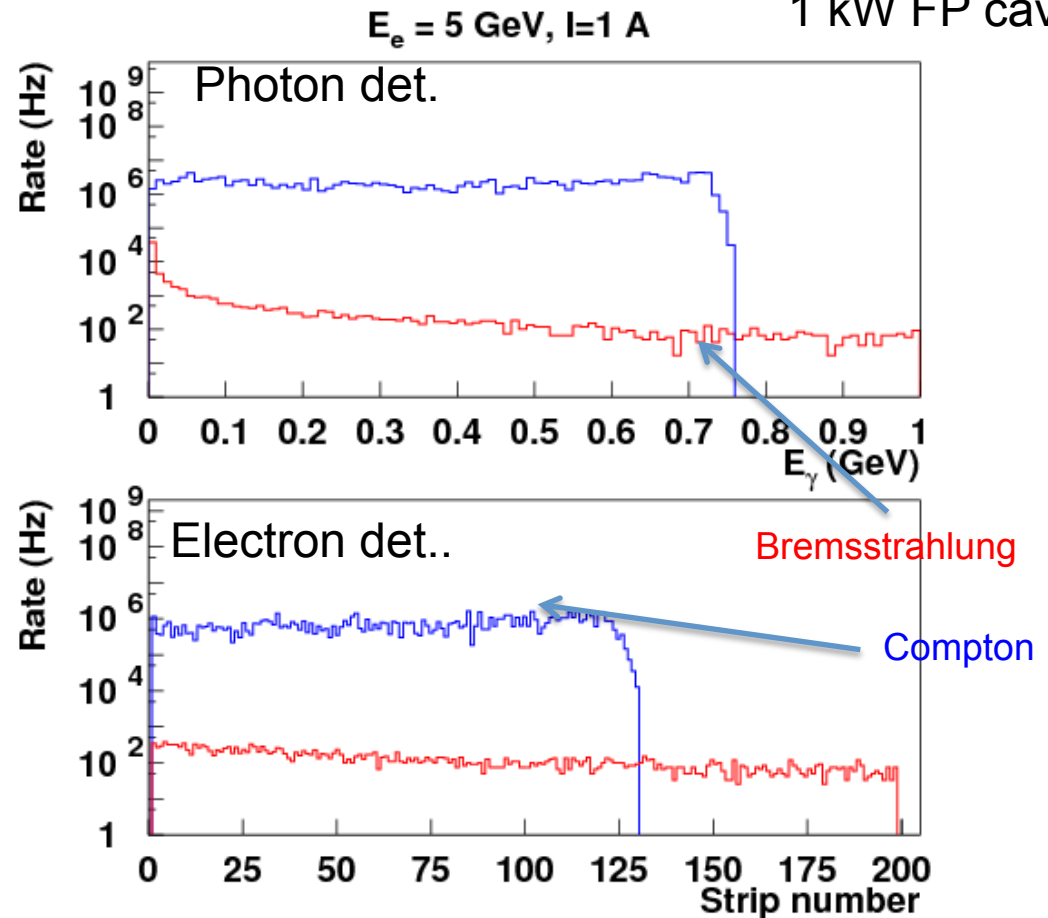
→ Modern CW machines, there is little to be gained in this manner

Conventional CW lasers may be useable if backgrounds can be controlled

→ FP cavity can provide significantly higher rates  
→ improved signal:background

