

A study of HERA I+II combined data at low Q^2

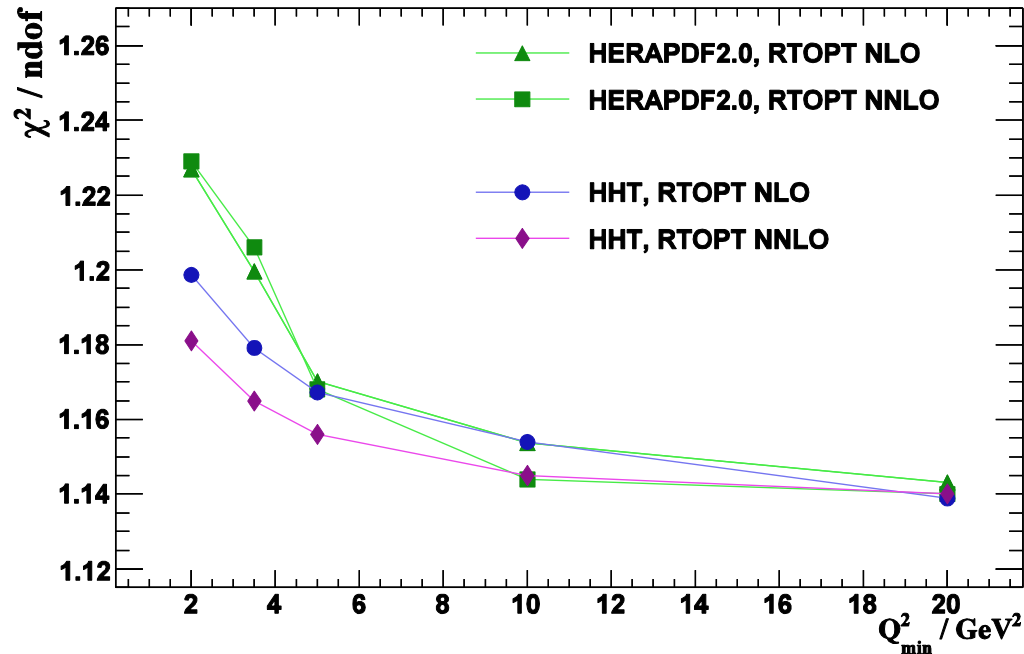
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A study of adding higher twist terms to the HERAPDF2.0 analysis of the HERA-I+II data for NLO and NNLO fits

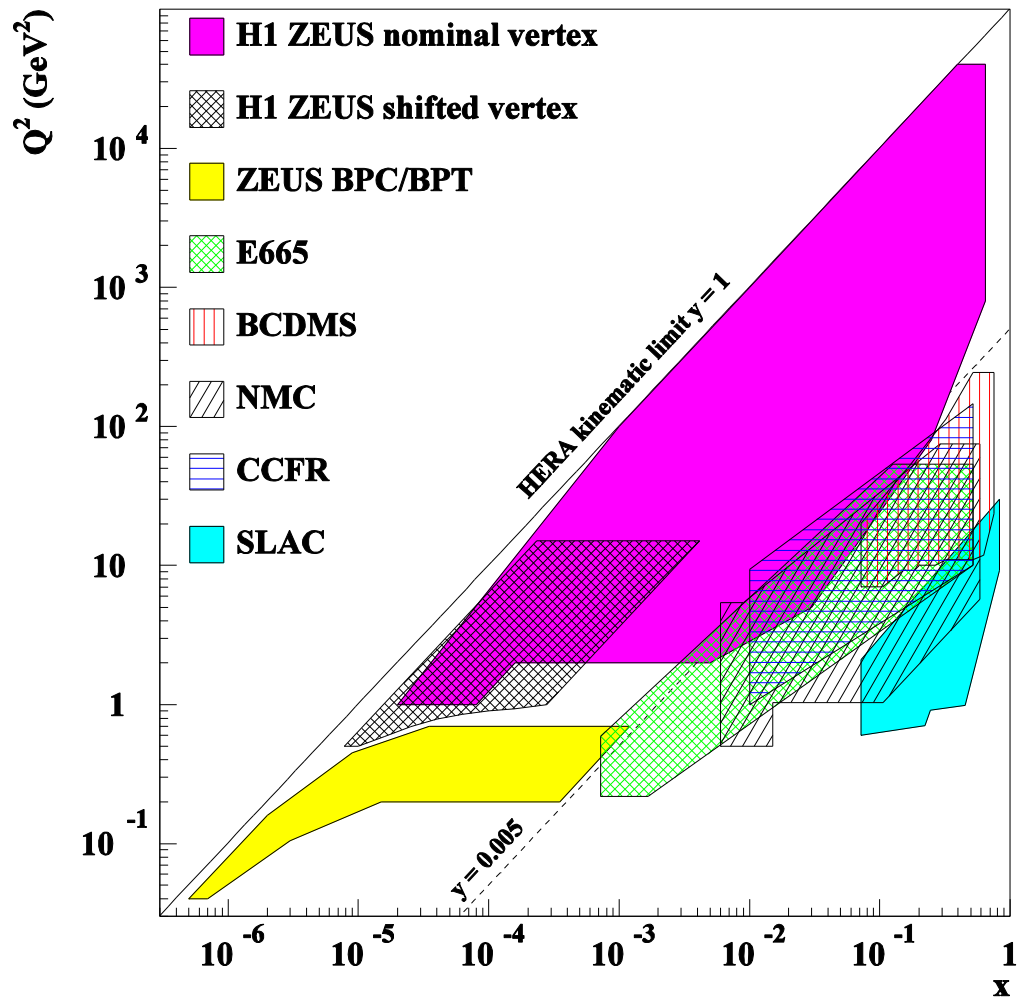
Higher twist terms are important in F_L for low Q^2 , which for HERA kinematics means at low- x

Such terms are significant in F_L for $2 < Q^2 < 50 \text{ GeV}^2$
But such an approach fails for $Q^2 < 2 \text{ GeV}^2$

