

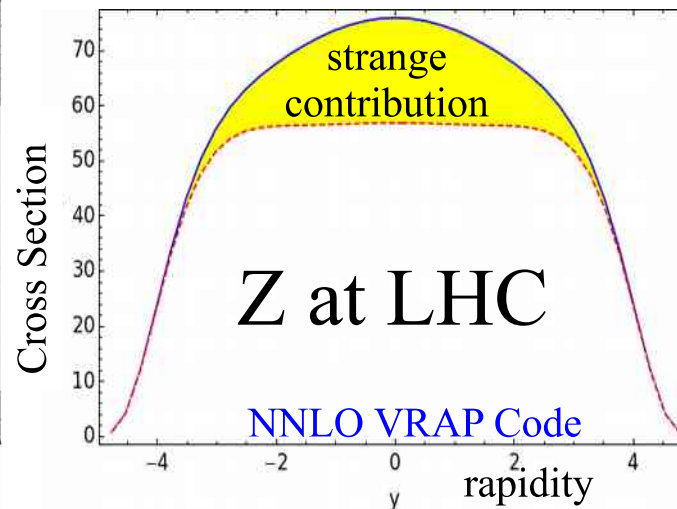
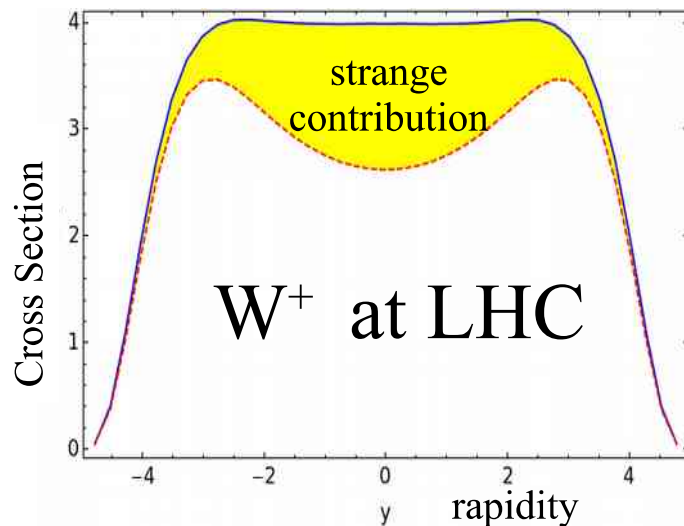
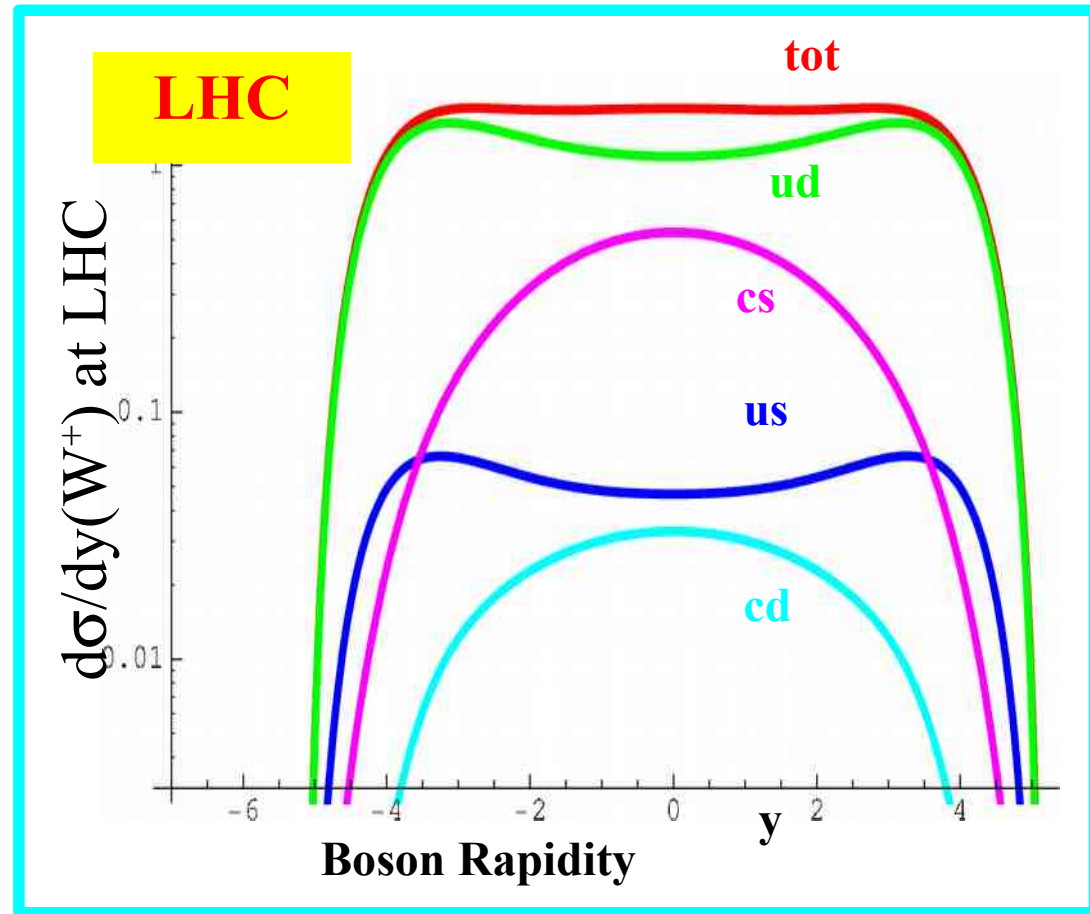
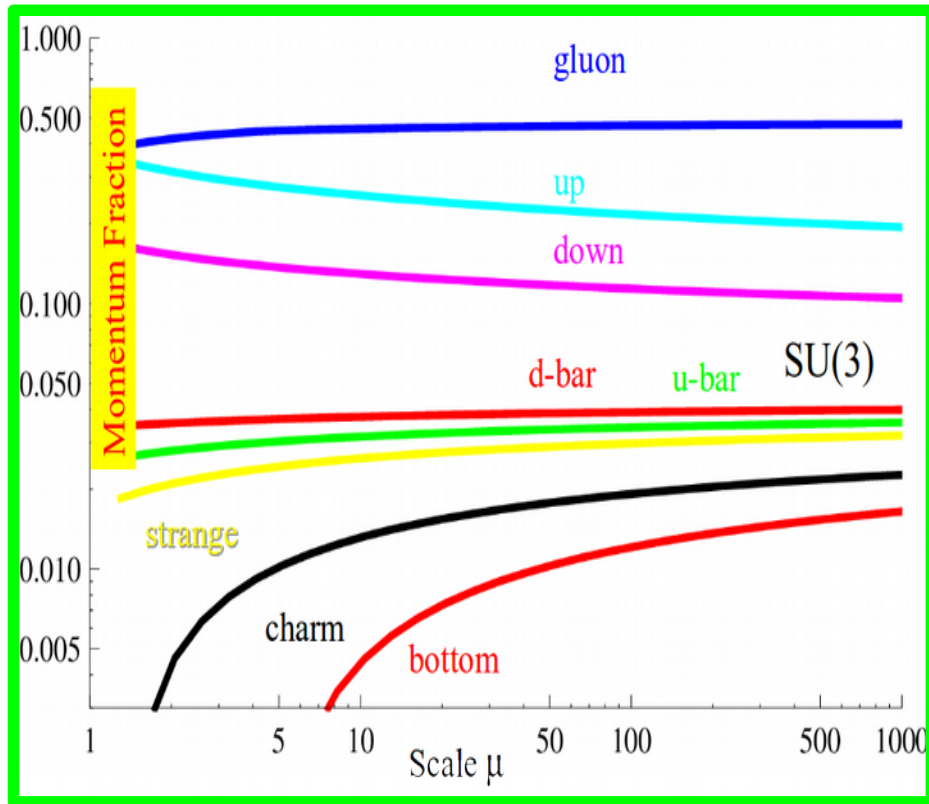
# Heavy flavors ( $s, c, b$ ) in LHC processes

Fred Olness

SMU

Thanks to:

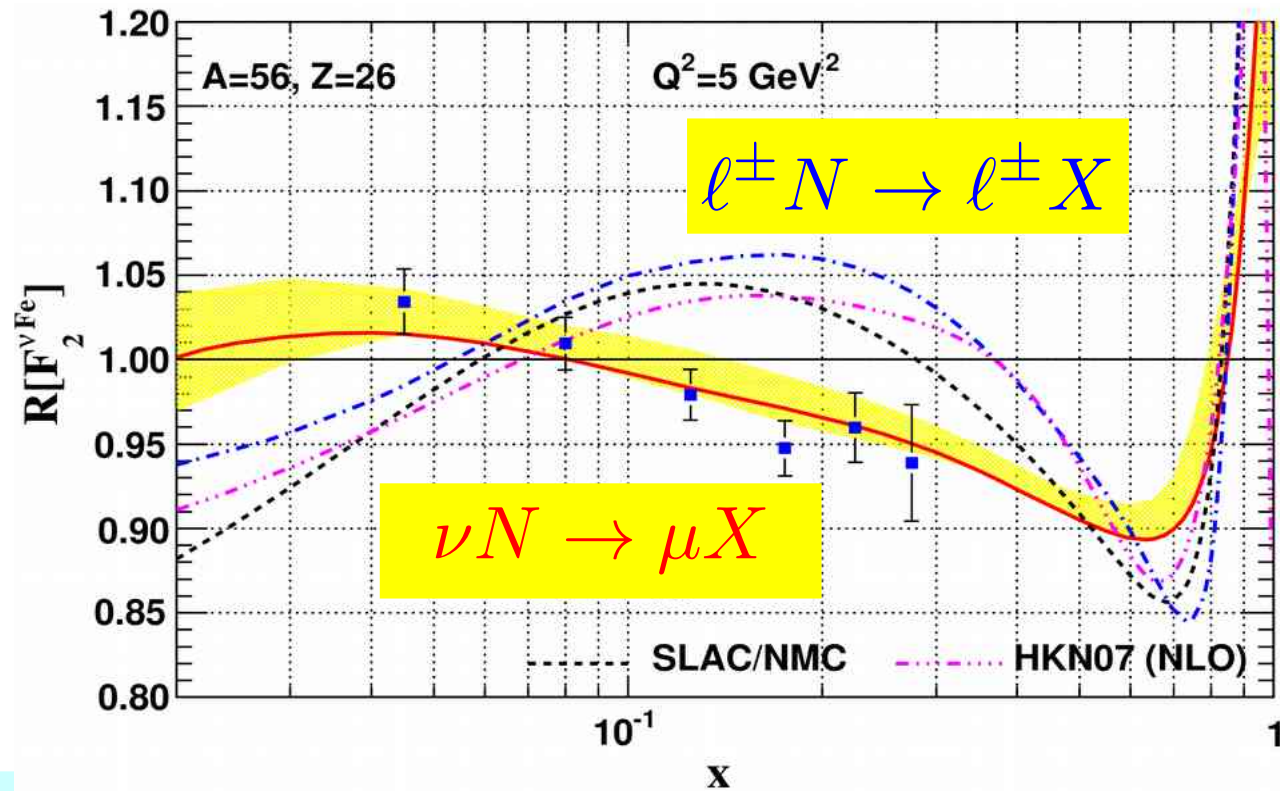
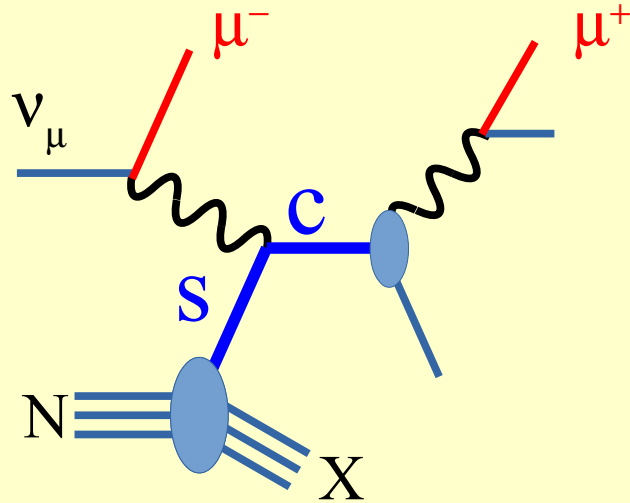
A. Kusina, F. Lyonnet, B. Clark, E. Godat, I. Schienbein, K. Kovarik, J.Y. Yu, T. Jezo, J.G. Morfin, J.F. Owens, P. Nadolsky, M. Guzzi, V. Radescu, C. Keppel



NNLO VRAP Code  
Anastasiou, Dixon, Melnikov, Petriello,  
Phys.Rev.D69:094008,2004.

