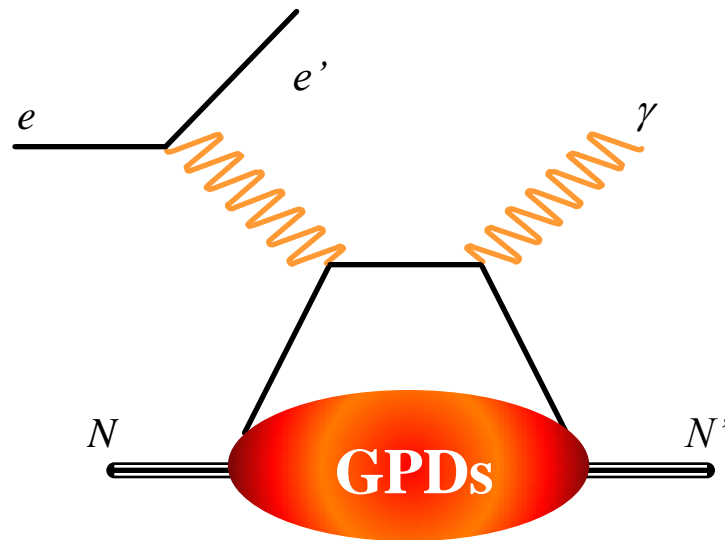


Experimental overview of DVCS and GPDs at Jefferson Lab

*Silvia Niccolai, IPN Orsay
for the CLAS Collaboration*

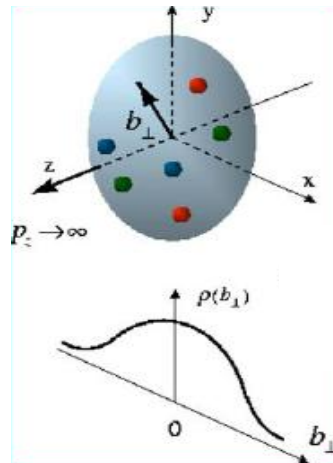
DIS16 – Hamburg (Germany) 13/4/2016



- Interest of GPDs
- GPDs and Deeply Virtual Compton Scattering
 - New DVCS results from Jefferson Lab
 - The JLab 12 GeV upgrade
 - Future JLab experiments on DVCS

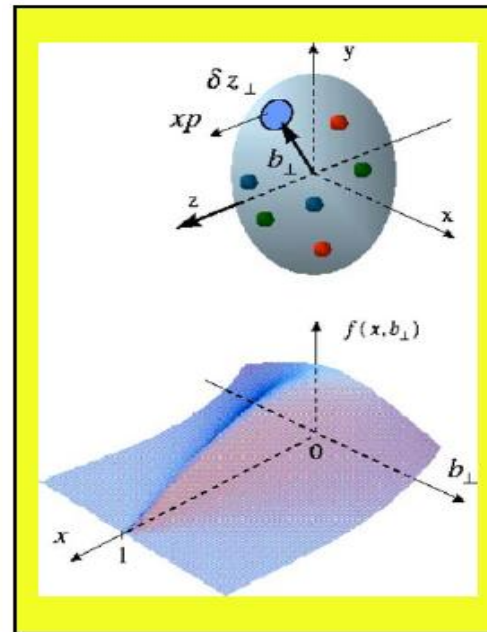


Electron-proton scattering to study nucleon structure



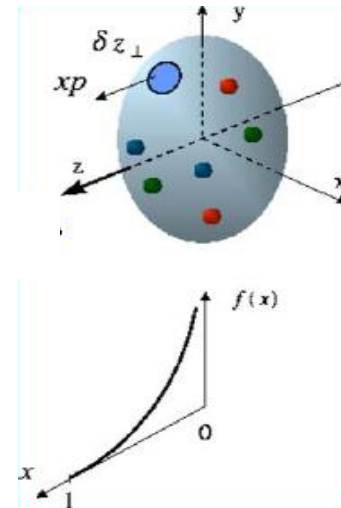
Form factors:
transverse quark
distribution in
 coordinate space

GPDs: $H, E, \tilde{H}, \tilde{E}$
Fully correlated quark
distributions in both coordinate
and momentum space



Accessible in
hard exclusive processes

High Q^2 **Final state known**



Parton distributions:
longitudinal
quark distribution
 in momentum space

The diagram shows an incoming electron (wavy line) and an incoming proton (solid line) interacting via a hard vertex (green oval) to produce an outgoing electron (wavy line) and an outgoing proton (solid line). The vertex is labeled with x .

$$\int H(x, \xi, t) dx = F_1(t) \quad (\forall \xi)$$

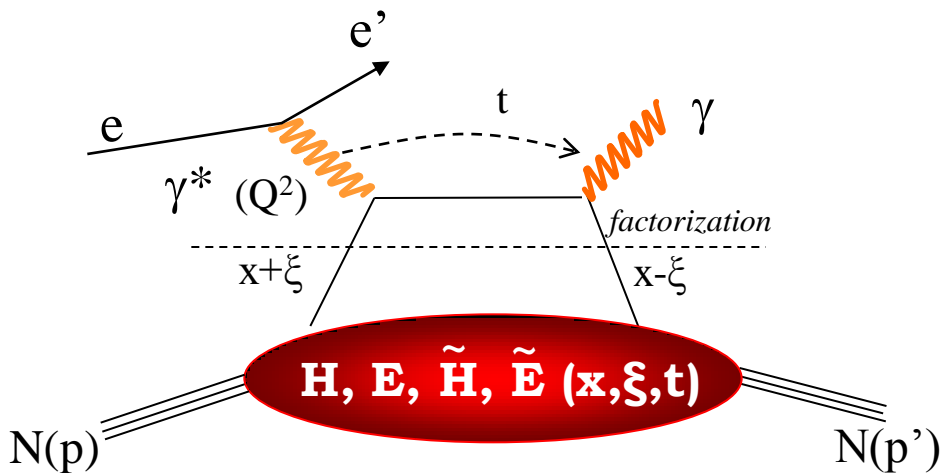
$$\int E(x, \xi, t) dx = F_2(t) \quad (\forall \xi)$$

The diagram shows an incoming electron (wavy line) and an incoming proton (solid line) interacting via a hard vertex (green oval) to produce an outgoing electron (wavy line) and an outgoing proton (solid line). The vertex is labeled with x .

$$H(x, 0, 0) = q(x),$$

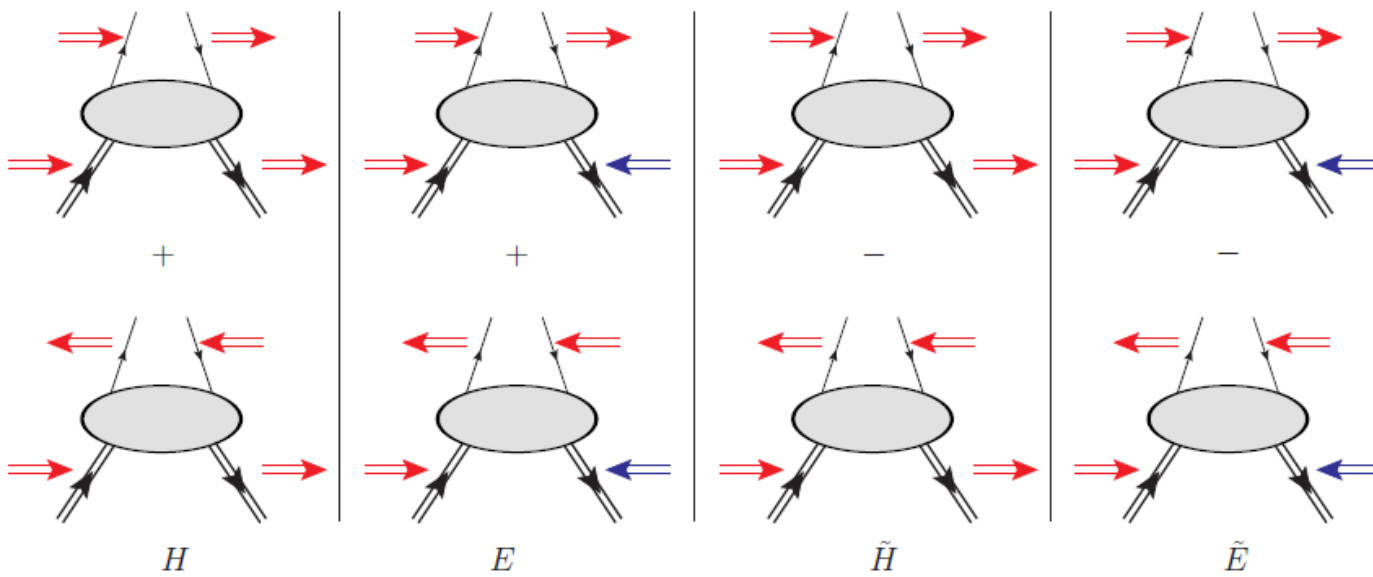
$$\tilde{H}(x, 0, 0) = \Delta q(x)$$

Deeply Virtual Compton Scattering and GPDs



- $Q^2 = -(e-e')^2$
- $x_B = Q^2/2M\nu$ $\nu = E_e - E_{e'}$
- $x+\xi, x-\xi$ longitudinal momentum fractions
- $t = \Delta^2 = (p-p')^2$
- $\xi \cong x_B/(2-x_B)$

« Handbag » factorization valid in the **Bjorken regime** (high Q^2 and ν , fixed x_B), $t \ll Q^2$



conserve nucleon spin

Vector: $H(x, \xi, t)$ Axial-Vector: $\tilde{H}(x, \xi, t)$

flip nucleon spin

Tensor: $E(x, \xi, t)$ Pseudoscalar: $\tilde{E}(x, \xi, t)$

GPDs: Fourier transforms of **non-local, non-diagonal QCD operators**

At leading order QCD, twist 2, chiral-even (quark helicity is conserved), quark sector
→ 4 GPDs for each quark flavor

Properties and “virtues” of GPDs

$$\left. \begin{aligned} \int H(x, \xi, t) dx &= F_1(t) \quad \forall \xi \\ \int E(x, \xi, t) dx &= F_2(t) \quad \forall \xi \\ \int \tilde{H}(x, \xi, t) dx &= G_A(t) \quad \forall \xi \\ \int \tilde{E}(x, \xi, t) dx &= G_P(t) \quad \forall \xi \end{aligned} \right\} \text{Link with FFs}$$

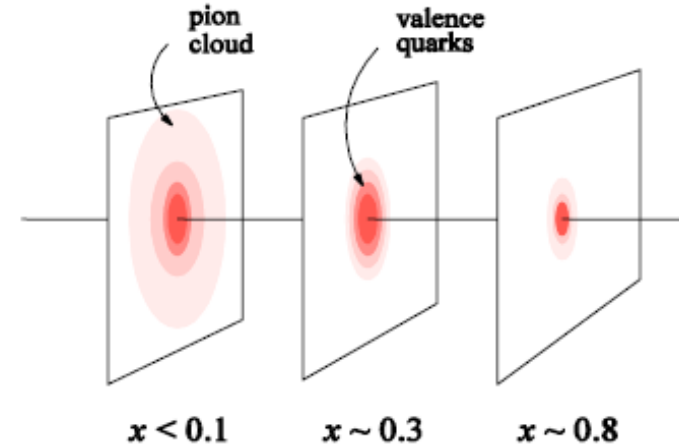
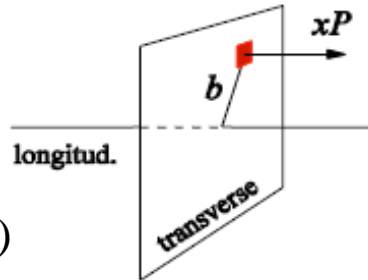
$$\left. \begin{aligned} H(x, 0, 0) &= q(x) \\ \tilde{H}(x, 0, 0) &= \Delta q(x) \end{aligned} \right\} \text{Forward limit: PDFs (not for E, } \tilde{E})$$

Nucleon tomography

$$q(x, b_{\perp}) = \int_0^{\infty} \frac{d^2 \Delta_{\perp}}{(2\pi)^2} e^{i\Delta_{\perp} b_{\perp}} H(x, 0, -\Delta_{\perp}^2)$$

$$\Delta q(x, b_{\perp}) = \int_0^{\infty} \frac{d^2 \Delta_{\perp}}{(2\pi)^2} e^{i\Delta_{\perp} b_{\perp}} \tilde{H}(x, 0, -\Delta_{\perp}^2)$$

M. Burkardt, PRD 62, 71503 (2000)



Quark angular momentum (Ji's sum rule)

$$\frac{1}{2} \int_{-1}^1 x dx (H(x, \xi, t=0) + E(x, \xi, t=0)) = J = \frac{1}{2} \Delta \Sigma + \Delta L$$

X. Ji, Phy.Rev.Lett.78,610(1997)

$$\text{Nucleon spin: } \frac{1}{2} = \underbrace{\frac{1}{2} \Delta \Sigma + \Delta L}_{\mathbf{J}} + \Delta G$$

Intrinsic spin of the quarks $\Delta \Sigma \approx 30\%$

Intrinsic spin on the gluons $\Delta G \approx 0\text{-}20\%$

Orbital angular momentum of the quarks ΔL ?

Accessing GPDs through DVCS

DVCS allows access to 4 complex GPDs-related quantities: **Compton Form Factors (ξ, t)**

$$T^{DVCS} \sim P \int_{-1}^{+1} \frac{GPDs(x, \xi, t)}{x \pm \xi} dx \pm i\pi GPDs(\pm \xi, \xi, t) + \dots$$

Only ξ and t are accessible experimentally

$$Re\mathcal{H}_q = e_q^2 P \int_0^{+1} \left(H^q(x, \xi, t) - H^q(-x, \xi, t) \right) \left[\frac{1}{\xi - x} + \frac{1}{\xi + x} \right] dx$$

$$Im\mathcal{H}_q = \pi e_q^2 \left[H^q(\xi, \xi, t) - H^q(-\xi, \xi, t) \right]$$

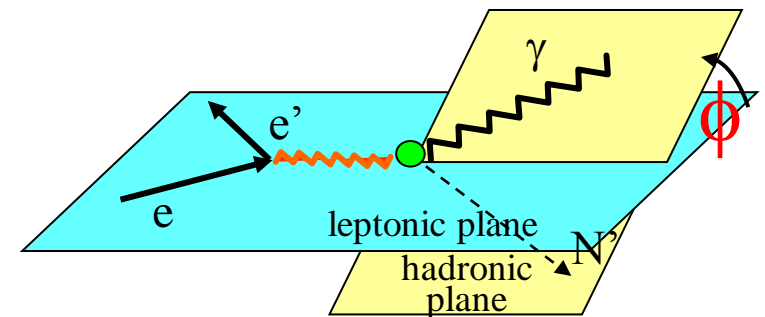
$$\sigma(eN \rightarrow eN\gamma) = \left| \begin{array}{c} \text{DVCS} \\ \text{Bethe-Heitler (BH)} \end{array} \right|^2$$

BH is calculable
(electromagnetic FFs)

$$\sigma \sim |T^{DVCS} + T^{BH}|^2 \rightarrow \text{Re}(CFFs) \quad (\text{also DSA})$$

$$\Delta\sigma = \sigma^+ - \sigma^- \propto I(DVCS \cdot BH) \rightarrow \text{Im}(CFFs)$$

$$A = \frac{\Delta\sigma}{2\sigma} \propto \frac{I(DVCS \cdot BH)}{|BH|^2 + |DVCS|^2 + I}$$



Sensitivity to CFFs of DVCS spin observables

$$A_{LU(UL)} = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} \propto \frac{s_{1,unp(UL)}^I \sin \phi}{c_{0,unp}^{BH} + c_{0,unp}^I + (c_{1,unp}^{BH} + c_{1,unp}^I) \cos \phi}$$

$$A_{LL} = \frac{\sigma^{++} + \sigma^{+-} - \sigma^{-+} - \sigma^{--}}{\sigma^{++} + \sigma^{+-} + \sigma^{-+} + \sigma^{--}} \propto \frac{c_{0,LP}^{BH} + c_{0,LP}^I + (c_{1,LP}^{BH} + c_{1,LP}^I) \cos \phi}{c_{0,unp}^{BH} + c_{0,unp}^I + (c_{1,unp}^{BH} + c_{1,unp}^I) \cos \phi}$$

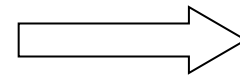
Twist-2
approximation
($-t \ll Q^2$)

$$(\xi = x_B / (2 - x_B) \quad k = -t / 4M^2)$$

Proton Neutron

Polarized beam, unpolarized target:

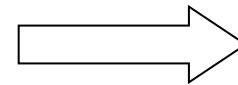
$$s_{1,unp}^I \sim \sin \phi \operatorname{Im}\{F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} - k F_2 \mathcal{E}\}$$



$$\begin{aligned} & \operatorname{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p, \mathcal{E}_p\} \\ & \operatorname{Im}\{\mathcal{H}_n, \tilde{\mathcal{H}}_n, \mathcal{E}_n\} \end{aligned}$$

Unpolarized beam, longitudinal target:

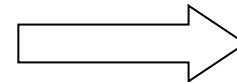
$$s_{1,UL}^I \sim \sin \phi \operatorname{Im}\{F_1 \tilde{\mathcal{H}} + \xi(F_1 + F_2)(\mathcal{H} + x_B / 2 \mathcal{E}) - \xi k F_2 \tilde{\mathcal{E}} + \dots\}$$



$$\begin{aligned} & \operatorname{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\} \\ & \operatorname{Im}\{\mathcal{H}_n, \mathcal{E}_n\} \end{aligned}$$

Polarized beam, longitudinal target:

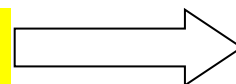
$$c_{1,LP}^I \sim (A + B \cos \phi) \operatorname{Re}\{F_1 \tilde{\mathcal{H}} + \xi(F_1 + F_2)(\mathcal{H} + x_B / 2 \mathcal{E}) \dots\}$$



$$\begin{aligned} & \operatorname{Re}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\} \\ & \operatorname{Re}\{\mathcal{H}_n, \mathcal{E}_n\} \end{aligned}$$