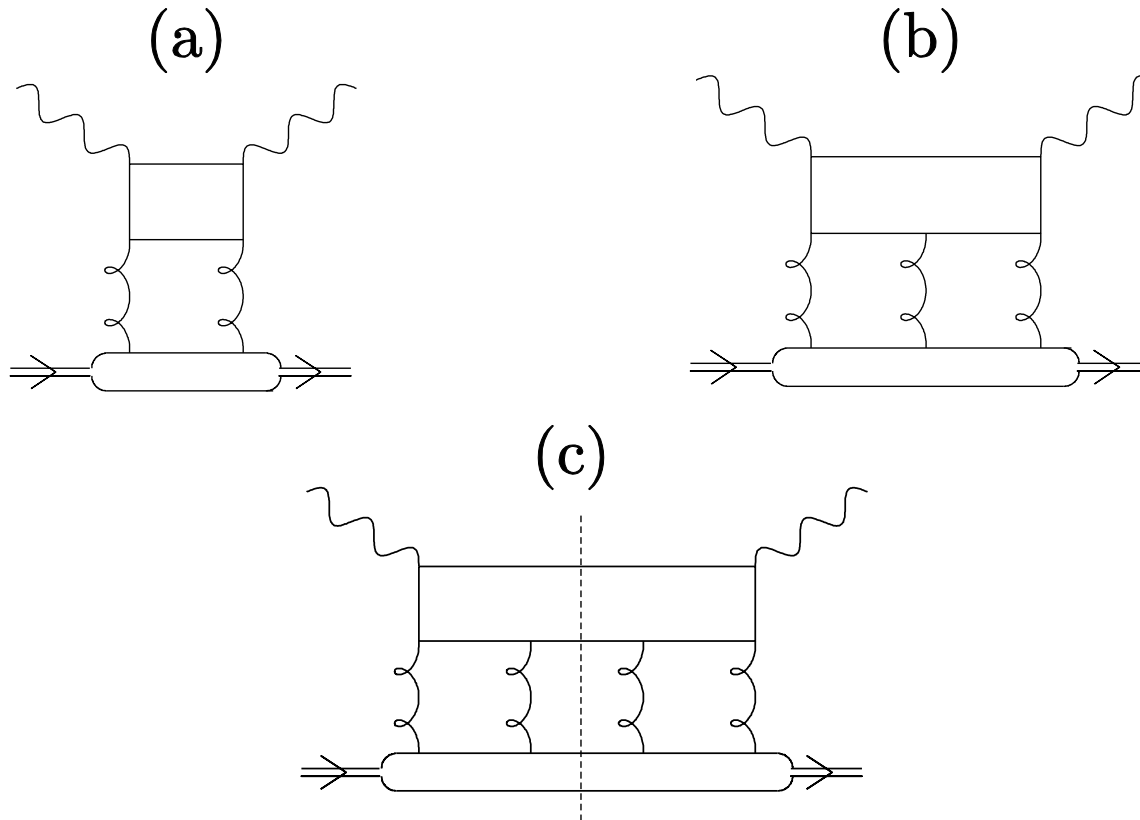
The χ^2/ndof of the HERAPDF2.0 NLO and NNLO fits deteriorate as the minimum value of Q^2 for data entering the fit is lowered



One way to improve this is to add higher twist terms - HHT analysis
 BUT NOTE- these are not the high-x, low Q^2 contributions that we usually
 associate with the terminology 'higher twist'
 Most groups exclude those contributions by a W cut, $W^2 > 12.5 \text{ GeV}^2$
 ALL HERA data is at much higher $W^2 > 300 \text{ GeV}^2$



HERA data at low Q^2
are also at low- x



We are now considering higher twist terms which act a low-x
 Their origin COULD be connected with the recombination of gluon ladders.
 Bartels, Golec-Biernat, Kowalski suggest that such higher twist terms would
 cancel between σ_L and σ_T in F_2 , but remain strong in F_L

Try the simplest of possible modification to the structure functions F_2 and F_L as calculated from HERAPDF2.0 formalism

$$F_{2,L} = F_{2,L} (1 + A_{2,L}^{\text{HT}}/Q^2)$$

We find that such a modification of F_L is favoured, whereas for F_2 it is not.

At NNLO the $\chi^2/\text{ndof} = 1363/1131$ for HERAPDF2.0

If A_2^{HT} is added this becomes 1357/1130 and $A_2^{\text{HT}} = 0.12 \pm 0.07 \text{ GeV}^2$

If A_L^{HT} is added this becomes 1316/1130 and $A_L^{\text{HT}} = 5.5 \pm 0.6 \text{ GeV}^2$

If both A_L^{HT} and A_2^{HT} are added the result is consistent with just adding A_L^{HT}

So now concentrating on just F_L , we call these fits HHT

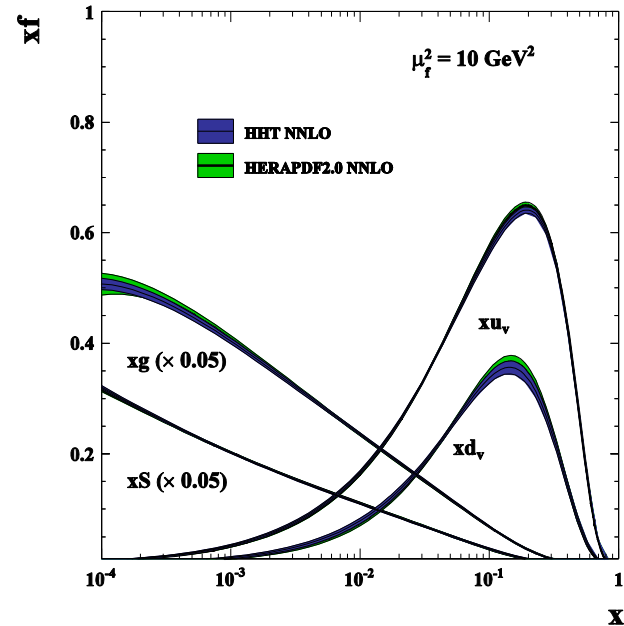
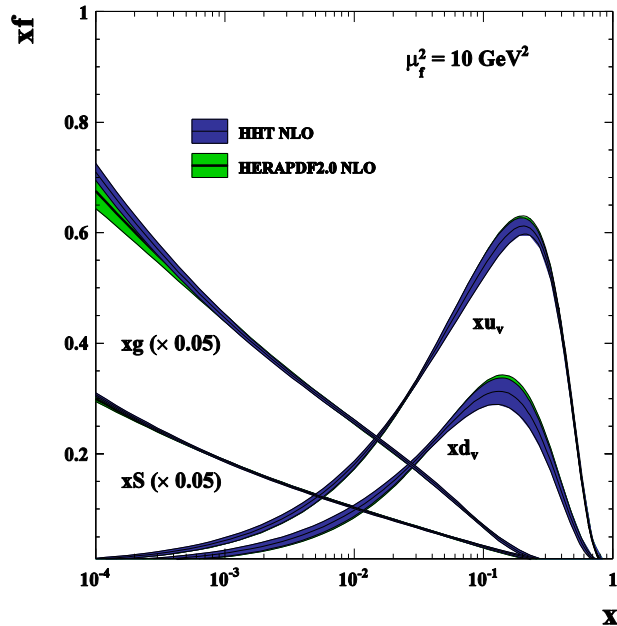
Fit at	with $Q_{\min}^2 = 3.5 \text{ GeV}^2$	HERAPDF2.0	HHT	$A_L^{\text{HT}}/\text{GeV}^2$	
NNLO	χ^2/ndof	1363/1131	1316/1130	5.5 ± 0.6	$\Delta\chi^2 = -47$
	χ^2/ndp for NC e^+p : $Q^2 \geq Q_{\min}^2$	451/377	422/377		
	χ^2/ndp for NC e^+p : $2.0 \text{ GeV}^2 \leq Q^2 < Q_{\min}^2$	41/25	32/25		
NLO	χ^2/ndof	1356/1131	1329/1130	4.2 ± 0.7	$\Delta\chi^2 = -28$
	χ^2/ndp for NC e^+p : $Q^2 \geq Q_{\min}^2$	447/377	431/377		
	χ^2/ndp for NC e^+p : $2.0 \text{ GeV}^2 \leq Q^2 < Q_{\min}^2$	46/25	46/25		