Kusina, Stavreva, Berge, Olness,  
Schienbein, Kovarik, Jezo, Yu, Park  
Phys.Rev. D85 (2012) 094028

$$\nu N \rightarrow \mu^+ \mu^- X$$



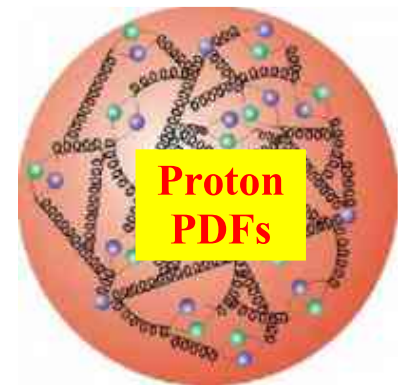
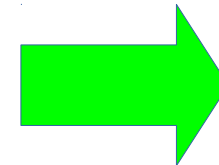
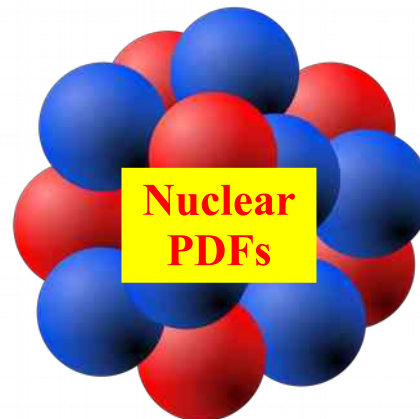
$$F_2^\nu \sim [d + s + \bar{u} + \bar{c}]$$

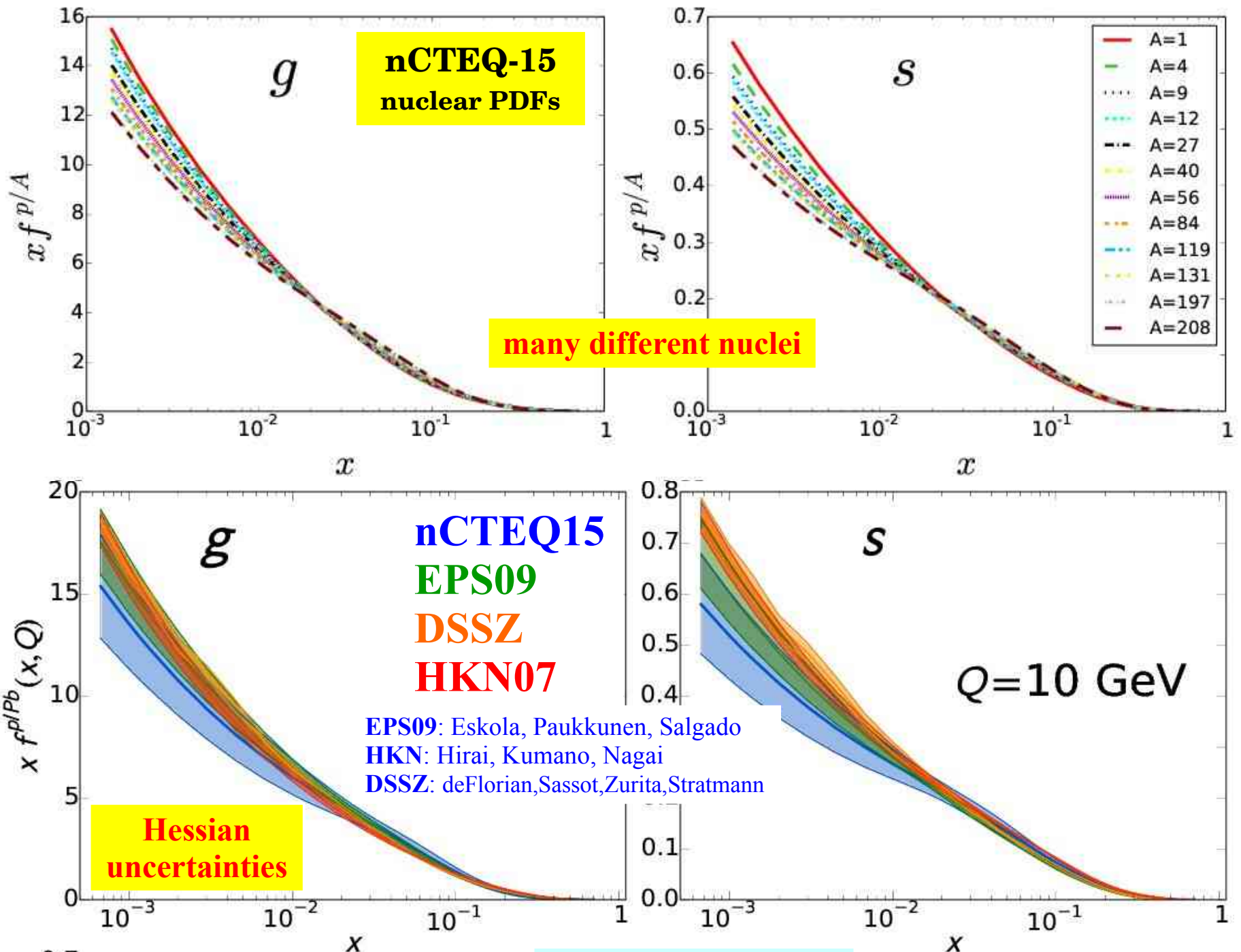
$$F_2^{\bar{\nu}} \sim [\bar{d} + \bar{s} + u + c]$$

$$F_3^\nu = 2[d + s - \bar{u} - \bar{c}]$$

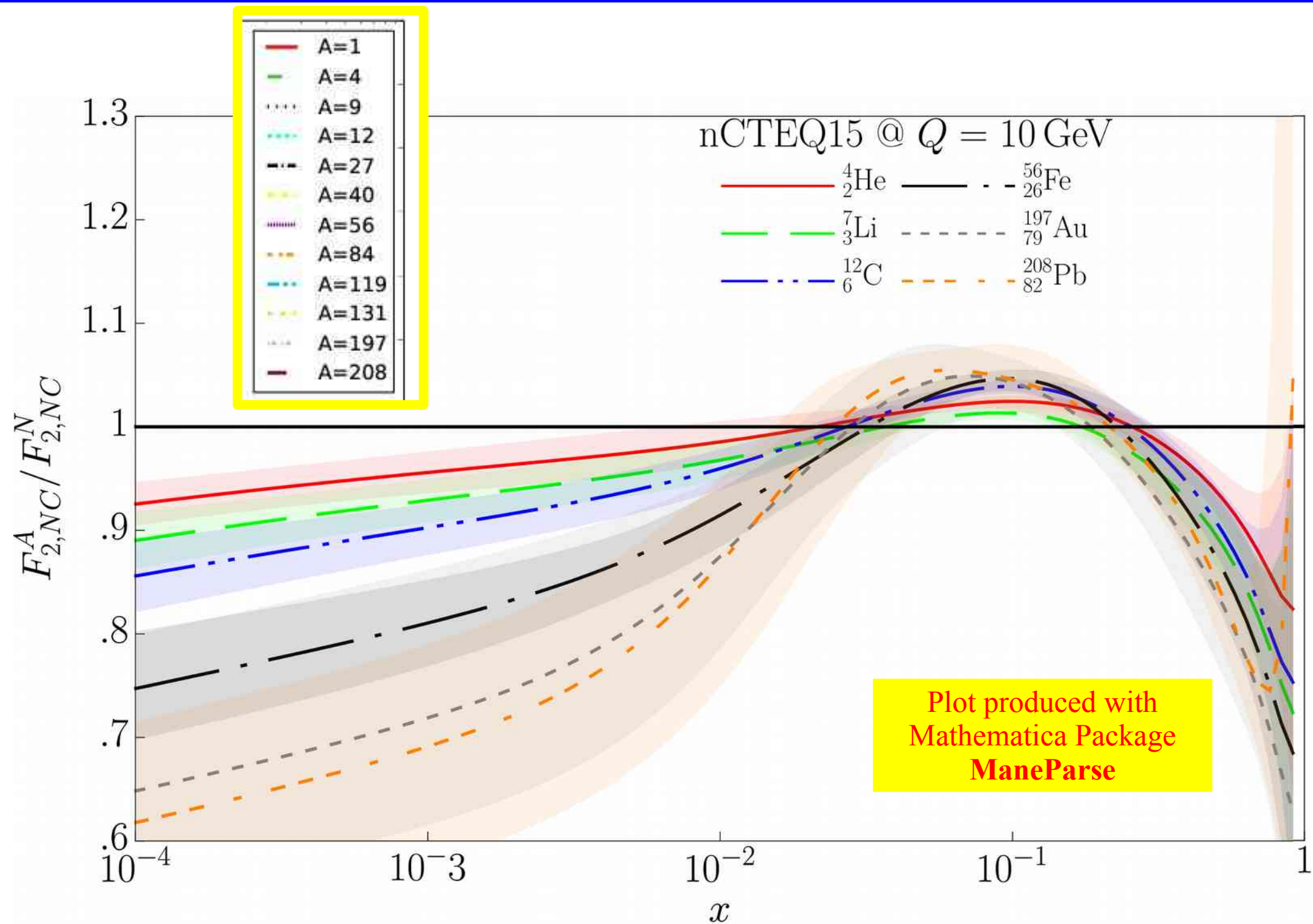
$$F_3^{\bar{\nu}} = 2[u + c - \bar{d} - \bar{s}]$$