Unpolarized beam, transverse target:

$$\Delta \sigma_{UT} \sim \cos \phi \sin(\phi_s - \phi) \operatorname{Im}\{k(F_2 \mathcal{H} - F_1 \mathcal{E}) + \dots\}$$



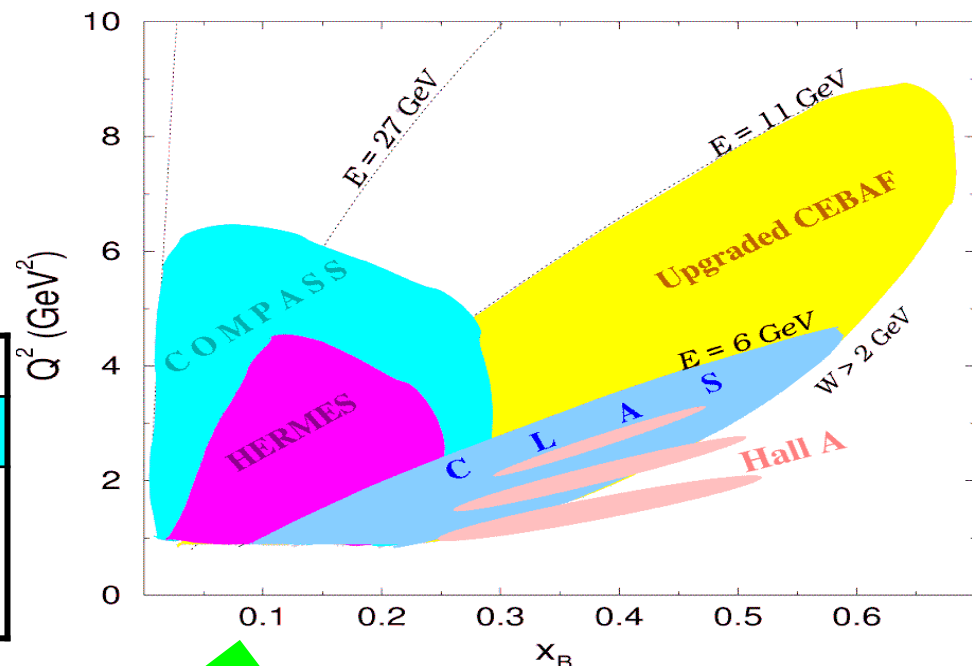
$$\begin{aligned} & \operatorname{Im}\{\mathcal{H}_p, \mathcal{E}_p\} \\ & \operatorname{Im}\{\mathcal{H}_n\} \end{aligned}$$

DVCS experiments worldwide

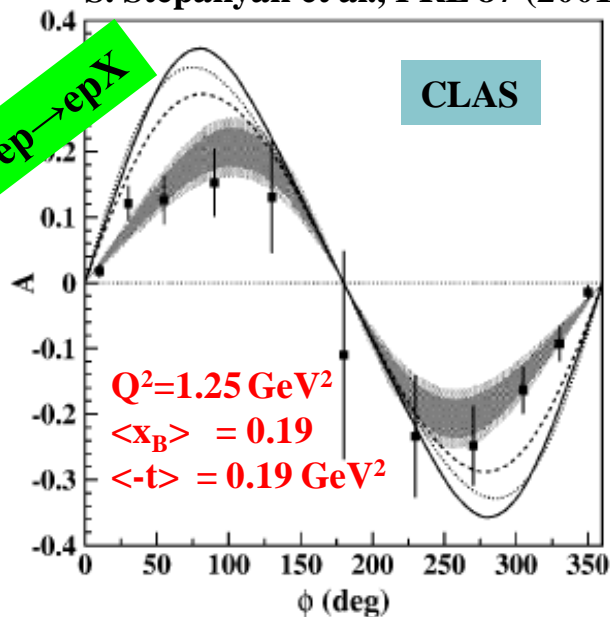
JLAB	
<i>Hall A</i>	<i>CLAS (Hall B)</i>
p,n-DVCS Beam-pol. CS	p-DVCS BSA,ITSA,DSA,CS

DESY	
<i>HERMES</i>	<i>H1/ZEUS</i>
p-DVCS BSA,BCA, tTSA,ITSA,DSA	p-DVCS CS,BCA

CERN
<i>COMPASS</i>
p-DVCS CS,BSA,BCA, tTSA,ITSA,DSA



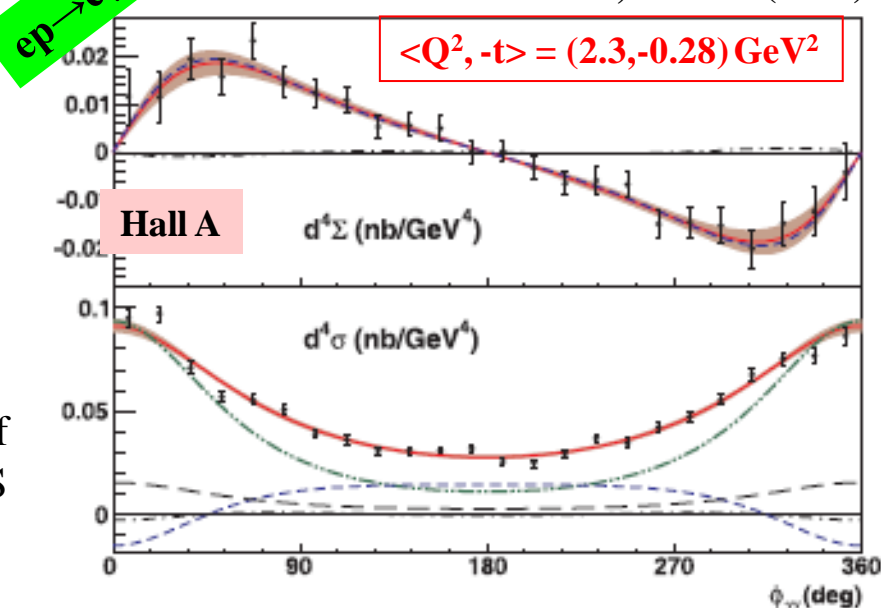
S. Stepanyan et al., PRL 87 (2001)



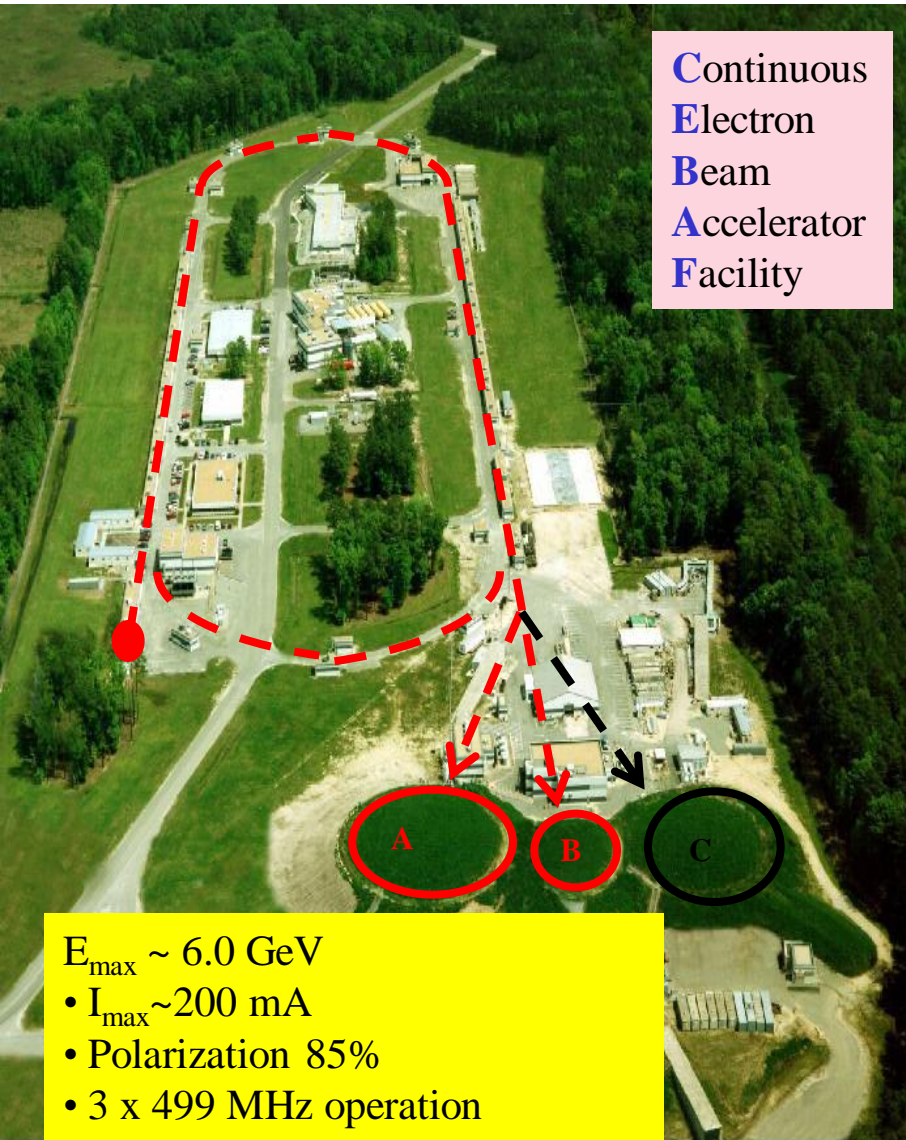
CLAS, HERMES:
first observation of
DVCS-BH
interference

Hall A: proof of
scaling for DVCS

C.M. Camacho et al., PRL 97 (2006)



JLab@6 GeV



$E_{\text{max}} \sim 6.0 \text{ GeV}$

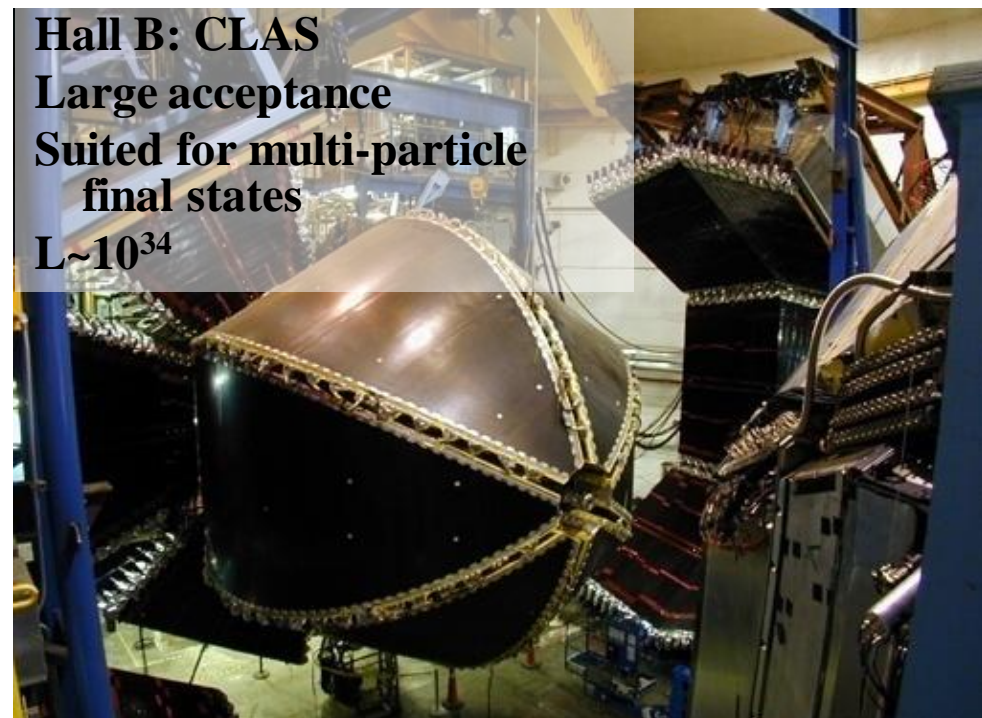
• $I_{\text{max}} \sim 200 \text{ mA}$

• Polarization 85%

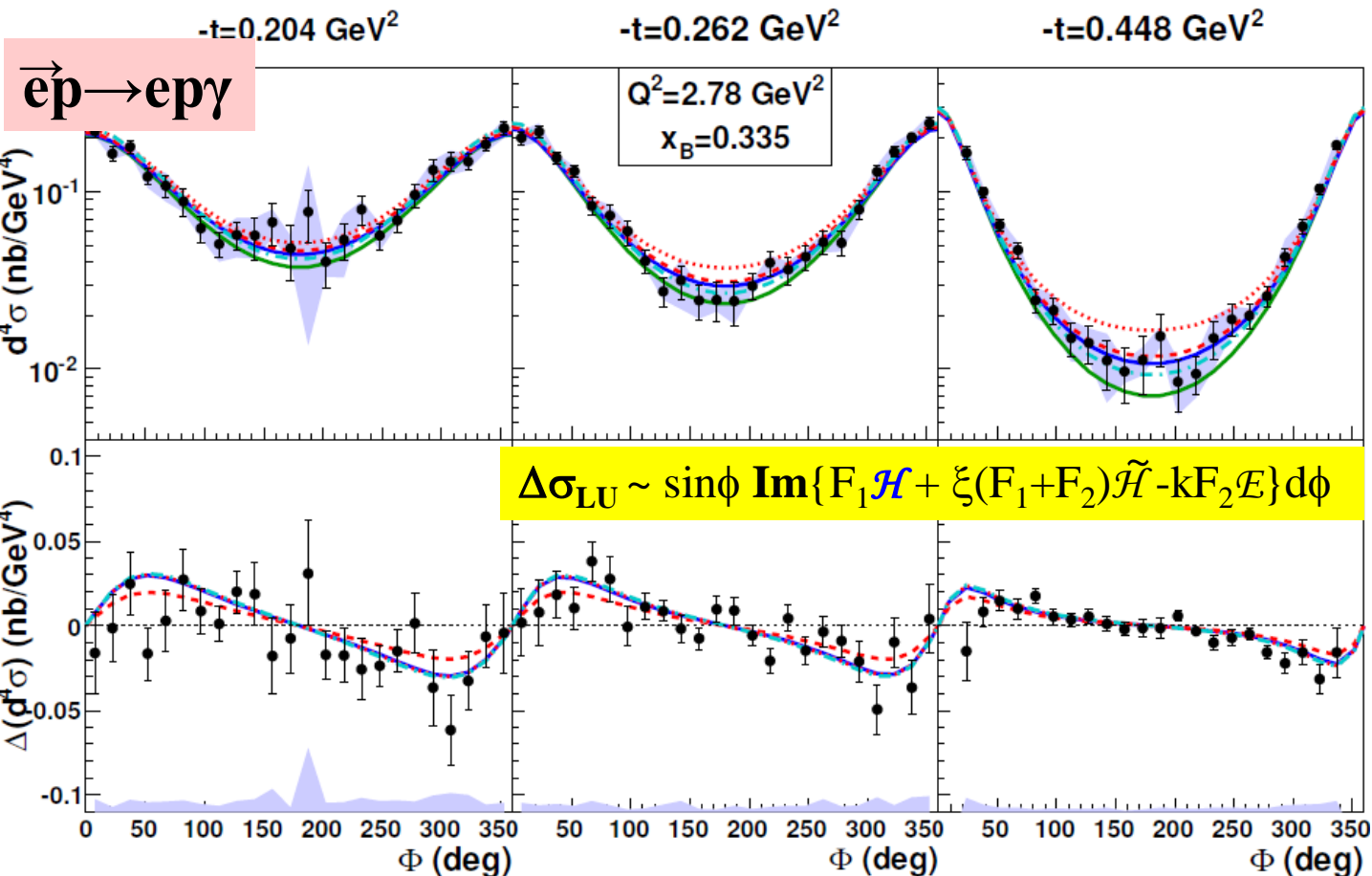
• 3 x 499 MHz operation

• Simultaneous delivery to 3 halls

• Shutdown in May 2012



CLAS: unpolarized and beam-polarized cross sections

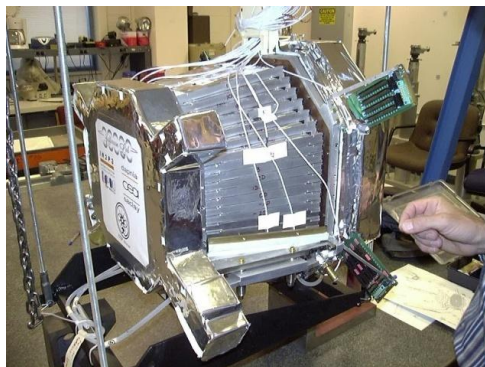


H.S. Jo et al., PRL
115, 212003 (2015)

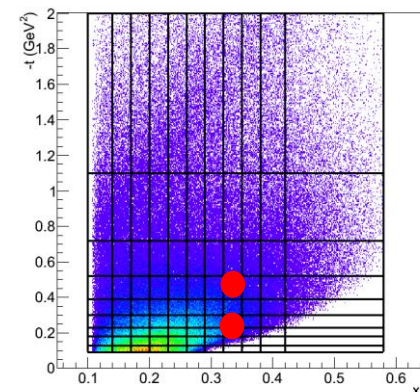
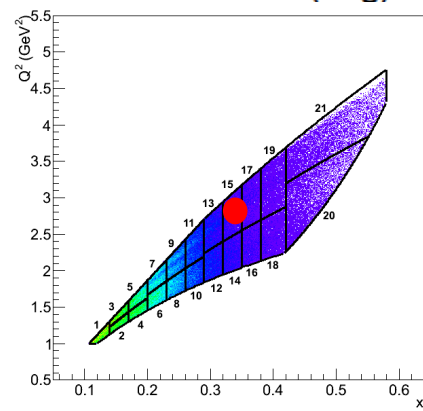
- Largest kinematic ever covered
- Two observables extracted