After HT is added the NNLO fit is better than the NLO fit

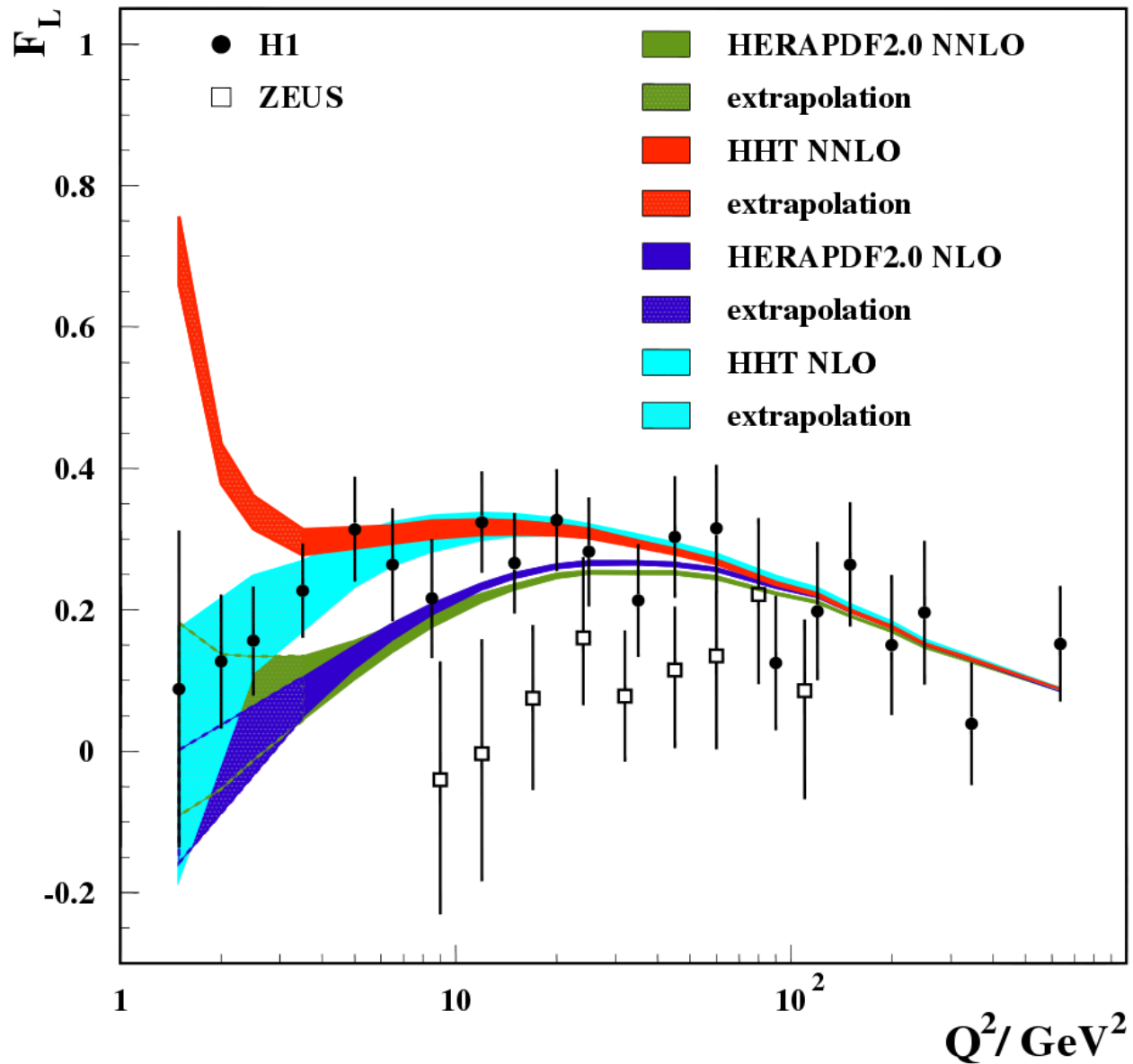
A substantial part of the improvement comes from the NCE+p 920 data

This persists even below the usual cut-off $Q_{\min}^2 = 3.5 \text{ GeV}^2$

NOTE: the HHT PDFs themselves barely change from HERAPDF2.0 – the higher twist modification does not affect high-scale LHC physics



The HHT fits give a larger F_L at low Q^2 for both NLO and NNLO



You might think that -since F_L is related to the gluon -

$$xG(x, Q^2) \approx \frac{3}{5} 5.9 \left[\frac{3\pi}{4\alpha_s} F_L(0.4x, Q^2) - \frac{1}{2} F_2(0.8x, Q^2) \right]$$

Simple LO relationship gives the idea

- an easier way to obtain larger F_L would be to drop the negative term in the gluon PDF parametrisation.

$$xg(x) = A_g x^{B_g} (1-x)^{C_g} - A'_g x^{B'_g} (1-x)^{C'_g},$$

So we did- we call this the alternative gluon (AG) parametrisation

This makes almost no difference for the NLO fits

Whereas it is strongly disfavoured for the NNLO fits.

At NNLO the fit wants a negative term in the gluon parametrization AND a higher twist term in F_L .

For HERAPDF2.0 AG the $\chi^2/\text{ndof} = 1389/1131$ cf 1363/1130 for the standard fit

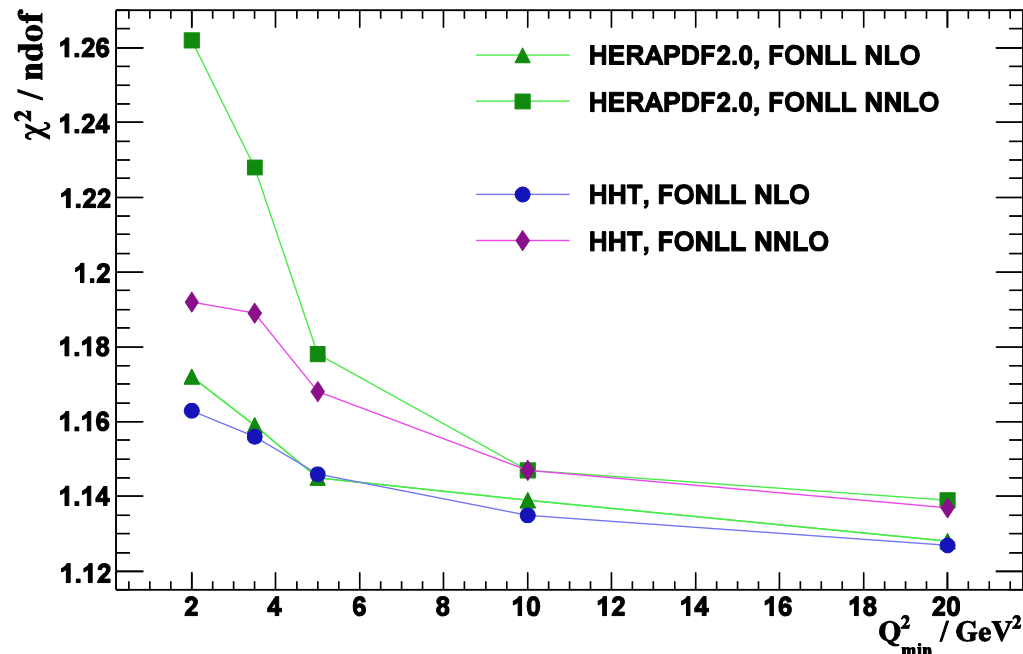
For HHT AG the $\chi^2/\text{ndof} = 1350/1130$ cf 1316/1130 for the standard fit

