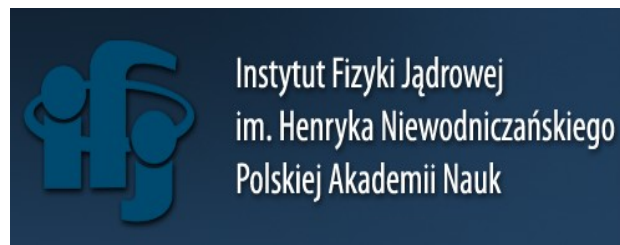




High Energy Factorization selected phenomenology results

Krzysztof Kutak



Based on results recently obtained in collaboration with:

M. Bury, P. Kotko, C. Marquet, E. Petreska, S. Sapeta, T. Salwa, M. Serino, A. van Hameren

High Energy Factorization

$$\frac{d\sigma}{dy_1 dy_2 d^2p_{1t} d^2p_{2t}} = \sum_{c,d} \int \frac{d^2k_{1t}}{\pi} \frac{d^2k_{2t}}{\pi} \frac{1}{16\pi^2(x_1 x_2 S)^2} |\overline{\mathcal{M}}_{g^*g^* \rightarrow cd}|^2 \delta^2(k_{1t} + k_{2t} - p_{1t} - p_{2t}) \mathcal{F}_A(x_1, k_{1t}^2) \mathcal{F}_B(x_2, k_{2t}^2) \frac{1}{1 + \delta_{cd}}$$

Originally written for total cross section

Gluons dominate in the t channel

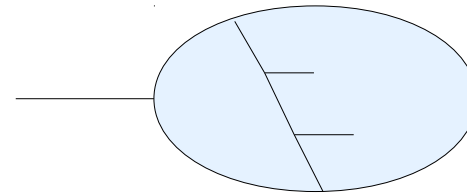
Obtained for heavy quarks in final state.

$$k_1^\mu = x_1 P_1^\mu + \bar{x}_1 P_2^\mu + k_{1t}^\mu \quad k_2^\mu = x_2 P_2^\mu + \bar{x}_2 P_1^\mu + k_{2t}^\mu$$

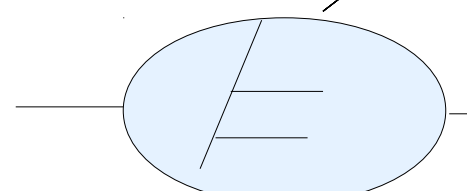
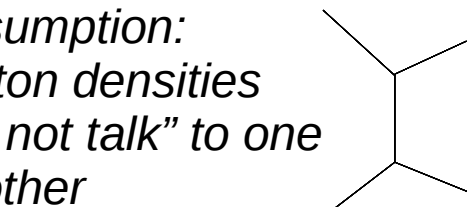
$$\bar{x}_1 = \frac{k_1^2 + k^2}{Sx_1}$$

$$\bar{x}_2 = \frac{k_2^2 + k^2}{Sx_2}$$

Assumption:
parton densities
"do not talk" to one
another



Decreasing longitudinal
momentum fractions
of off-shell partons



Decreasing longitudinal
momentum fractions
of off-shell partons

Gribov, Levin, Ryskin '81

Ciafaloni, Catani, Hautman '93

$$|\mathcal{M}_{ab \rightarrow cd}|^2 = \frac{2x_1 k_1^{\mu_1} k_1^{\nu_1}}{k_1^2} \frac{2x_2 k_2^{\mu_2} k_2^{\nu_2}}{k_2^2} \mathcal{M}_{ab \rightarrow cd, \mu_1 \nu_1} \mathcal{M}_{ab \rightarrow cd, \mu_2 \nu_2}^*$$

Does not take into account MPI
as formulated in DGLAP i.e.
emissions from independent chains

Helicity method based method for any process

KK, Kotko, van Hameren '13

Off-shell matrix elements

One consider embedding off-shell amplitude in onshell and introduces eikonal lines

Kotko, KK, van Hameren 2013,
KK, Salwa, van Hameren 2013

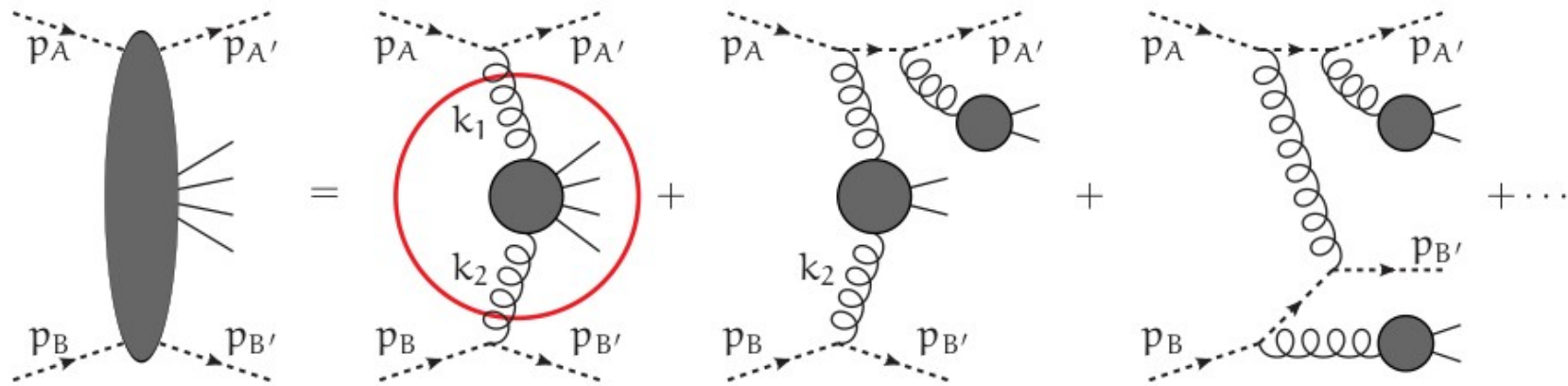


Diagram showing a fermion eikonal line (wavy line) connecting vertex j to vertex i . The equation is:

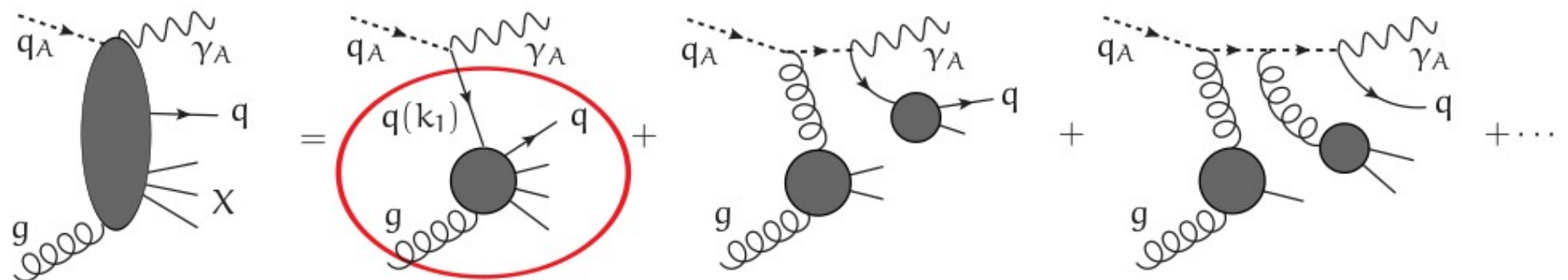
$$= -i \delta_{i,j} u(p_1)$$

Diagram showing a gluon eikonal line (coiled line) connecting vertex j to vertex i . The equation is:

$$= -i T_{i,j}^a p_1^\mu$$

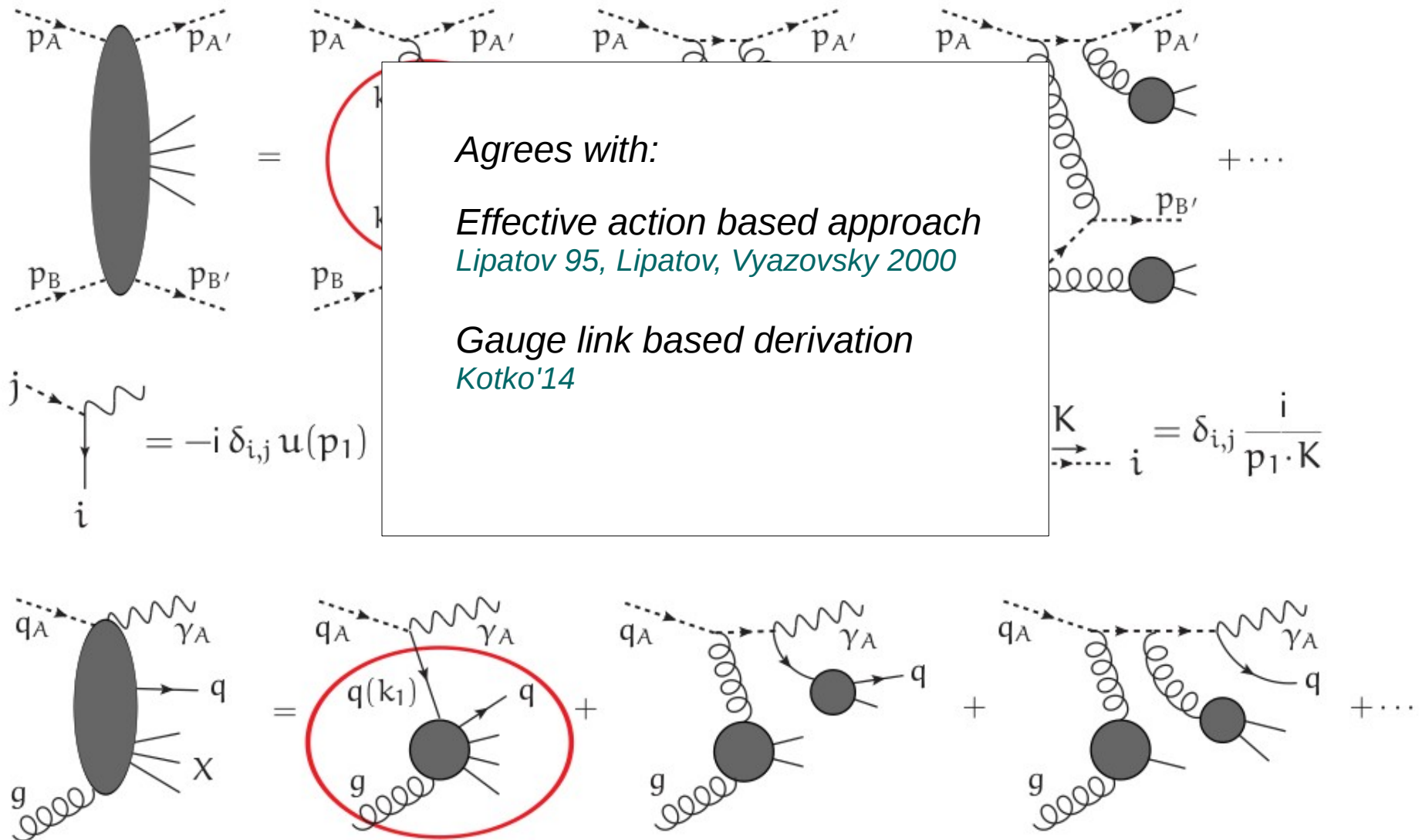
Diagram showing a scalar eikonal line (dashed line) connecting vertex j to vertex i . The equation is:

$$= \delta_{i,j} \frac{i}{p_1 \cdot K}$$

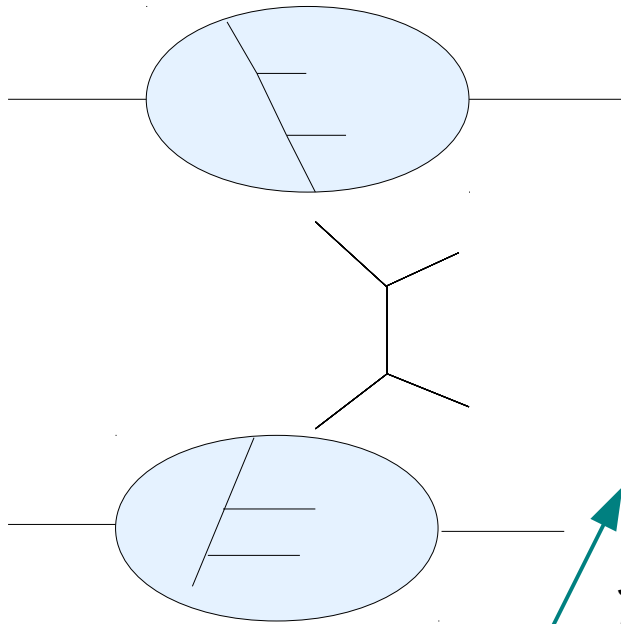


Off-shell matrix elements

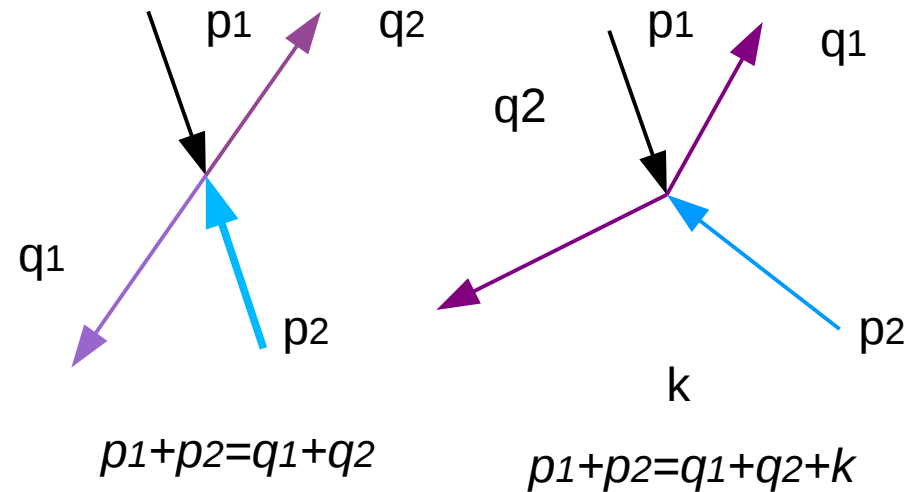
Kotko, KK, van Hameren 2013,
KK, Salwa, van Hameren 2013



Hybrid high energy factorization



Strongly decreasing
Longitudinal momentum
fractions of off-shell partons



ME + parton densities in kt factorization

Theory

*Gribov, Levin, Ryskin '81
Ciafaloni, Catani, Hautman '93
Collins, Ellis '93*

Phenomenology

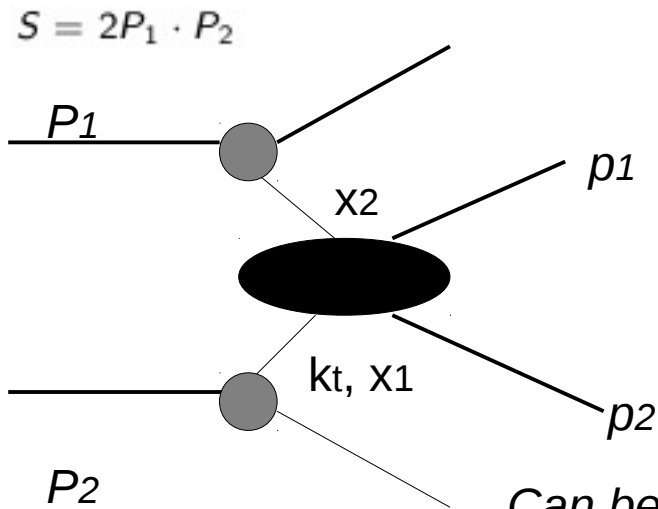
*Baranov, Hautmann, Jung, Maciuła,
Marquet, Motyka, KK, Kotko, Lipatov,
van Hameren, Saleev, Szczurek,
Zotov,...*

*Helicity based methods for ME
Kotko, K.K, van Hameren, '12*

Hybrid formula for cross section

$$\frac{d\sigma}{dy_1 dy_2 dp_{1t} dp_{2t} d\Delta\phi} = \sum_{a,c,d} \frac{p_{t1} p_{t2}}{8\pi^2 (x_1 x_2 S)^2} |\overline{\mathcal{M}}_{ag \rightarrow cd}|^2 x_1 f_{a/A}(x_1, \mu^2) \mathcal{F}_{g/B}(x_2, k^2, \mu^2) \frac{1}{1 + \delta_{cd}}$$

Deak, Jung, KK, Hautmann '09

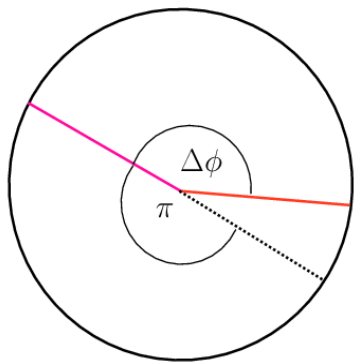


- logs of x
- logs of hard scale

Can be obtained from CGC after neglecting nonlinearities
In that limit gluon density is so called dipole gluon density.

Laidet, Iancu '13

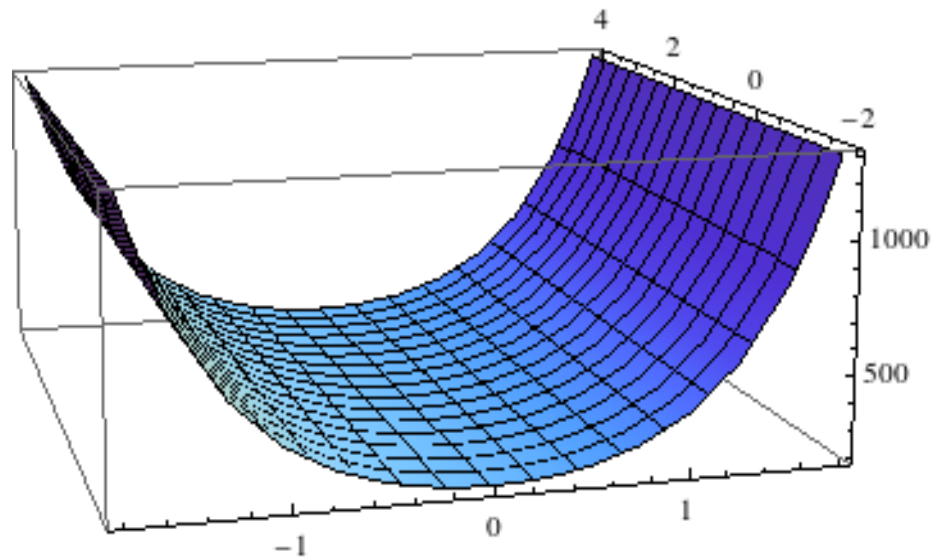
The refinement of formula above for dilute-dense: *talk of Cyrille Marquet*



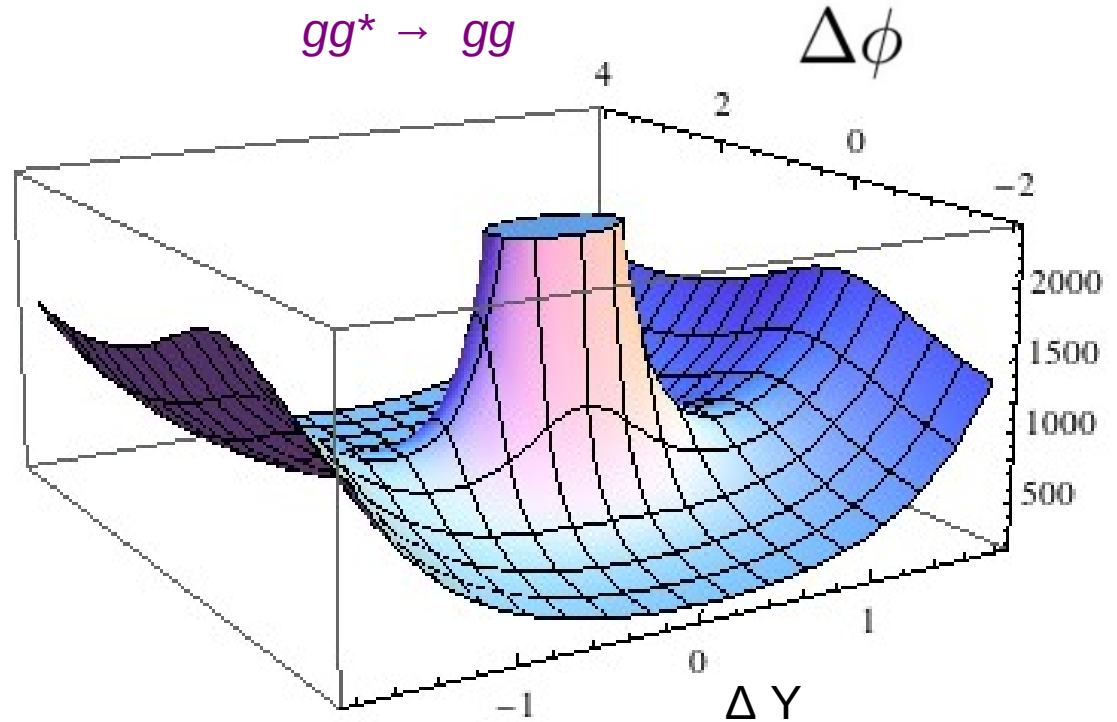
In azimuthal plane

Example of matrix element for hybrid factorization

$gg \rightarrow gg$



$gg^* \rightarrow gg$



Two on-shell partons $2 \rightarrow 2$ ME

One off-shell parton $2 \rightarrow 2$ ME

Bad approximation by collinear physics of region where angle between produced jets is small

Numerical tools for HEF

AVHLIB

<http://bitbucket.org/hameren/avhlib> (A. van Hameren)

- complete Monte Carlo program for tree-level calculations
- any process within the Standard Model
- any initial-state partons on-shell or off-shell
- employs numerical Dyson-Schwinger recursion to calculate helicity amplitudes
- automatic phase space optimization

AMP4HEF

<http://bitbucket.org/hameren/amp4hef> (A. van Hameren, M. Bury, K. Bilko, H. Milczarek)