KM10 model fits Hall A
2006 data using
« anomalous » \tilde{H}

- Data taken in **2005**, e1-dvcs
- Beam energy $\sim 5.75 \text{ GeV}$
- Beam polarization $\sim 80\%$
- Target **LH₂**
- **Inner Calorimeter (IC)**



21 Q^2 - x_B bins, 9 t bins, 24 ϕ bins



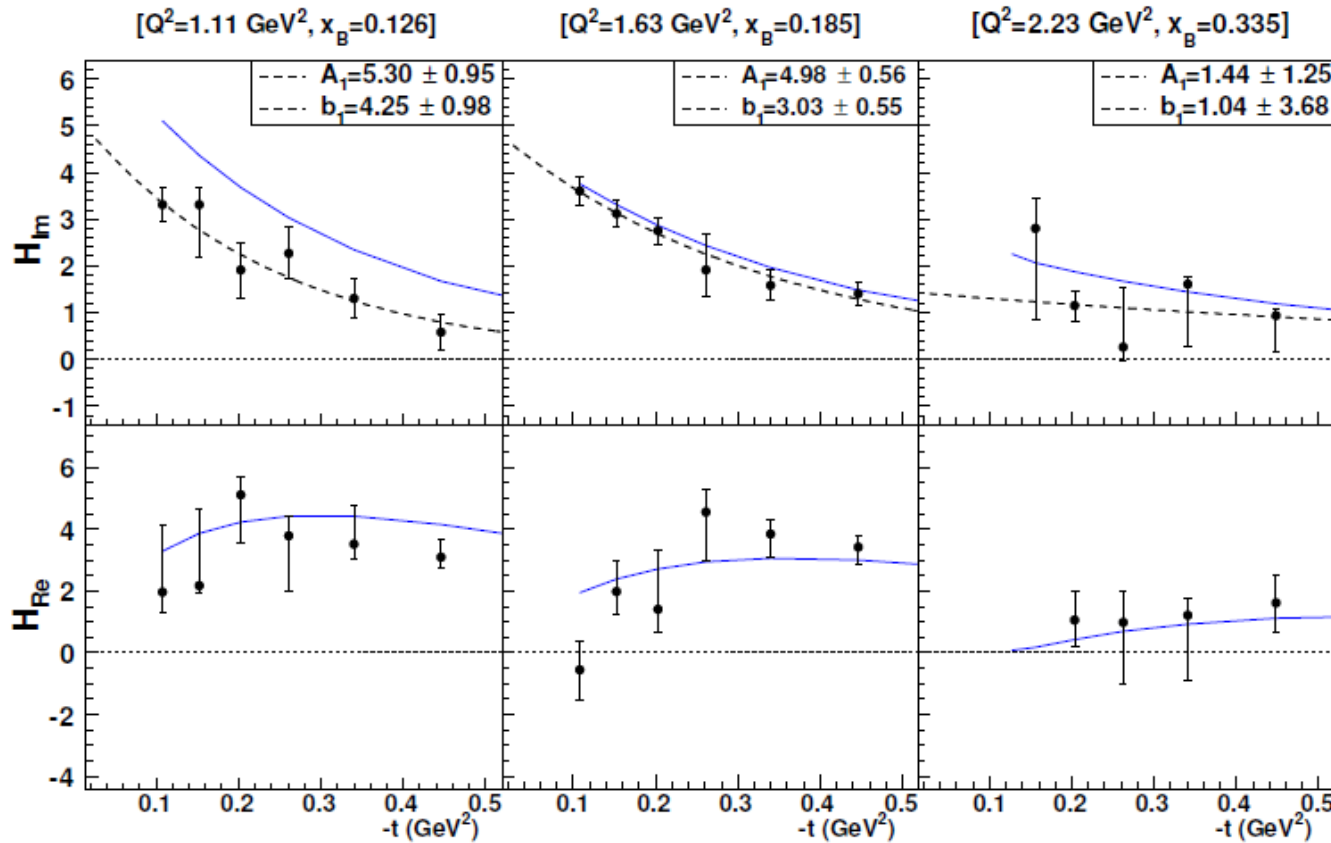
Extraction of Compton Form Factors from DVCS observables

GPDs cannot directly be extracted from DVCS observables, one can access
Compton Form Factors:

$$\begin{aligned} 8 \text{ CFF} \left\{ \begin{aligned} \text{Re}(\mathcal{H}) &= P \int_0^1 dx [H(x, \xi, t) - H(-x, \xi, t)] C^+(x, \xi) \\ \text{Re}(\mathcal{E}) &= P \int_0^1 dx [E(x, \xi, t) - E(-x, \xi, t)] C^+(x, \xi) \\ \text{Re}(\tilde{\mathcal{H}}) &= P \int_0^1 dx [\tilde{H}(x, \xi, t) + \tilde{H}(-x, \xi, t)] C^-(x, \xi) \\ \text{Re}(\tilde{\mathcal{E}}) &= P \int_0^1 dx [\tilde{E}(x, \xi, t) + \tilde{E}(-x, \xi, t)] C^-(x, \xi) \\ \text{Im}(\mathcal{H}) &= H(\xi, \xi, t) - H(-\xi, \xi, t) \\ \text{Im}(\mathcal{E}) &= E(\xi, \xi, t) - E(-\xi, \xi, t) \\ \text{Im}(\tilde{\mathcal{H}}) &= \tilde{H}(\xi, \xi, t) - \tilde{H}(-\xi, \xi, t) \\ \text{Im}(\tilde{\mathcal{E}}) &= \tilde{E}(\xi, \xi, t) - \tilde{E}(-\xi, \xi, t) \end{aligned} \right. \\ \text{with } C^\pm(x, \xi) = \frac{1}{x - \xi} \pm \frac{1}{x + \xi} \end{aligned}$$

M. Guidal: Model-independent fit, at fixed Q^2 , x_B and t of DVCS observables
8 parameters (the CFFs), loosely bound (± 5 x VGG prediction)
M. Guidal, Eur. Phys. J. A 37 (2008) 319 & many other papers...

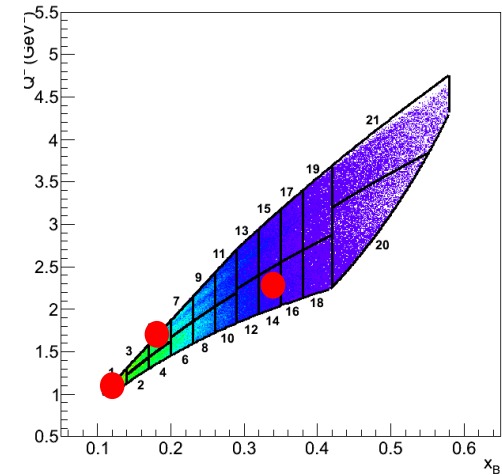
Extraction of CFFs from e1dvcs pol. and unpol. cross sections



*CFF fits by M. Guidal
(H and \tilde{H} only)*

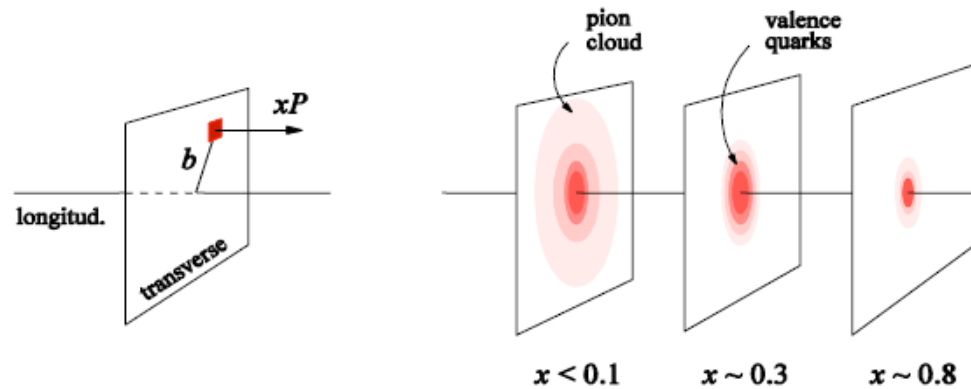
--- Ae^{-bt} fit

— VGG predictions



$$q(x, b_{\perp}) = \int_0^{\infty} \frac{d^2 \Delta_{\perp}}{(2\pi)^2} e^{i\Delta_{\perp} b_{\perp}} H(x, 0, -\Delta_{\perp}^2)$$

$Im(\mathcal{H}_p)$, flatter t slope at high x_B : faster quarks (valence) at the core of the nucleon, slower quarks (sea) at its periphery
→ **PROTON TOMOGRAPHY**

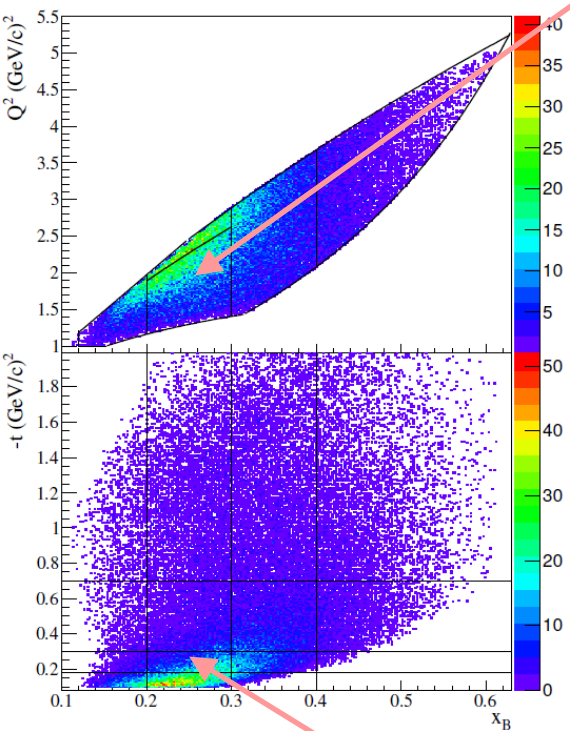


CLAS: DVCS on longitudinally polarized target

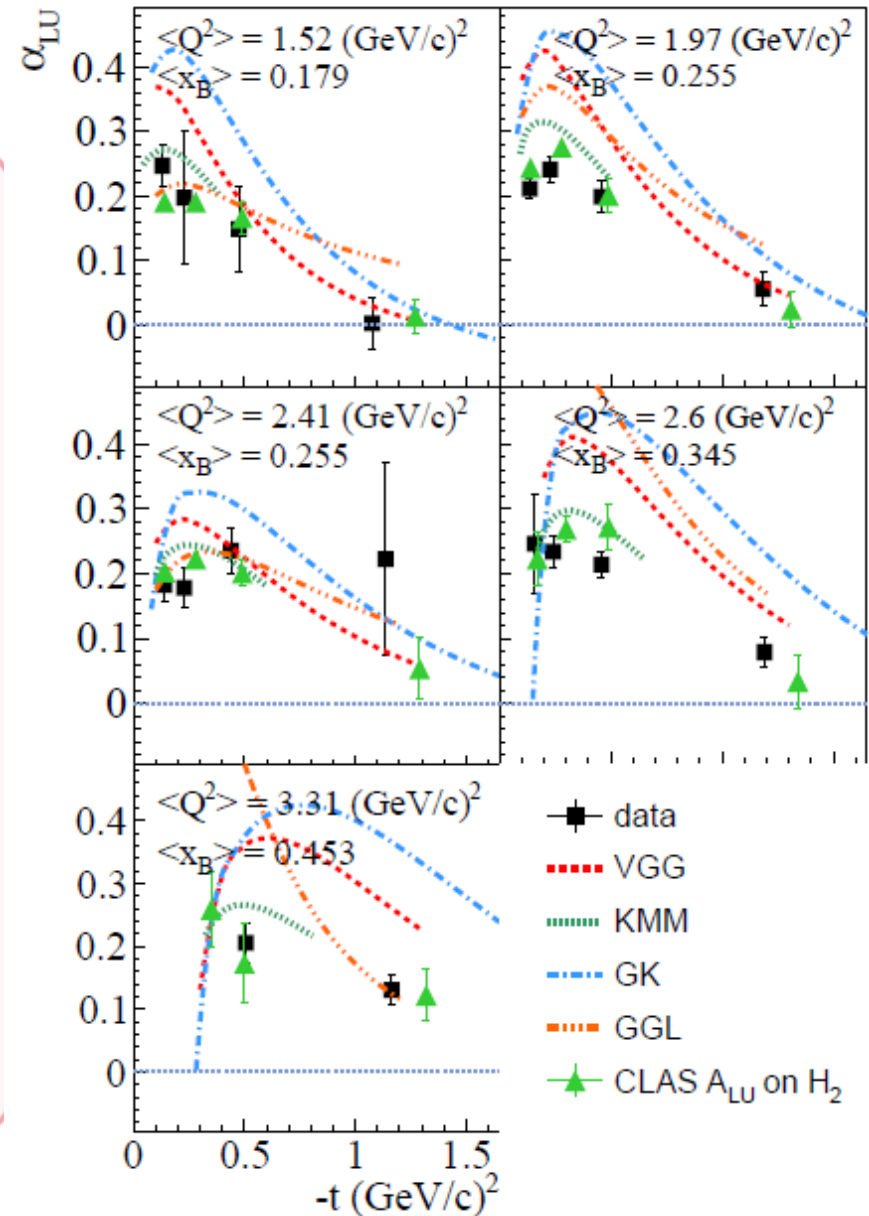
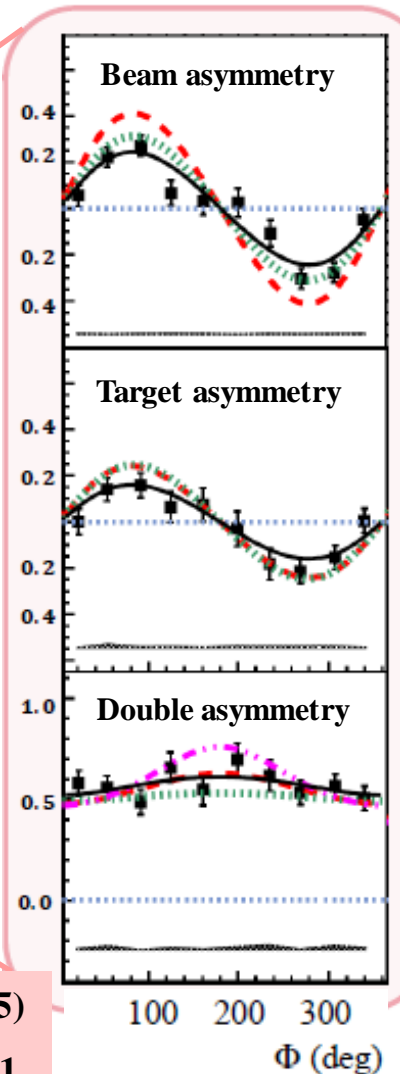
$$\vec{e}\vec{p} \rightarrow e\vec{p}\gamma$$

$$\text{BSA} \sim \text{Im}\{\mathcal{H}_p\}$$

- Data taken in **2009, eg1-dvcs**
- Beam energy **~ 5.9 GeV**
- **CLAS + IC** to detect forward photons
- Target: **longitudinally polarized NH_3** ($P \sim 80\%$)
- **3 DVCS observables**