Depends on nuclear corrections

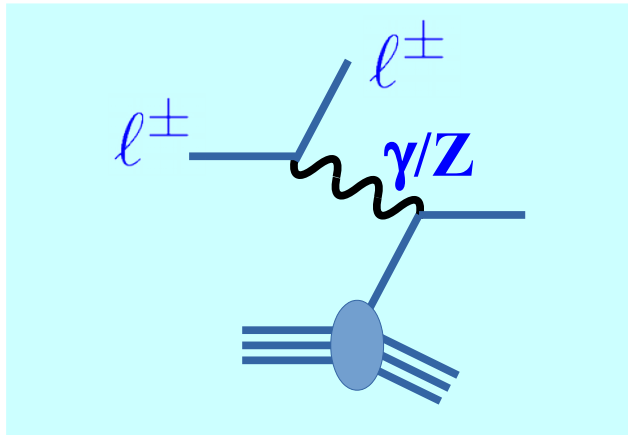








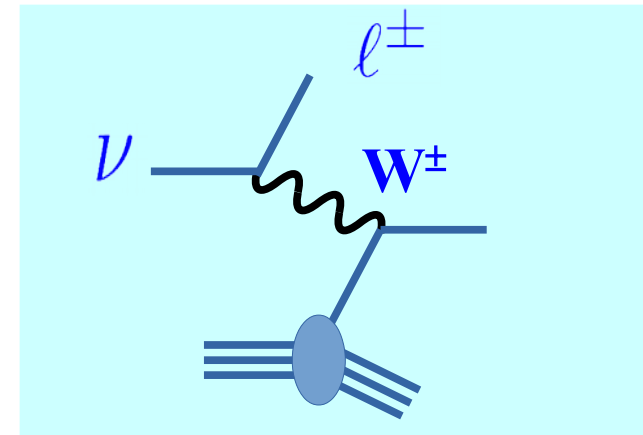
## Charged Lepton DIS



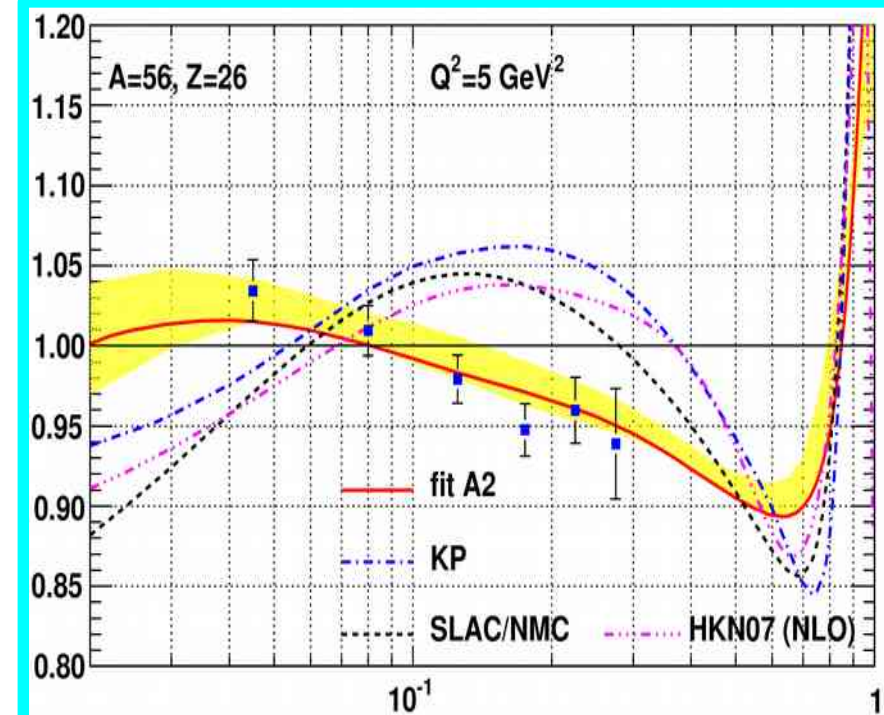
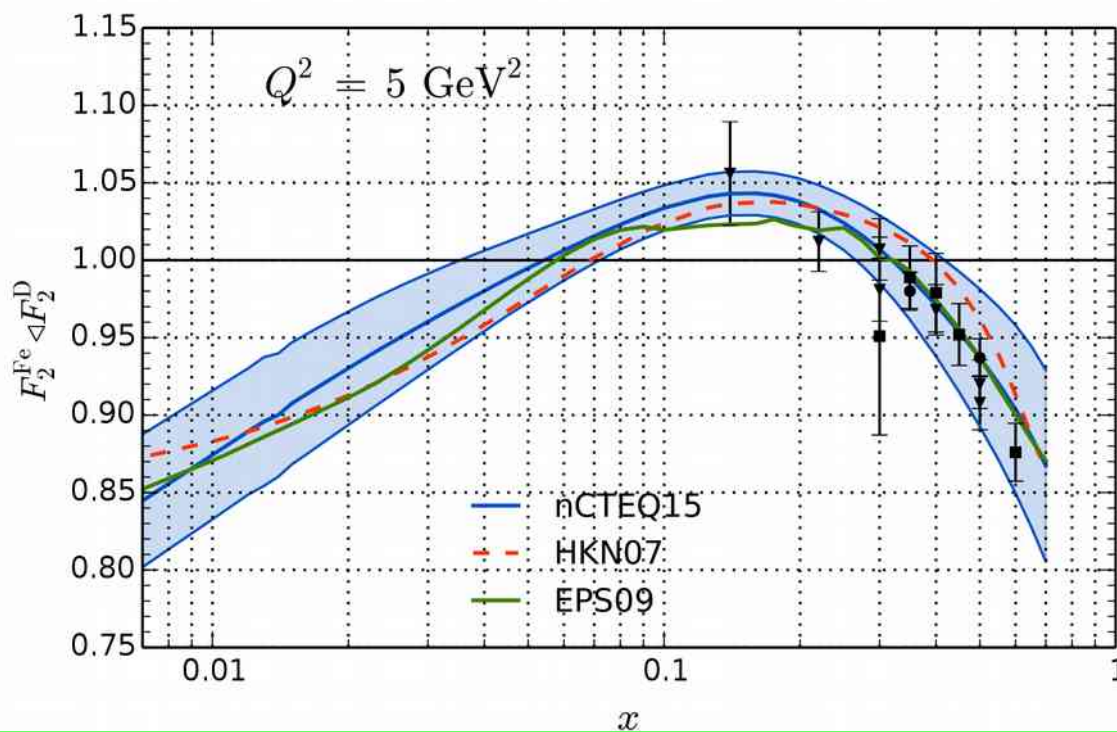
Very different  
nuclear  
corrections

$$\ell^\pm N \rightarrow \ell^\pm X$$

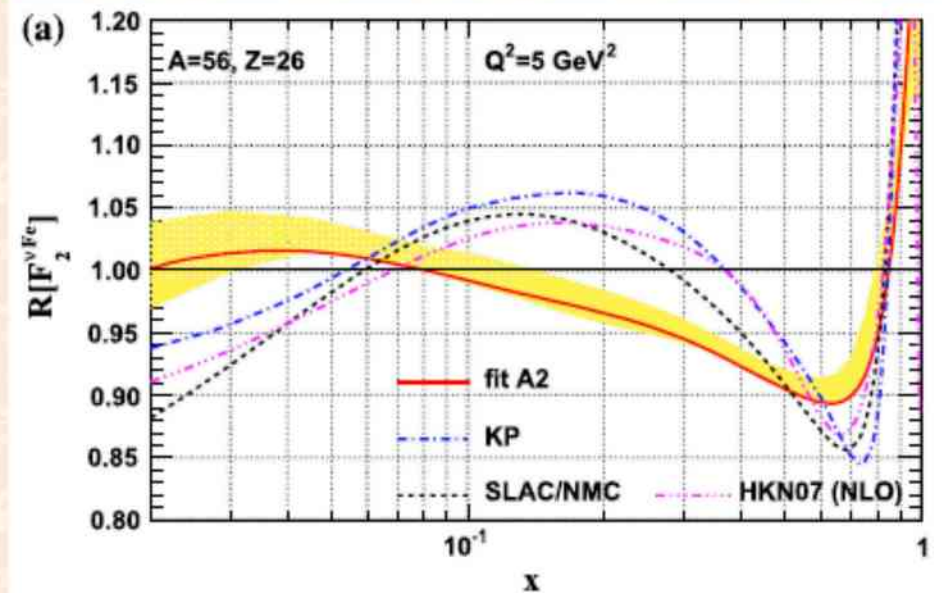
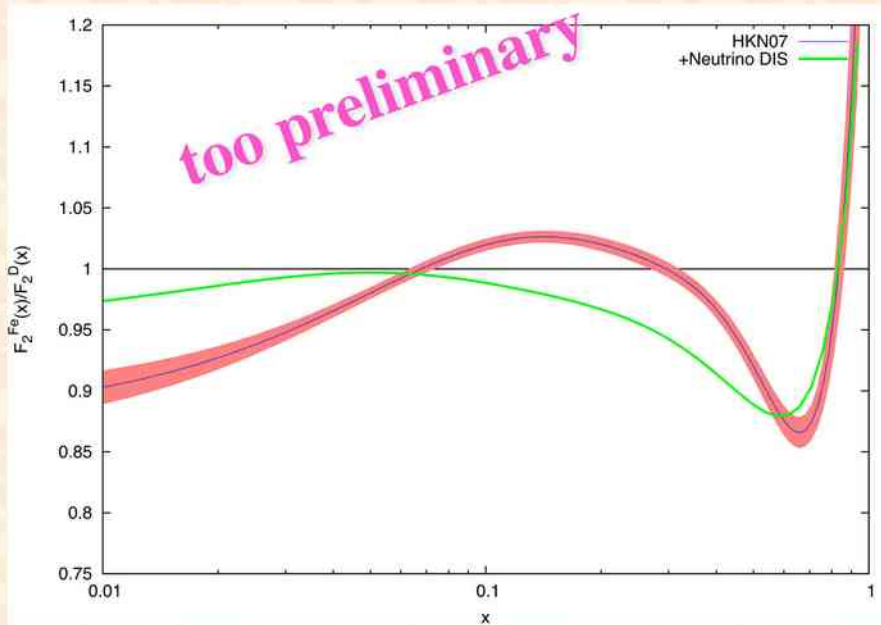
## Neutrino DIS



$$\nu N \rightarrow \ell^\pm X$$



# Our research in progress (M. Hirai, SK, K. Saito)

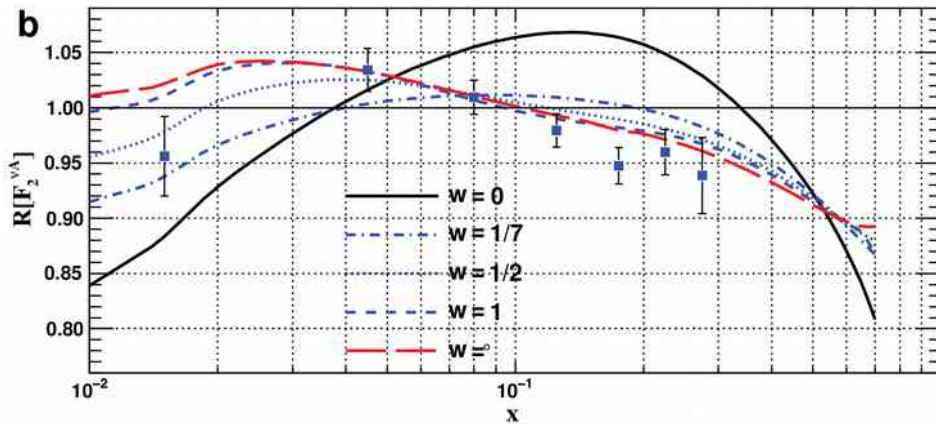


We are getting a similar modification to the nCTEQ one.