- only provides tree-level matrix elements (or color-ordered helicity amplitudes)
- employs BCFW recursion to calculate color-ordered helicity amplitudes
- available processes (plus those with fewer on-shell gluons):

$$\begin{array}{lll} \emptyset \rightarrow g g + 4g & \emptyset \rightarrow \bar{q} q + 3g & \emptyset \rightarrow \bar{q}^* q + 3g \\ \emptyset \rightarrow g^* g + 4g & \emptyset \rightarrow g^* + \bar{q} q + 2g & \emptyset \rightarrow \bar{q} q^* + 3g \\ \emptyset \rightarrow g^* g^* + 4g & & \end{array}$$

*Easy to use in
Fortran and C++*

LxJet

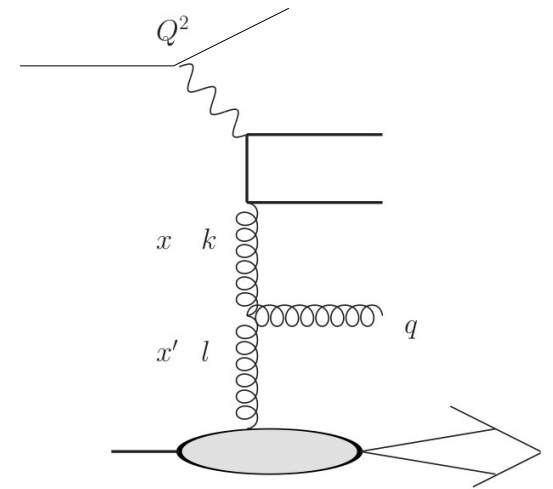
<http://annapurna.ifj.edu.pl/~pkotko/LxJet.html> (P. Kotko)

- hybrid high energy factorization suitable for forward jets,
- implemented helicity tree-level amplitudes for all channels for dijets and three jets
- recursive relation for color ordered tree-level amplitudes with single off-shell leg for arbitrary number of gluons
- currently the native phase space generator is up to three final state partons

Evolution

BFKL reummation of logs $1/x$. Dependence on x and k_t

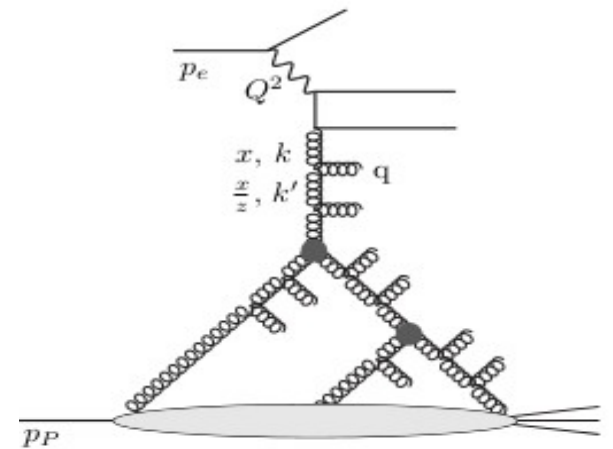
$$\mathcal{F} = \mathcal{F}_0 + K \otimes \mathcal{F}$$



Balitsky-Kovchegov reummation of logs $1/x$ + rescatterings.
Dependence on x , k_t and shape of the target

$$\mathcal{F} = \mathcal{F}_0 + K \otimes \mathcal{F} - \frac{1}{R^2} V \otimes \mathcal{F}^2$$

*BK formulated for large nucleus
often used as a model for proton*



Numerical solutions and phenomenological applications:

Lublinsky, Levin, Maor;

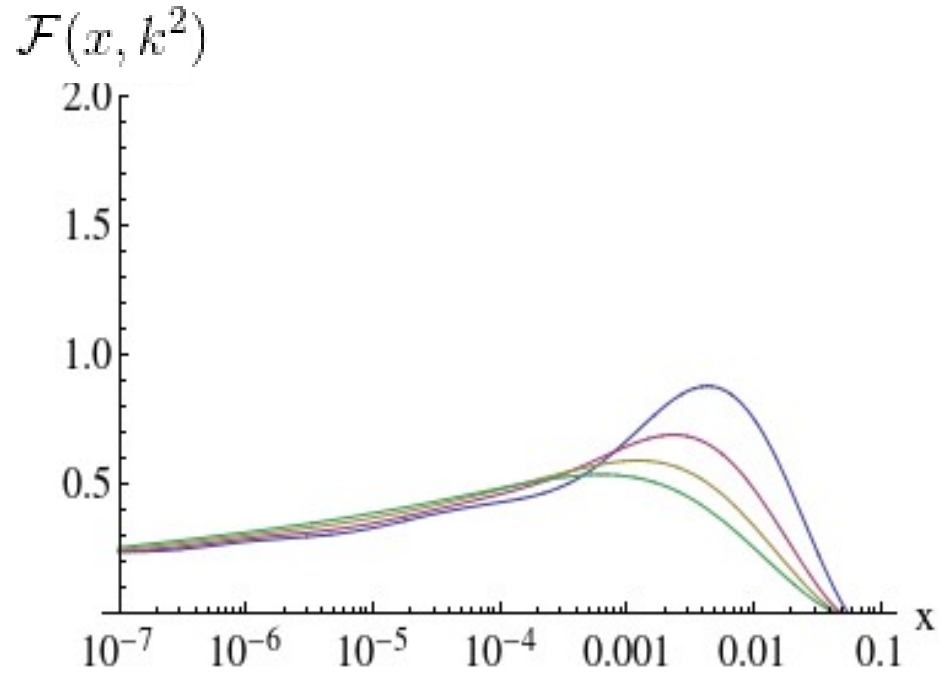
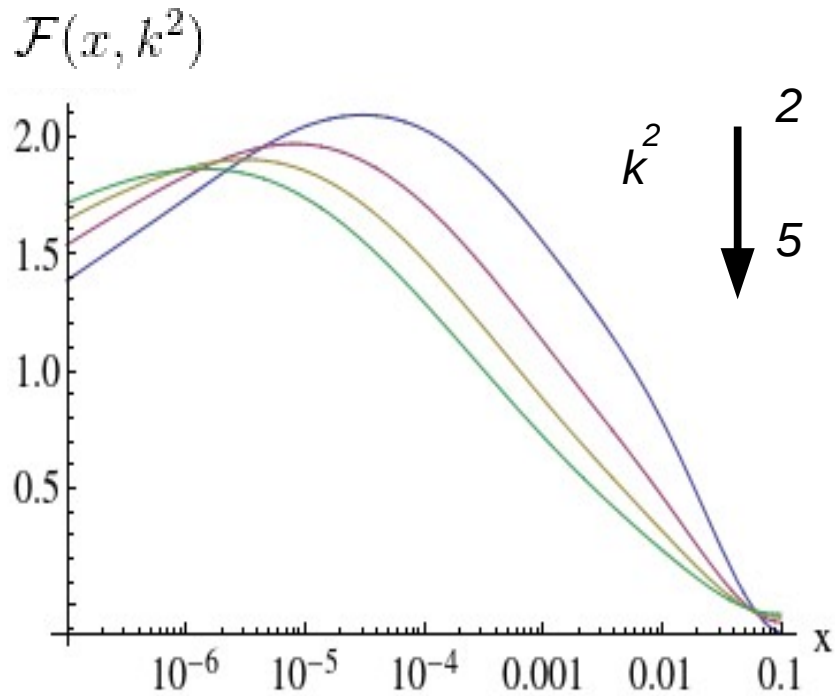
Golec-Biernat, Motyka, Stasto;

KK, Kwiecinski;

Marquet, Soyez'

Albacete, Armesto, Milhano, Salgado,.....

Glue in p vs. glue in Pb



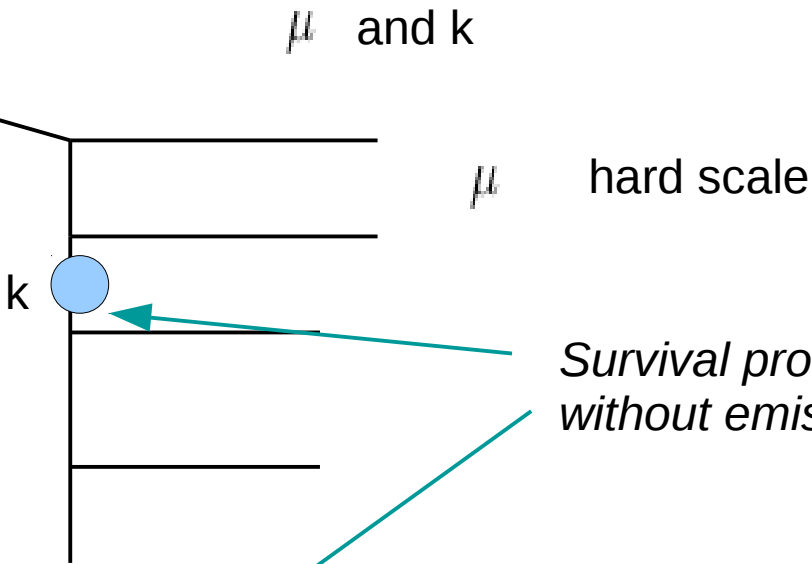
Maximum → emergence of saturation scale

Introducing hard scale dependence

The relevance in low x physics at linear level recognized by:

Catani, Ciafaloni, Fiorani, Marchesini;
Kimber, Martin, Ryskin;
Collins, Jung

Survival probability of the gap without emissions



Survival probability of the gap without emissions

Kimber, Martin, Ryskin procedure '01:

$$T_s(\mu^2, k^2) = \exp \left(- \int_{k^2}^{\mu^2} \frac{dk'^2}{k'^2} \frac{\alpha_s(k'^2)}{2\pi} \sum_{a'} \int_0^{1-\Delta} dz' P_{a'a}(z') \right)$$

$$\Delta = \frac{\mu}{\mu+k}$$

$$\mathcal{F}(x, k^2, \mu^2) \sim \partial_{\lambda^2} (T(\lambda^2, \mu^2) x g(x, \lambda^2)) |_{\lambda^2=k^2}$$

KK procedure '14 :

Motivated by paper by [Mueller, Yuan, Xiao '13](#)

$$\mathcal{F}(x, k^2, \mu^2) := \theta(\mu^2 - k^2) T_s(\mu^2, k^2) \frac{xg(x, \mu^2)}{xg_{hs}(x, \mu^2)} \mathcal{F}(x, k^2) + \theta(k^2 - \mu^2) \mathcal{F}(x, k^2)_{11}$$