Rates and backgrounds: MEIC

Green laser  
1 kW FP cavity

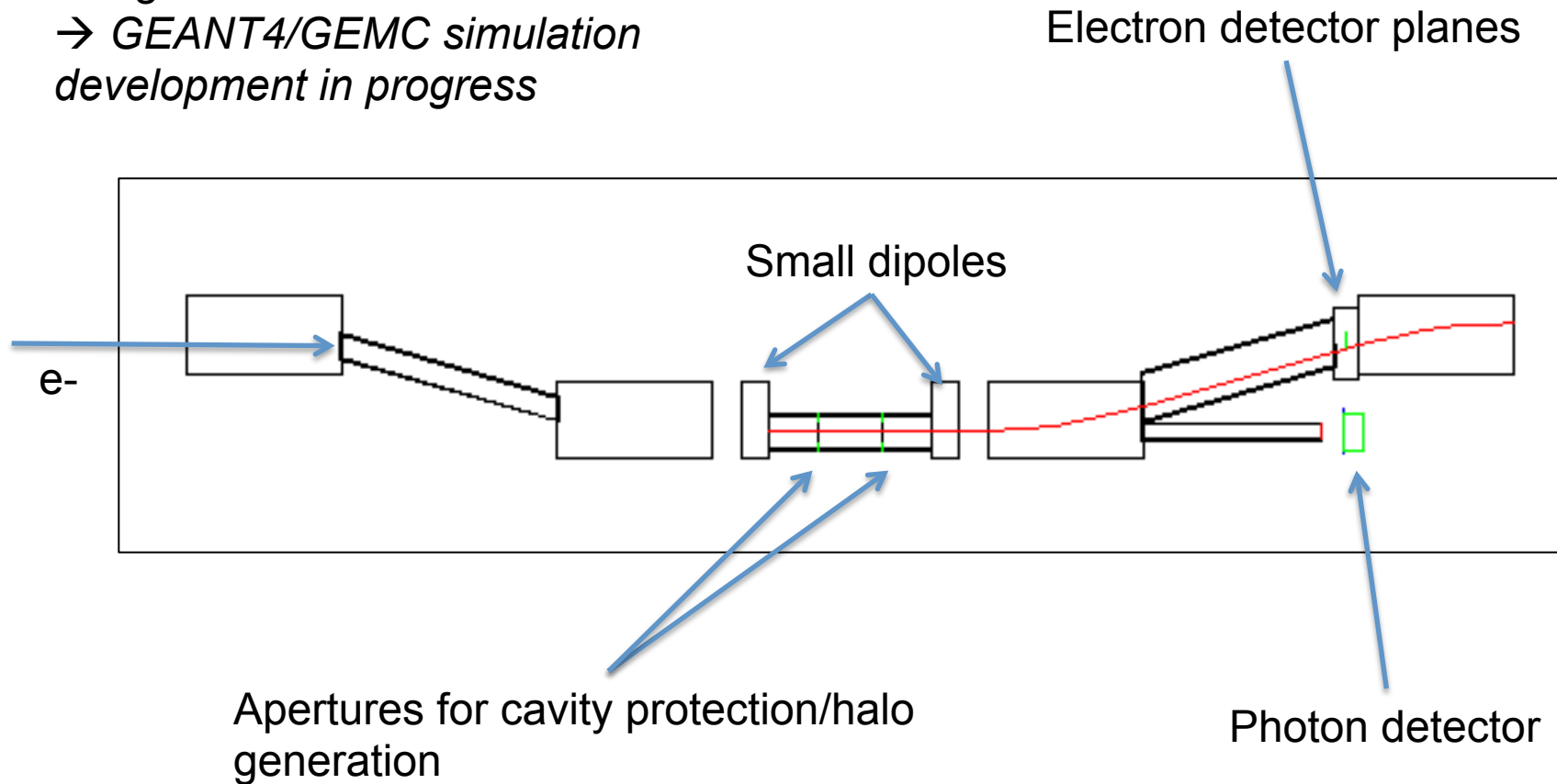


Beamline vac. =  $10^{-9}$  Laser-beam crossing angle  $\sim 2.6$  deg.

# Simulation of Rates and Backgrounds

Initial estimates performed  
using GEANT3

→ *GEANT4/GEMC simulation  
development in progress*



Realistic beam sizes from machine design used

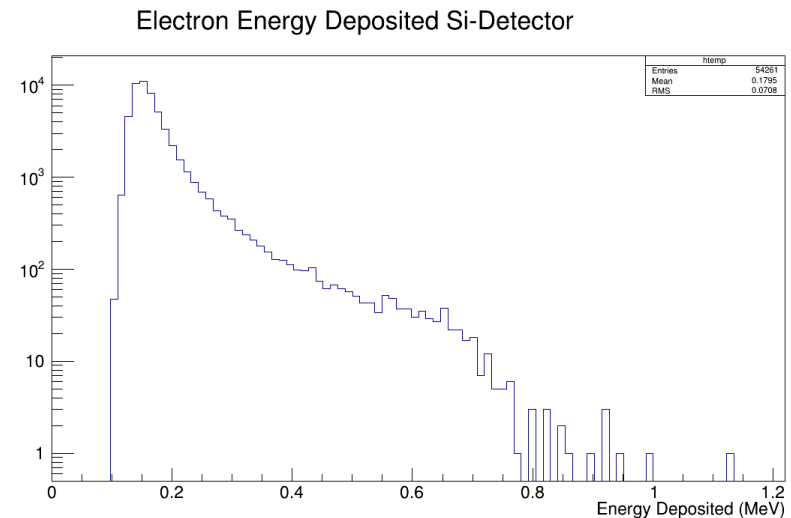
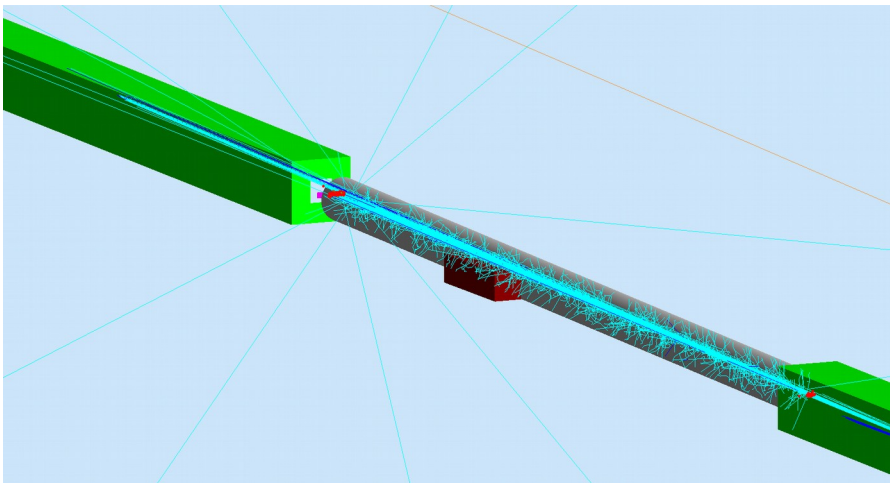


# Simulation of Rates and Backgrounds

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→ *GEANT4/GEMC simulation  
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GEANT4 Compton simulation will use same framework (GEMC) as Central  
Detector simulation → fully integrated into one common simulation



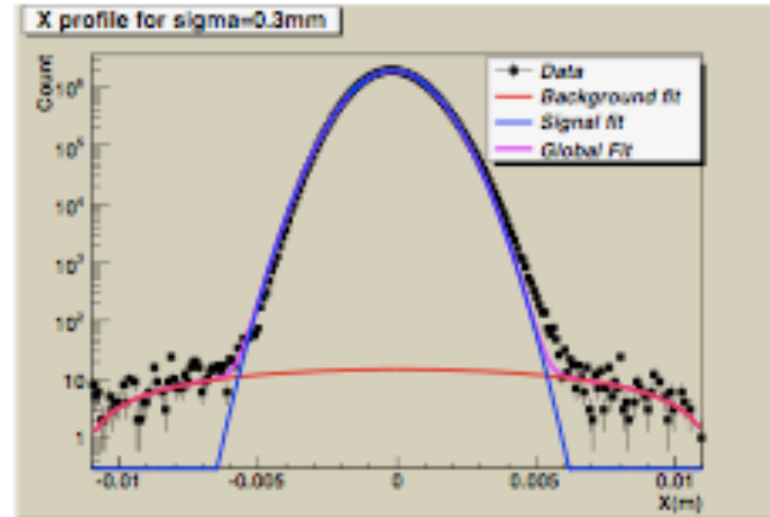
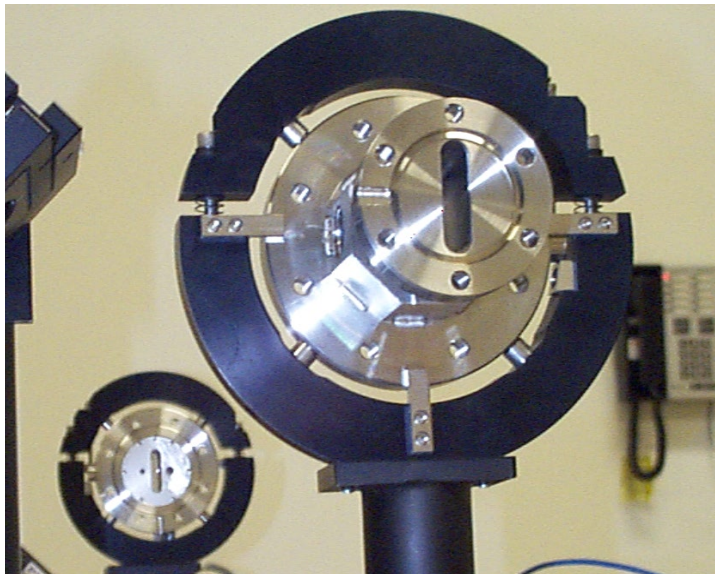
*Josh Hoskins, University of Manitoba*