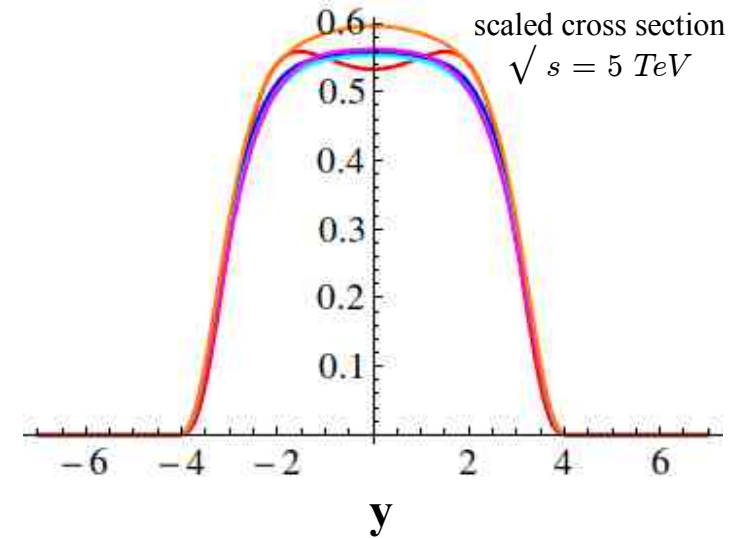


Interpolate between  $\nu N$  and  $\ell^\pm N$

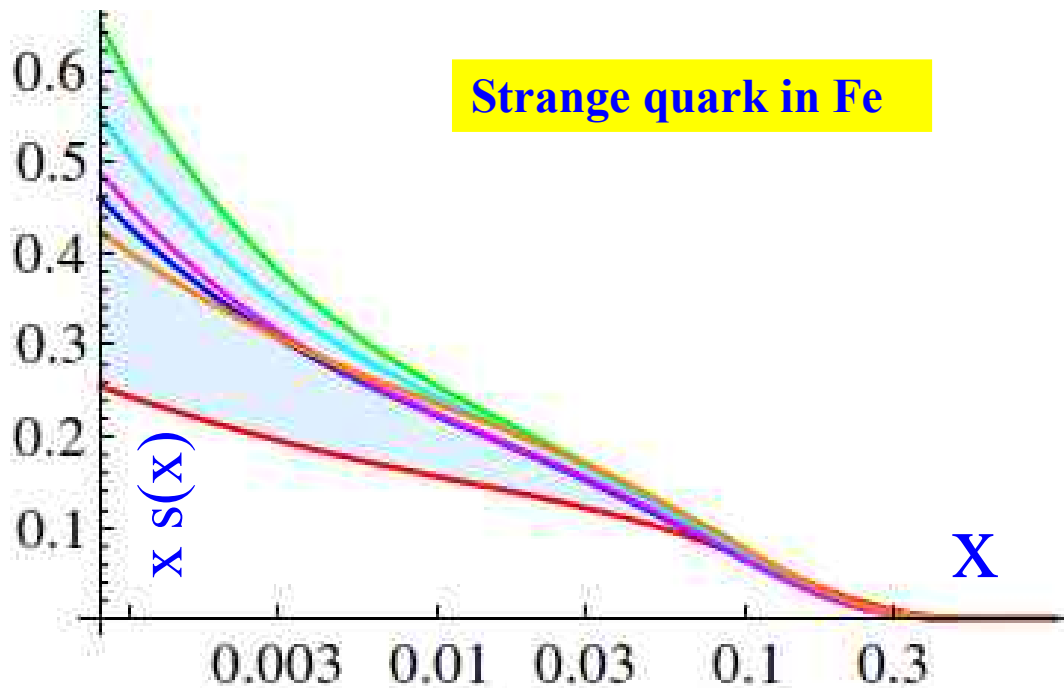
Nuclear Correction



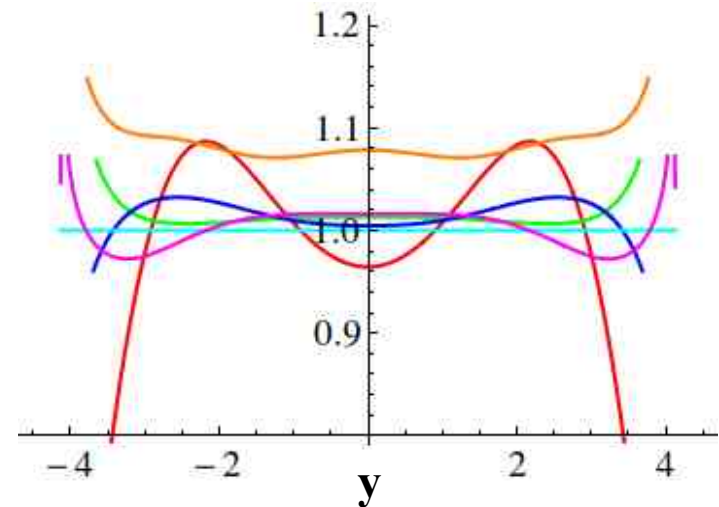
$W^+$  Production:  $d\sigma/dy$



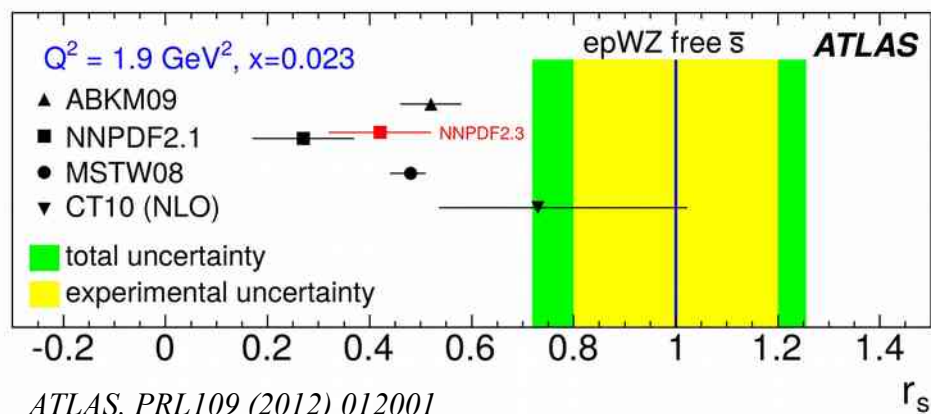
Strange quark in Fe



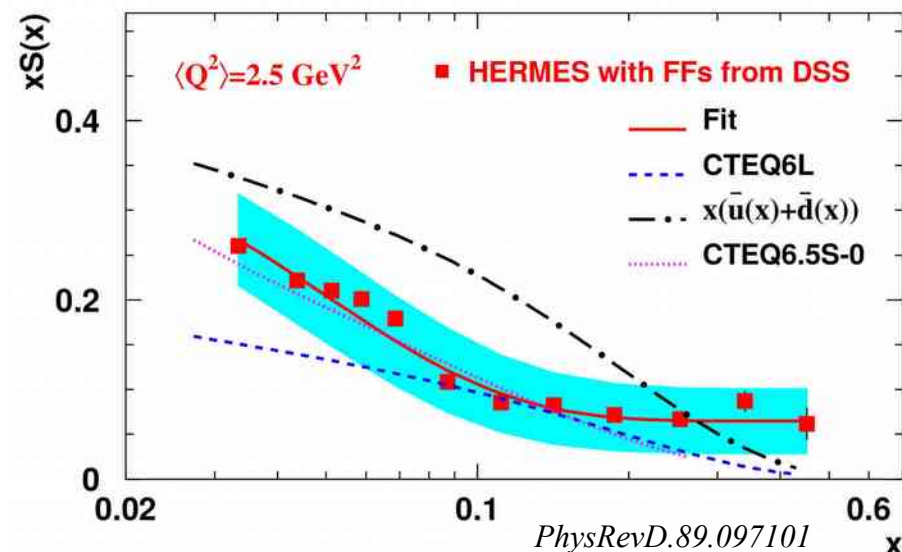
$W^+$  Production: ratio







$$r_s = 0.5 \frac{s + \bar{s}}{\bar{d}}$$



$$\kappa_s = \frac{\int_0^1 x(s + \bar{s}) dx}{\int_0^1 x(\bar{u} + \bar{d}) dx}$$

CMS:  $\kappa_s = 0.52$

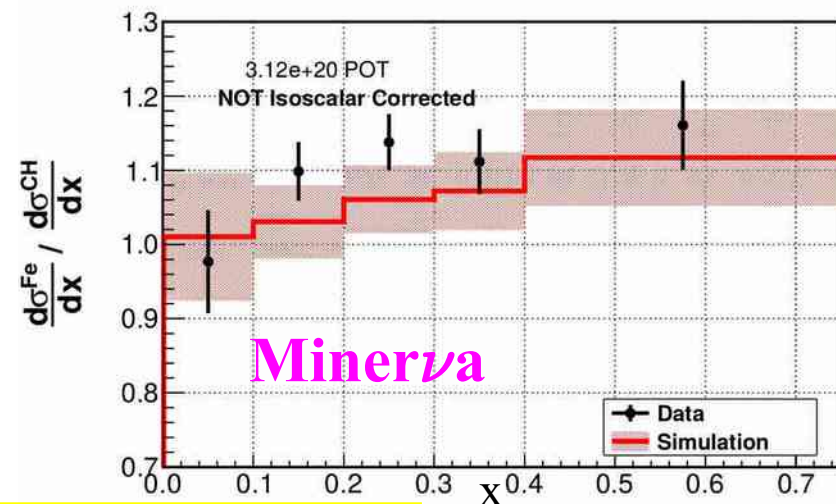
PRD 90, 032004 (2014)

NOMAD:  $\kappa_s = 0.591$

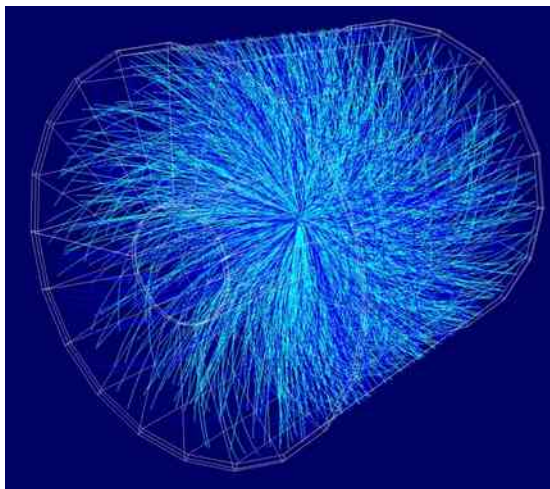
NPB876, 339 (2013).