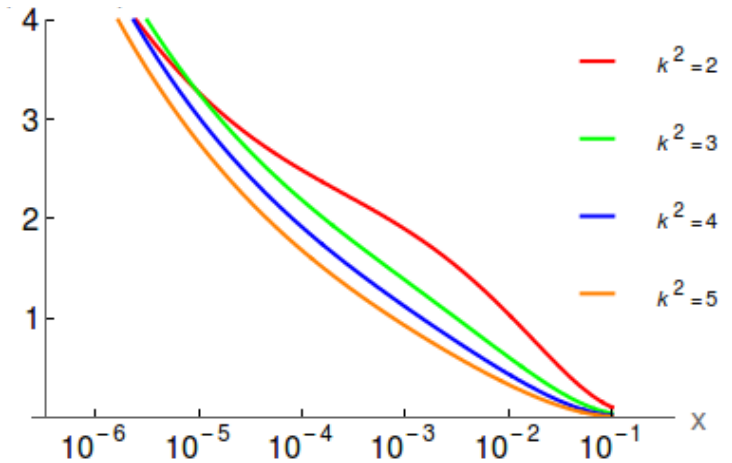
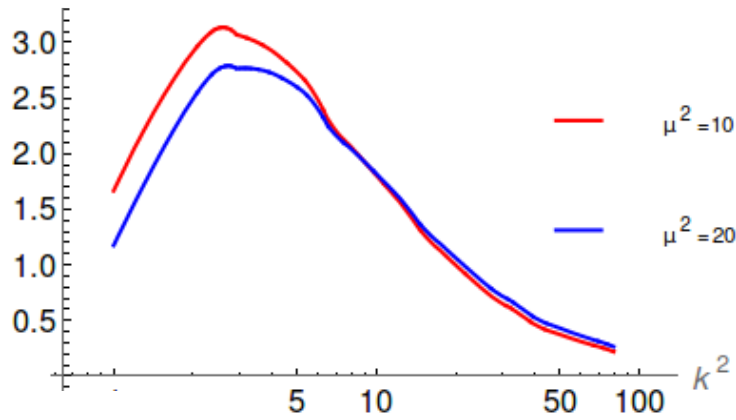
Introducing hard scale dependence

KMR for CTEQ 10 NLO

Fundamental difference.
x dependence

$$\mathcal{F}(x, k^2, \mu^2)$$

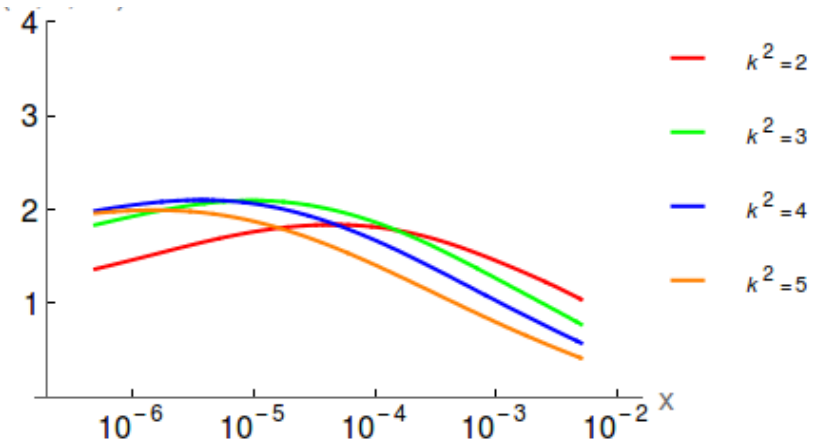
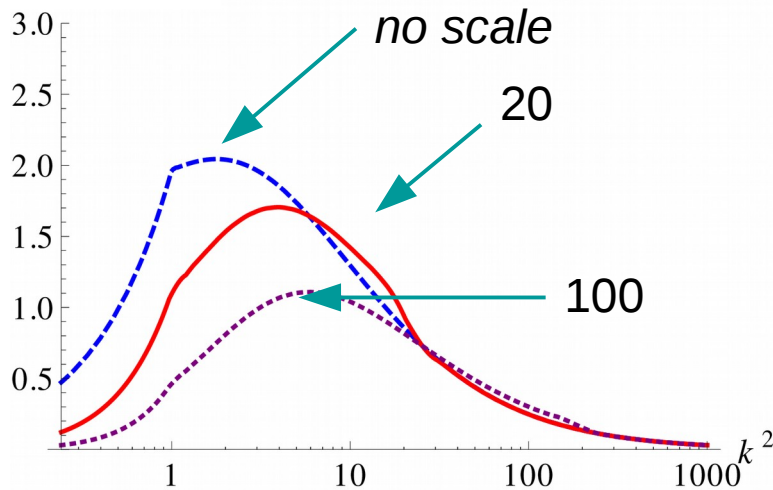
$$\mathcal{F}(x, k^2, \mu^2)$$



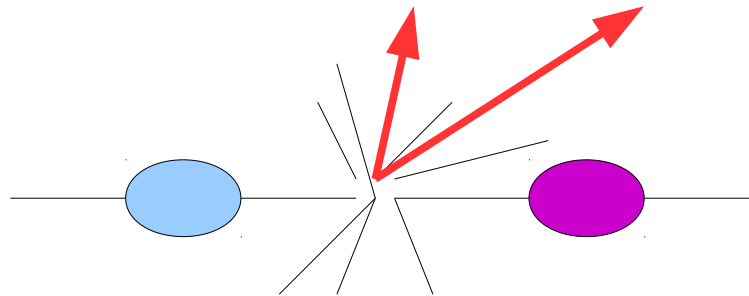
$$\mathcal{F}(x, k^2, \mu^2)$$

KShardscale

$$\mathcal{F}(x, k^2, \mu^2)$$

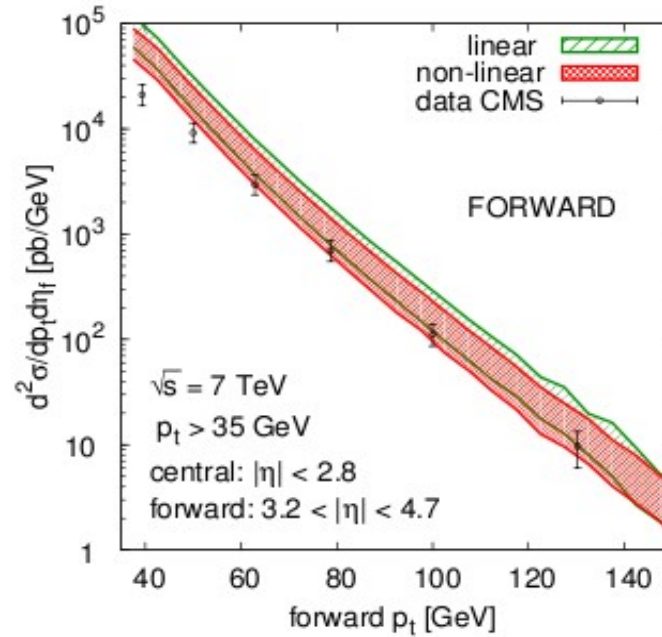
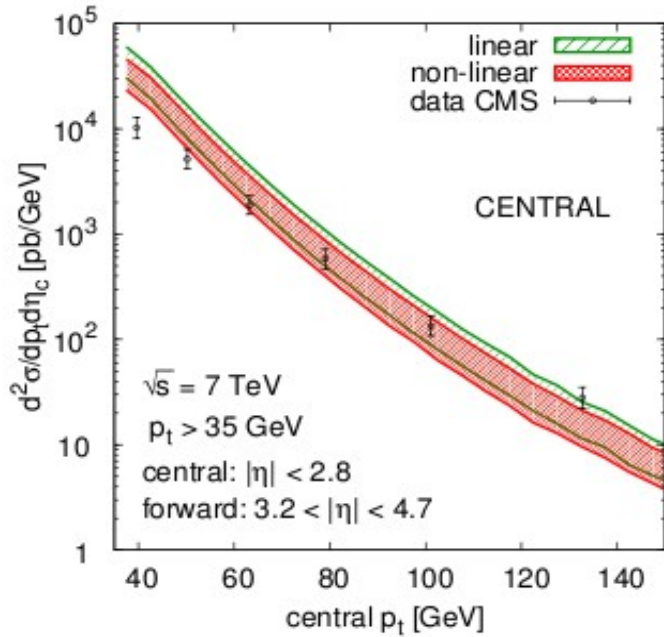


Central-forward di-jets



Di-jets p_t spectra

S.Sapeta. KK ,12



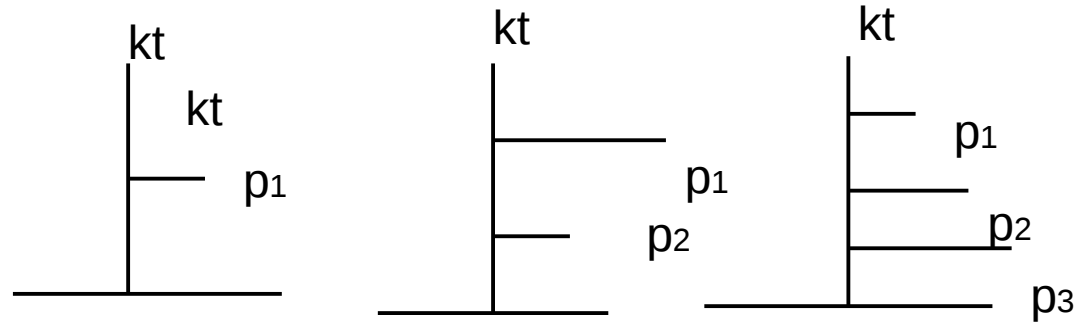
Reasonable agreement.

No usage of traditional parton shower

Gluon emissions are unordered in p_t and add up to $k_t = |p_1 + p_2 + \dots + p_n|$

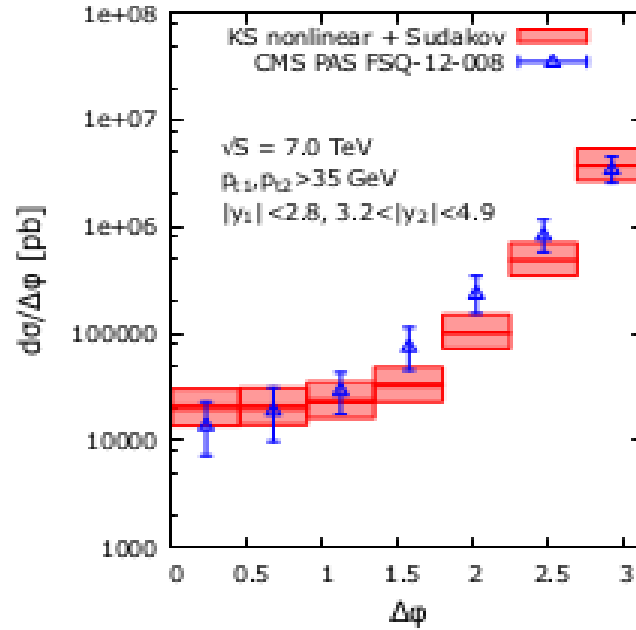
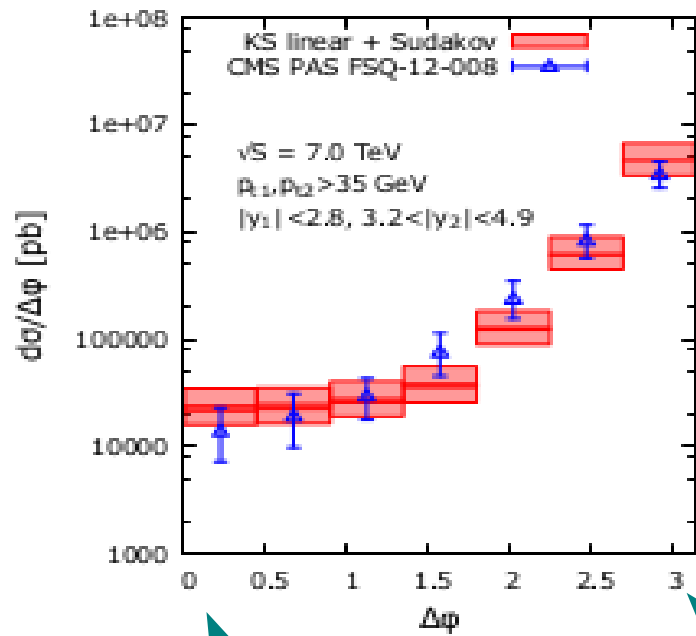
During evolution time incoming gluon becomes off-shell