5 Q^2 - x_B bins, 4 t bins, 10 ϕ bins



S. Pisano et al., PRD 91, 052014 (2015)

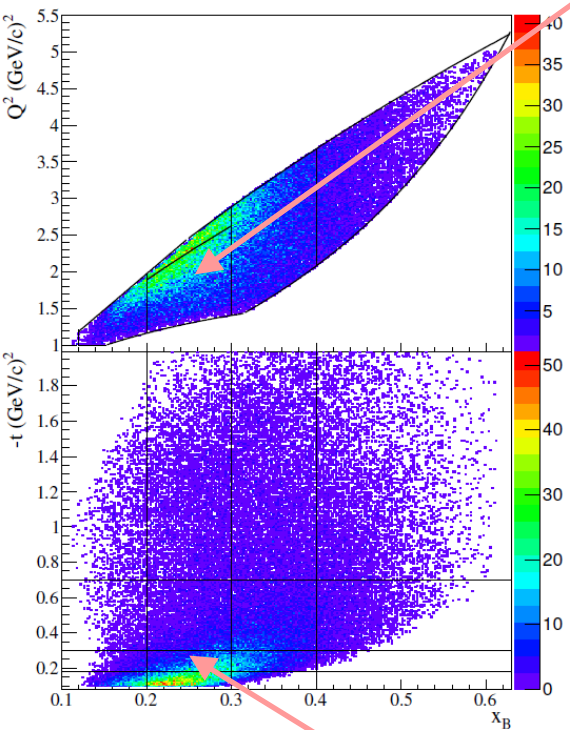
E. Seder et al., PRL 114 (2015) 032001

CLAS: DVCS on longitudinally polarized target

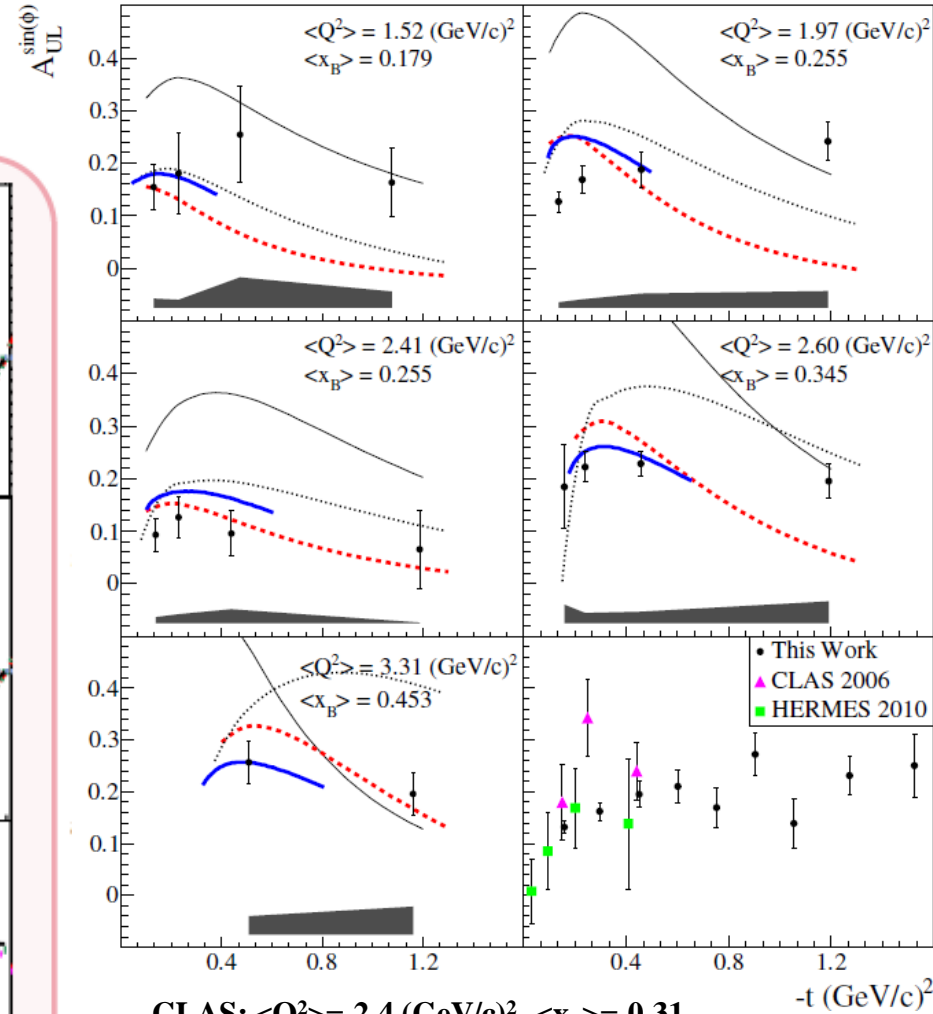
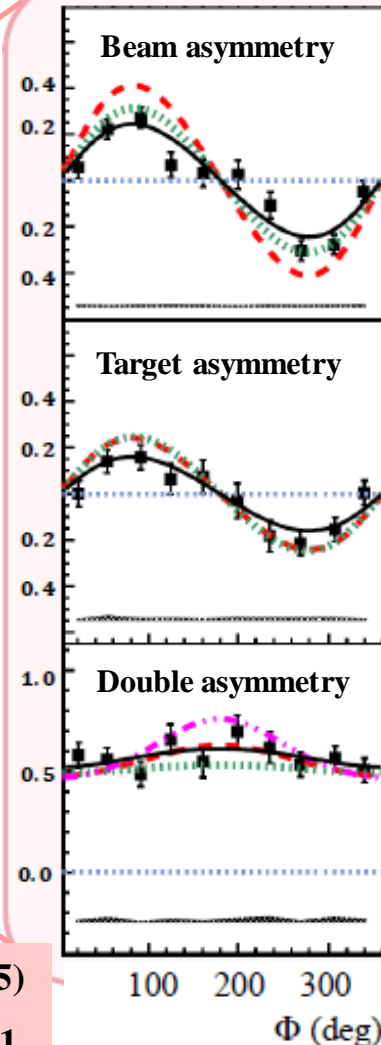
$$\vec{e}\vec{p} \rightarrow e\gamma$$

- Data taken in **2009, eg1-dvcs**
- Beam energy **~5.9 GeV**
- **CLAS + IC** to detect forward photons
- Target: **longitudinally polarized NH₃** (P~80%)
- **3 DVCS observables**

$$\text{TSA} \sim \text{Im}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$$



5 Q^2 - x_B bins, 4 t bins, 10 ϕ bins



CLAS: $\langle Q^2 \rangle = 2.4 \text{ (GeV/c)}^2$, $\langle x_B \rangle = 0.31$

HERMES: $\langle Q^2 \rangle = 2.459 \text{ (GeV/c)}^2$, $\langle x_B \rangle = 0.096$

CLAS2006: $\langle Q^2 \rangle = 1.82 \text{ (GeV/c)}^2$, $\langle x_B \rangle = 0.28$

- Improved statistics x10 at low $-t$
- Extended kinematic coverage

S. Pisano et al., PRD 91, 052014 (2015)

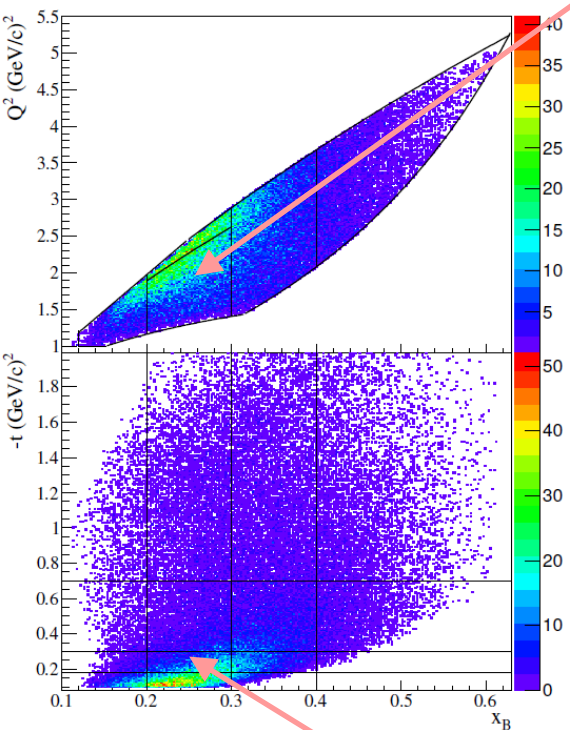
E. Seder et al., PRL 114 (2015) 032001

CLAS: DVCS on longitudinally polarized target

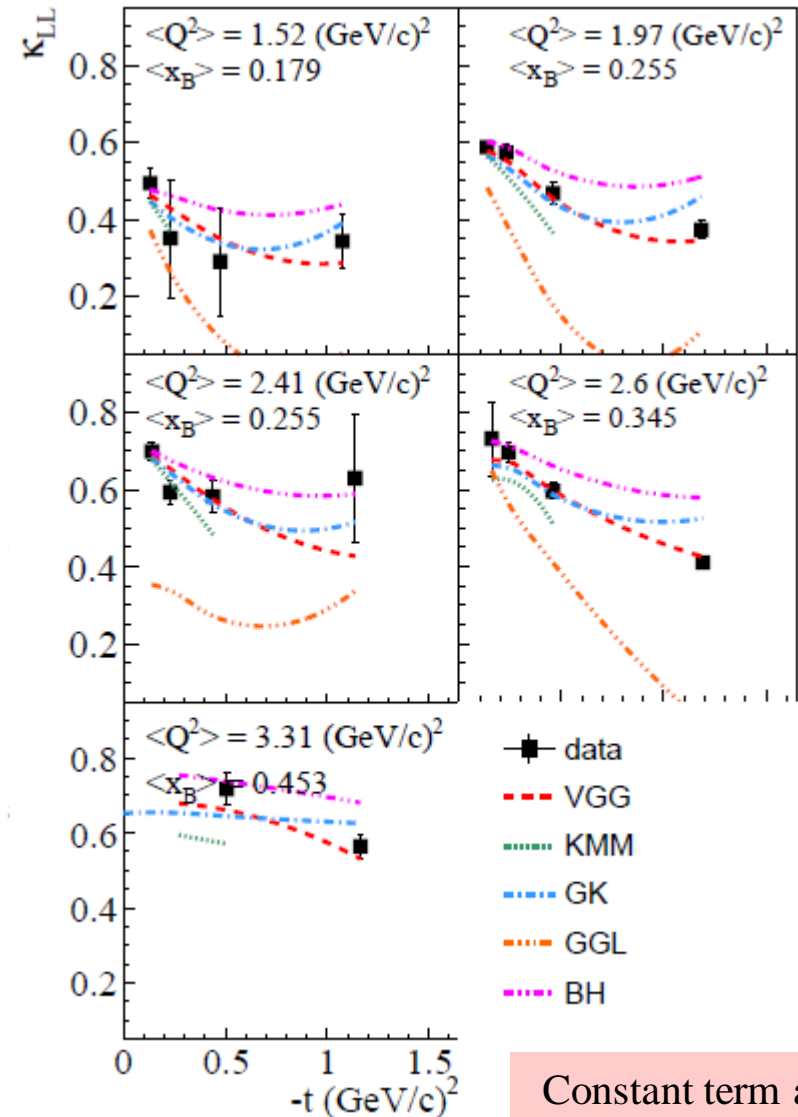
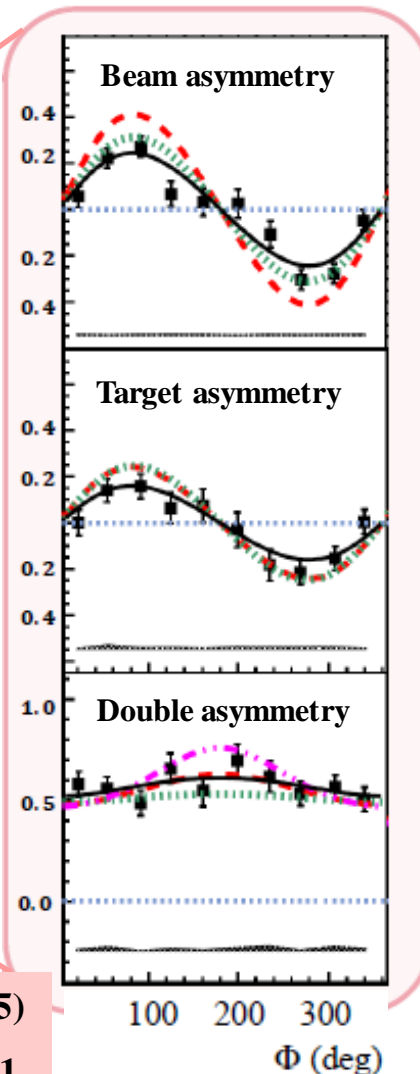
$$\vec{e}p \rightarrow e\gamma$$

- Data taken in **2009, eg1-dvcs**
- Beam energy **~5.9 GeV**
- **CLAS + IC** to detect forward photons
- Target: **longitudinally polarized NH₃** (P~80%)
- **3 DVCS observables**

$$\text{DSA} \sim \text{Re}\{\mathcal{H}_p, \tilde{\mathcal{H}}_p\}$$



5 Q^2 - x_B bins, 4 t bins, 10 ϕ bins



Constant term and $\cos\phi$ term are dominated by **BH**

S. Pisano et al., PRD 91, 052014 (2015)

E. Seder et al., PRL 114 (2015) 032001

Extraction of CFFs from eg1-dvcs TSA, BSA, DSA

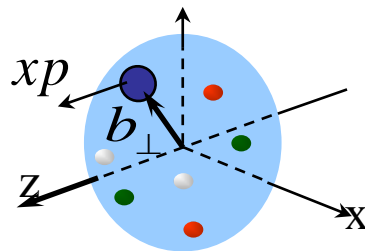
CFFs fitting code by M. Guidal
(7 CFFs included in the fit)

$Im\mathcal{H}$ has steeper t-slope than $Im\tilde{\mathcal{H}}$: is axial charge more “concentrated” than the electromagnetic charge?
→ **PROTON TOMOGRAPHY**

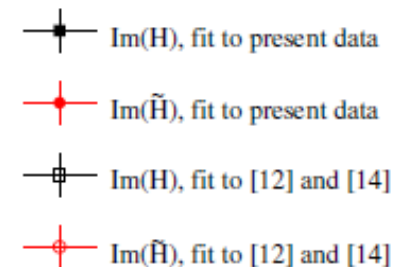
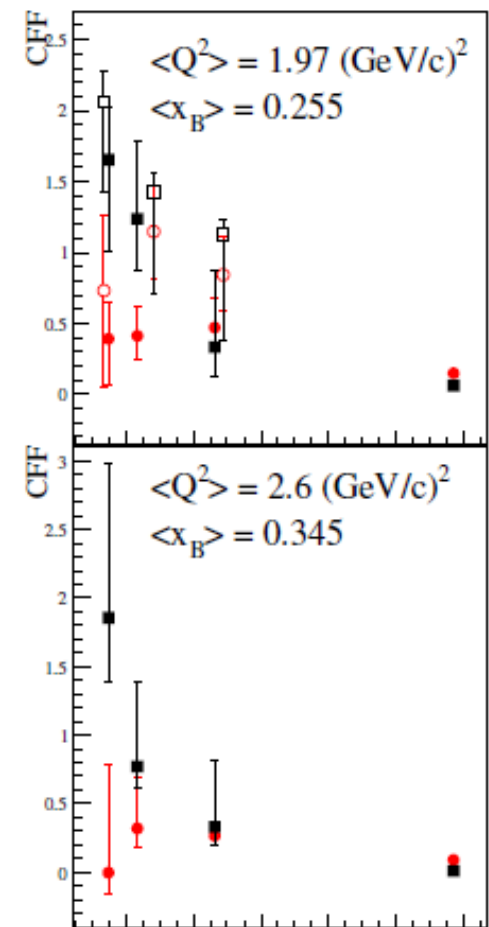
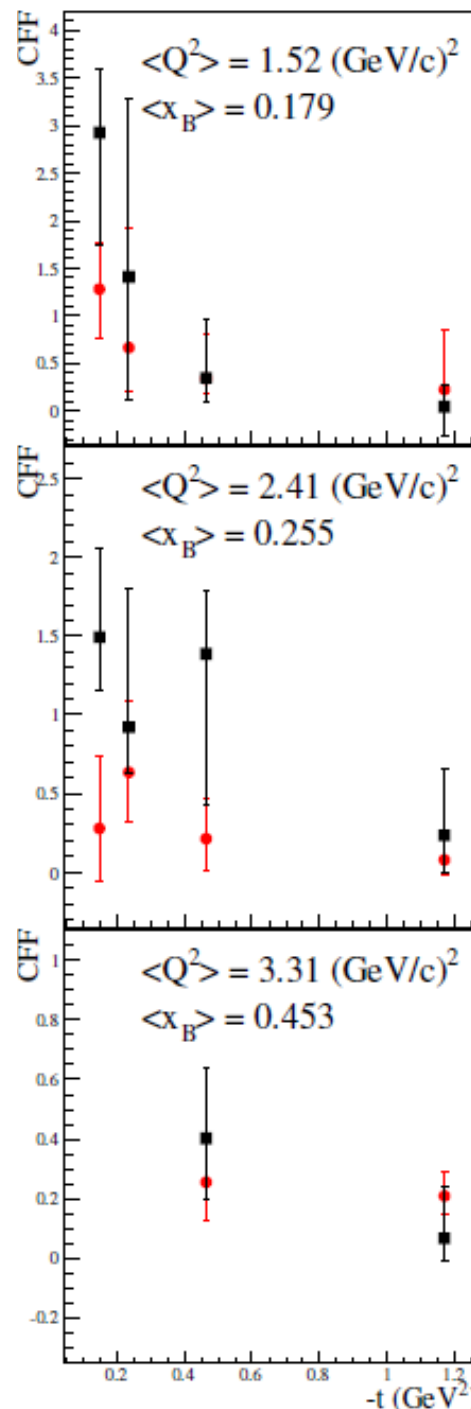
$$\Delta q(x, b_{\perp}) = \int_0^{\infty} \frac{d^2 \Delta_{\perp}}{(2\pi)^2} e^{i\Delta_{\perp} b_{\perp}} \tilde{H}(x, 0, -\Delta_{\perp}^2)$$

$$\int H(x, \xi, t) dx = F_1(t)$$

$$\int \tilde{H}(x, \xi, t) dx = G_A(t)$$



S. Pisano et al., PRD 91, 052014 (2015)



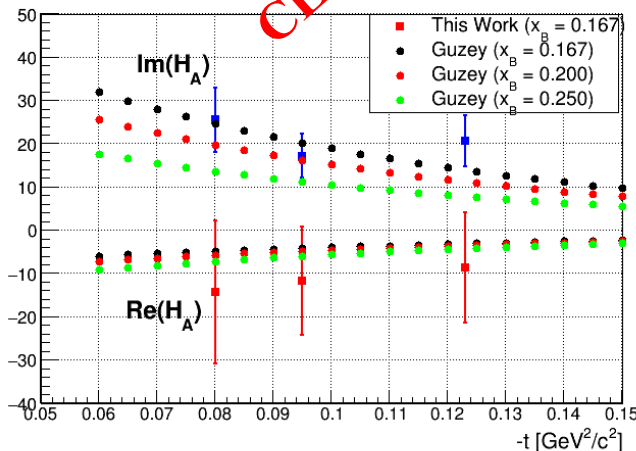
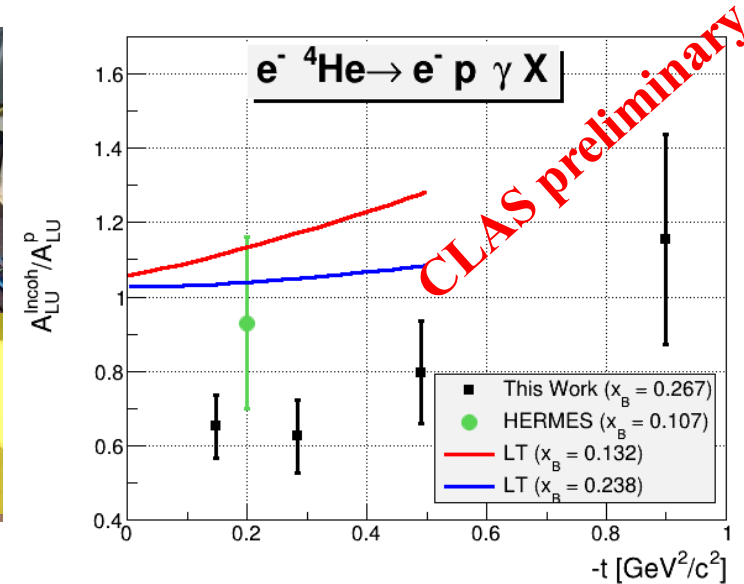
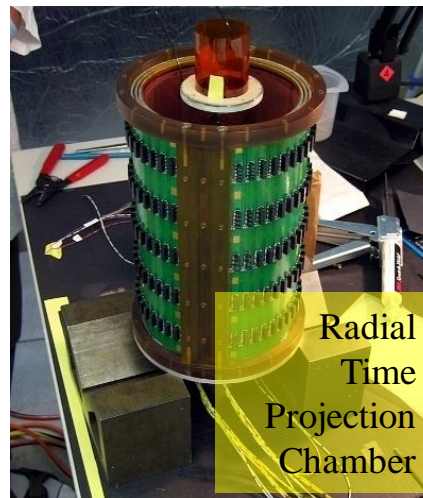
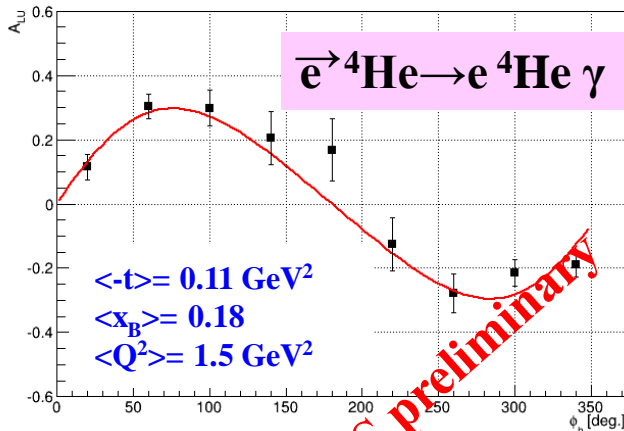
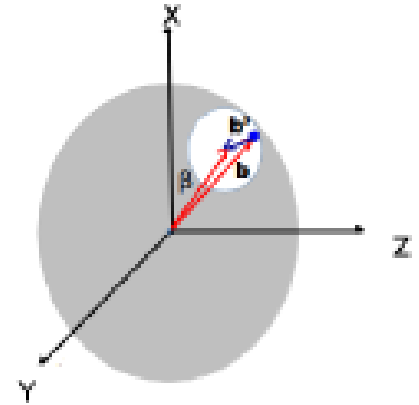
[12] CLAS BSA (Girod *et al.*)
[14] CLAS TSA (Chen *et al.*)