NuTeV:  $\kappa_s = 0.58$

PRL99 (2007) 192001

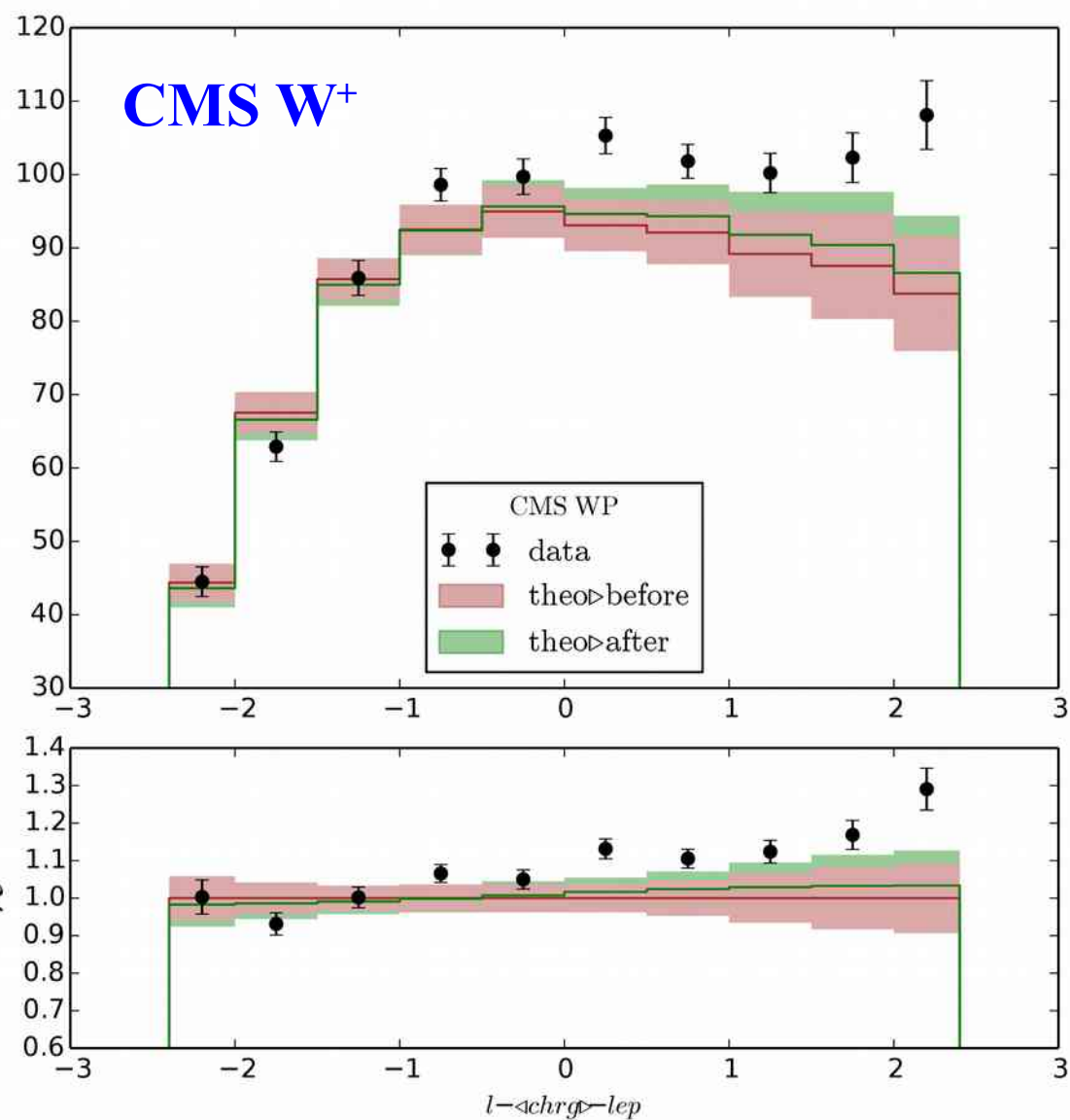
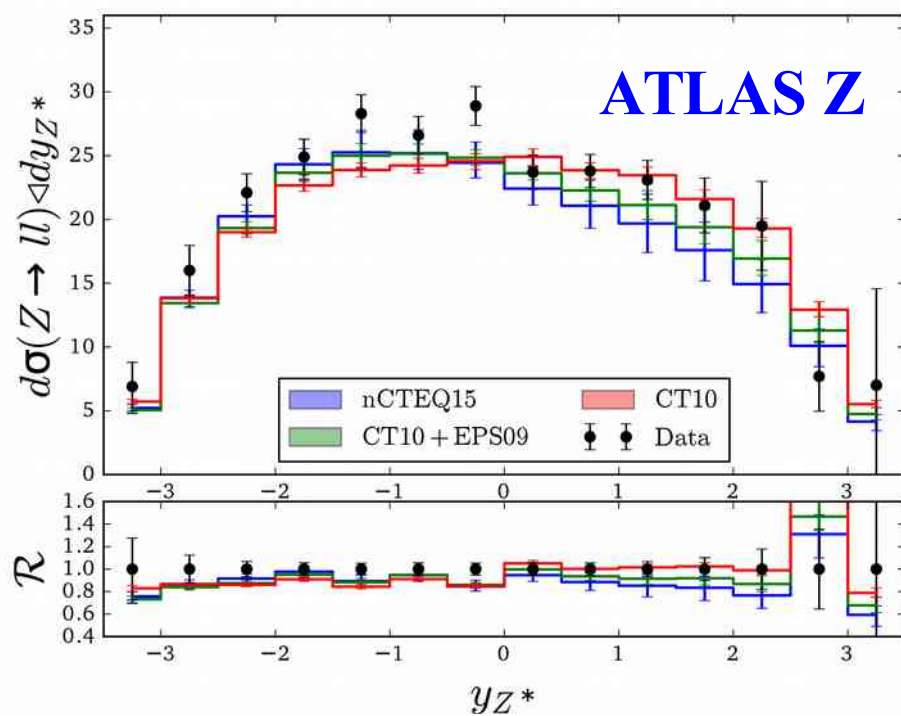


more data coming

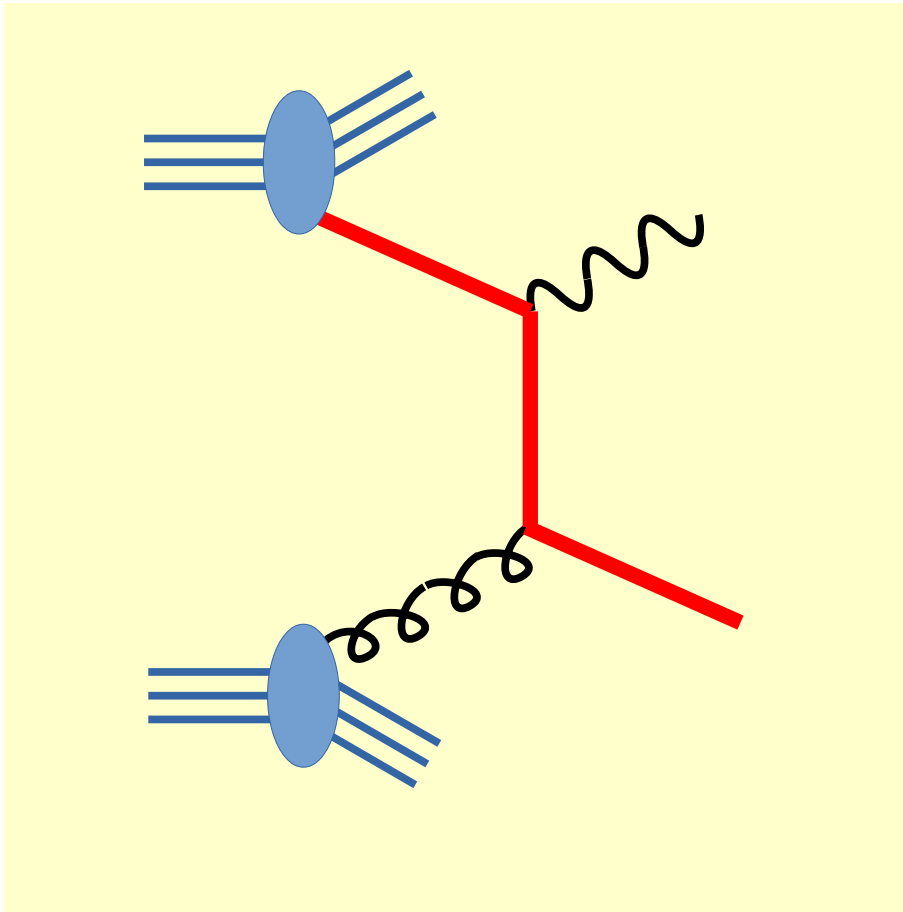


PDF re-weighting with CMS data

Compare with rapidity distribution



$$p + p \rightarrow V + Q$$



### Charged Current

$$s \ g \rightarrow c \ W$$

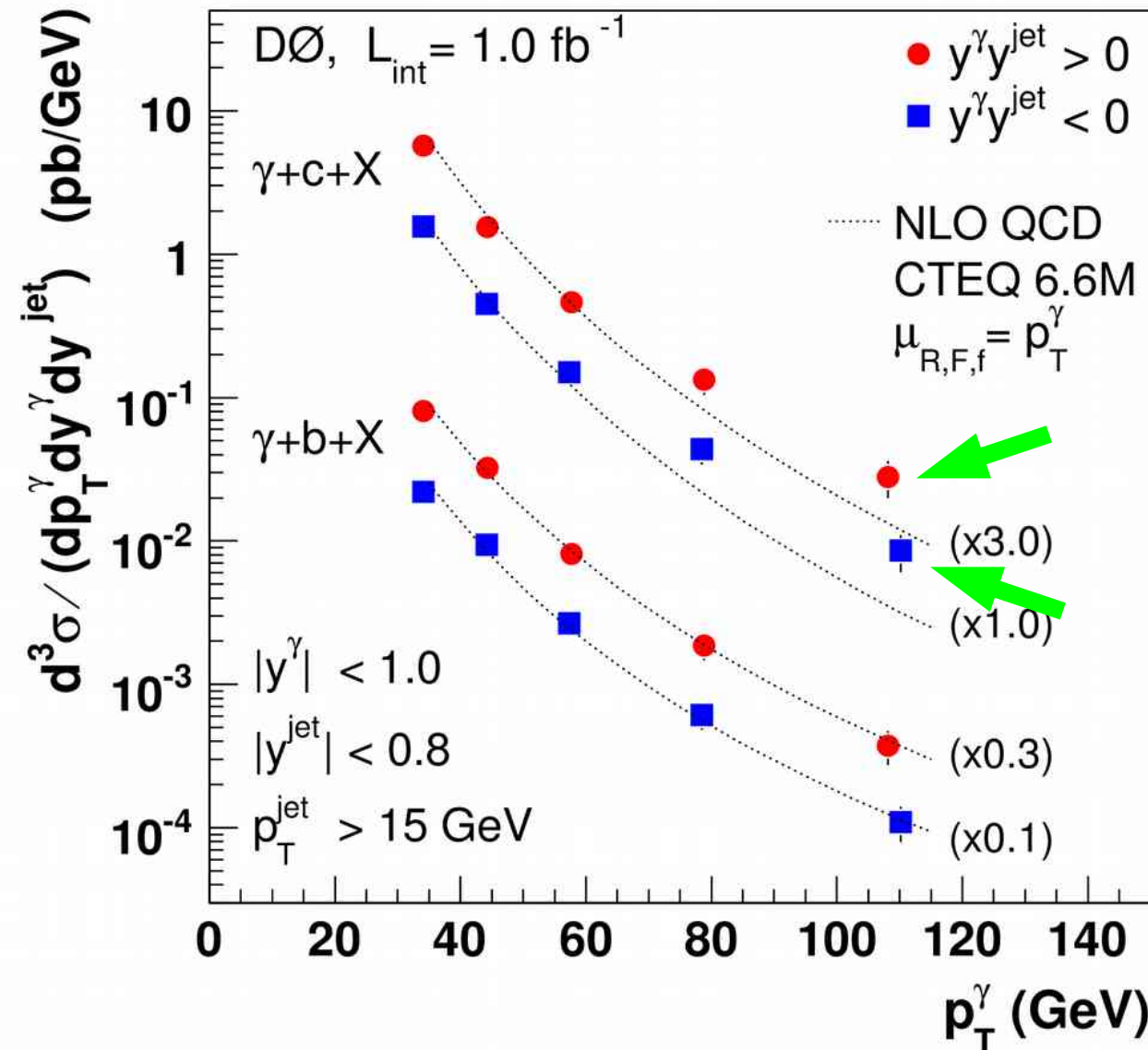
$$c \ g \rightarrow b \ W$$

### Neutral Current

$$c \ g \rightarrow c \ \gamma/Z$$

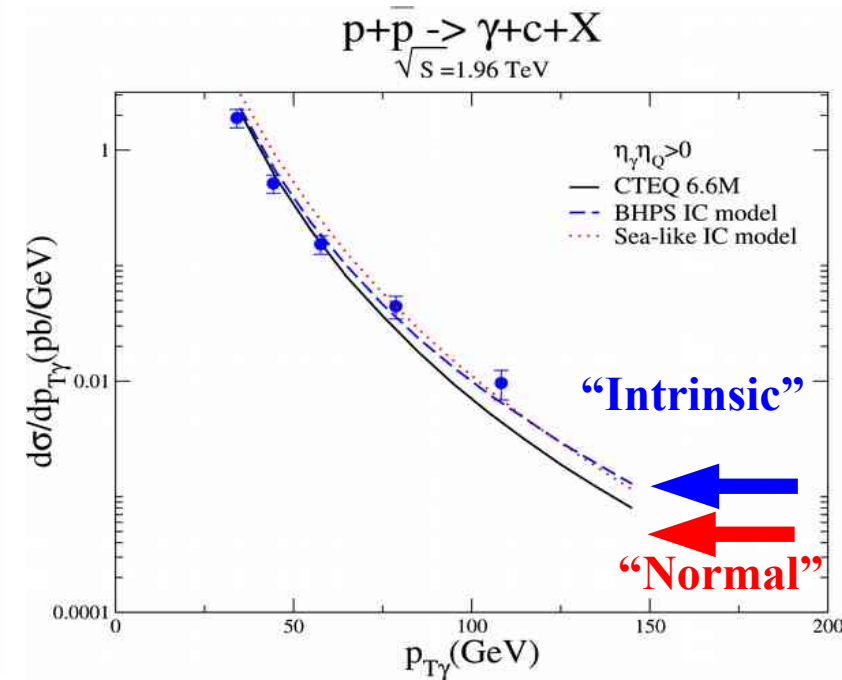
$$b \ g \rightarrow b \ \gamma/Z$$

*good place to find both extrinsic & intrinsic PDFs*



Excess in Charm,  
NOT Bottom

only at high PT



T. Stavreva, I. Schienbein, F. Arleo, K. Kovarik, F. Olness,  
J.Y. Yu, J.F. Owens, JHEP 1101 (2011) 152



# Short Cut: how to add “intrinsic” charm to any PDF

DGLAP Evolution equations ...