These two contributions clearly affect the fit in different ways

Another consideration is that we know that the rate of decrease χ^2/ndof with increasing Q_{\min}^2 differs with the heavy flavour scheme used AND with the order in α_s to which F_L is evaluated

So let's take a look at FONLL

For FONLL-C at NNLO a higher twist term in F_L brings a substantial decrease in the χ^2/ndof with a similar value of $A_L^{\text{HT}} = 6.0 \pm 0.7 \text{ GeV}^2$ to that for the RTOPT scheme. For FONLL-B at NLO a higher twist term in F_L brings almost no decrease in χ^2/ndof . This is probably related to the order in α_s to which F_L is evaluated



For FONLL-C/RTOPT at NNLO, F_L is evaluated to $O(\alpha_s^2)/O(\alpha_s^3)$

For FONLL-B/RTOPT at NLO, F_L is evaluated to $O(\alpha_s)/O(\alpha_s^2)$

The value of F_L at $O(\alpha_s)$ is relatively large in any scheme and thus there is little need for higher twist.

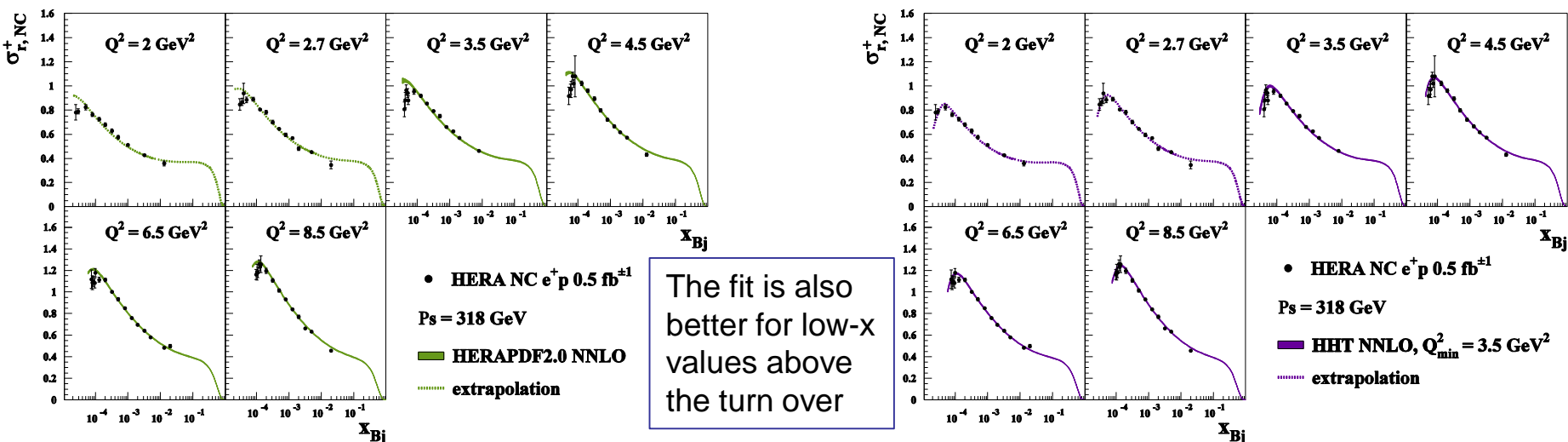
However as soon as F_L is evaluated to $O(\alpha_s^2)$ or higher the need for higher twist appears

So now let's look at why the HHT fits do so well

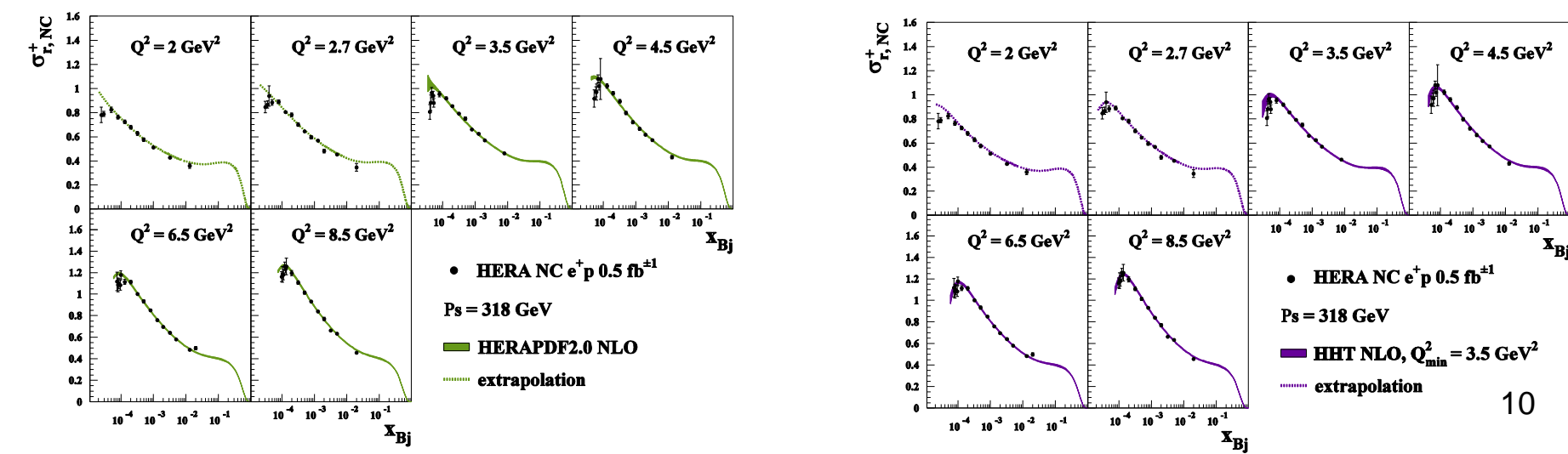
It is because they describe the turn over of the cross section at low x, Q2 much better

$$\sigma_{red} = F_2 - y^2/Y_+ F_L$$

The data clearly wants a larger F_L and this is what the higher twist term provides