# Beam Halo and Backgrounds

Halls A and C use CW, Fabry-Perot cavities

→ Both systems have mirrors ~5 mm from the beam

→ Small apertures protect mirrors from beam excursions, really bad beam properties

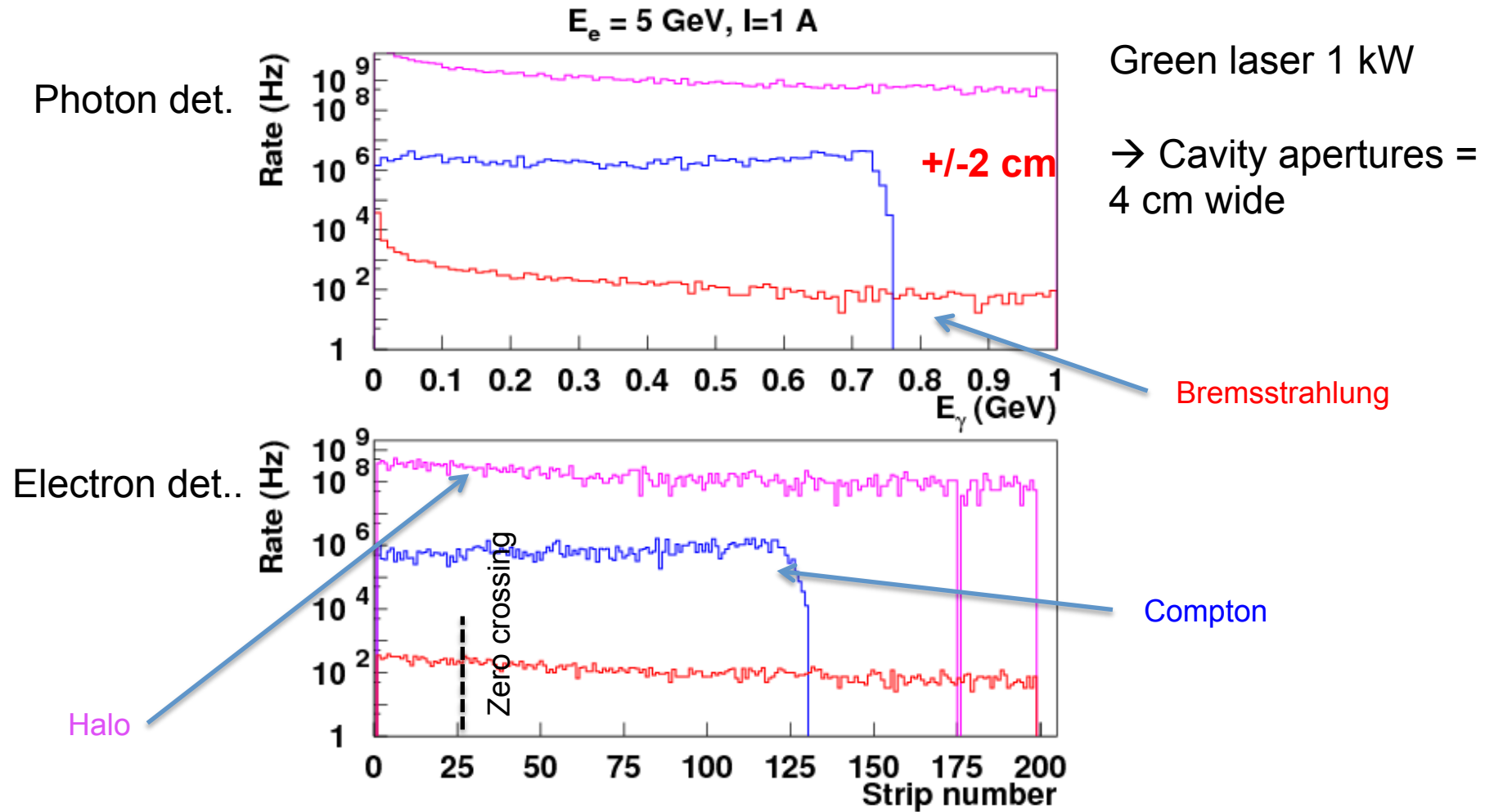


*Yves Roblin and Arne Freyberger*  
*JLAB-TN-06-048*

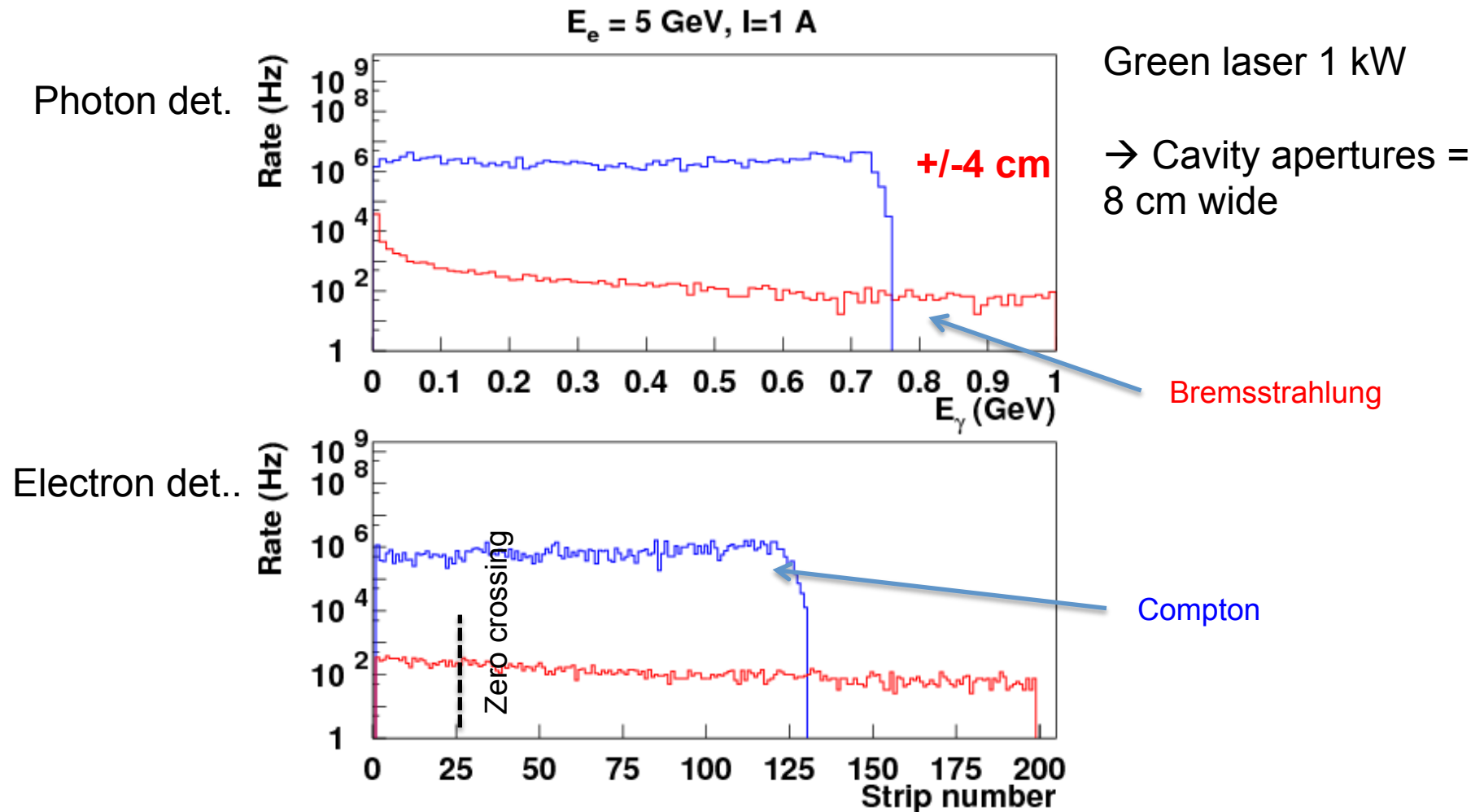
Same protective apertures can lead to backgrounds due to interactions with beam halo

Use of FP cavity at MEIC depends on understanding halo

# Laser and Backgrounds - Halo



# Laser and Backgrounds - Halo



# Projected Rates and Measurements Times

Energy (GeV)	Current (A)	1 pass laser (10 W)		FP cavity (1 kW)	
		Rate (MHz)	Time (1%)	Rate (MHz)	Time (1%)
3 GeV	3	26.8	161 ms	310	14 ms
5 GeV	3	16.4	106 ms	188	9 ms
10 GeV	0.72	1.8	312 ms	21	27 ms

1 pass laser crossing angle = 0.3 degrees, FP cavity = 2.6 degrees  
Time for 1% (statistics) measurement assumes 70% polarization  
Rates integrated from asymmetry zero-crossing

Extremely high rates when using FP cavity means that detectors (electron and photon) will have to operate in integrating mode in that case

# Electron Detector Requirements

- Segmentation → allows determination of the beam polarization with high precision by fitting the spectrum
- High rate capability
  - Scattered electron rates will be very large
  - Typical “strip” detectors have relatively slow response times after amplification → large dead time
  - Integrating mode?
- Radiation hard
  - Dose rates will be on the order of 7-25 krad/hour
  - Example: Silicon signal/noise smaller by factor of 2 after 3 Mrad
- Diamond micro-strip detector default choice for JLEIC Compton
  - Excellent performance during Q-Weak experiment in Hall C
  - Exposed to 10 Mrad dose with no degradation in performance