Crucial effect of higher order corrections

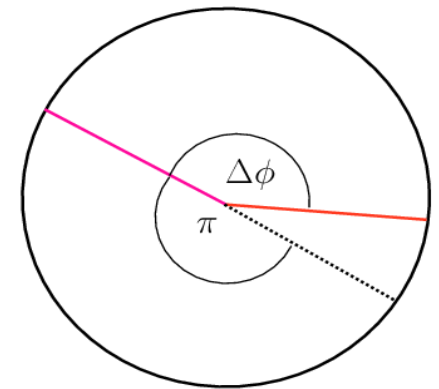


Decorelations inclusive scenario forward-central

van Hameren,, Kotko, K.K, Sapeta '14



$p_{T1}, p_{T2} > 35$, leading jets
 $|y_1| < 2.8, 3.2 < |y_2| < 4.7$
No further requirement on jets

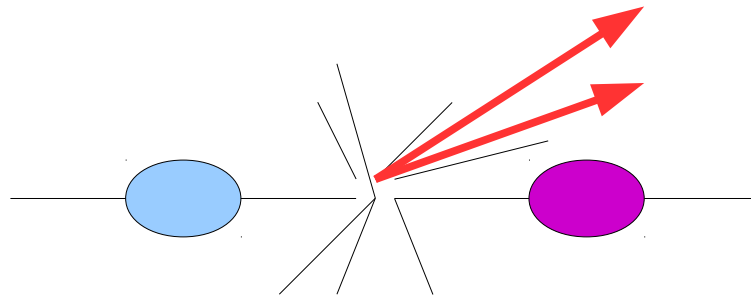


Sudakov effects by reweighting
implemented in LxJet Monte Carlo
P. Kotko

Observable suggested to
study BFKL effects
Sabio-Vera, Schwensen '06

Studied also context of RHIC
Albacete, Marquet '10

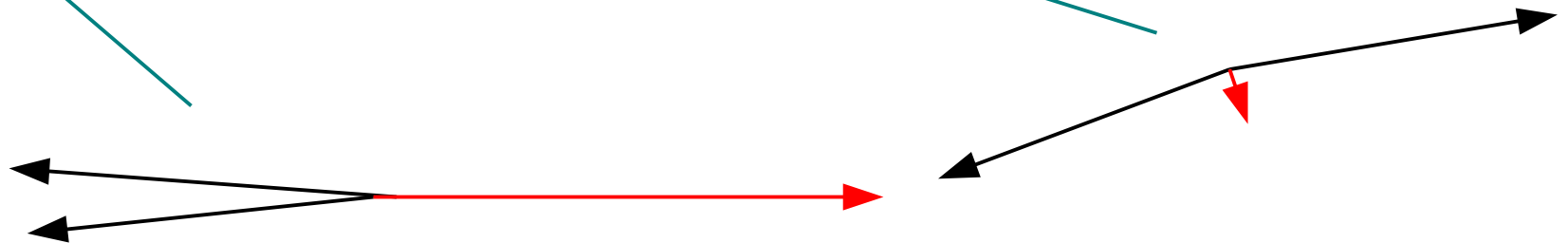
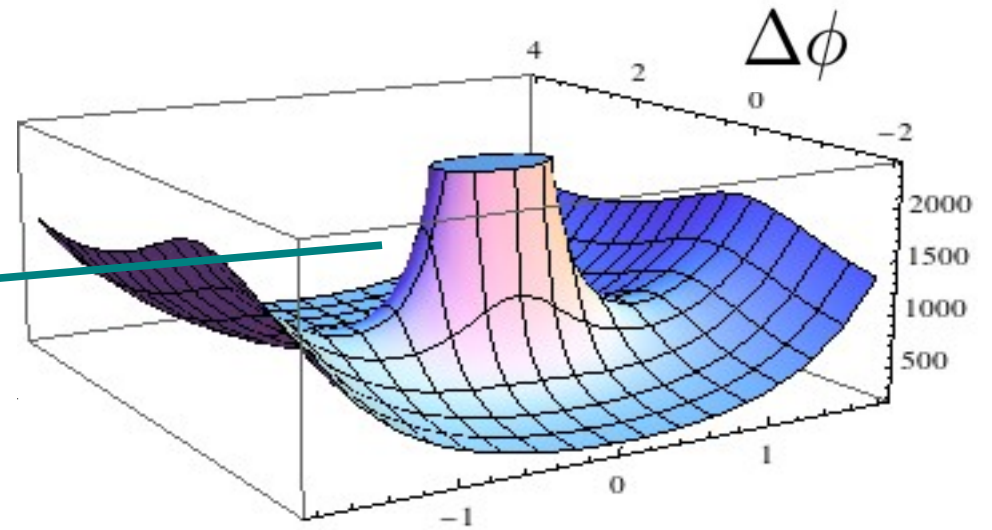
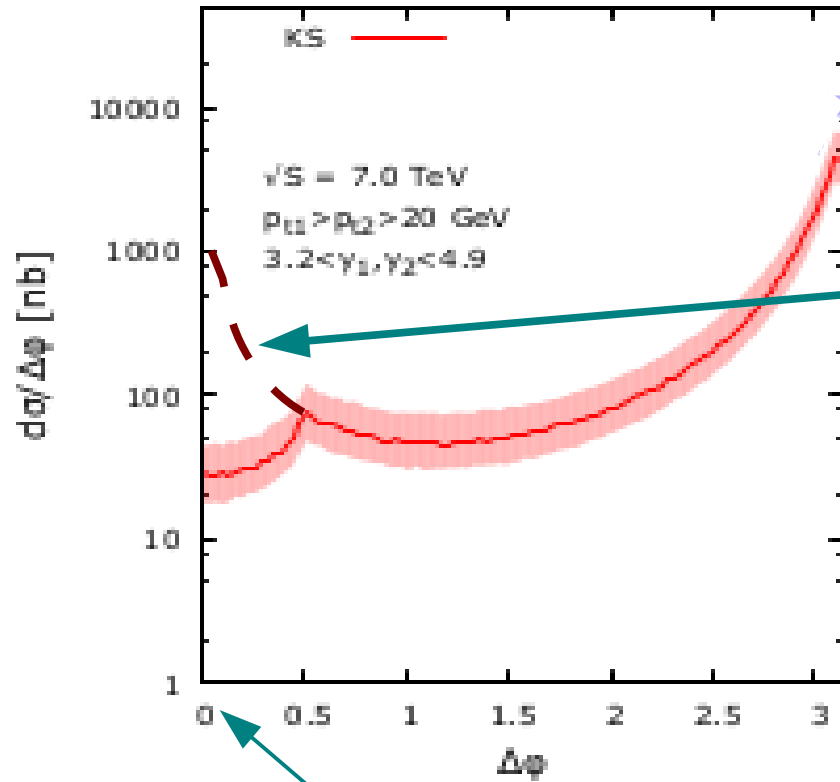
Forward-forward di-jets



Results for decorrelations $p+p$

Divergence regularized
by jet algorithm

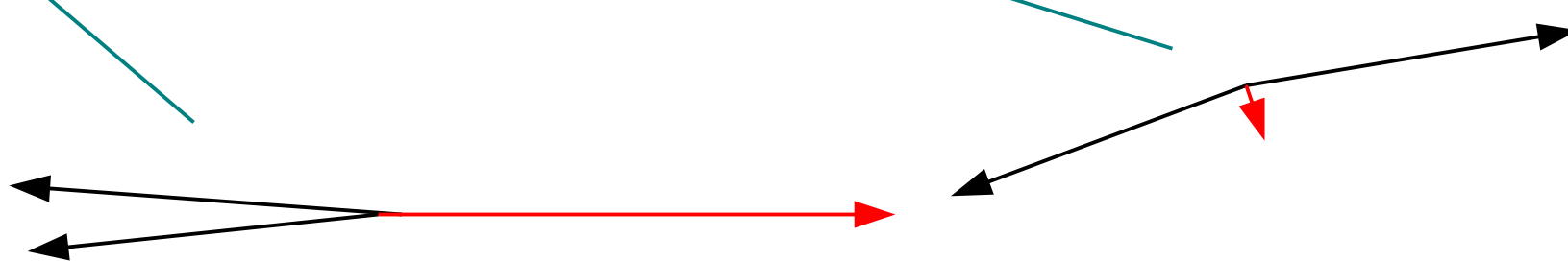
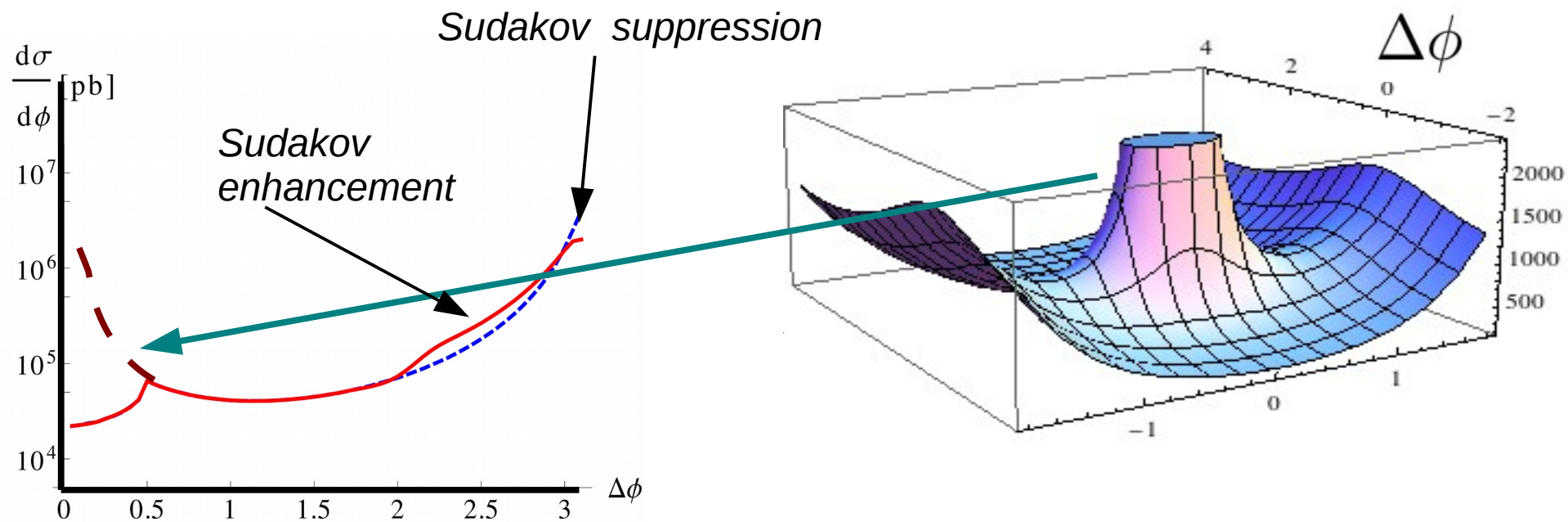
A. van Hameren, Kotko, KK, Marquet, Sapeta '14