DVCS on nuclei: the CLAS eg6 experiment

- Data taken in the fall 2009
- CLAS+IC+**RTPC**+ ^4He target; $E \sim 6.065$ GeV
- **Coherent** and **incoherent** DVCS: **nuclear GPDs**, **EMC effect**

^4He is a spin-0 nucleus: at twist-2 **only one CFF** in DVCS BSA

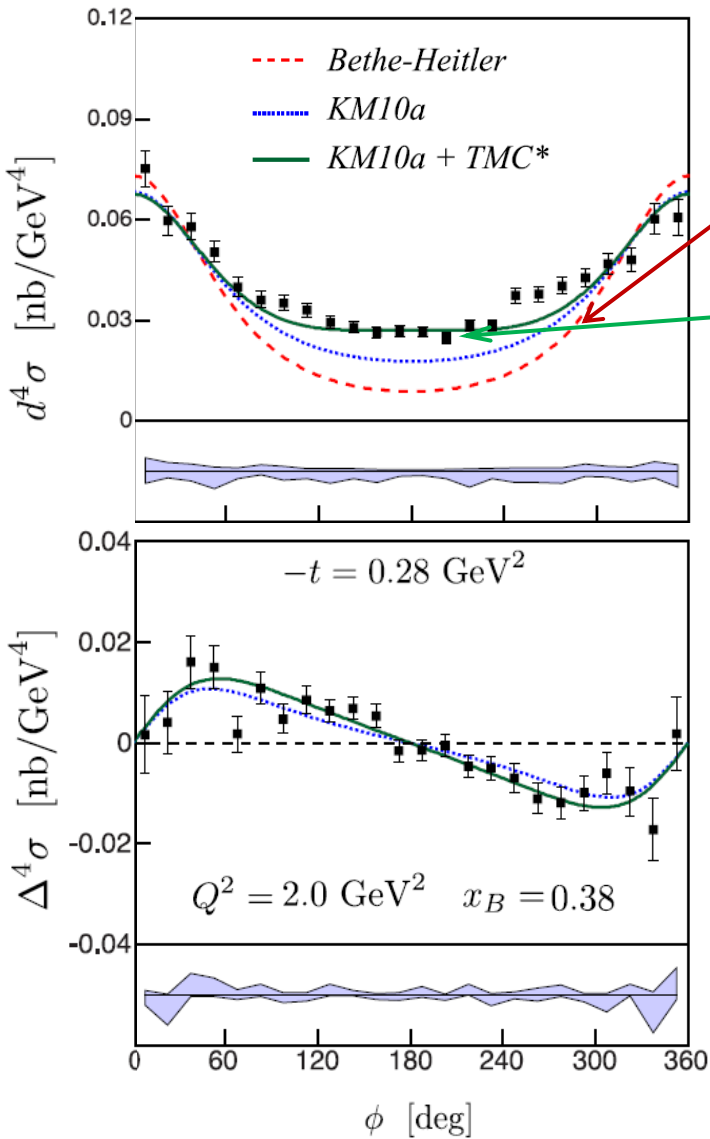
$$A_{\text{LU}}^{^4\text{He}}(\varphi) = \frac{\alpha_0(\varphi) F_A(t) \Im m[\mathcal{H}_A]}{\alpha_1(\varphi) F_A^2(t) + \alpha_2(\varphi) F_A(t) \Re e[\mathcal{H}_A] + \alpha_3(\varphi) \Re e[\mathcal{H}_A]^2 + \alpha_3(\varphi) \Im m[\mathcal{H}_A]^2}$$



- **Small $-t$:** asymmetry for ^4He lower than the bound proton one
- **High $-t$:** the two asymmetries tend to become compatible

Work by M. Hattawy, IPNO & ANL

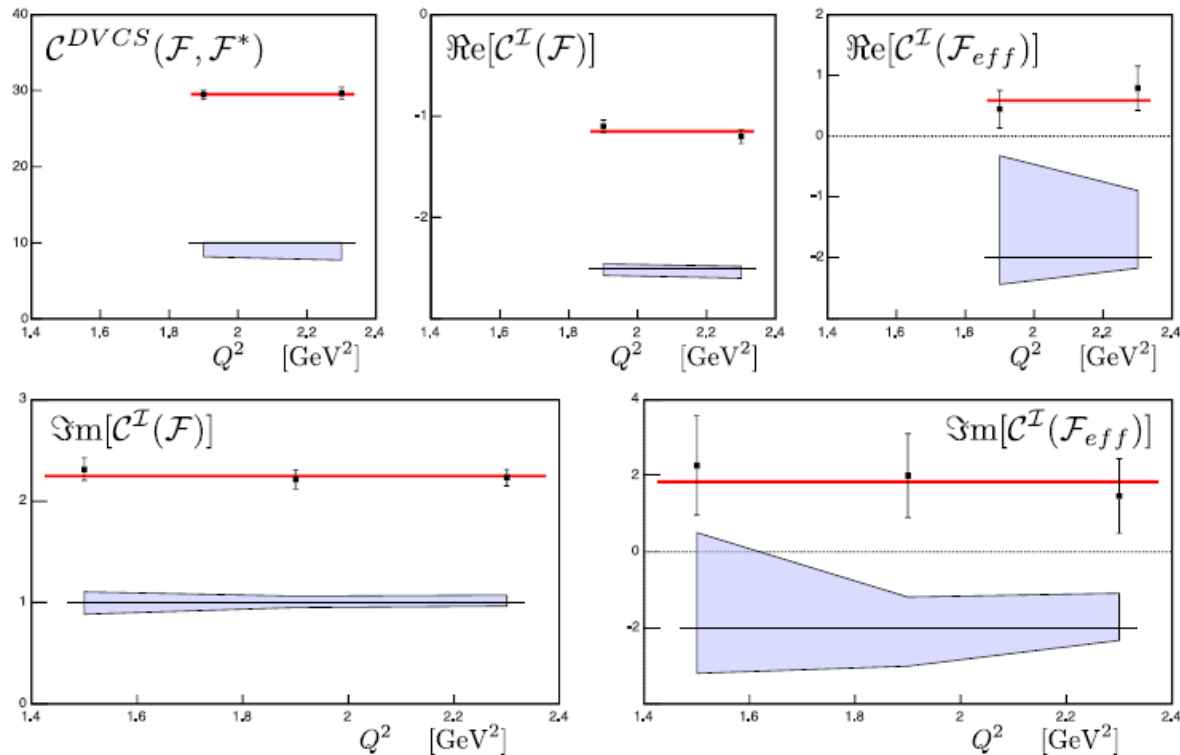
DVCS on the proton in Hall A



Re-analysis of E00-110 (2001)

M. Defurne et. al., Phys. Rev. C92, 055202 (2015)

- Significant deviation from Bethe-Heitler (BH)
- Both $I(\text{BH} \cdot \text{DVCS})$ and DVCS^2 contribute to the cross section
- Twist-4 corrections (TMC) may be necessary to describe the data



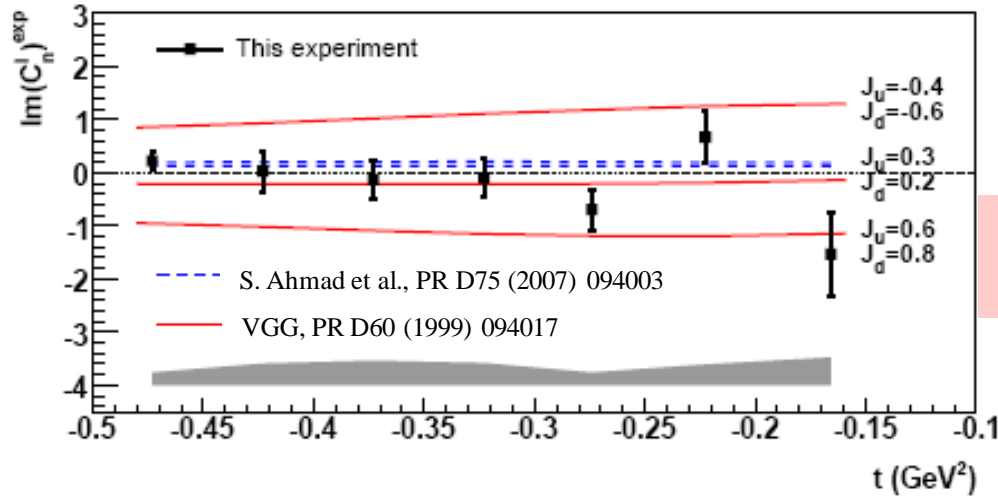
Beam-energy separation at constant Q^2 , x_B and t :
experiment E07-007 (Analysis ongoing)

$$\sigma = |BH|^2 + I(BH \cdot DVCS) + |DVCS|^2$$

$\sim E_b^3$ $\sim E_b^2$

DVCS on the neutron in Hall A

M. Mazouz et al., PRL 99 (2007) 242501



Proton and neutron GPDs (and CFFs) are **linear combinations** of quark GPDs

$$\mathcal{H}_p(\xi, t) = \frac{4}{9} \mathcal{H}_u(\xi, t) + \frac{1}{9} \mathcal{H}_d(\xi, t)$$

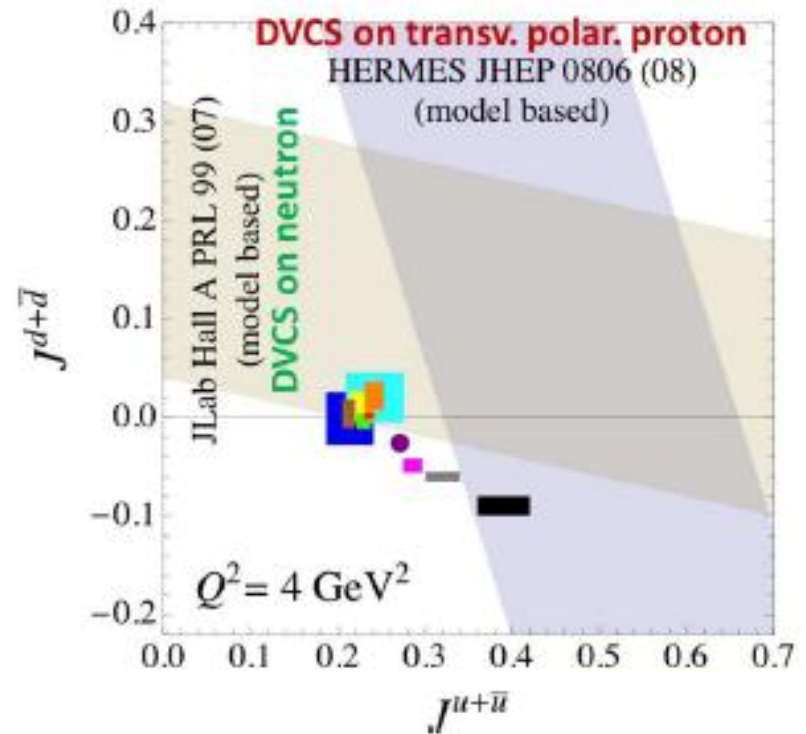
$$\mathcal{H}_n(\xi, t) = \frac{1}{9} \mathcal{H}_u(\xi, t) + \frac{4}{9} \mathcal{H}_d(\xi, t)$$

A **combined analysis** of DVCS observables for **proton and neutron targets** is necessary to perform a **quark-flavor separation** of the GPDs

$$\Delta\sigma_{LU} \sim \sin\phi \operatorname{Im}\{F_1\mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2\mathbf{E}\}$$

$$\frac{1}{2} \int_{-1}^1 x dx (H(x, \xi, t=0) + E(x, \xi, t=0)) = J$$

• **E03-106: First-time measurement of $\Delta\sigma_{LU}$ for nDVCS, model-dependent extraction of J_u, J_d**



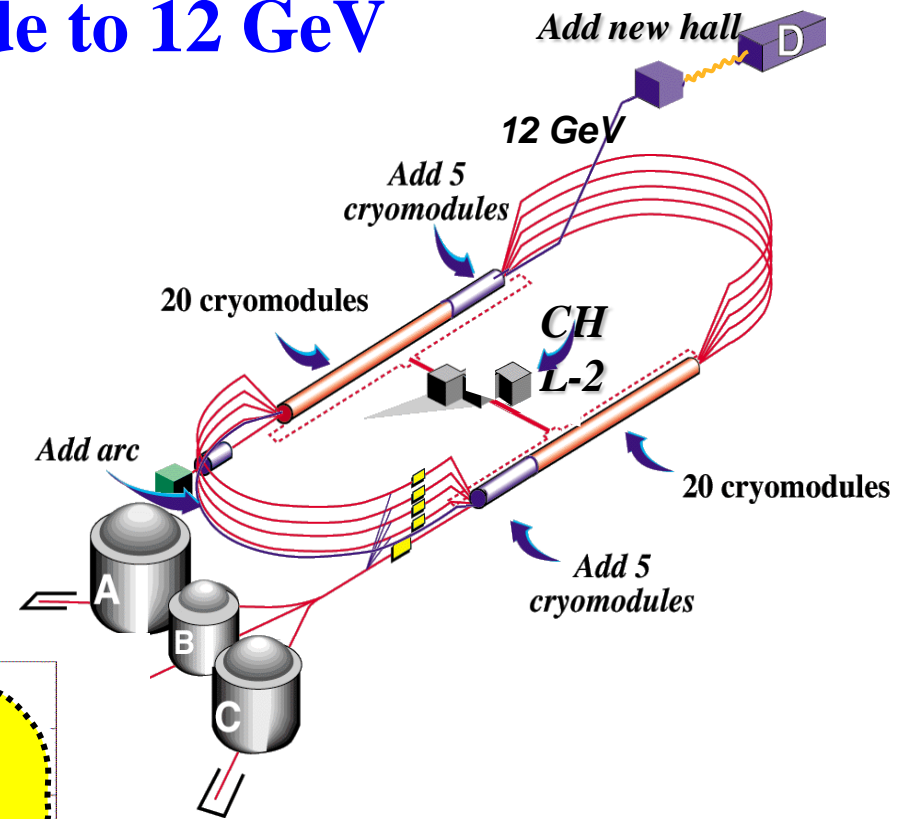
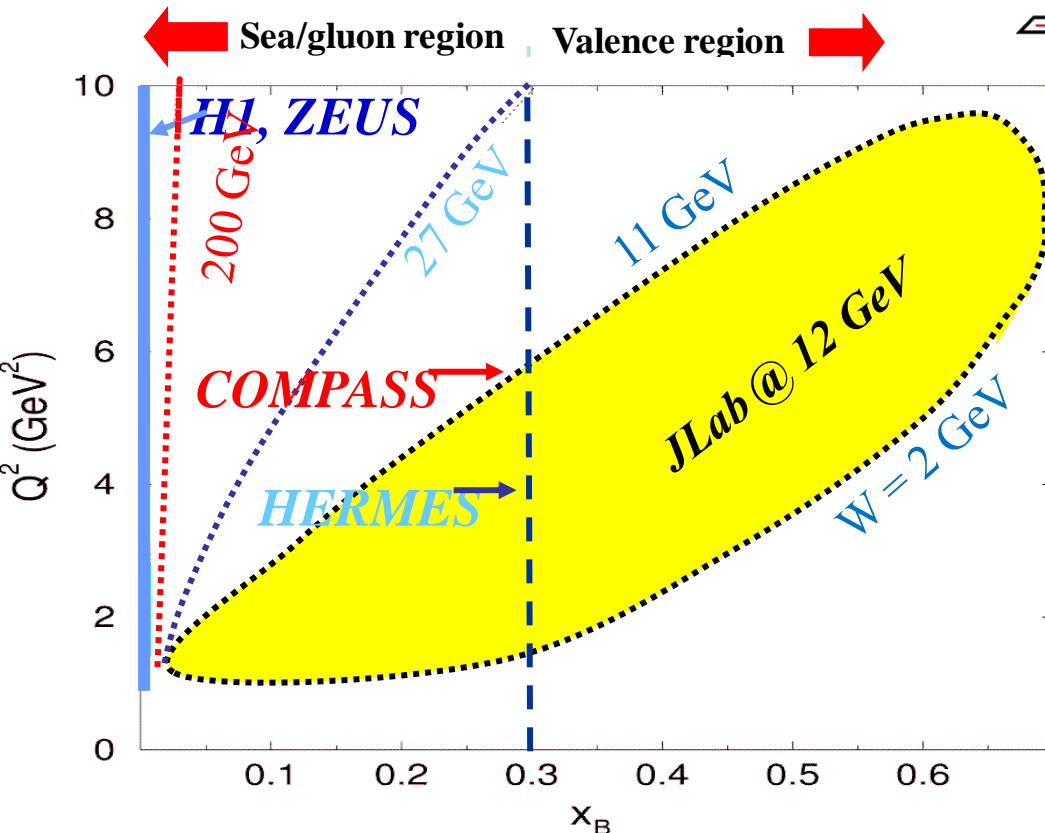
• **E08-025: Beam-energy « Rosenbluth » separation of nDVCS CS using an LD2 target**
 • Data taken in fall 2010, **analysis ongoing**