# Baseline JLEIC Compton electron detector

- Diamond strip detector
  - At least 5 cm long
  - 200 strips
  - 4 planes
- Roman pot required
  - Need for RF and detector shielding
  - Cooling required?
  - Detector far enough from beam to use roman pot (  $\sim 1.5$  cm = 15 sigma )
  - More convenient access to detector
  - Easier placement of electronic close to detector
  - “Small profile” design needed

# Compton R&D

- Current efforts\* focused on detailed simulations related to:
  - Backgrounds, including halo and synchrotron
  - Effect of using electron detector in Roman pot
  - Backgrounds from interactions in upstream IP
- Roman pot design also under investigation – conventional design likely not practical
- In later stages, hope to build “test stand” at JLab to try out Roman pots, and also potentially different detectors
- Operation in integration mode?
- Additional R&D could be performed on the laser system
  - Investigate fast laser polarization flipping – both for single-pass and Fabry-Perot cavity



# Summary

- Excellent progress in design of chicane for combined Low- $Q^2$  tagger and Compton polarimeter
- Compton polarimeter design in progress, although baseline concept mature
  - Emphasis on electron detection → easiest avenue to achieve high precision
  - One-pass laser and high-gain Fabry-Perot cavity laser solutions both look feasible – choice will be dictated by need for “fast” measurements
- Baseline design will include “simple” photon detection scheme
  - More complex photon schemes (pair spectrometer?) also possible

# Low $Q^2$ tagger/Compton Design team

- JLab:
  - Fanglei Lin, Vasiliy Morozov, Alexandre Camsonne, Pawel Nadel-Turonski, Dave Gaskell, Kijun Park
- SLAC:
  - Mike Sullivan
- Duke:
  - Zhiwen Zhao
- ODU:
  - Charles Hyde
- U. Manitoba
  - Juliette Mammei, Josh Hoskins