Results for decorrelations $p+p$

Divergence regularized
by jet algorithm

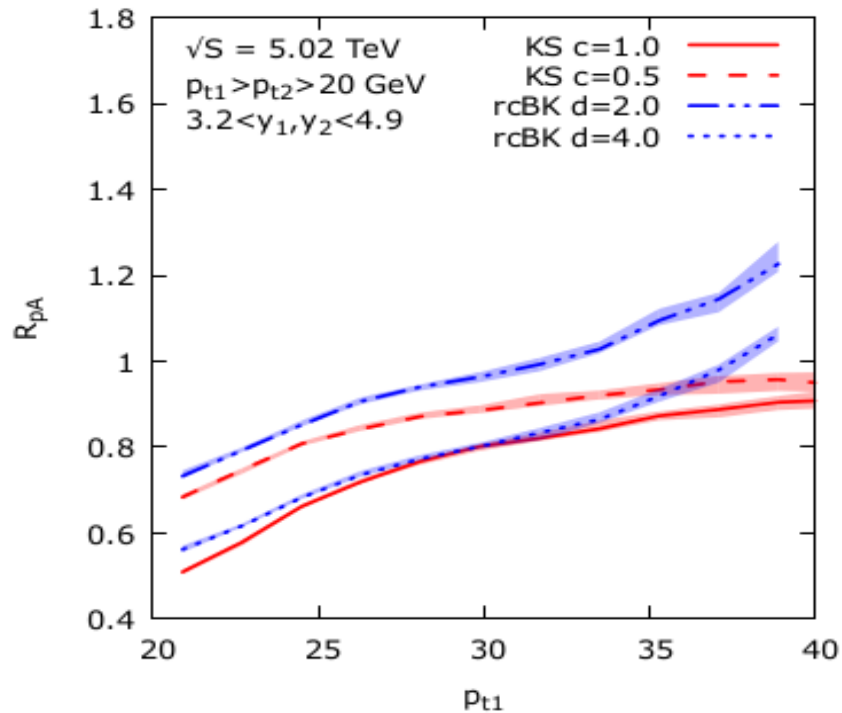
KK '14



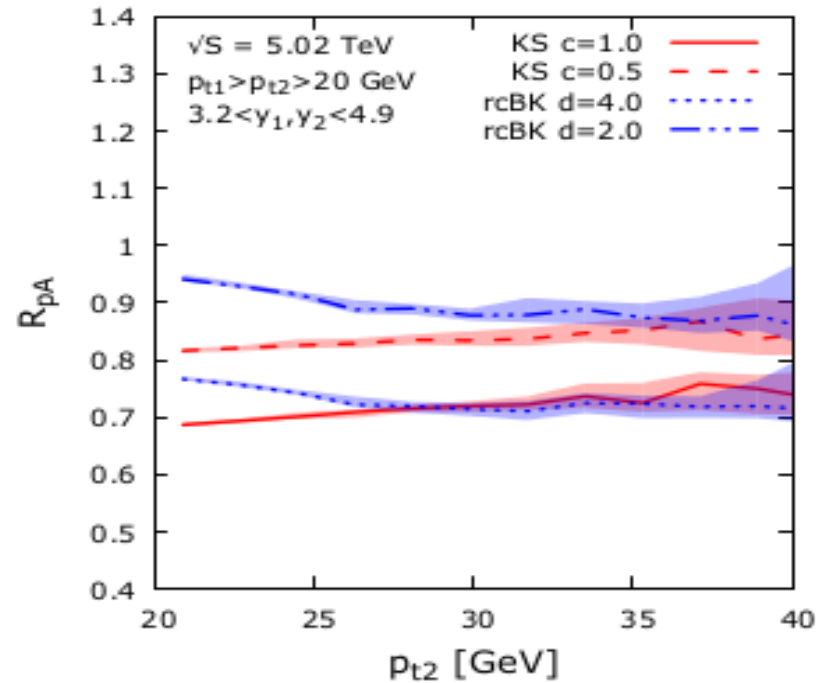
Forward-forward dijets

$$R_{pA} = p+Pb / p+p$$

A. van Hameren, Kotko, KK, Marquet, Sapeta '14

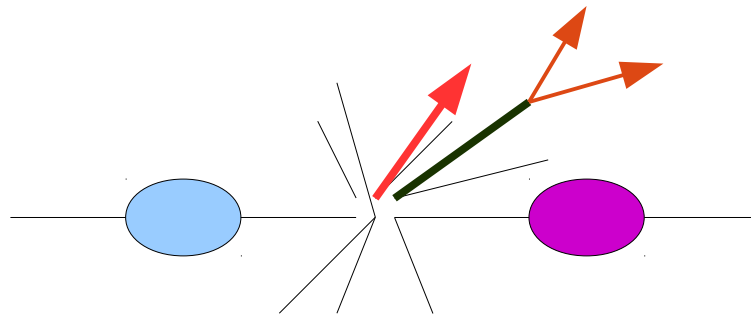


rcBK: above unity at large p_t
KS: reaches unity at large p_t



Studies of sub-leading jet gives more pronounced signal of nonlinear effects.