JLab upgrade to 12 GeV

$E = 2.2, 4.4, 6.6, 8.8, 11$ GeV
for the Halls A, B, C

Beam polarization $> 80\%$

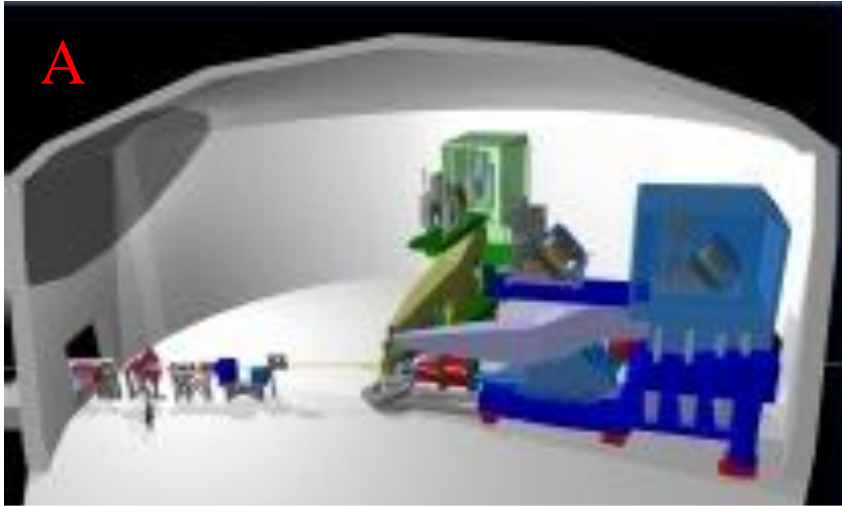
Upgrade completed in 2014



Study of high x_B domain
requires high luminosity

The 12-GeV upgrade is
well matched to studies in
the valence-quark regime

New capabilities in Halls A, B & C



**High Resolution Spectrometer (HRS) pair
and specialized large installation experiments**



CLAS12: large acceptance, high luminosity



**Super High Momentum Spectrometer (SHMS)
at high luminosity and forward angles**

DVCS experiments at 11 GeV have been approved for each of these **three halls**.

Complementary programs:

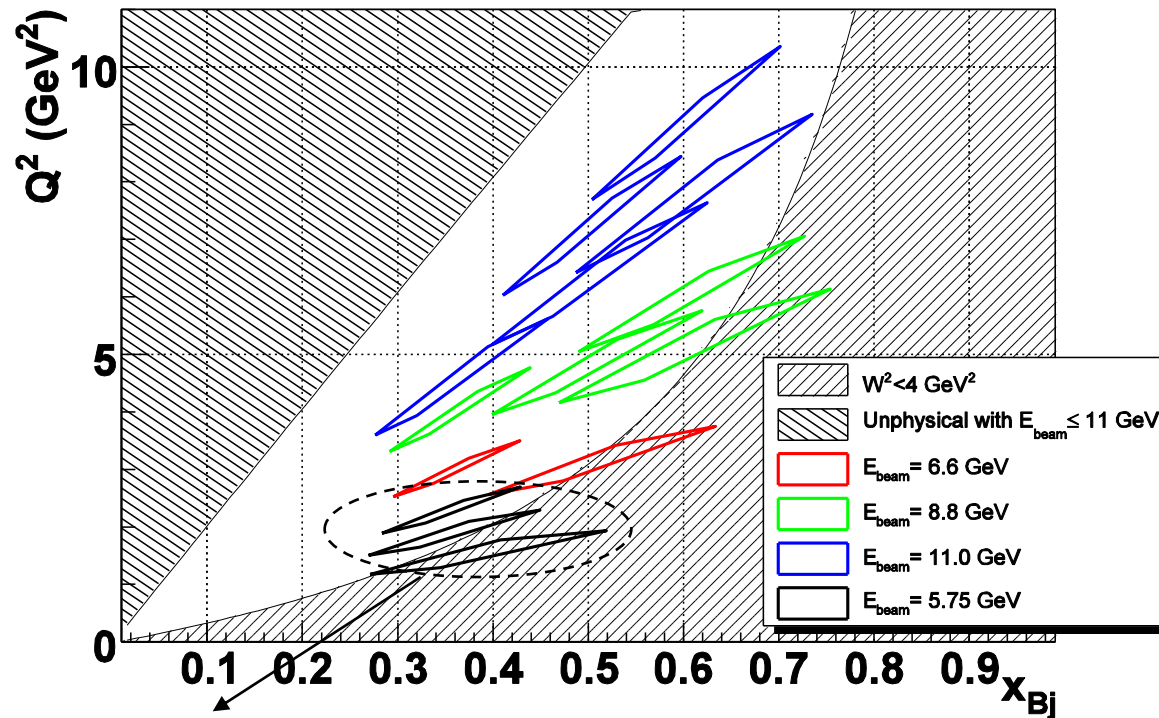
- different kinematic coverage
- different precisions/resolutions
- focus on different observables

E12-06-114: pDVCS at 11 GeV in Hall A

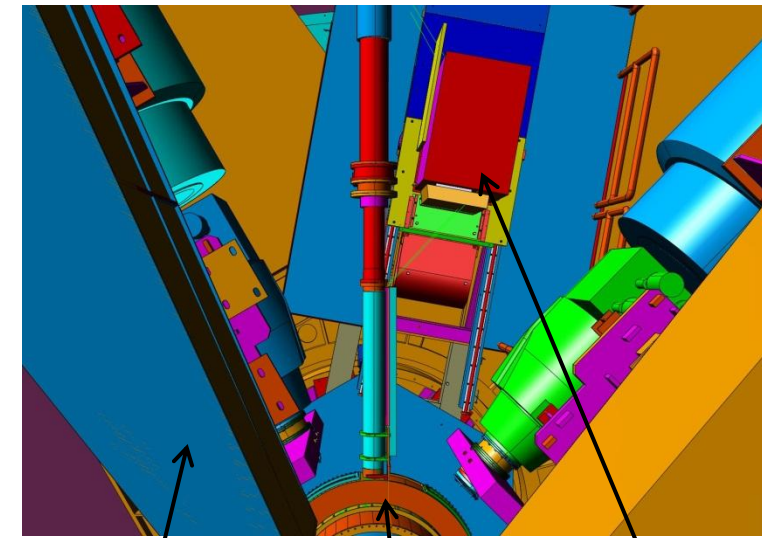
- Absolute cross-section measurements
- Test of scaling: Q^2 dependence of $d\sigma$ at fixed x_{Bj}
- Increased kinematical coverage

JLab12 with 3, 4, 5 pass beam (6.6, 8.8, 11.0 GeV)

DVCS measurements in Hall A/JLab



JLab @ 6 GeV



L-HRS

Scattering
chamber

PbF2
calorimeter

**1st experiment to run after
the 12-GeV upgrade**

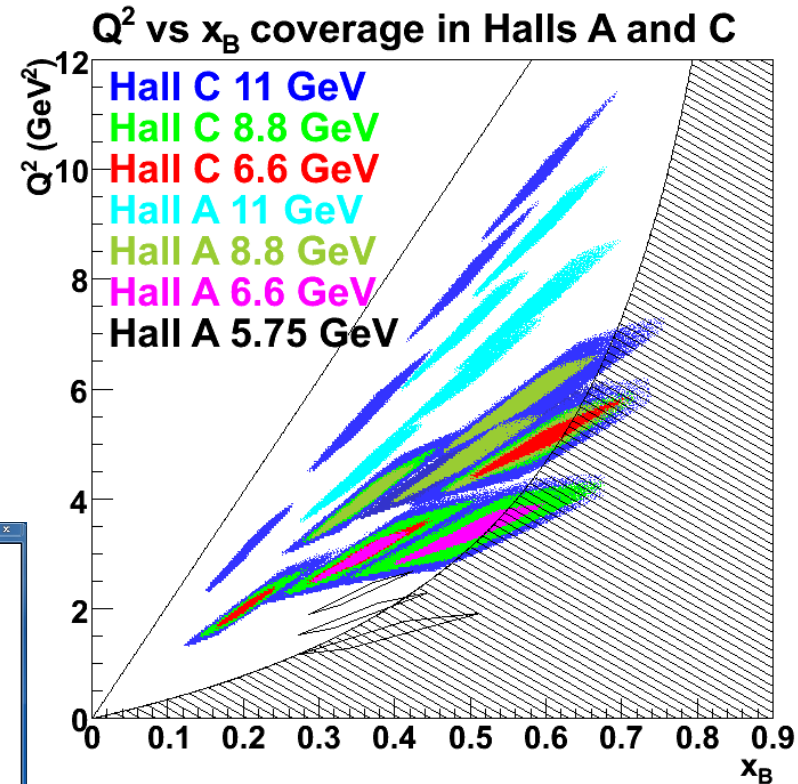
**Started in fall 2014,
continuing in 2016**

E12-13-010: pDVCS at 11 GeV in Hall C

- Energy separation of the DVCS cross section
- Higher Q^2 : measurement of higher twist contributions
- Low- x_B extension (thanks to sweeping magnet)

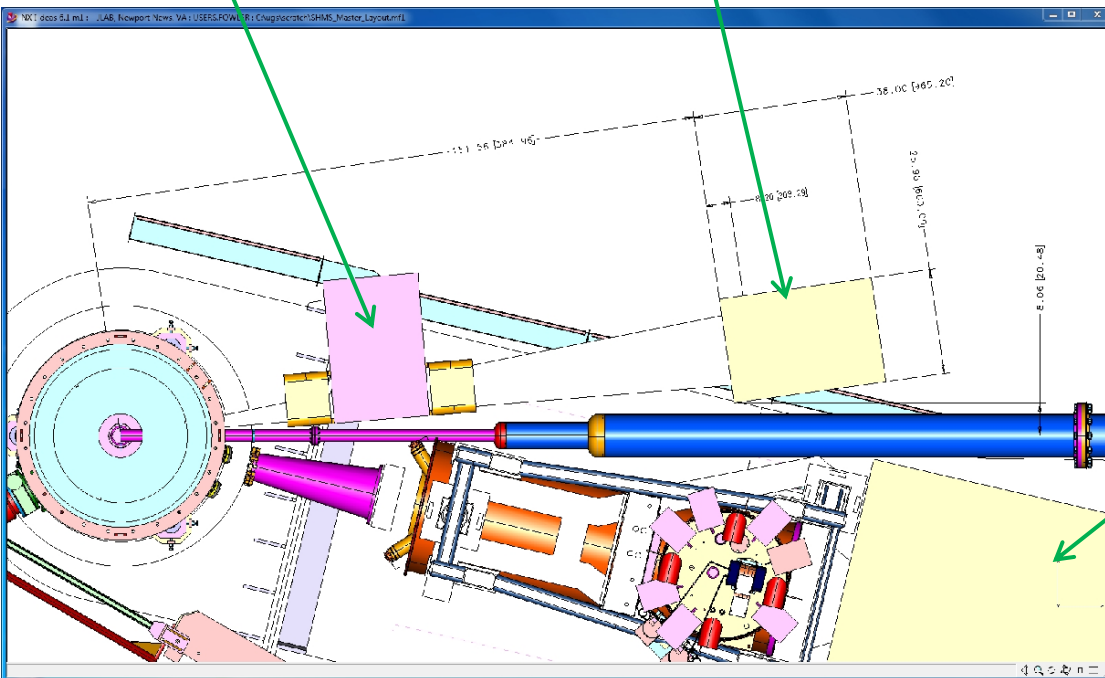
Sweeping
magnet

1116-block PbWO_4
calorimeter



Hall C
HMS

Tentative running:
~ 2019-20



DVCS BSA and TSA with CLAS12 & 11 GeV beam

85 days of beam time

Liquid hydrogen target

$P_{\text{beam}} = 85\%$

$L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}$

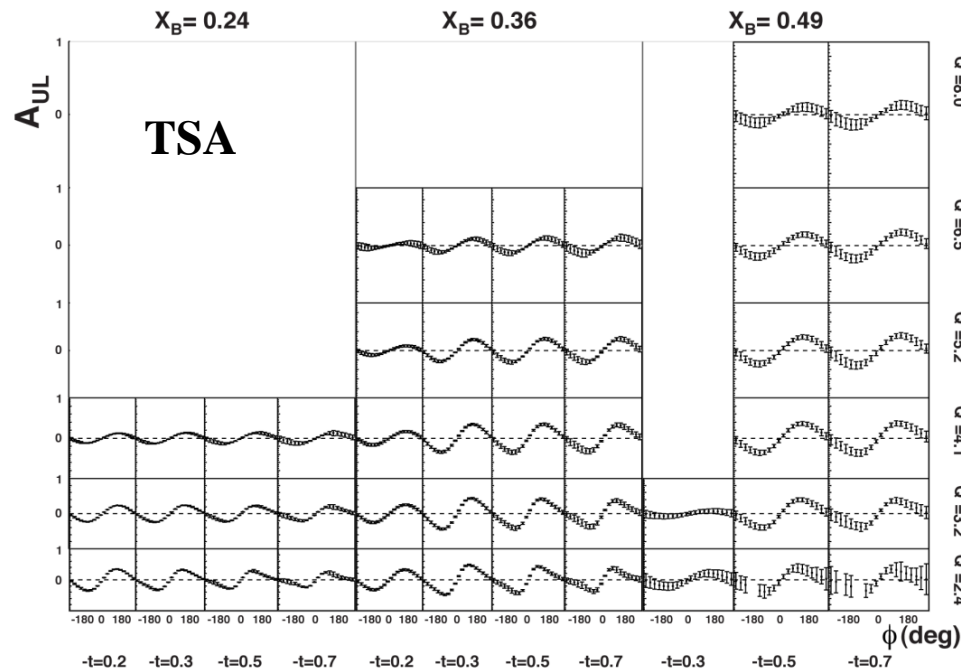
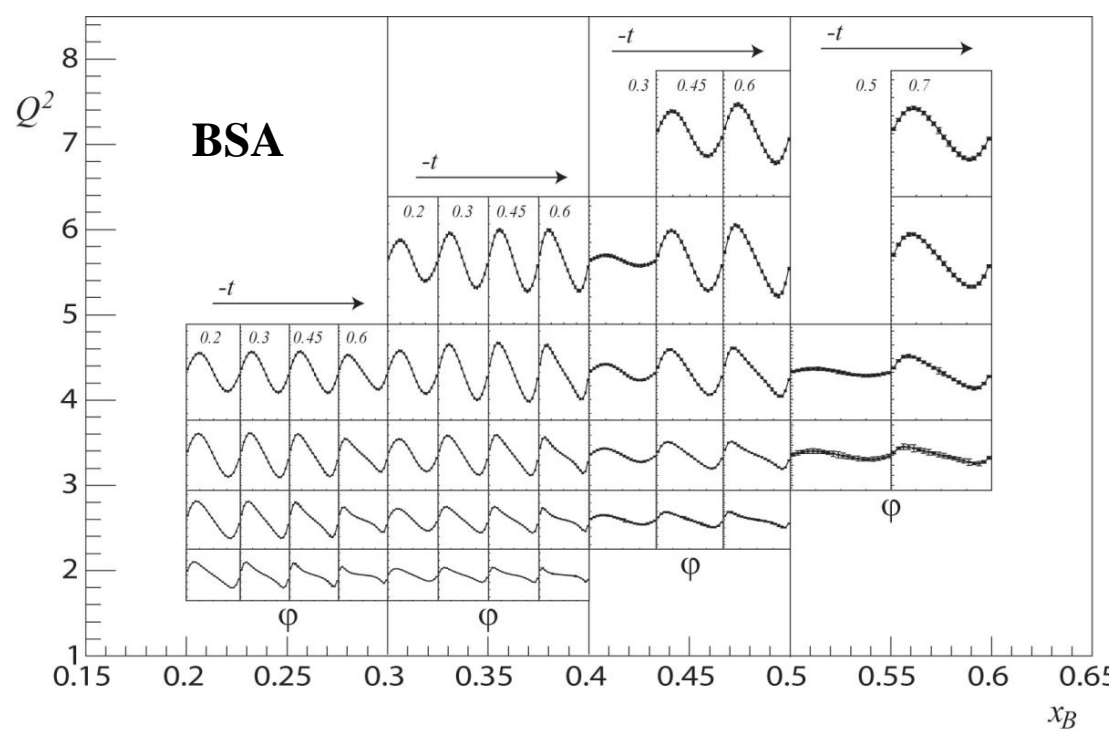
$1 < Q^2 < 10 \text{ GeV}^2, 0.1 < x_B < 0.65$