# Extra

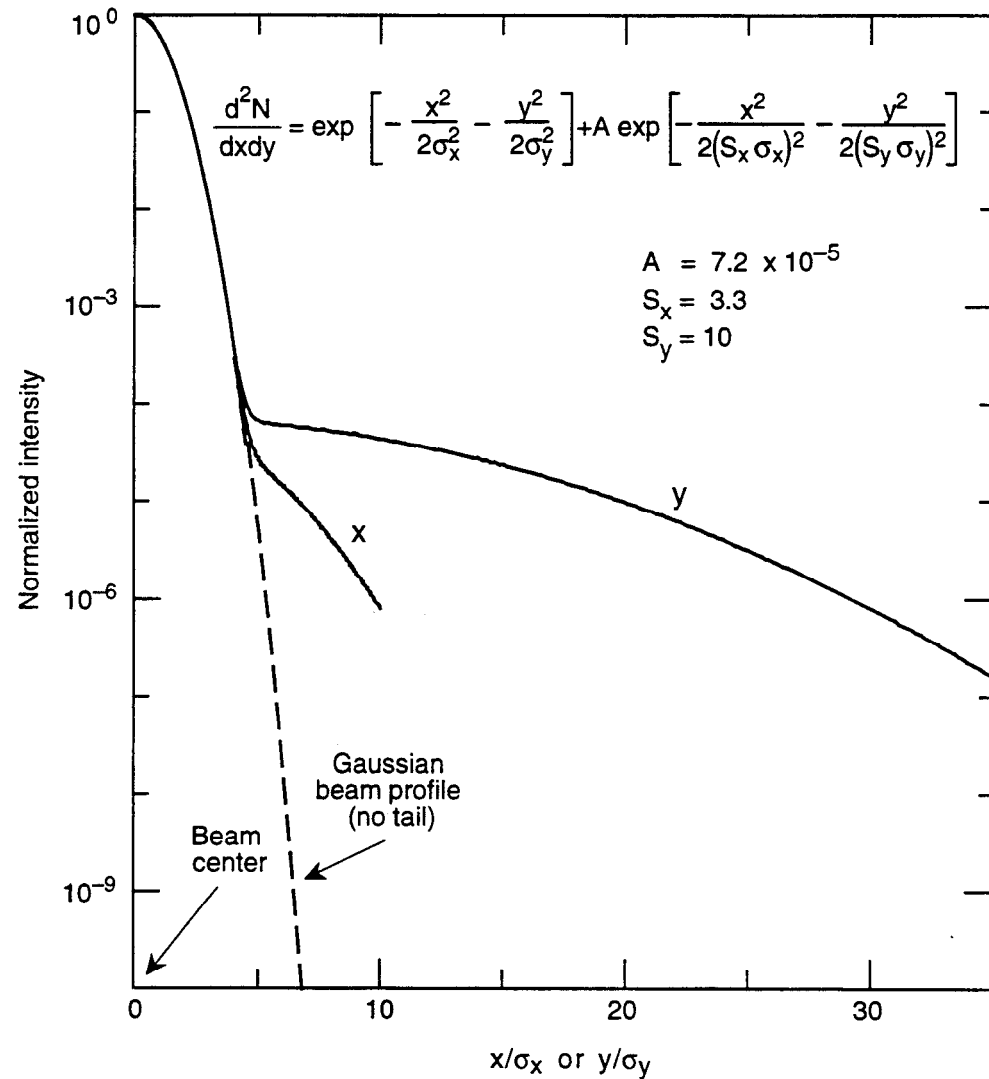
# Simulations - Halo

GEANT3 simulation uses description of beam halo from PEP-II design report (SLAC-R-418 p. 113)

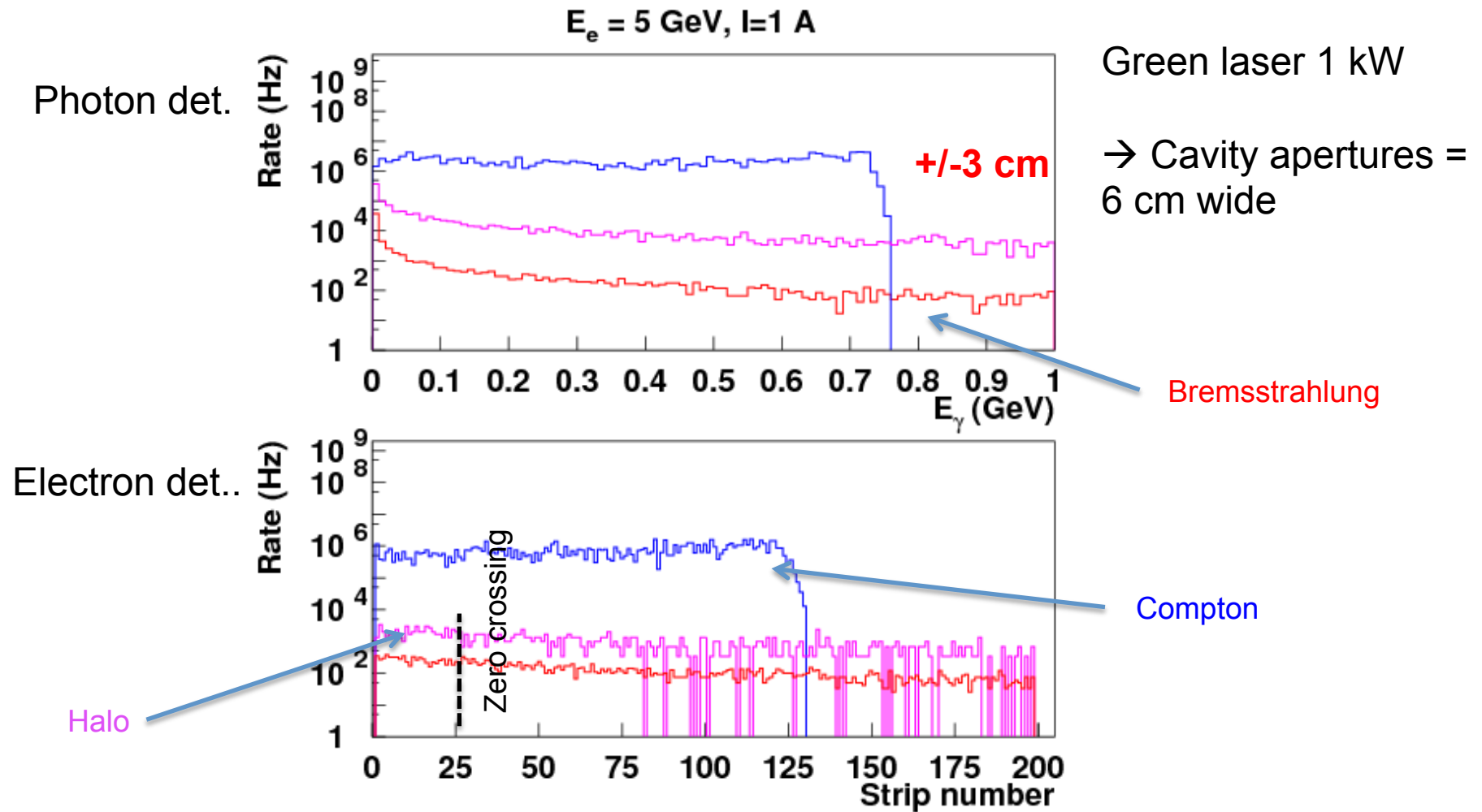
Halo flux is about 0.25% of total beam flux