including **ordinary**  $Q_0$  and **intrinsic**  $Q_1$  heavy quark

$$\begin{aligned}\dot{g} &= P_{gg} \otimes g + P_{gq} \otimes q + P_{gQ} \otimes Q_0 + \cancel{P_{gQ} \otimes Q_1}, \\ \dot{q} &= P_{qg} \otimes g + P_{qq} \otimes q + P_{qQ} \otimes Q_0 + \cancel{P_{qQ} \otimes Q_1}, \\ \dot{Q}_0 + \dot{Q}_1 &= P_{Qg} \otimes g + P_{Qq} \otimes q + P_{QQ} \otimes Q_0 + P_{QQ} \otimes Q_1.\end{aligned}$$

neglect

## Equations decouple:

Intrinsic component evolves independently  
Scale set by  $m_Q$   
Adjust normalization by simple rescaling

$$\dot{Q}_1 = P_{QQ} \otimes Q_1.$$

$$c_1(x) = \bar{c}_1(x) \propto x^2 [6x(1+x) \ln x + (1-x)(1+10x+x^2)]$$

BHPS Model: Brodsky, Hoyer, Peterson, Sakai

*Phys.Lett. B93 (1980) 451-455*

$$c_1(x) = \bar{c}_1(x) \propto x^2 [6x(1+x) \ln x + (1-x)(1+10x+x^2)]$$

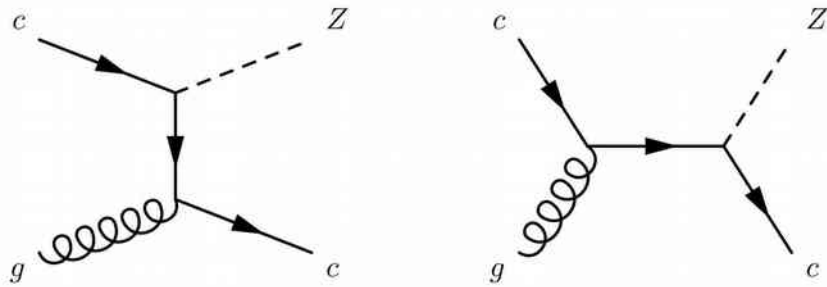
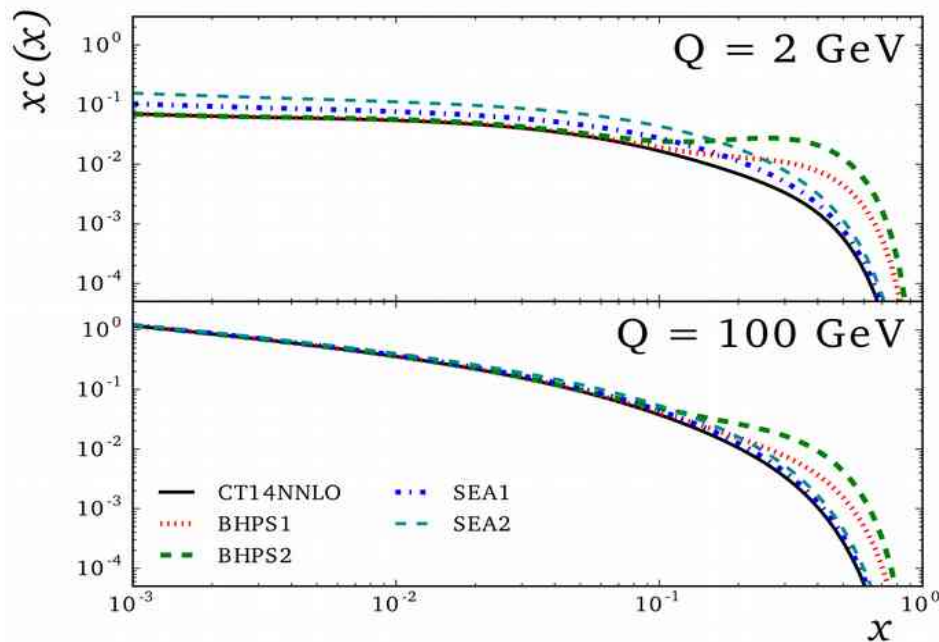
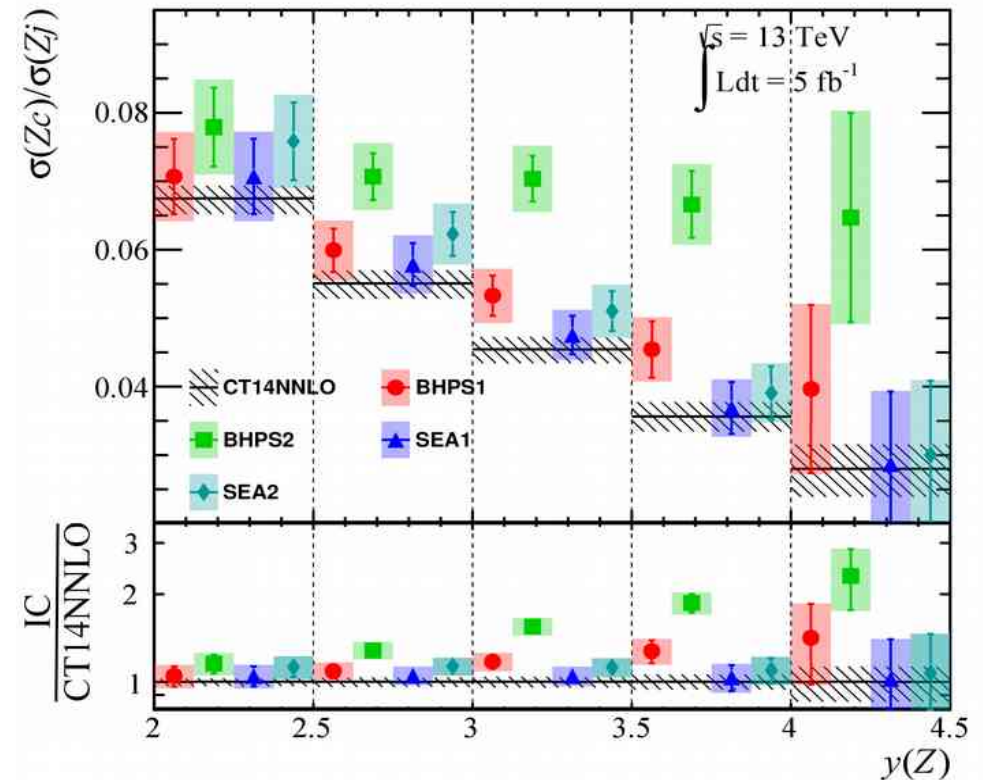


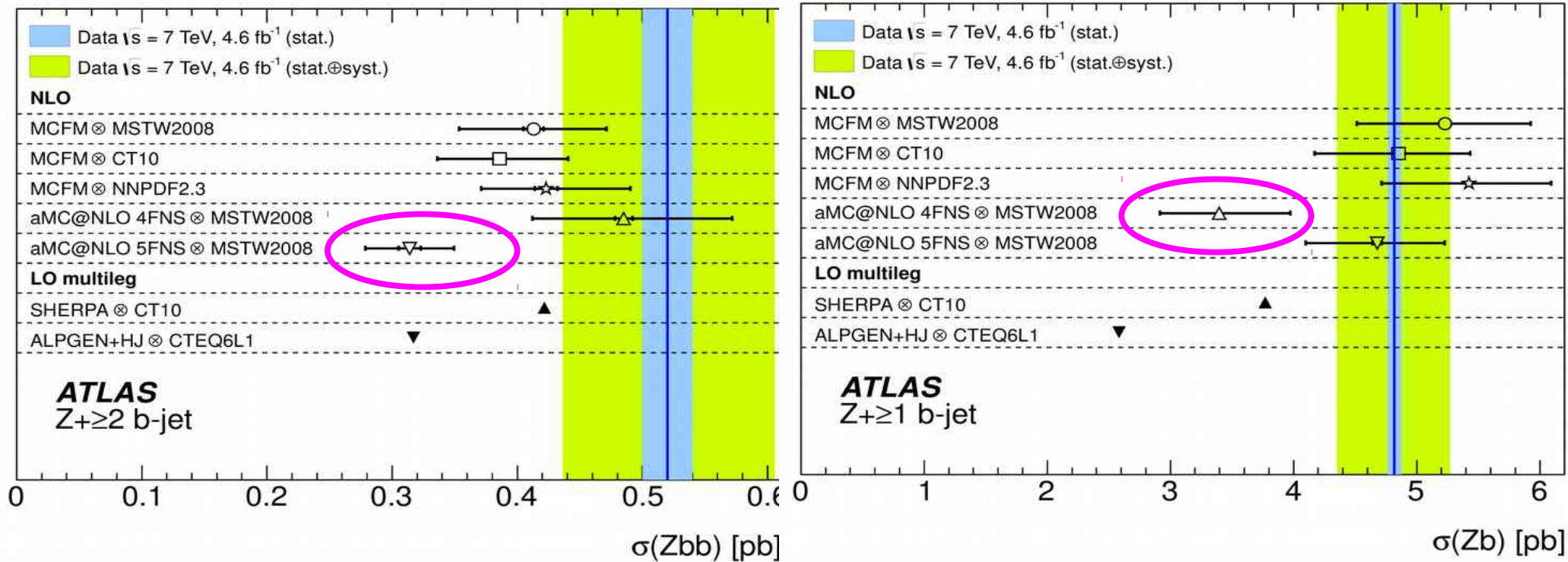
FIG. 1. Leading-order Feynman diagrams for  $gc \rightarrow Zc$ .