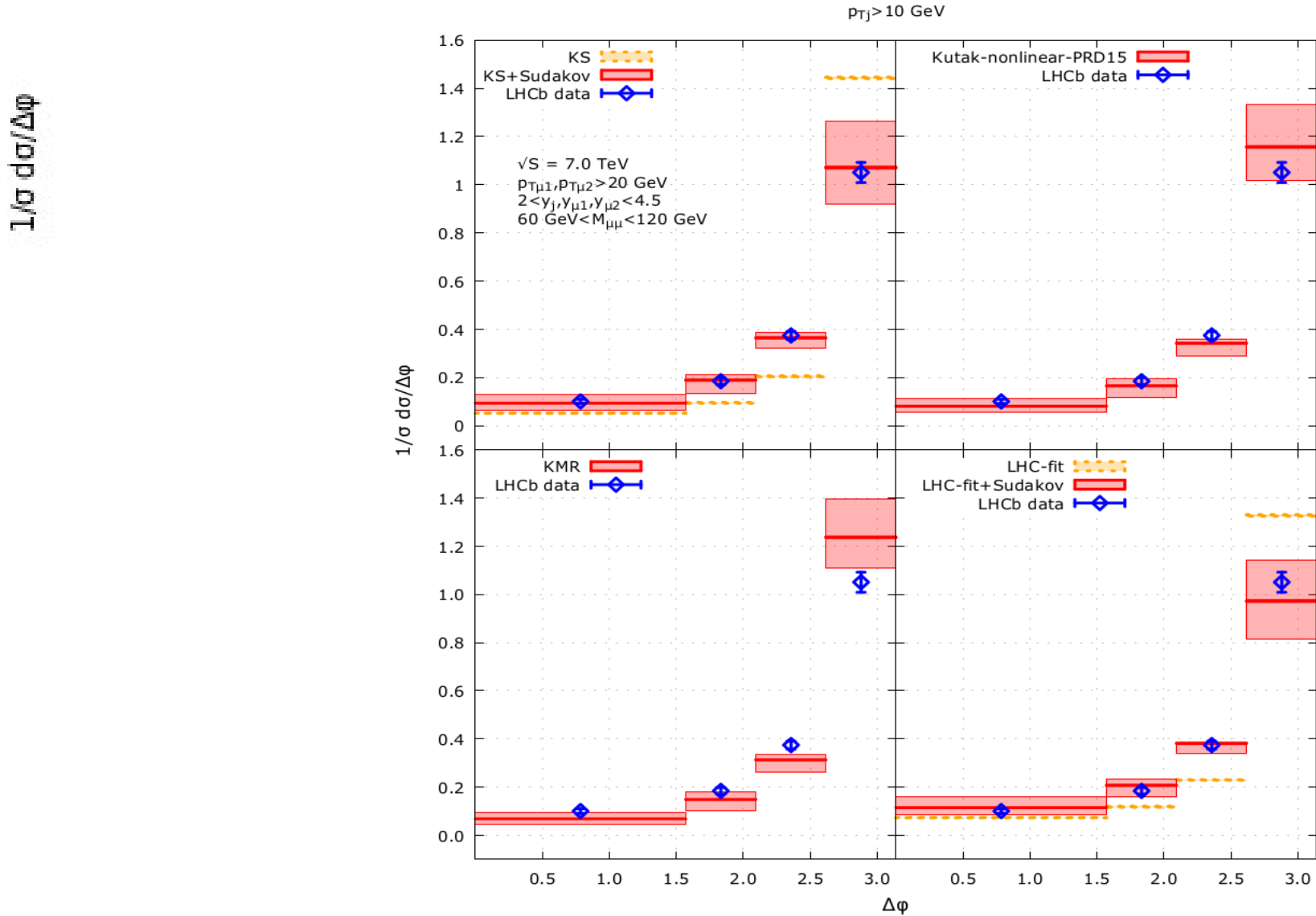
$Z + \text{jet}$



Decorelations in $Z0 + \text{jet}$

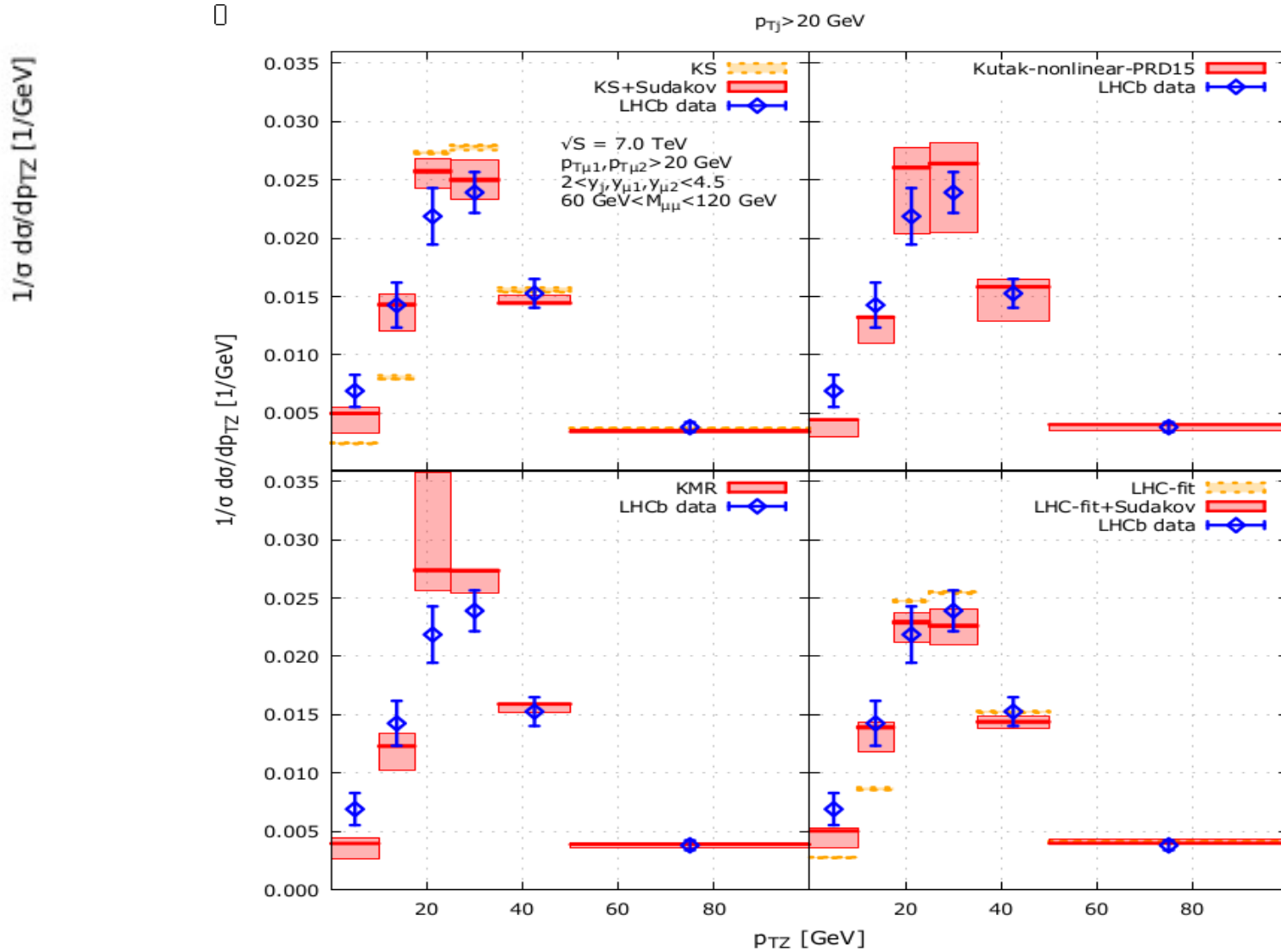
Kotko, KK van Hameren '15

$$g^*q \rightarrow q\mu^+\mu^-, g^*\bar{q} \rightarrow \bar{q}\mu^+\mu^-$$



Z0 pt spectrum

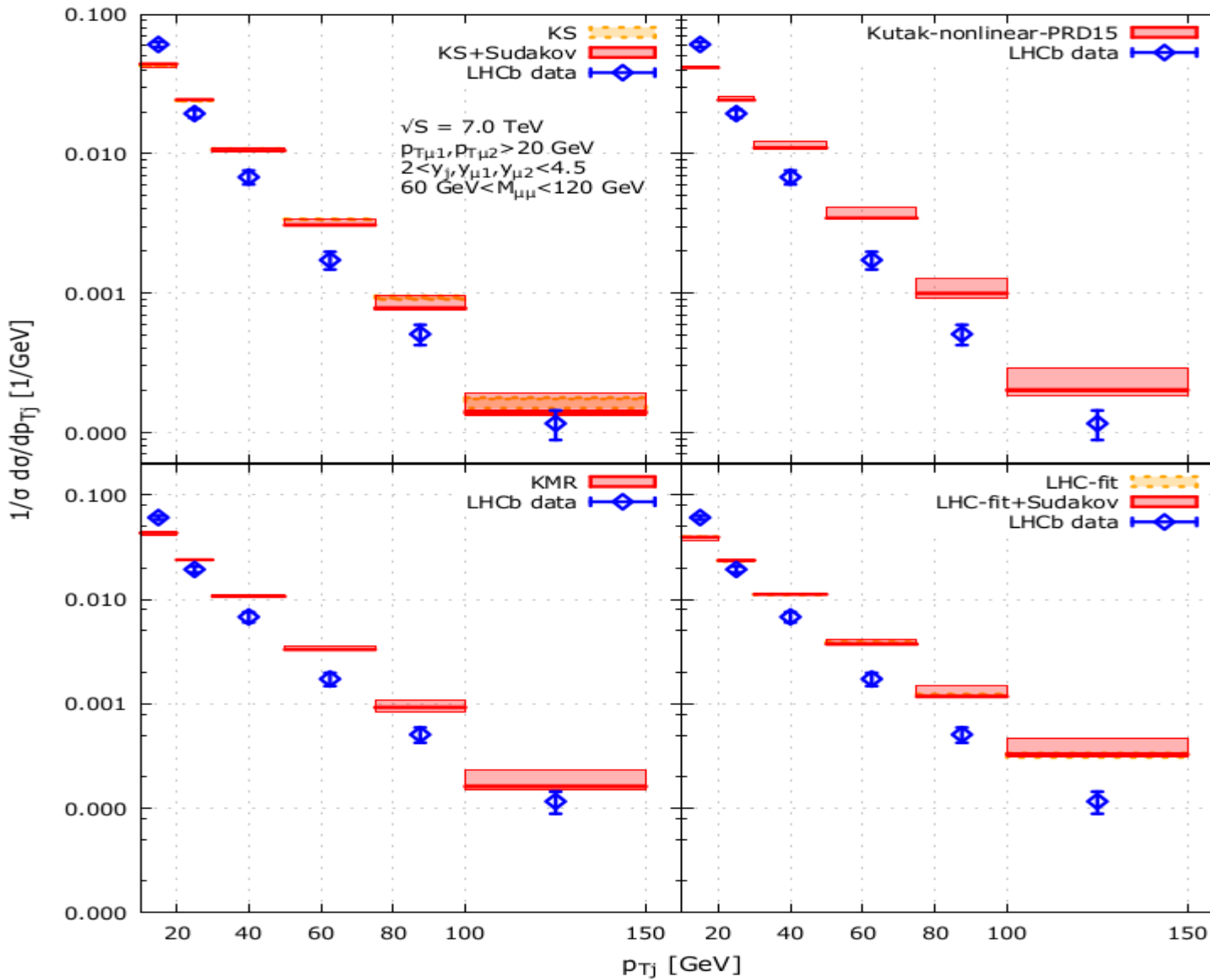
$$g^*q \rightarrow q\mu^+\mu^-, g^*\bar{q} \rightarrow \bar{q}\mu^+\mu^-$$



*Colorless final state.
No color rescatterings.
Description OK*

Message: calculation with hard scale dependent gluon closer to data

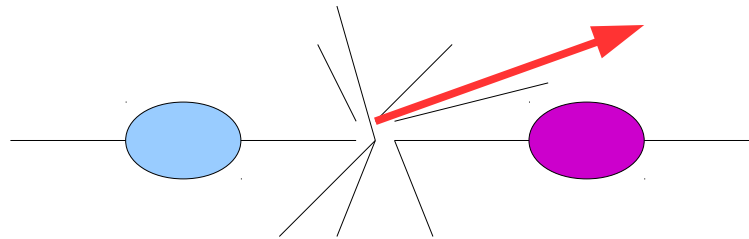
Jet pt spectrum



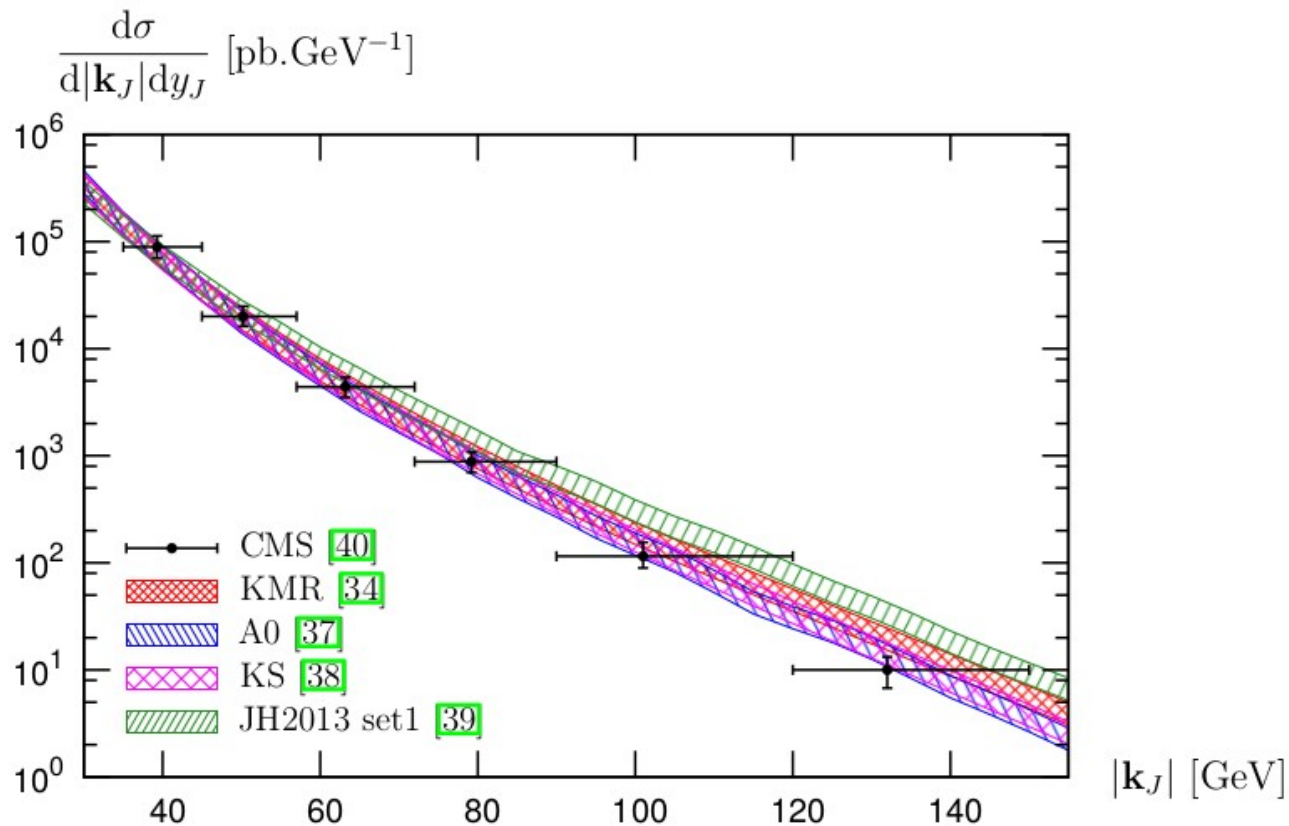
Good tendency but probably effects of final state interactions not taken into account.