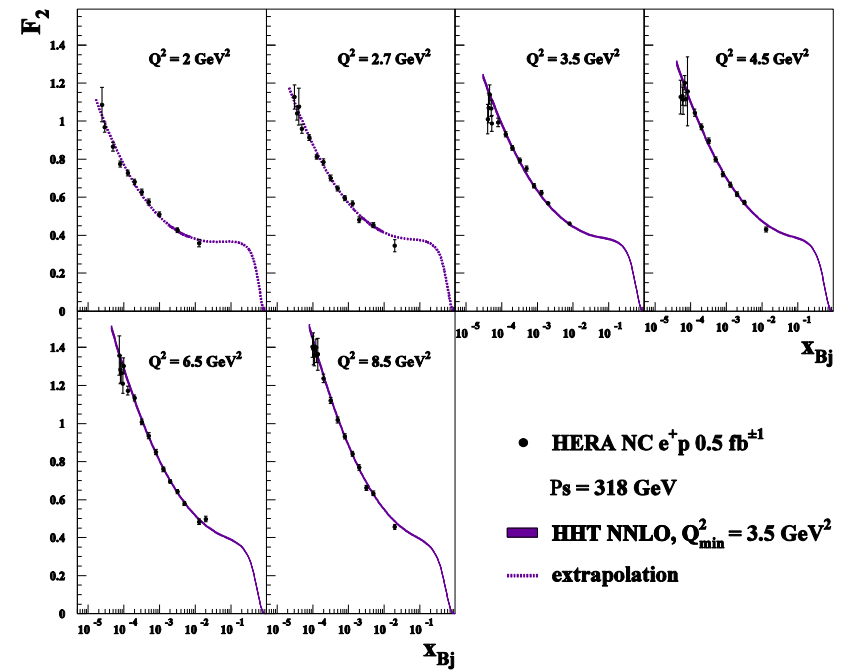
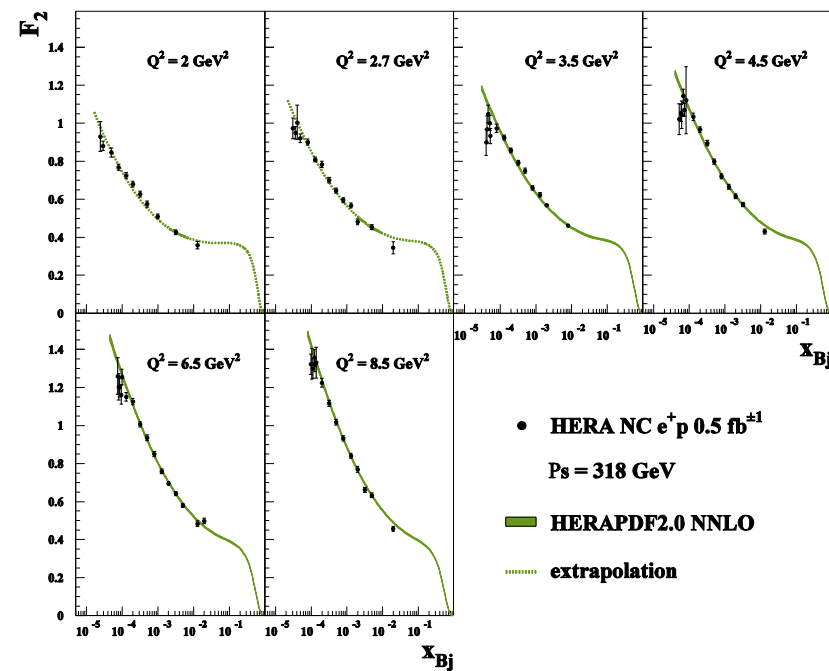
You can also see that NNLO does better than NLO



It is also interesting to look at F_2 , where the data points are extracted as

$$F_2^{\text{extracted}} = F_2^{\text{predicted}} \sigma_{\text{red}}^{\text{measured}} / \sigma_{\text{red}}^{\text{predicted}}$$

Since F_2 is the dominant part of the reduced cross section this is a reasonable procedure



This essentially means that we get F_2 by correcting σ_{red} with our predicted F_L

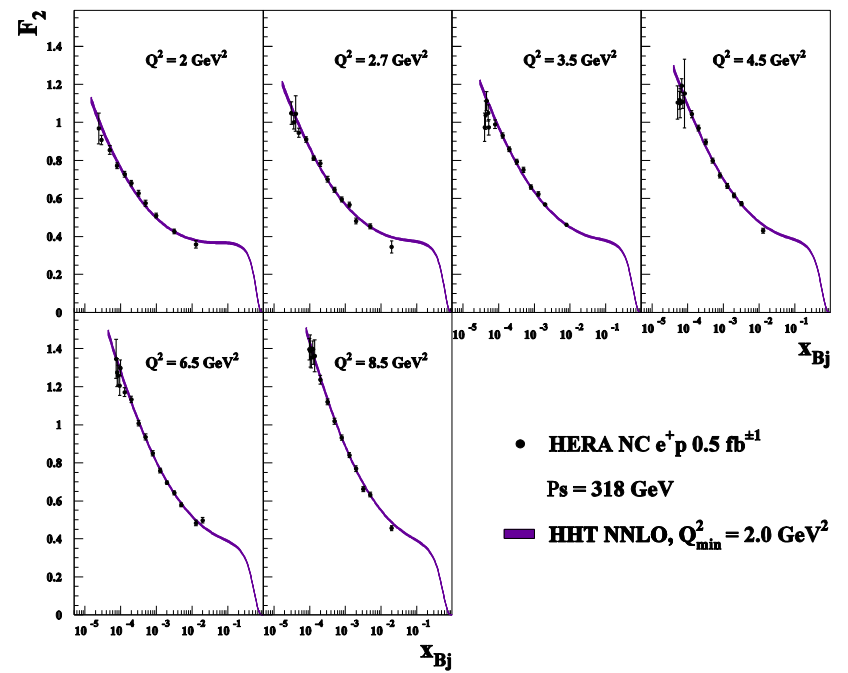
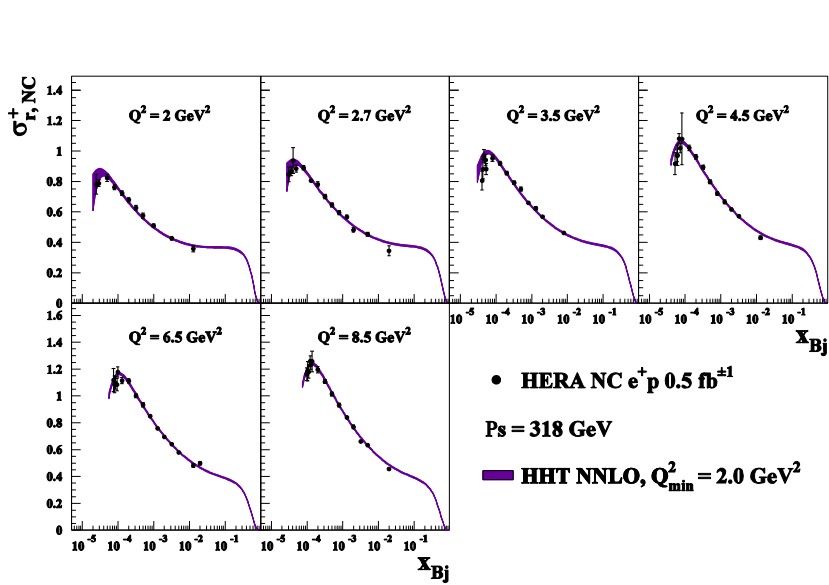
$$F_2 = \sigma_{\text{red}} + y^2/Y_+ F_L$$

If our predicted F_L is too small the F_2 will also be too small and this is what we see in HERAPDF2.0 F_2 at low x, Q^2 . **The extracted F_2 takes a turn over!**

This is not what the pQCD F_2 predictions say.