$-t_{\text{min}} < -t < 2.5 \text{ GeV}^2$

Statistical error: 1% to 10%

on $\sin\phi$ moments

Systematic uncertainties: $\sim 6\text{-}8\%$



120 days of beam time

NH_3 long. polarized target

$P_{\text{beam}} = 85\%, P_{\text{target}} = 80\%$

$L = 2 \cdot 10^{35} \text{ cm}^{-2}\text{s}^{-1}$

$1 < Q^2 < 10 \text{ GeV}^2, 0.1 < x_B < 0.65$

$-t_{\text{min}} < -t < 2.5 \text{ GeV}^2$

Statistical error: 2% to 15%

on $\sin\phi$ moments

Systematic uncertainties: $\sim 6\text{-}8\%$

DVCS BSA and TSA with CLAS12 & 11 GeV beam

85 days of beam time

Liquid hydrogen target

$P_{\text{beam}} = 85\%$

$L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}$

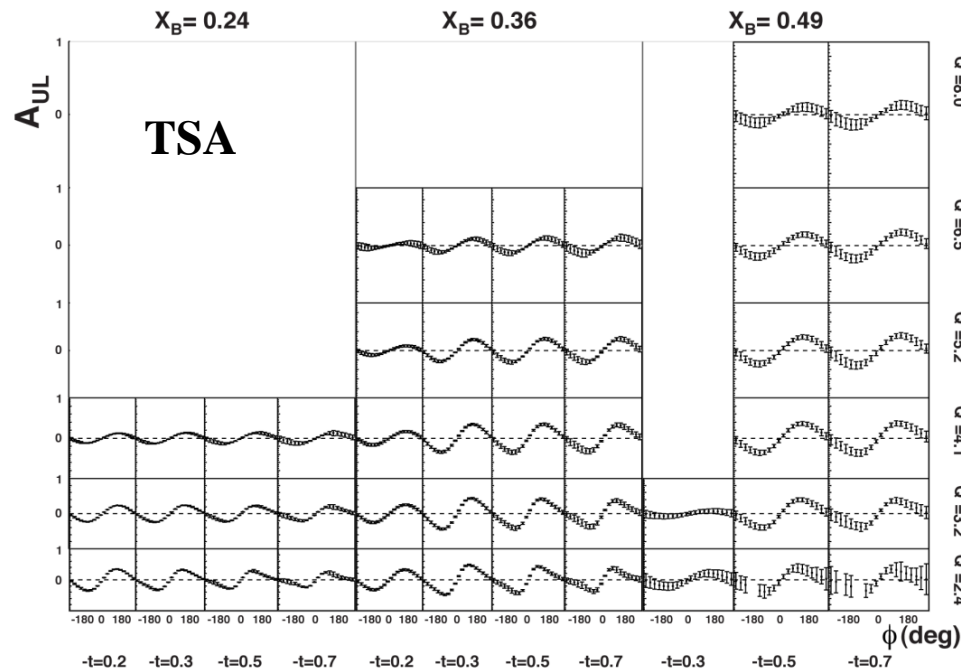
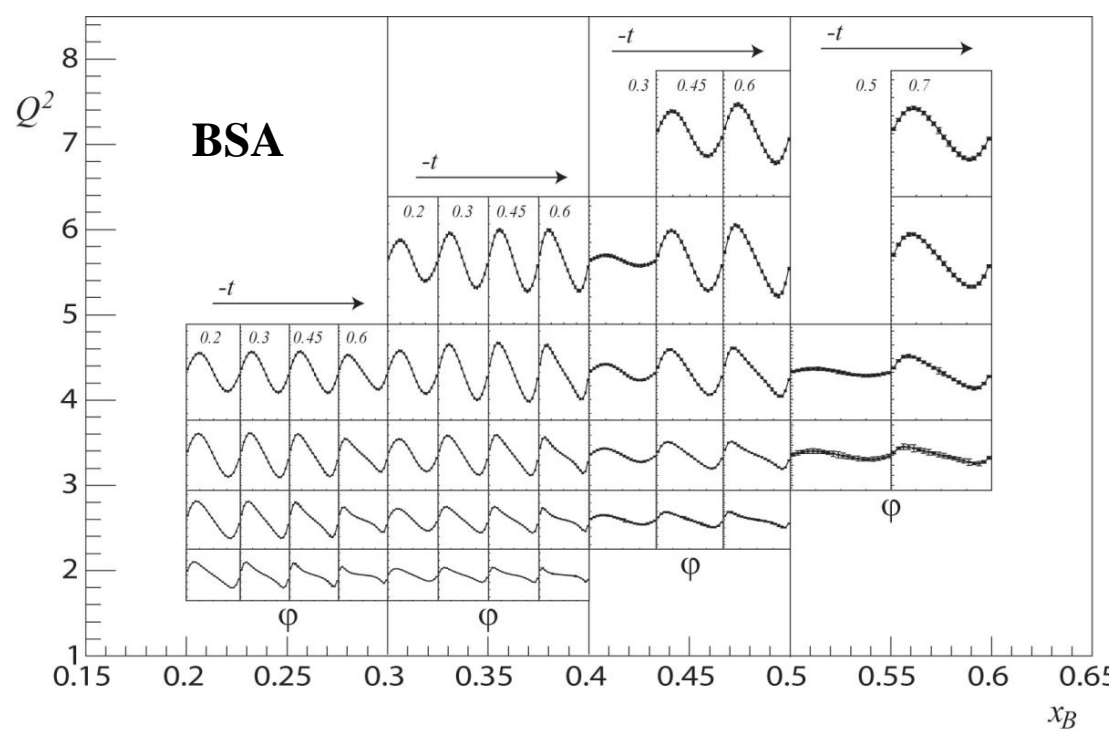
$1 < Q^2 < 10 \text{ GeV}^2, 0.1 < x_B < 0.65$

$-\mathbf{t}_{\text{min}} < -\mathbf{t} < 2.5 \text{ GeV}^2$

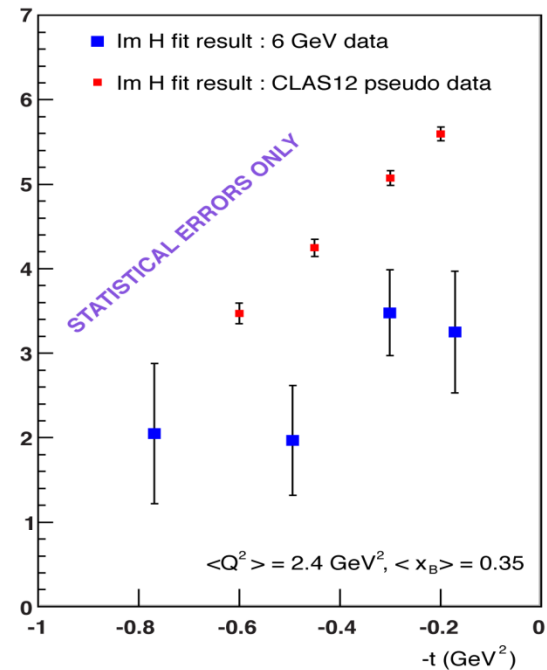
Statistical error: 1% to 10%

on $\sin\phi$ moments

Systematic uncertainties: $\sim 6\text{-}8\%$



Impact of CLAS12 DVCS-BSA data on model-independent fit to extract $\text{Im}(H)$



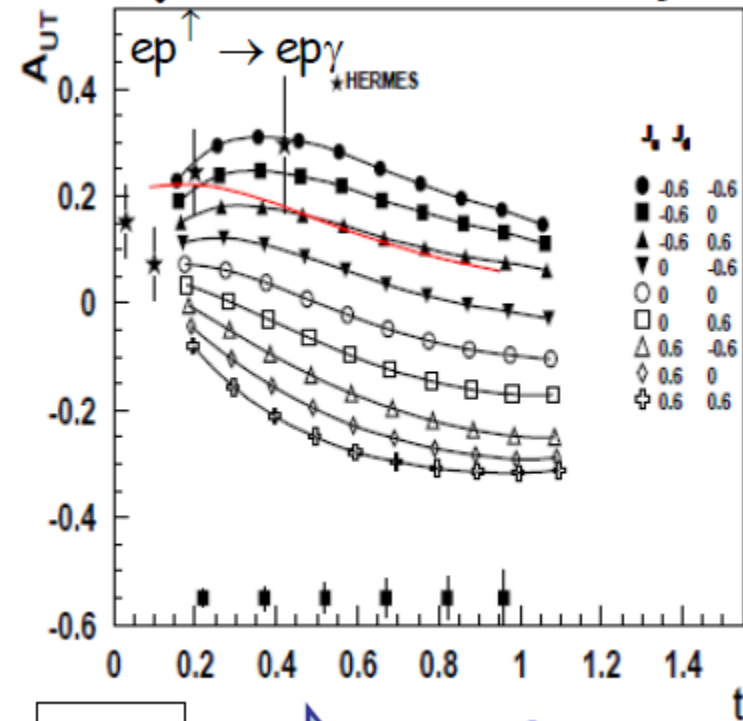
CLAS12: p-DVCS *transverse* target-spin asymmetry

100 days of beam time

Beam pol. = 80% ; **target pol. (HDIce) = 60%** ; Luminosity = $5 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$

$1 < Q^2 < 10 \text{ GeV}^2$, $0.06 < x_B < 0.66$, $-t_{\min} < -t < 1.5 \text{ GeV}^2$

Projections for $Q^2 = 2.5 \text{ GeV}^2$, $x_B = 0.2$

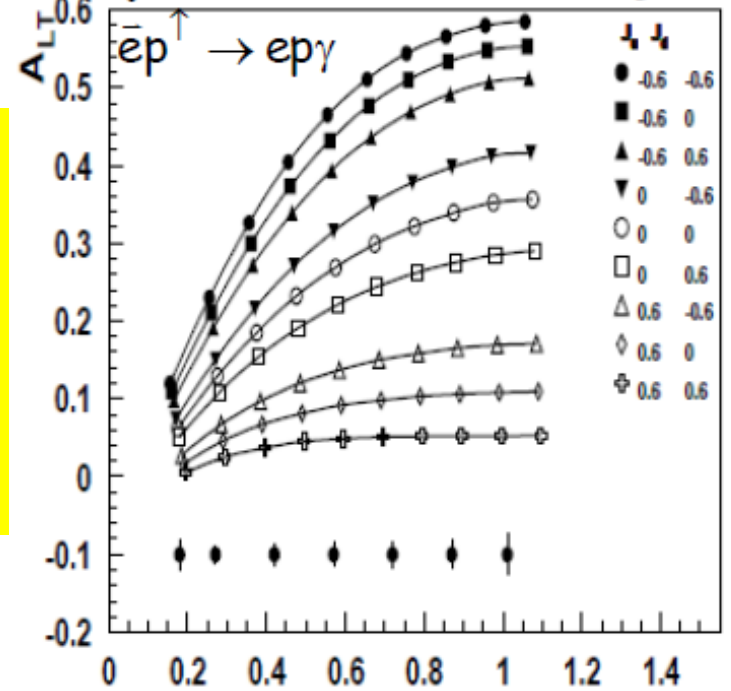


Transverse-target
spin asymmetry
for p-DVCS is
highly sensitive
to the **u-quark**
contributions to
proton spin.

$$\Delta\sigma_{UT} \longrightarrow \text{Im}\{\mathcal{H}_p, \mathcal{E}_p\}$$

JLab PAC:
high-impact
experiment

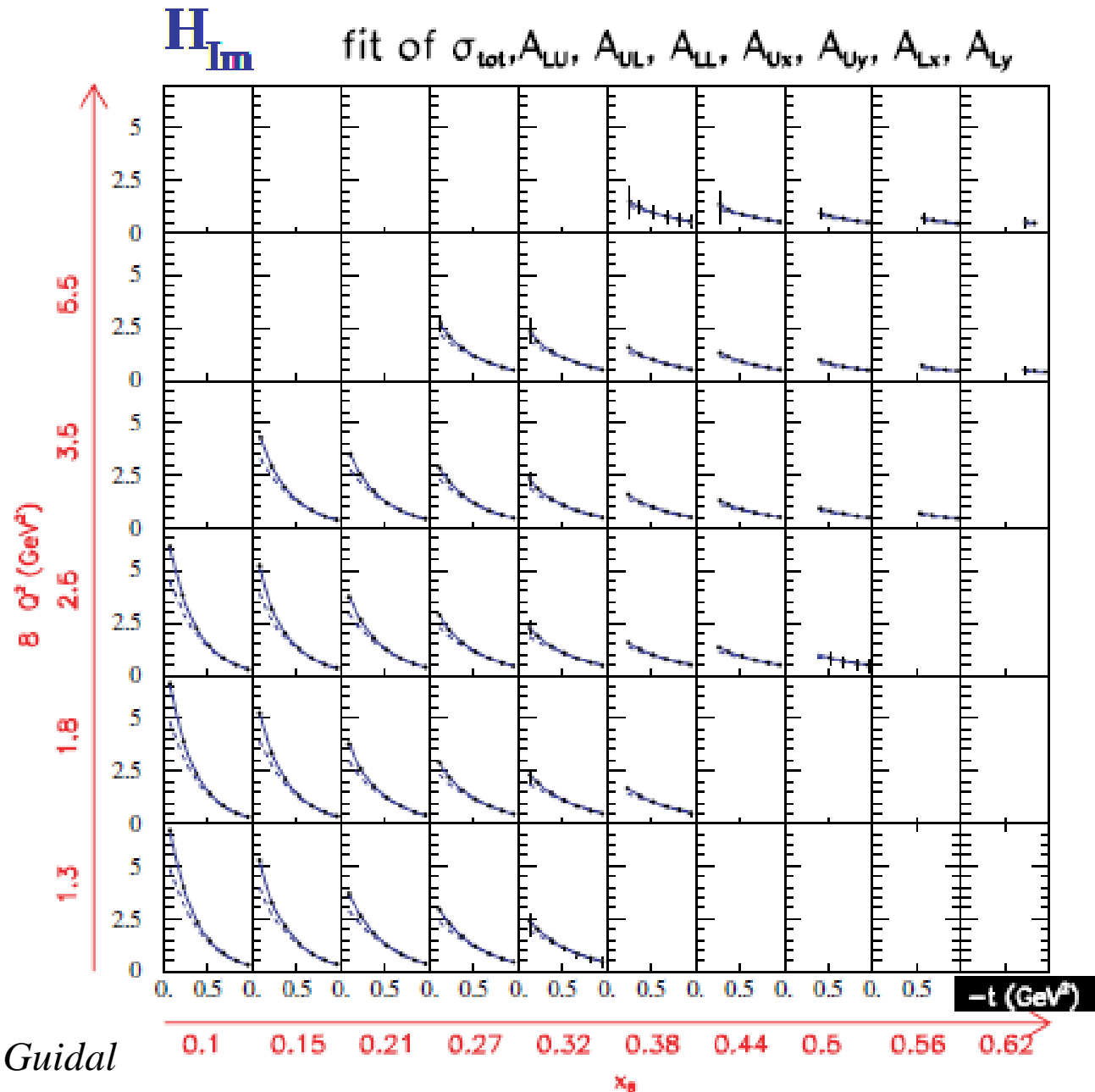
Projections for $Q^2 = 2.5 \text{ GeV}^2$, $x_B = 0.2$



$$\Delta\sigma_{LT} \longrightarrow \text{Re}\{\mathcal{H}_p, \mathcal{E}_p\}$$

Proposal conditionally approved by PAC39

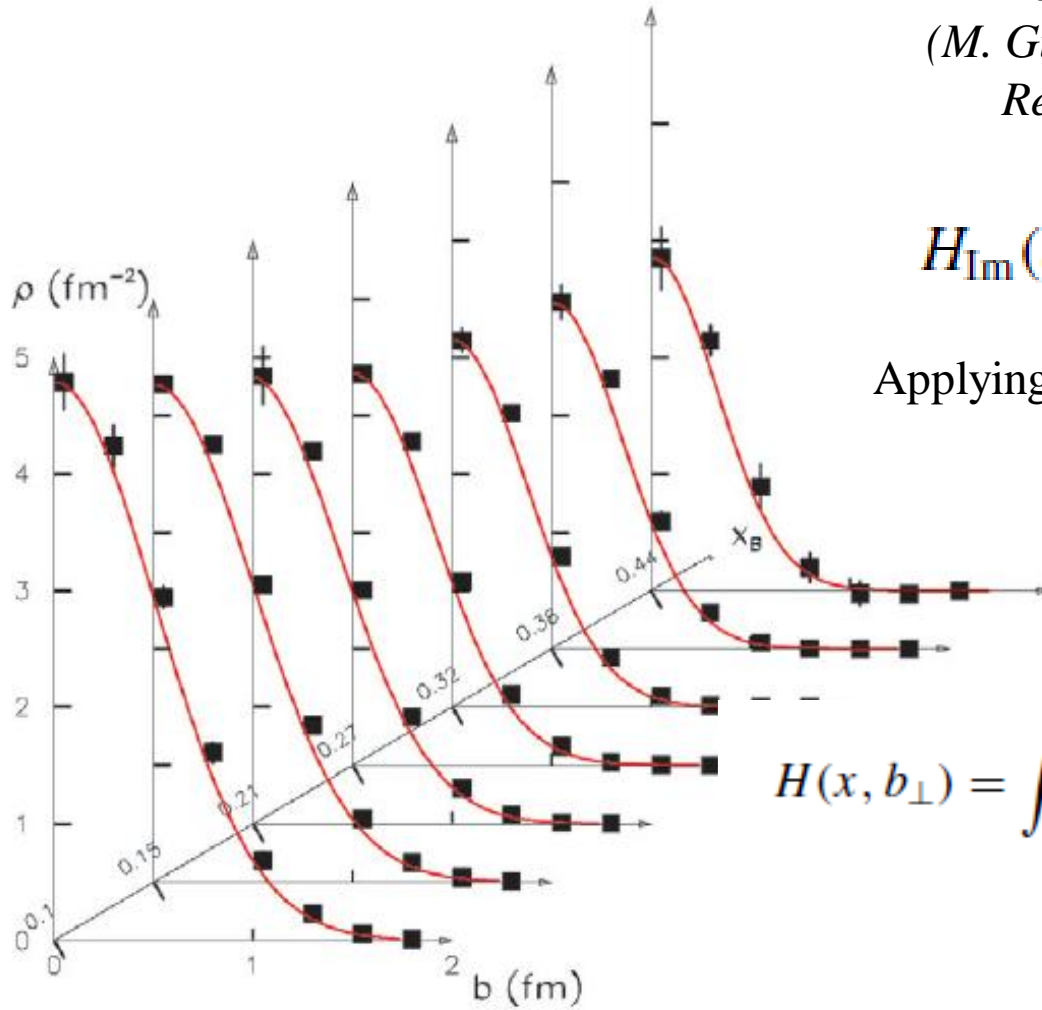
Projections for CLAS12 for H_{Im}



From CFFs to spatial densities

How to go from momentum coordinates (t)
to space-time coordinates (b) ?