BHPS1: 0.6%  
BHPS2: 2%

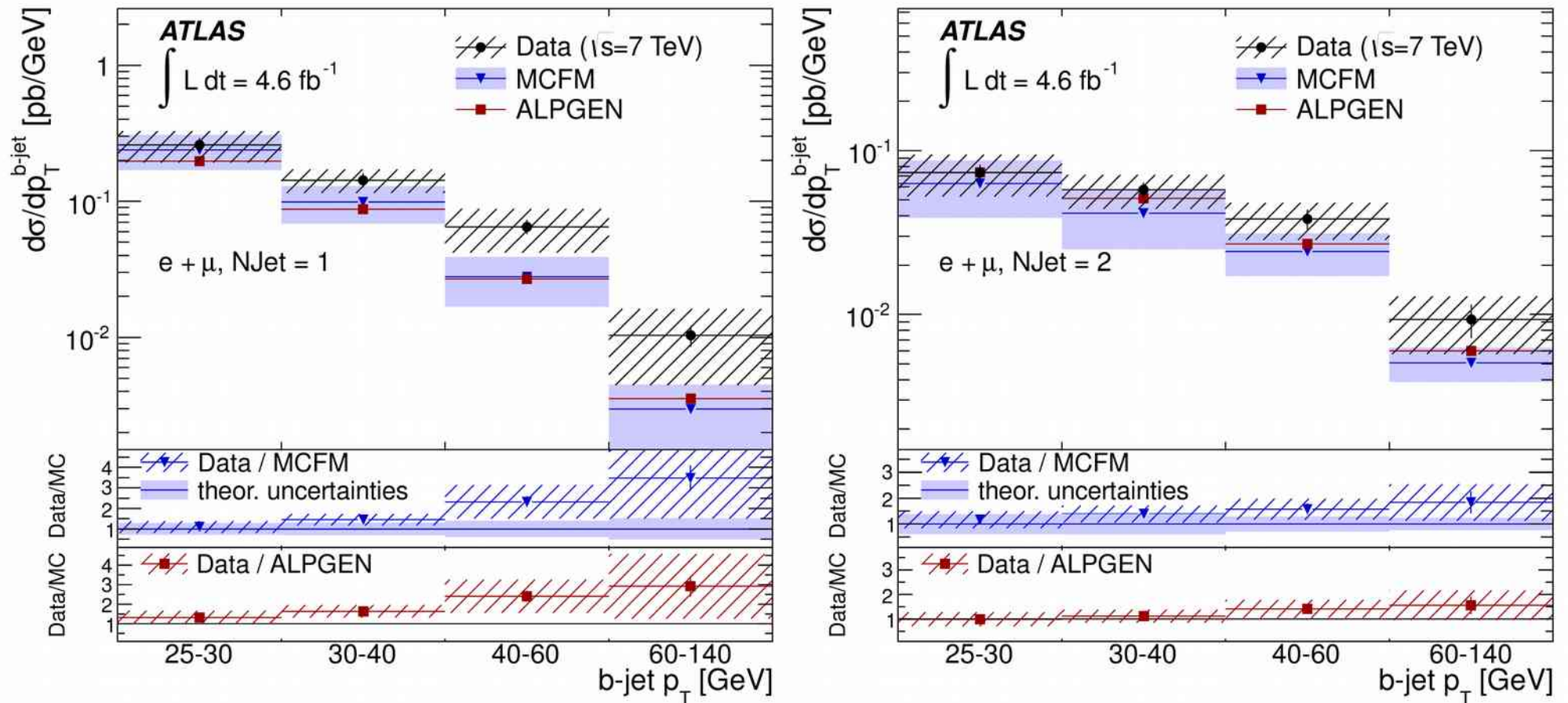


**A direct probe of the intrinsic charm content of the proton**  
Tom Boettcher, Philip Ilten, Mike Williams. [arXiv:1512.06666](https://arxiv.org/abs/1512.06666)



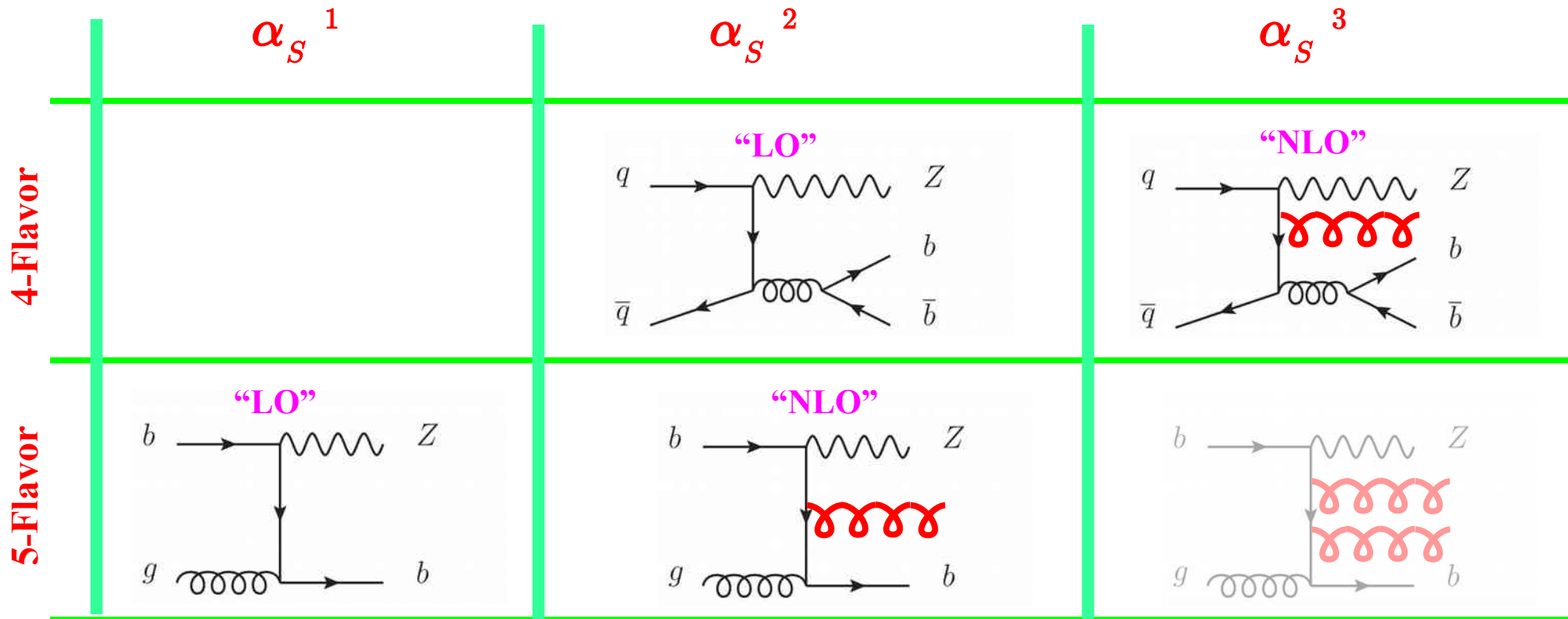
Measurement of differential production cross-sections for a Z boson in association with b-jets in 7 TeV proton-proton collisions with the ATLAS detector. JHEP10(2014)141

The agreement of the amc@nlo cross-section prediction with data differs in the  $Z+\geq 1$  b-jet and  $Z+\geq 2$  b-jets cases, with the former better described by the 5FNS prediction and the latter better described by the 4FNS prediction. Even at NLO, scale uncertainties dominate and currently limit any sensitivity to different PDF sets. Descriptions of the shapes of the differential cross-sections are generally good within uncertainties for both LO and NLO predictions. For angular distributions in the  $Z+\geq 1$  b-jet selection, where the fixed-order NLO prediction is observed to break down, the differential shapes in data are well modelled by LO multi-legged predictions.

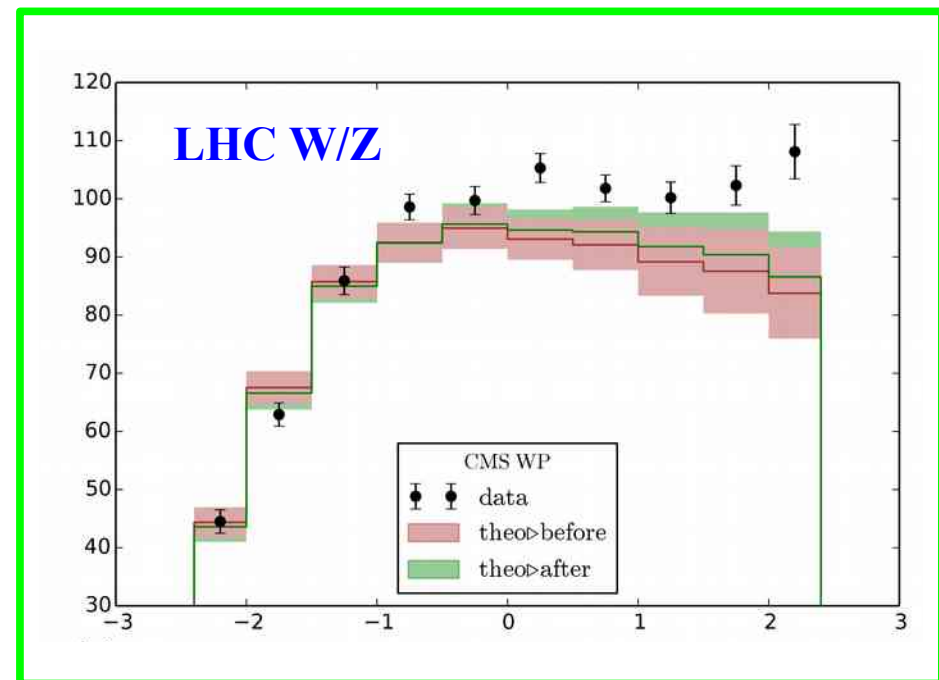
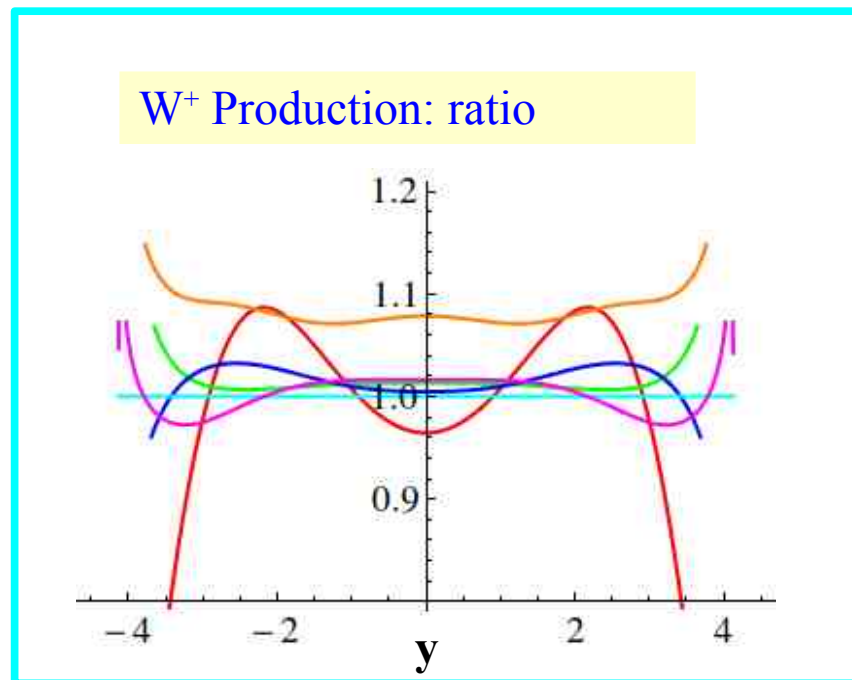
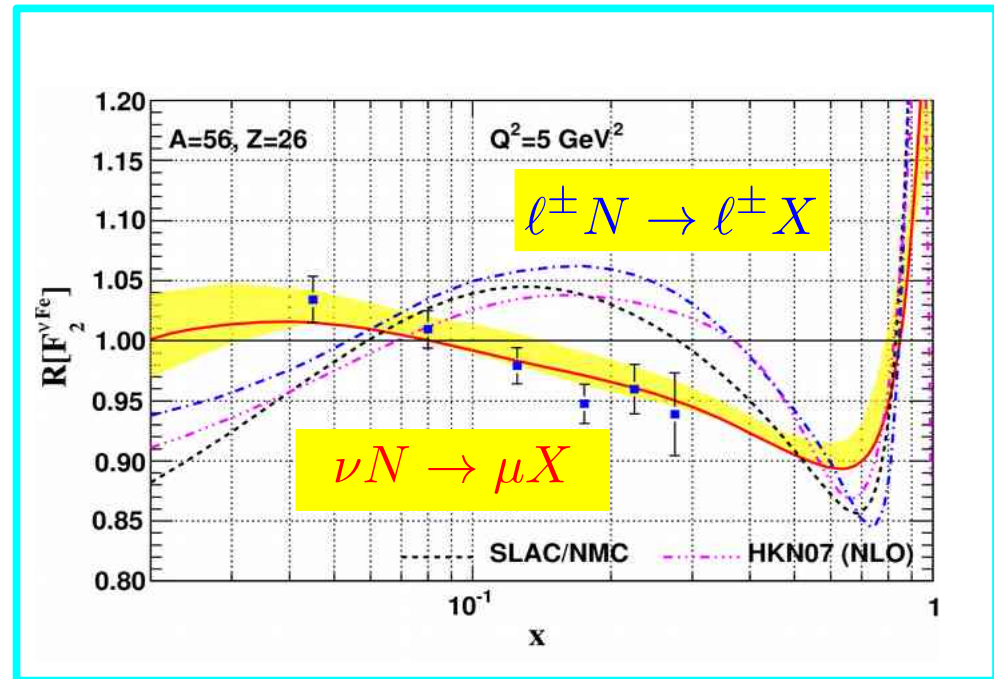
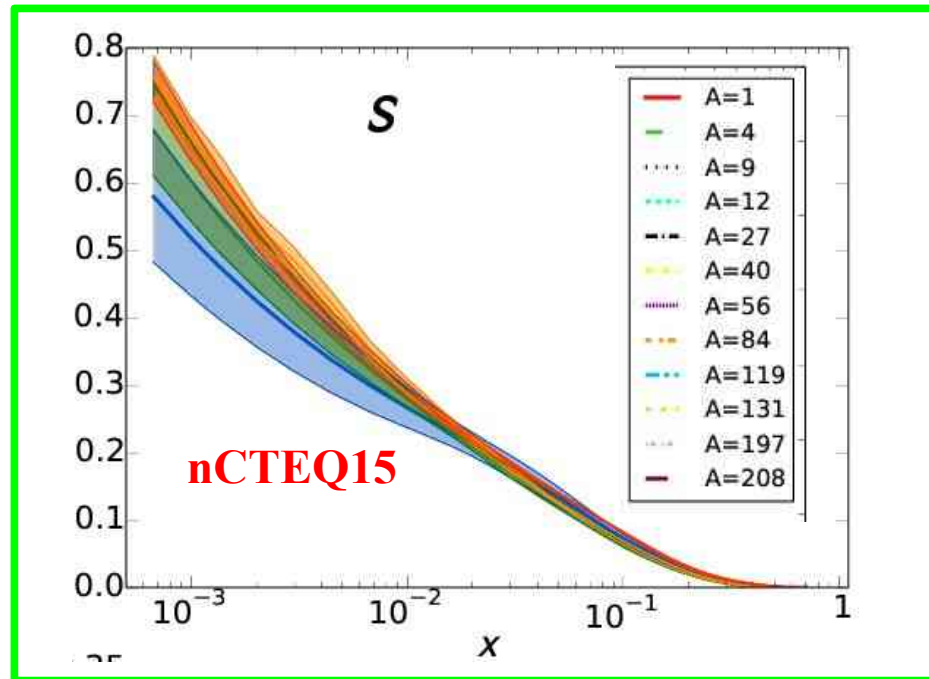