If we use the HHT predictions for F_L then the F_2 extracted is much closer to the F_2 predictions— and note these F_2 predictions are very similar for HERAPDF2.0 and HHT because they depend ONLY on the very similar PDFs.

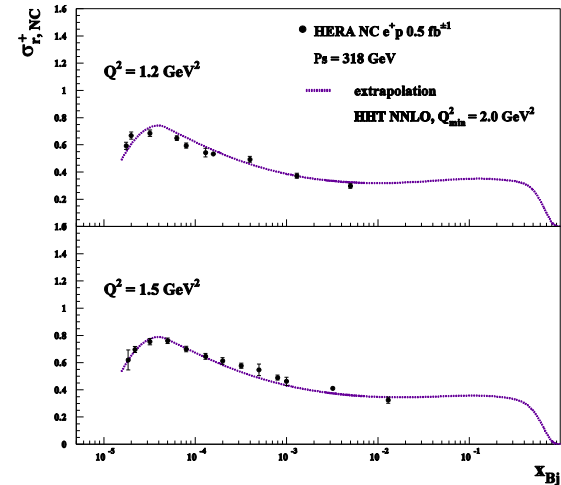
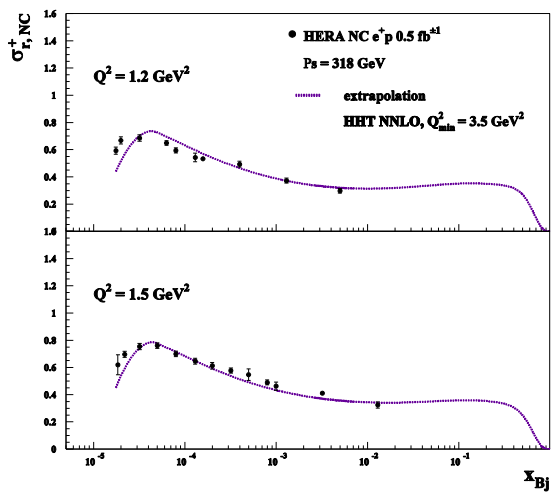
(The picture is similar but not quite so good for NLO- see back-up)



Looking at the extrapolations of our fits below $Q_{\min}^2 = 3.5 \text{ GeV}^2$ made us bold enough to extend the fit down to $Q_{\min}^2 = 2.0 \text{ GeV}^2$

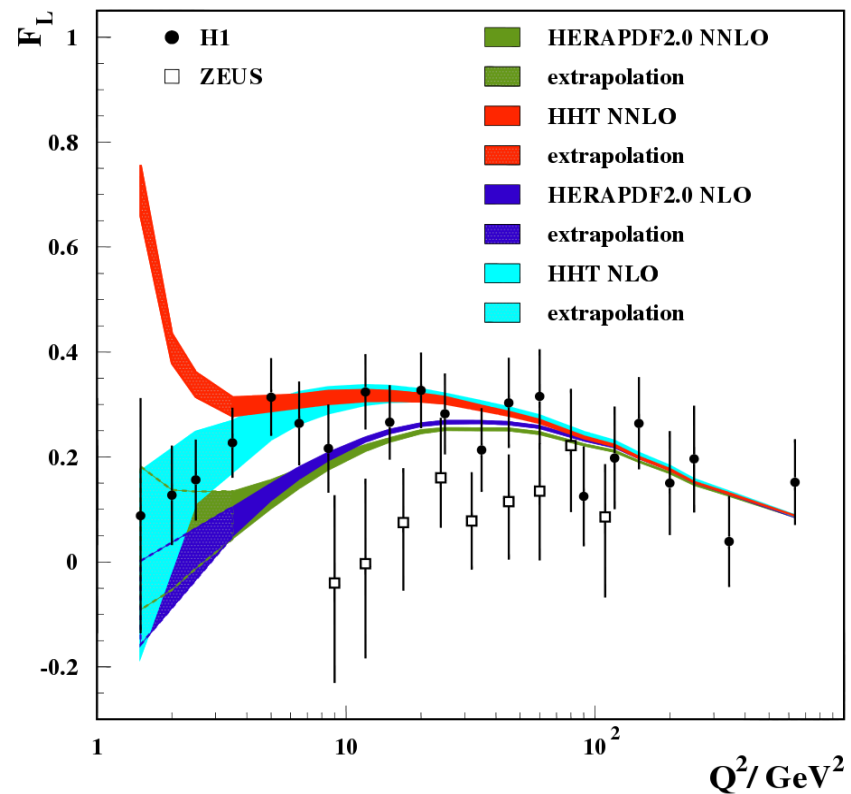
Fit at	with $Q_{\min}^2 = 2.0 \text{ GeV}^2$	HERAPDF2.0	HHT	$A_L^{\text{HT}}/\text{GeV}^2$
NNLO	χ^2/ndof	1437/1171	1381/1170	5.2 ± 0.7
	χ^2/ndp for NC e^+p : $Q^2 \geq Q_{\min}^2$	486/402	457/402	
	χ^2/ndp NC e^+p : $Q_{\min}^2 \leq Q^2 < 3.5 \text{ GeV}^2$	31/25	26/25	
NLO	χ^2/ndof	1433/1171	1398/1170	4.0 ± 0.6
	χ^2/ndp for NC e^+p : $Q^2 \geq Q_{\min}^2$	487/402	466/402	
	χ^2/ndp NC e^+p : $Q_{\min}^2 \leq Q^2 < 3.5 \text{ GeV}^2$	40/25	31/25	

Not much changes for the NNLO fit and the NLO fit improves a little
See back-up



So we got even bolder and looked at lower Q^2 - by backward evolution

But beware...is this actually reasonable?
 What does FL itself look like?



NNLO HHT F_L prediction is becoming untamed at low Q^2 — this approach cannot be pushed too far.
 This comes from NNLO coefficient functions and the $1/Q^2$ term just makes it worse

Another interesting way to look at this is by looking at plots of F_2 and F_L at fixed W as a function of Q^2 (This is the Golec-Biernat Wusthoff dipole model way of looking at it)

First look at the upper three curves for F_2

Compare the HHT F_2 extracted points to the F_2 predictions – the description is good. Then compare the HERAPDF2.0 F_2 extracted points to the F_2 predictions the description is not so good.