Backgrounds due to halo can contribute in 2 locations

1. Direct strike of electron detector
2. Interactions with FP cavity apertures

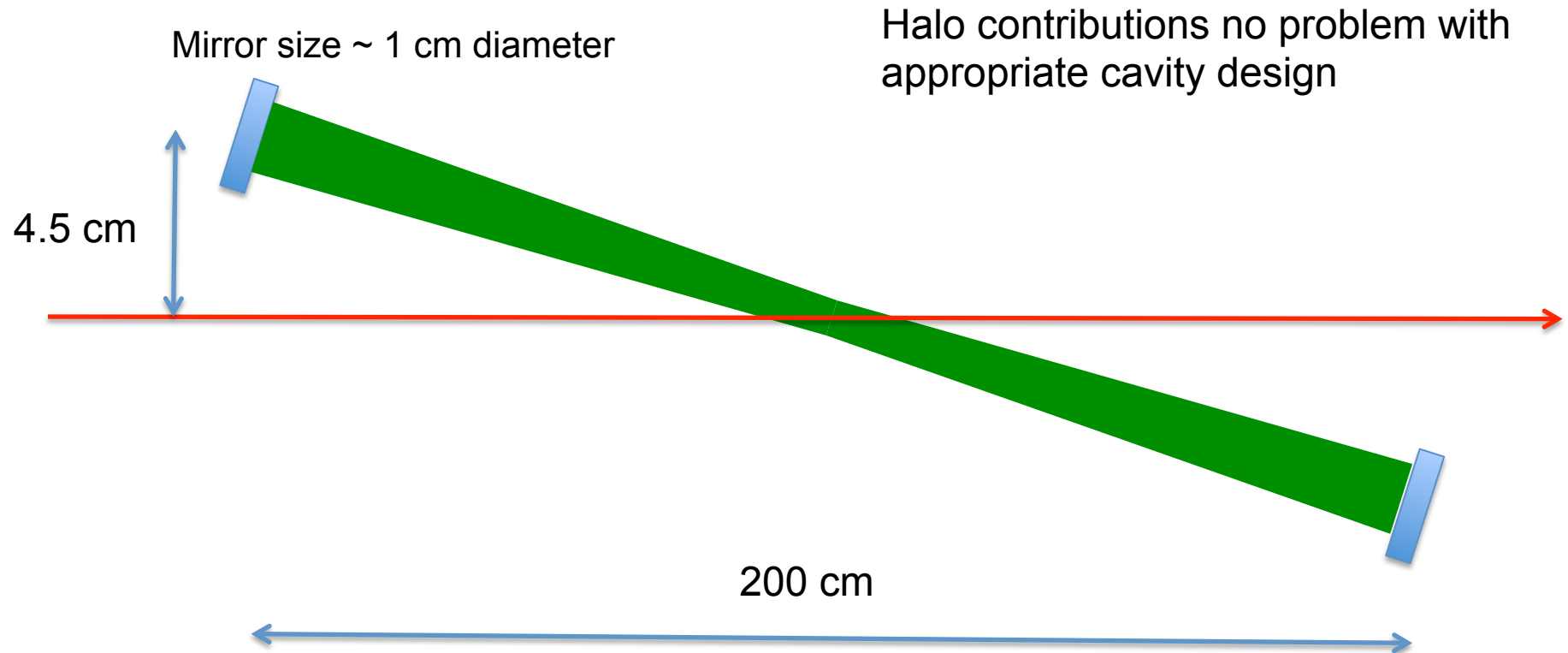


# Laser and Backgrounds - Halo



Compton edge 4 cm from beam, zero crossing = 2 cm from beam

# Fabry-Perot Cavity Design

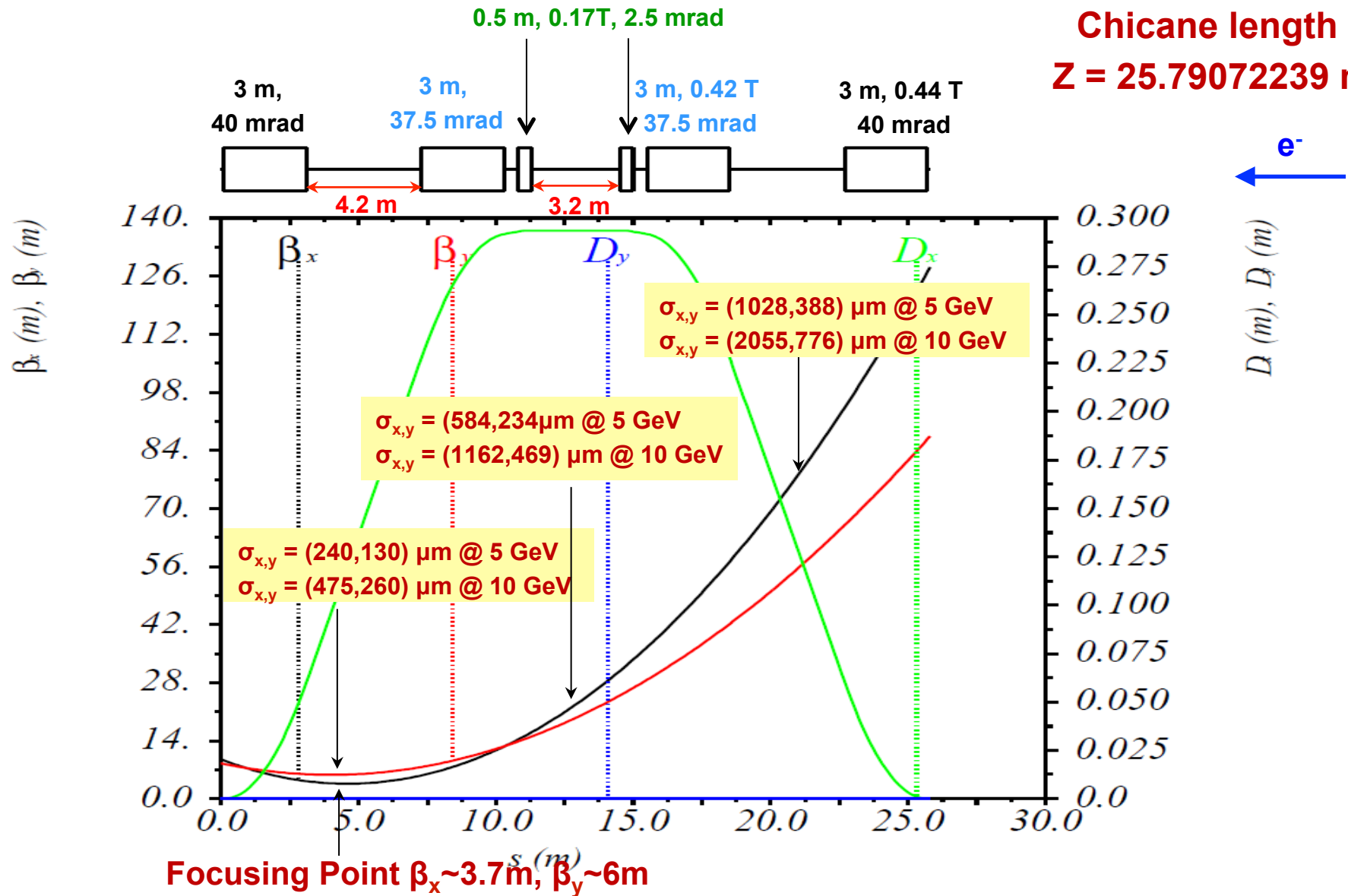


Electron-laser crossing angle = 2.58 degrees  