Measurement of the cross-section for W boson production in association with b-jets in pp collisions at  $\sqrt{s} = 7$  TeV with the ATLAS detector. JHEP 06 (2013) 084

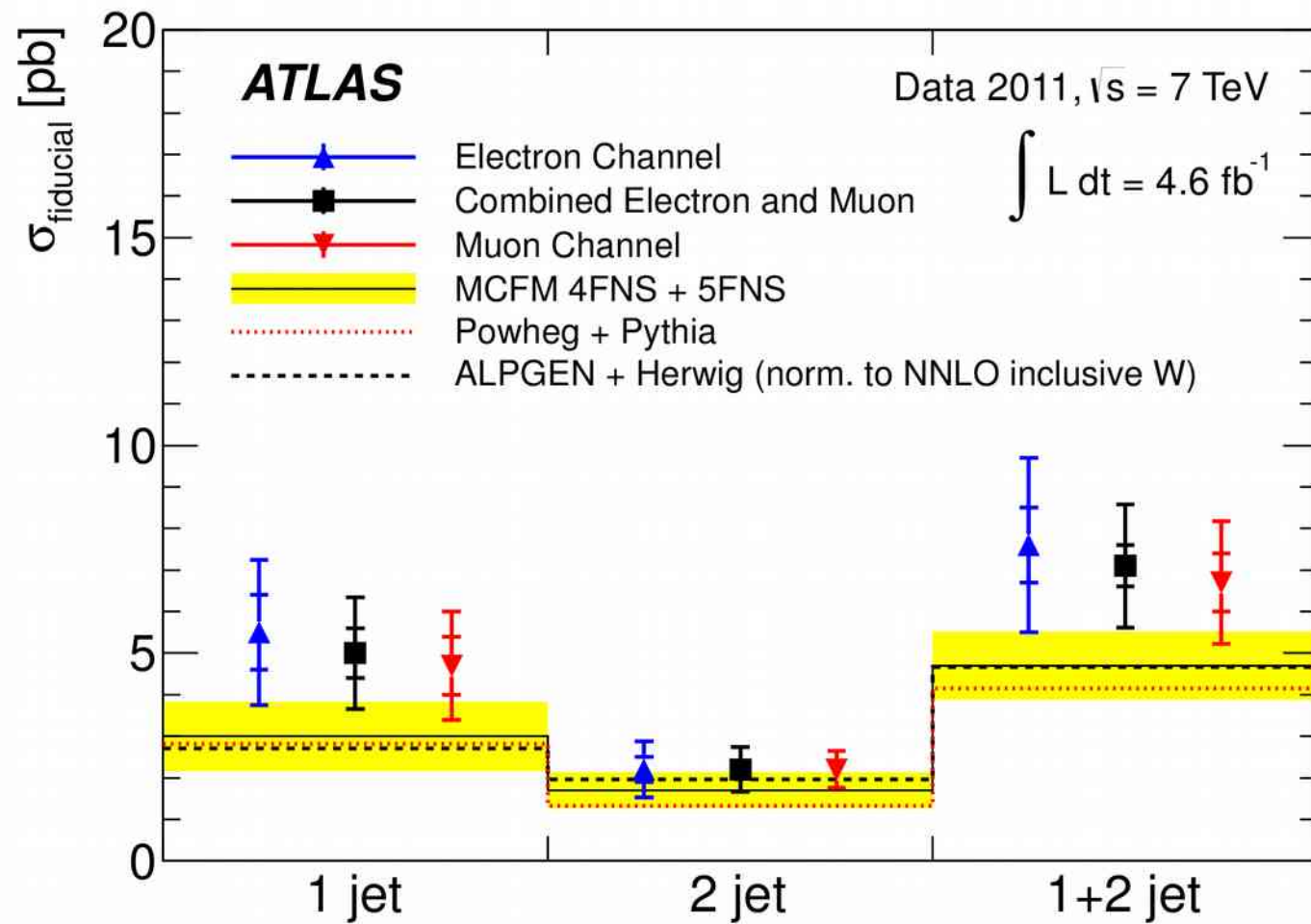




The agreement of the amc@nlo cross-section prediction with data differs in the  $Z + \geq 1$  b-jet and  $Z + \geq 2$  b-jets cases, with the former better described by the 5FNS prediction and the latter better described by the 4FNS prediction. Even at NLO, scale uncertainties dominate and currently limit any sensitivity to different PDF sets. Descriptions of the shapes of the differential cross-sections are generally good within uncertainties for both LO and NLO predictions. For angular distributions in the  $Z + \geq 1$  b-jet selection, where the fixed-order NLO prediction is observed to break down, the differential shapes in data are well modeled by LO multi-legged predictions.

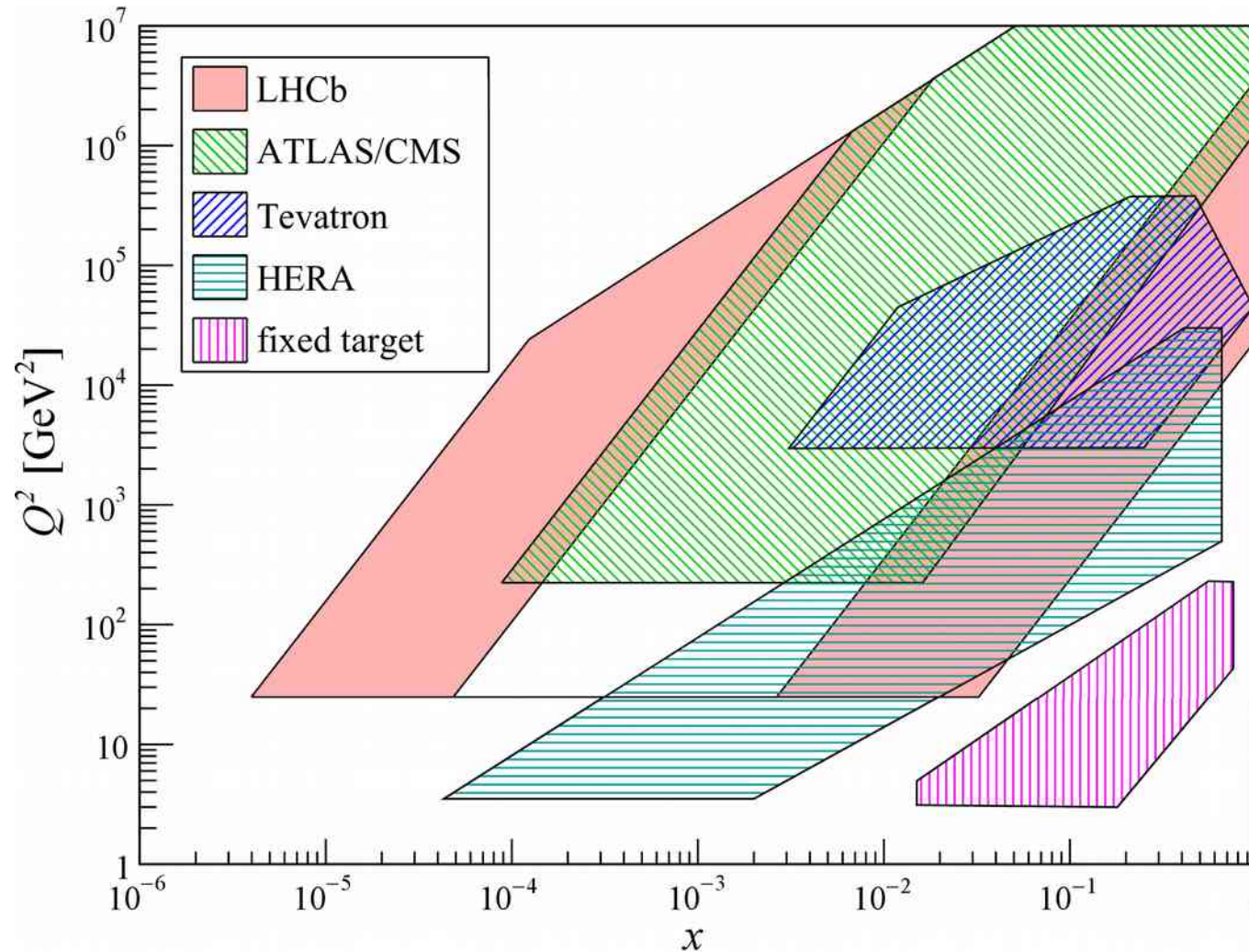






Measurement of the cross-section for W boson production in association with b-jets in pp collisions at  $\sqrt{s} = 7$  TeV with the ATLAS detector. JHEP 06 (2013) 084





A direct probe of the intrinsic charm content of the proton

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arXiv:1512.06666