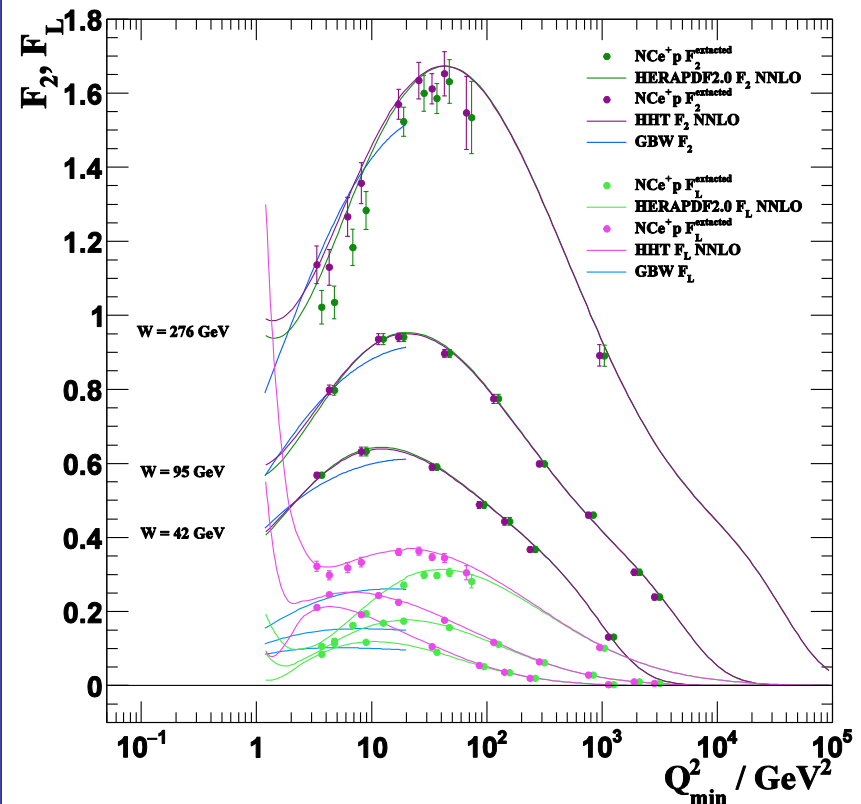
This is essentially what we saw in the F_2 curves on slide 11 but it emphasizes that the discrepancy comes at low x . Only the top curve $W=276$ GeV involves data at really low x

$$x = Q^2 / (W^2 + Q^2)$$

Now look at the lower three curves for F_L

The predictions for HHT go crazy at very low Q^2 .

In fact this upturn happens in HERAPDF as well- and it is starting to happen in F_2 . It is a feature of the low- x coefficient functions



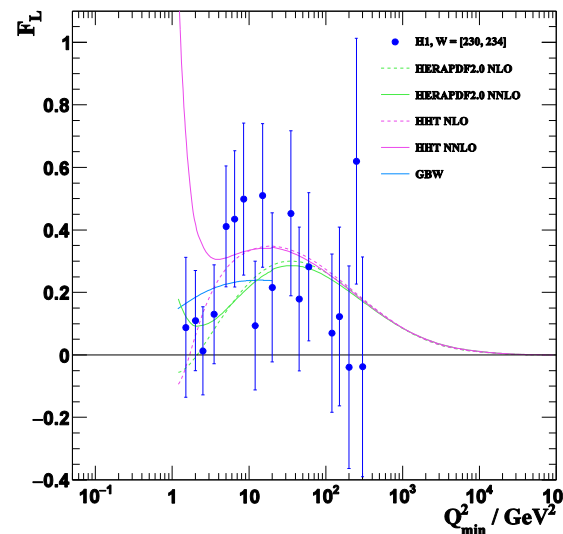
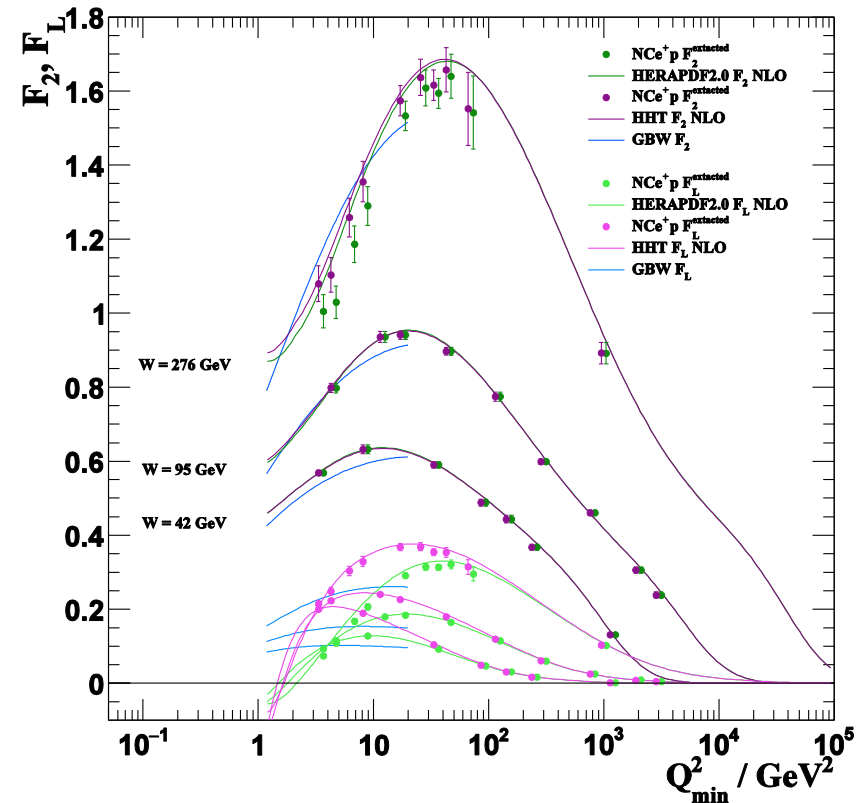
Here the extracted F_L points are got from $F_L^{\text{extracted}} = F_L^{\text{predicted}} \sigma_{\text{red}}^{\text{measured}} / \sigma_{\text{red}}^{\text{predicted}}$. Since F_L is not the dominant part of the reduced cross section these cannot be considered as measurements and they simply follow the predictions

It is not just the NNLO F_L which is becoming unacceptable at low Q^2 , the NLO predictions also have problems. They are becoming negative. This is not allowed for a structure function (as opposed to a PDF)

The GBW predictions at both NNLO and NLO are also compared to the extracted data points in these figures. They are broadly compatible with the HHT predictions for F_2 for $Q^2 < 10 \text{ GeV}^2$

Finally we look at the F_L predictions for HERAPDF2.0 and HHT at NNLO as compared to the H1 direct measurements at $W = 232 \text{ GeV}$.