Mirror radius of curvature = 120 cm  
Laser size at cavity center  $(\sigma_x, \sigma_y) = 151.4 \text{ } \mu\text{m}$

Cavity gains of 1000-5000 easily achievable

# Chicane Design (baseline)

Chicane length  
 **$Z = 25.79072239 \text{ m}$**

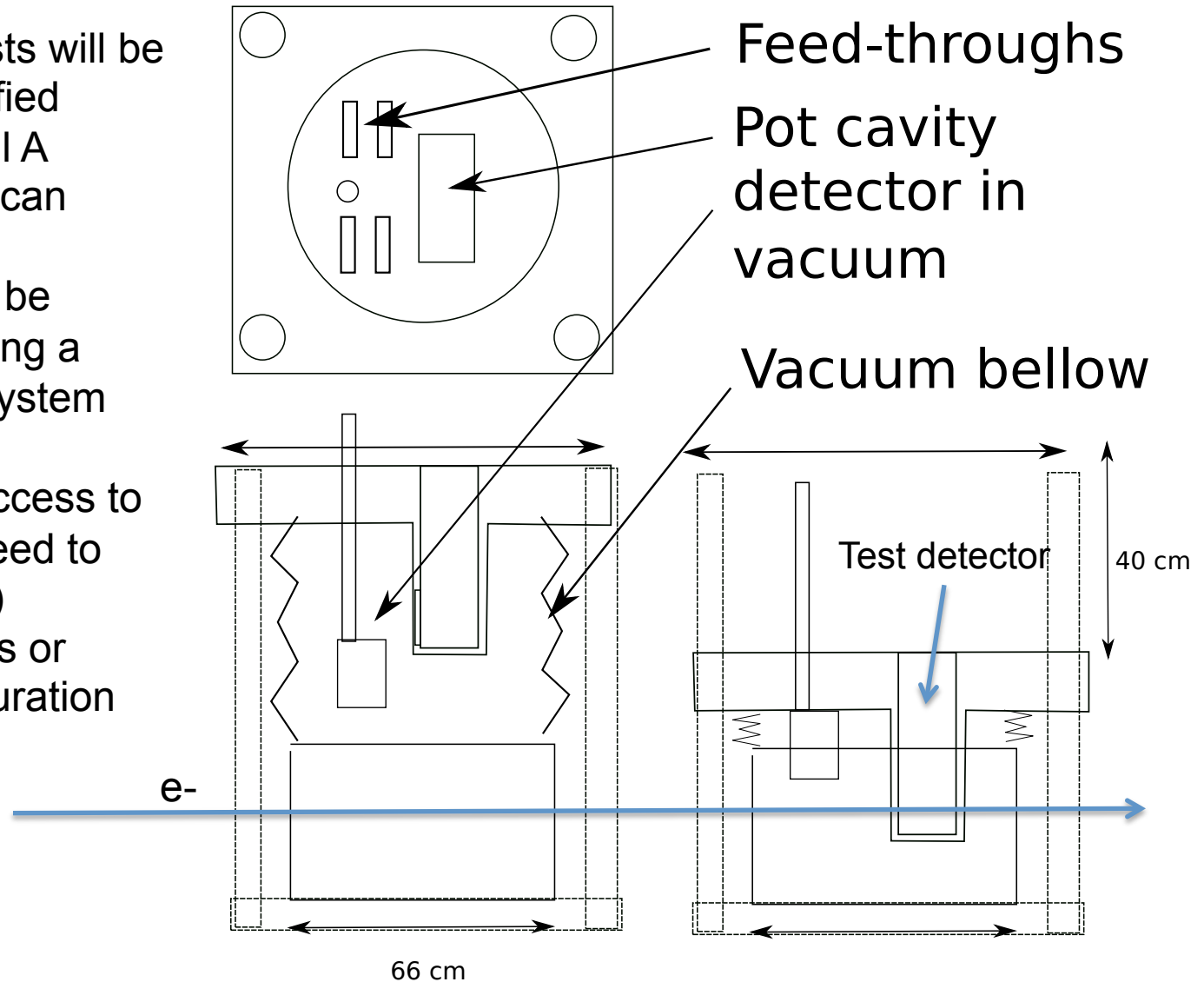


# Roman Pot

Initial detector tests will be done with a modified version of the Hall A electron detector can