(M. Guidal, H. Moutarde, M. Vanderhagen,
Rept.Prog.Phys. 76 (2013) 066202)



$$H_{\text{Im}}(\xi, t) \equiv H(\xi, \xi, t) - H(-\xi, \xi, t)$$

Applying a model-dependent “deskewing” factor:

$$\frac{H(\xi, 0, t)}{H(\xi, \xi, t)}$$

$$H(x, b_{\perp}) = \int_0^{\infty} \frac{d\Delta_{\perp}}{2\pi} \Delta_{\perp} J_0(b_{\perp} \Delta_{\perp}) H(x, 0, -\Delta_{\perp}^2)$$

Burkardt (2000)

Projections for CLAS12

E12-11-003: BSA for DVCS *on the neutron* with CLAS12

$$(H, E)_u(\xi, \xi, t) = \frac{9}{15} [4(H, E)_p(\xi, \xi, t) - (H, E)_n(\xi, \xi, t)]$$

$$(H, E)_d(\xi, \xi, t) = \frac{9}{15} [4(H, E)_n(\xi, \xi, t) - (H, E)_p(\xi, \xi, t)]$$

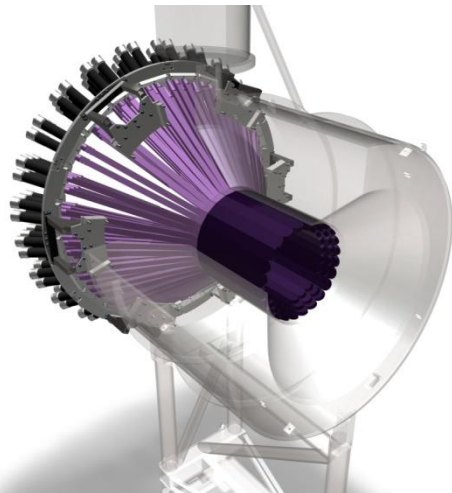
$$\Delta\sigma_{\text{LU}} \sim \sin\phi \operatorname{Im}\{F_1\mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2\mathcal{E}\}d\phi$$

The most sensitive observable to the GPD \mathcal{E}

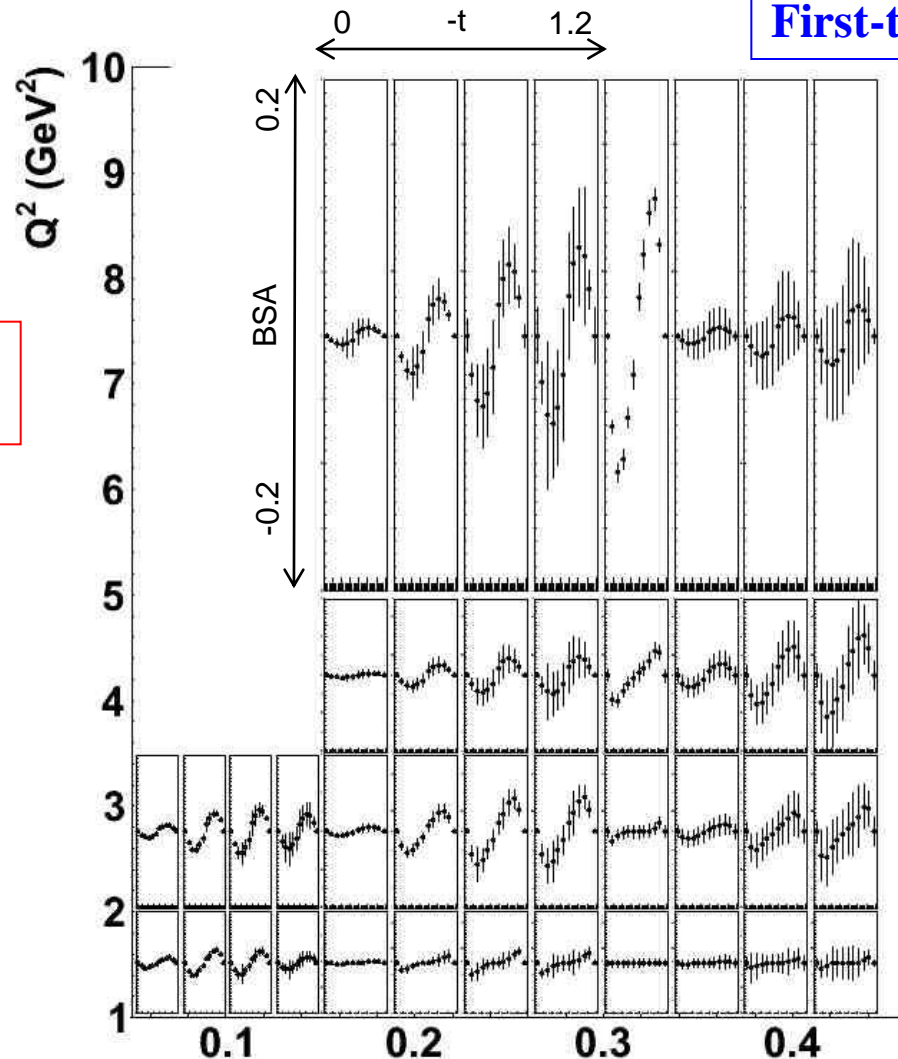
$ed \rightarrow e(p)n\gamma$

CLAS12 +
Forward Calorimeter +
Neutron Detector

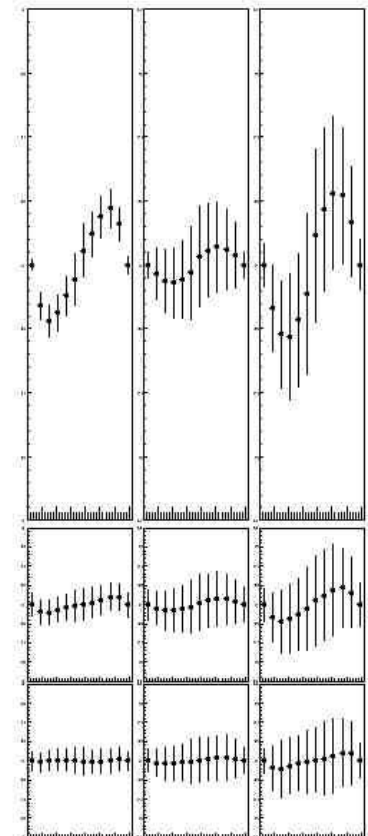
80 days of data taking
 $L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}/\text{nucleon}$



Construction completed
(IPN Orsay)



First-time measurement



JLab PAC:
high-impact
experiment

x_B

E12-11-003: BSA for DVCS *on the neutron* with CLAS12

$$(H, E)_u(\xi, \xi, t) = \frac{9}{15} [4(H, E)_p(\xi, \xi, t) - (H, E)_n(\xi, \xi, t)]$$

$$(H, E)_d(\xi, \xi, t) = \frac{9}{15} [4(H, E)_n(\xi, \xi, t) - (H, E)_p(\xi, \xi, t)]$$

$$\Delta\sigma_{\text{LU}} \sim \sin\phi \operatorname{Im}\{F_1\mathcal{H} + \xi(F_1+F_2)\tilde{\mathcal{H}} - kF_2\mathcal{E}\}d\phi$$

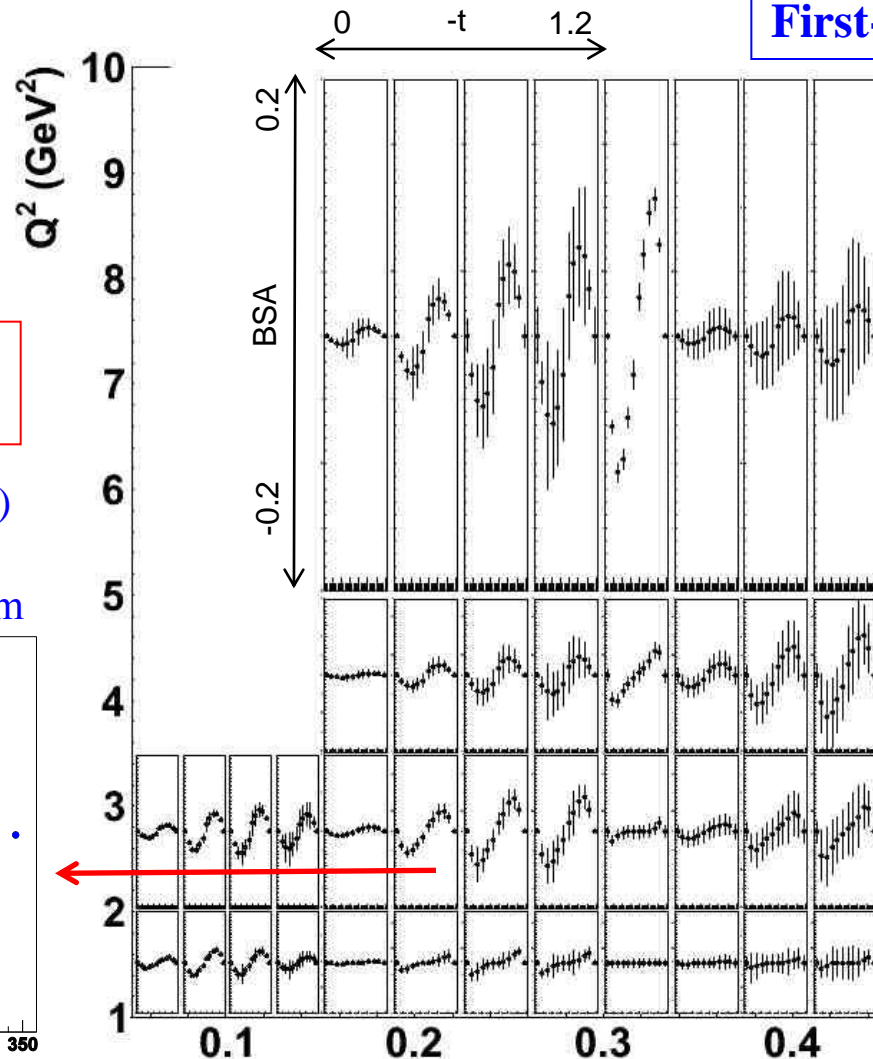
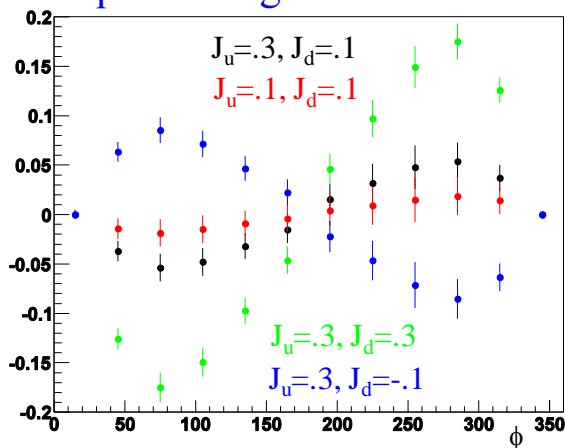
The most sensitive observable to the GPD \mathcal{E}

$ed \rightarrow e(p)n\gamma$

CLAS12 +
Forward Calorimeter +
Neutron Detector

80 days of data taking
 $L = 10^{35} \text{ cm}^{-2}\text{s}^{-1}/\text{nucleon}$

Model predictions (VGG)
for different values of
quarks' angular momentum



First-time measurement

JLab PAC:
high-impact
experiment

x_B

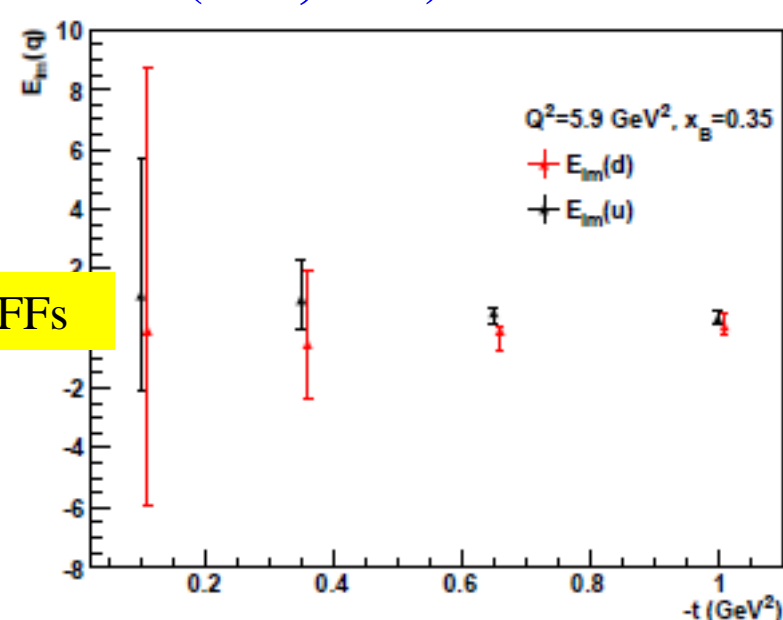
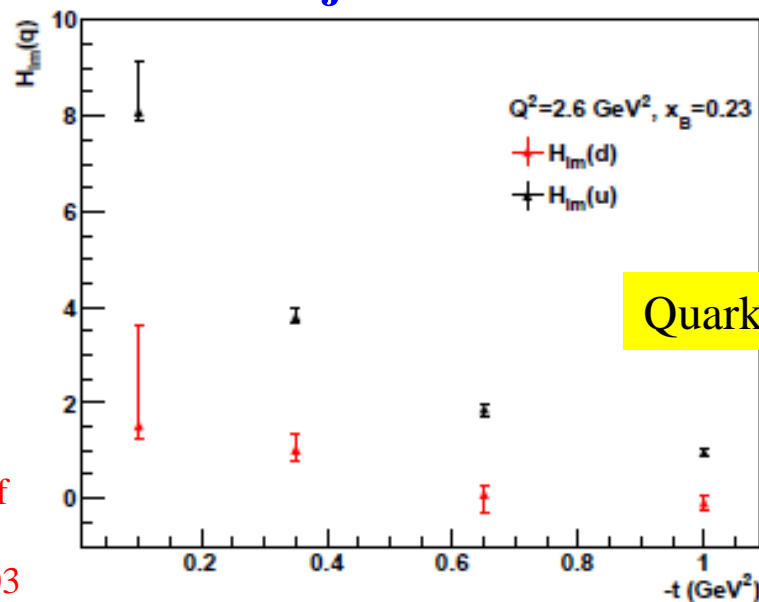
DVCS: global view

Observable (target)	Target	Sensitivity to CFFs	Completed experiments	12-GeV experiments
$\Delta\sigma_{beam}(p)$	Unpolarized hydrogen	$\Im m\mathcal{H}_p$	Hall A, CLAS	Hall A, CLAS12, Hall C
BSA(p)	Unpolarized hydrogen	$\Im m\mathcal{H}_p$	HERMES, CLAS	CLAS12
TSA(p)	Long. pol. NH3	$\Im m\mathcal{H}_p, \Im m\mathcal{H}_p$	HERMES, CLAS	CLAS12
DSA(p)	Long. pol. NH3	$\Re\mathcal{H}_p, \Re\mathcal{H}_p$	HERMES, CLAS	CLAS12
ϵ TSA(p)	Transv. pol. protons	$\Im m\mathcal{H}_p, \Im m\mathcal{E}_p$	HERMES	CLAS12
$\Delta\sigma_{beam}(n)$	Unpolarized deuterium	$\Im m\mathcal{E}_n$	Hall A	
BSA(n)	Unpolarized deuterium	$\Im m\mathcal{E}_n$		CLAS12
TSA(n)	Long. pol. ND3	$\Im m\mathcal{H}_n$		PR12-15-004
DSA(n)	Long. pol. ND3	$\Re\mathcal{H}_n$		PR12-15-004

Projections for flavor separation ($\Im m\mathcal{H}, \Im m\mathcal{E}$)

Fit using all the projected pDVCS asymmetries of the CLAS12 program

Fit using the projected nDVCS asymmetries of PR12-15-004 and E12-11-003



Conclusions

- GPDs are a unique tool to explore the **internal dynamics of the nucleon**:
 - **3D** quark/gluon **imaging** of the nucleon
 - **orbital angular** momentum carried by quarks
- Their extraction from experimental data is **very difficult**:
 - there are **4 GPDs for each quark flavor**
 - they depend on **3 variables**, only two (ξ , t) experimentally accessible via DVCS
- ✓ Recently-developed fitting methods allow to **extract CFFs from DVCS observables**. Need to measure several **p-DVCS** and **n-DVCS observables** over a **wide phase space**
- ✓ A wealth of **new results** on various DVCS observables is coming from recent **CLAS** and **Hall-A** experiments (on the proton, deuterium and ^4He targets)
- ✓ First **tomographic interpretations** of the quarks in the **proton**:
 - ✓ **valence quarks** are concentrated in its **center**, **sea quarks** at its **periphery**
 - ✓ **axial charge** more concentrated than the **electric** one
- The 12-GeV-upgraded JLab will be **the only facility** to perform DVCS experiments **in the valence region**, for Q^2 up to 11 GeV
- DVCS experiments on both **proton** and **deuterium** targets (polarized and unpolarized) are planned for **3 of the 4 Halls at JLab@12 GeV: quarks' spatial densities, quark-flavor separation, quarks' orbital angular momentum...**
- **Beyond DVCS: double DVCS (x dependence), TCS, exclusive meson production, ...**