The data are able to exclude the extreme behaviour of the HHT prediction for $Q^2 < 2.0 \text{ GeV}^2$



Summary

A study of adding higher twist terms to the HERAPDF2.0 analysis of the HERA-I+II data for NLO and NNLO fits

Such terms are significant in F_L for $2 < Q^2 < 50 \text{ GeV}^2$

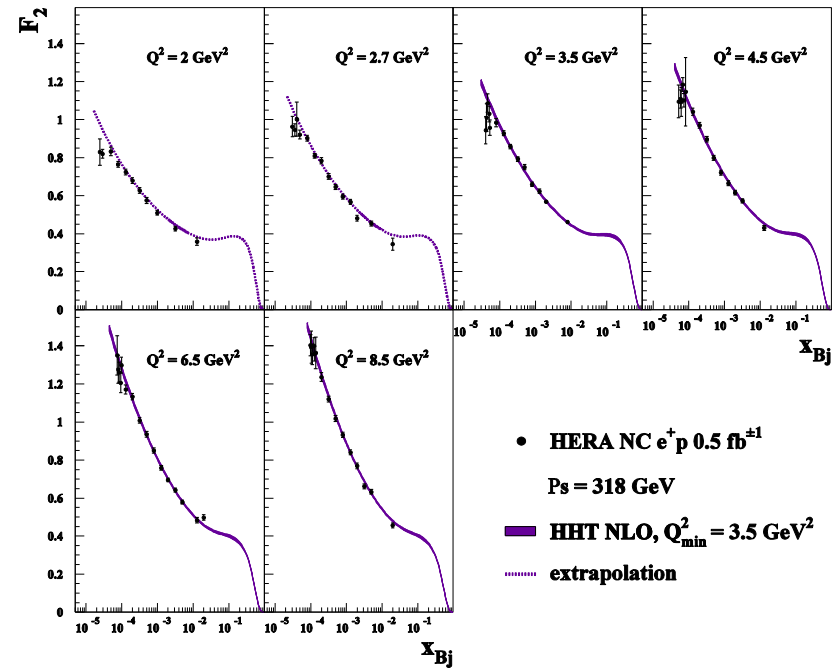
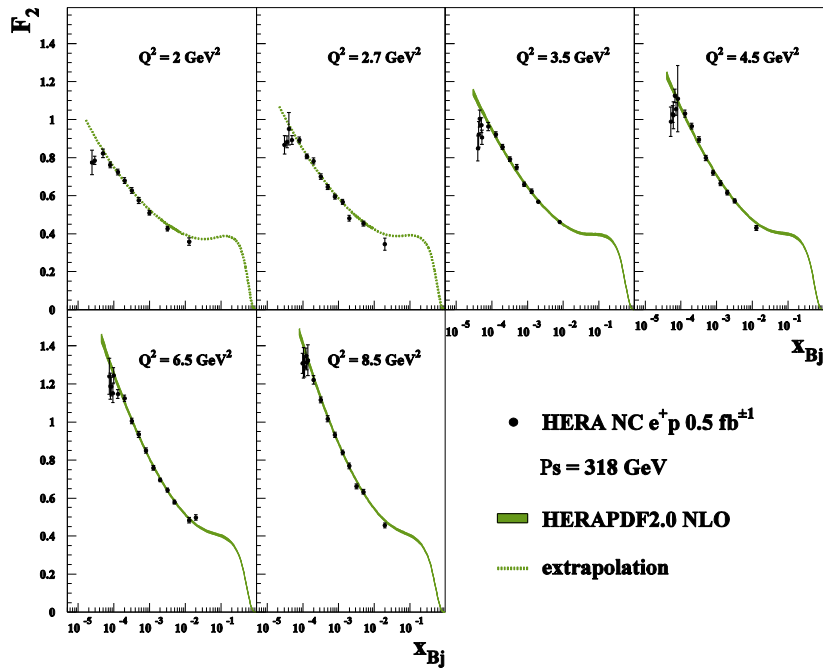
Improves the χ^2 significantly, and makes NNLO fits clearly better than NLO

Does not change the HERAPDF2.0 NLO or NNLO significantly- no change at higher Q^2

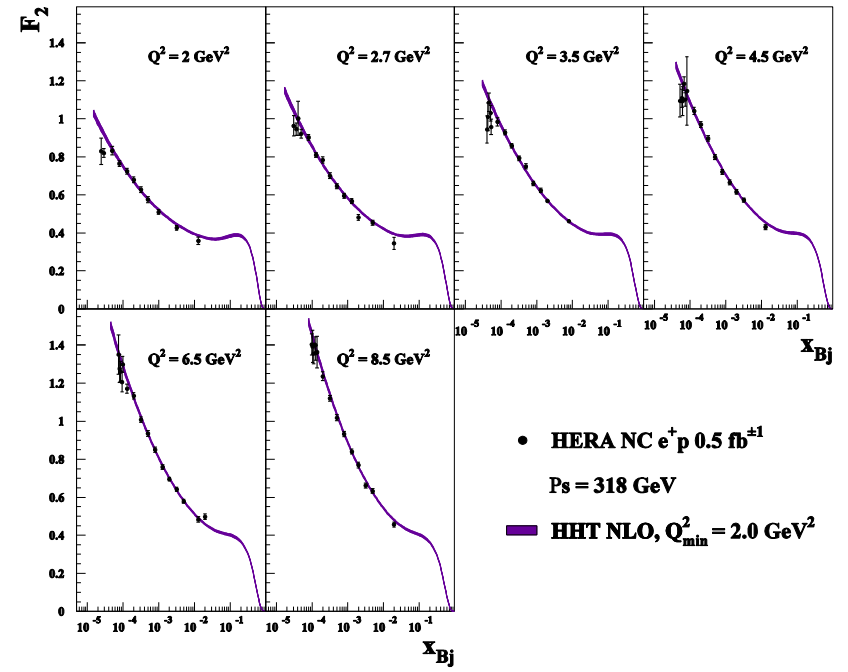
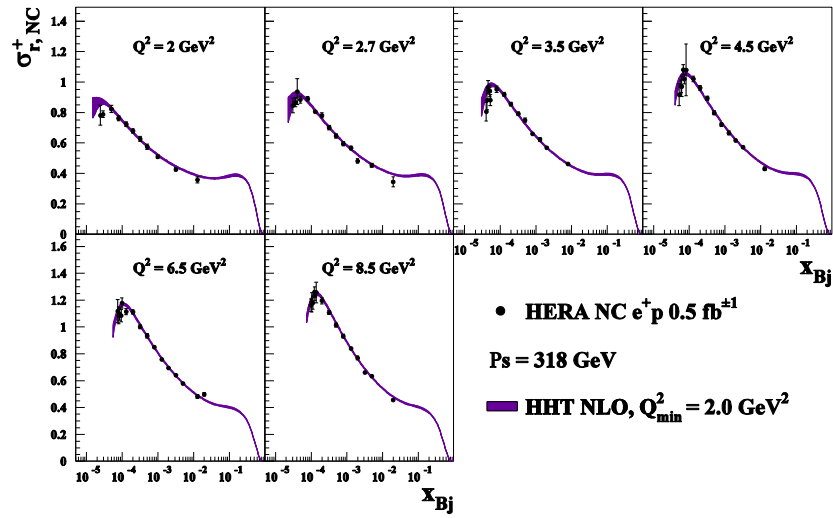
Higher twist terms are important for low Q^2 , which for HERA kinematics means at low- x

But such a simple approach fails for $Q^2 < 2 \text{ GeV}^2$

Back-up



And at NLO –the F_2 down to
 $Q_{\min}^2 = 3.5$



And at NLO down to $Q_{\min}^2 = 2.0$