Later tests would be facilitated by adding a Roman Pot-like system

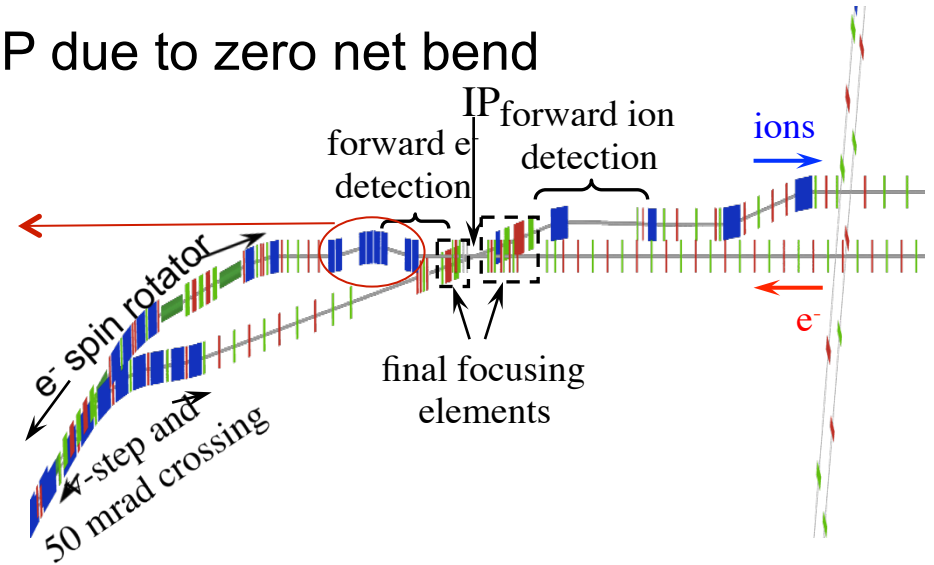
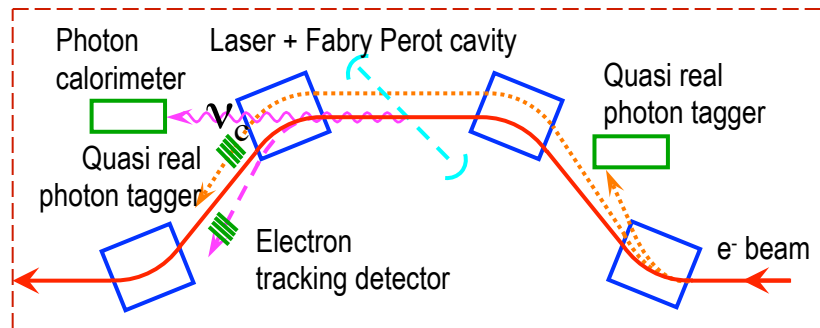
- Allow easier access to detector (no need to break vacuum)
- Swap detectors or change configuration rapidly



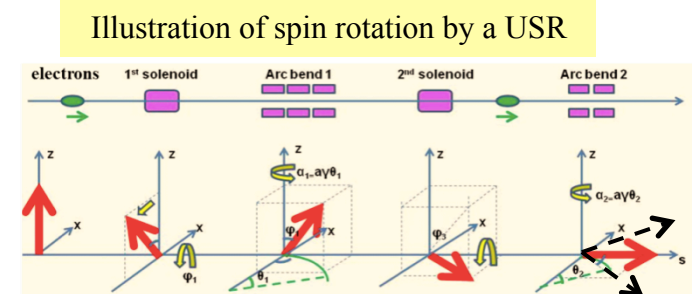
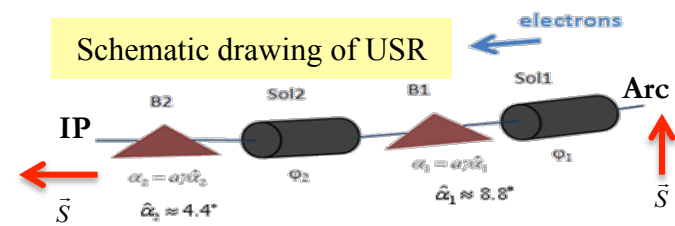
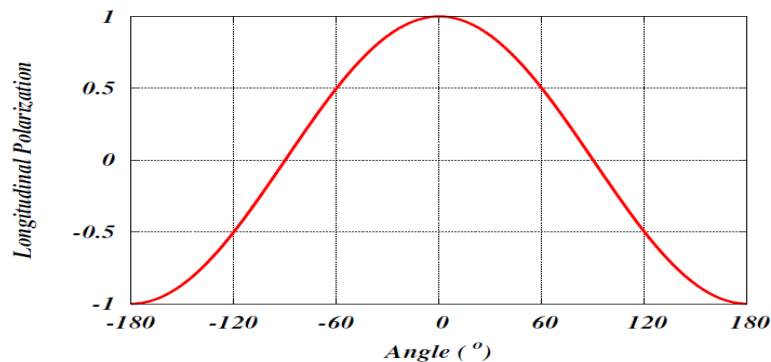


# Polarization Measurement

- Compton polarimetry:
  - same polarization at laser as at IP due to zero net bend



- Spin dancing (using spin rotators):
  - Experimentally optimize (calibrate) longitudinal polarization at IP



# Hall C Compton Electron Detector

Diamond microstrips used to detect scattered electrons

- Radiation hard: exposed to 10 MRad without significant signal degradation
- Four 21mm x 21mm planes each with 96 horizontal 200  $\mu\text{m}$  wide microstrips.
- Rough-tracking based/coincidence trigger suppresses backgrounds