Message: possibly larger rescatterings as compared to color neutral Z0 therefore data is not described

Single inclusive forward jet



Single inclusive p_t jet spectra



Decloue, Szymanowski, Wallon '15

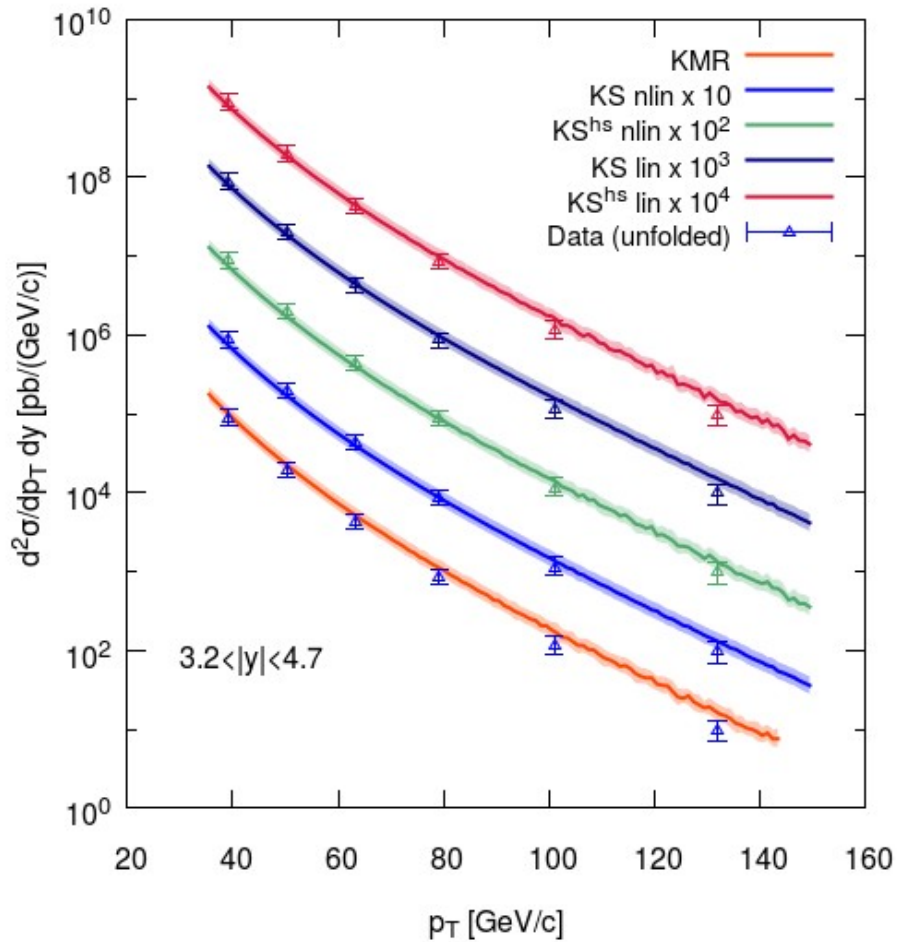
$$\frac{d\sigma}{dy_1 dp_{1t}} = \frac{1}{2} \frac{\pi p_{1,t}}{(x_1 x_2 S)^2} \sum_{a,b,c} \overline{|\mathcal{M}_{ab \rightarrow c}|^2} x_1 f_{a/A}(x_1, \mu^2) \mathcal{F}_{b/B}(x_2, p_{1t}^2, \mu^2)$$

Dumitru, Hayashigski, Jalilian-Marian '05

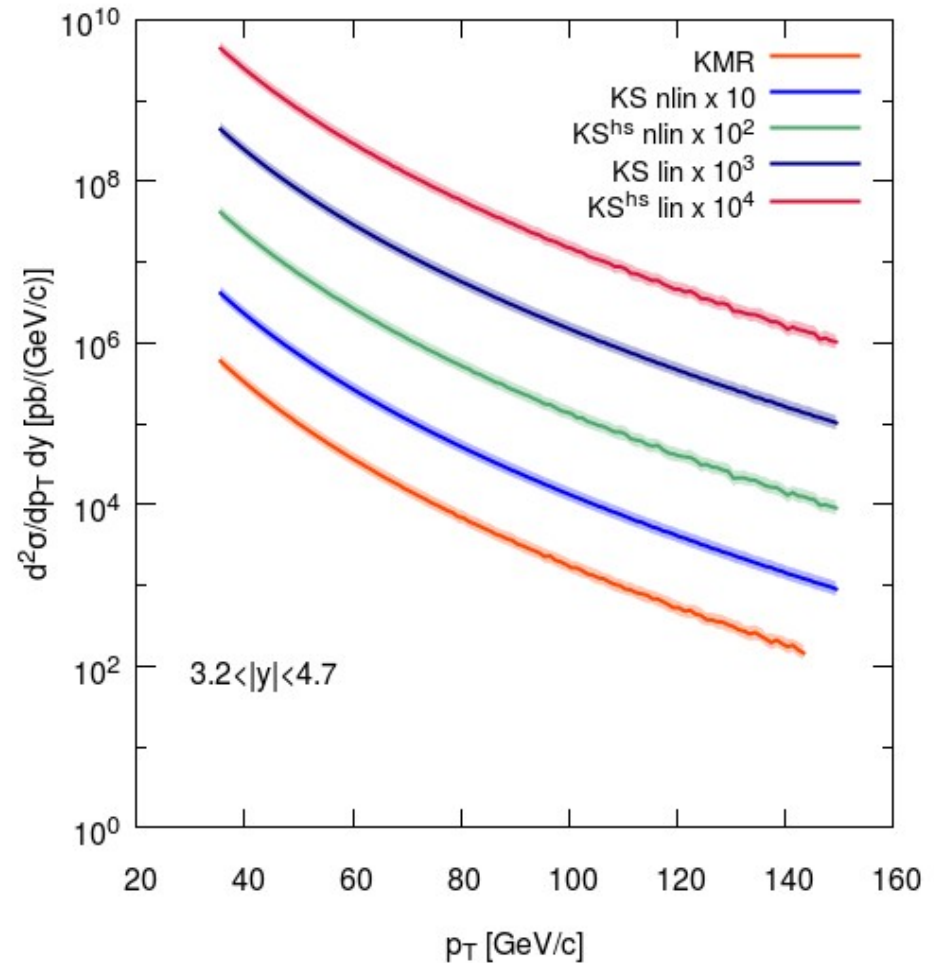
Single inclusive p_T jet spectra

Preliminary *KK, Bury, Sapeta*

CMS, $pp \rightarrow \text{jet}_{\text{fwd}} + X$, $\sqrt{S} = 7$ TeV, $L_{\text{int}} = 3.14 \text{ pb}^{-1}$

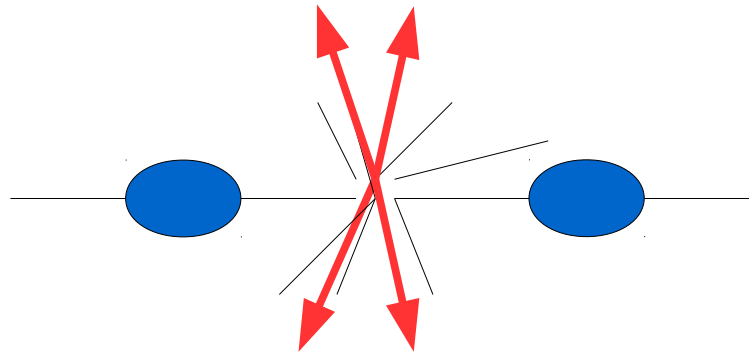


$pp \rightarrow \text{jet}_{\text{fwd}} + X$, $\sqrt{S} = 13$ TeV



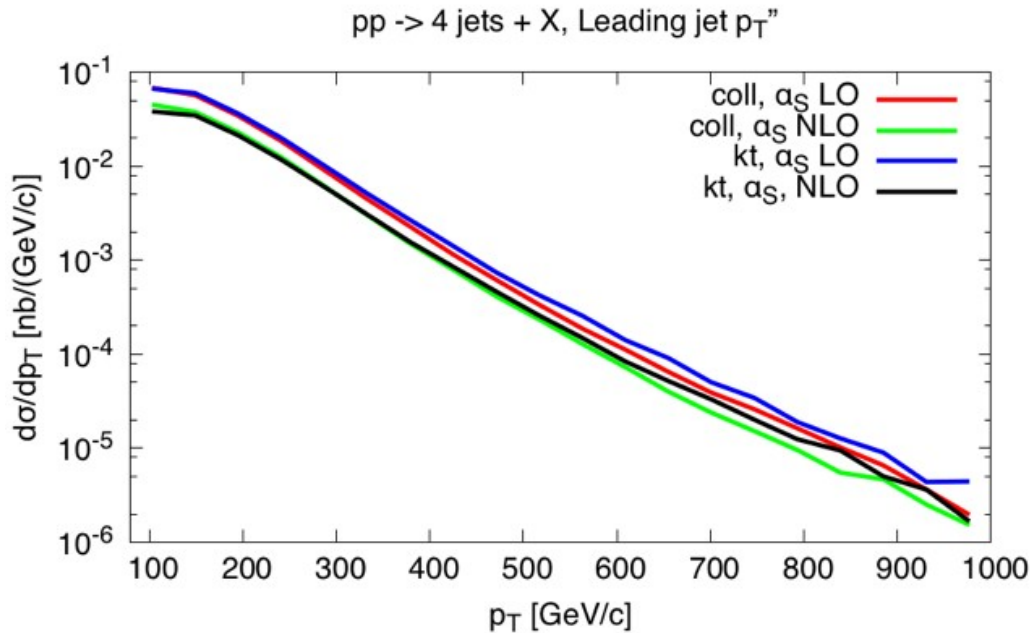
$|3.2| < y < |4.7|$

4 inclusive central jets



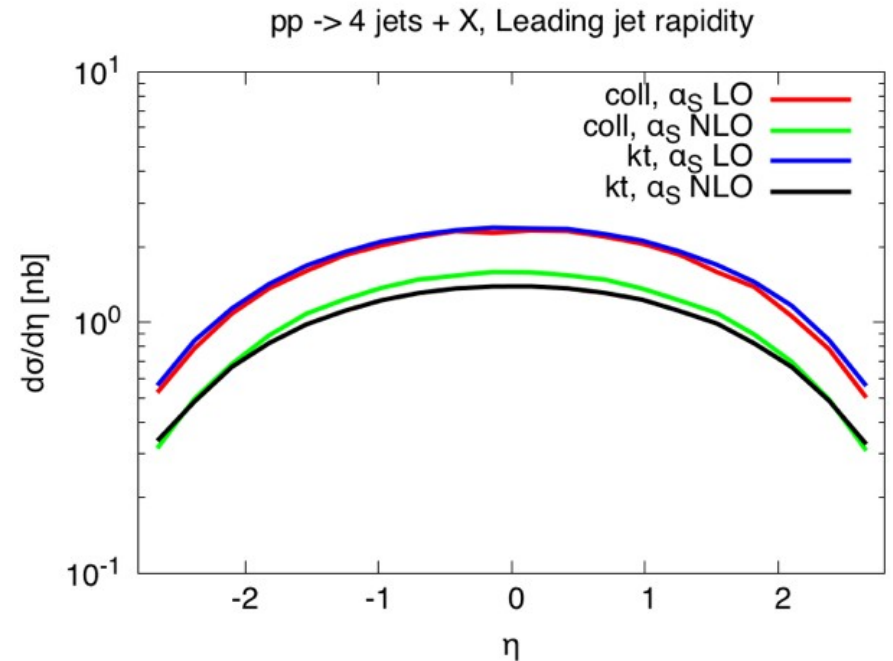
HEF and central jets

pt spectra of leading jets



preliminary

KK, M.Serino, A. van Hameren



- *on both sides unintegrated pdfs.*
- *all channels take into account i.e. off-shell quarks and off-shell gluons*
- *gluon density obtained using the KMR prescription.*
- *the underlying collinear pdf is at NLO accuracy*

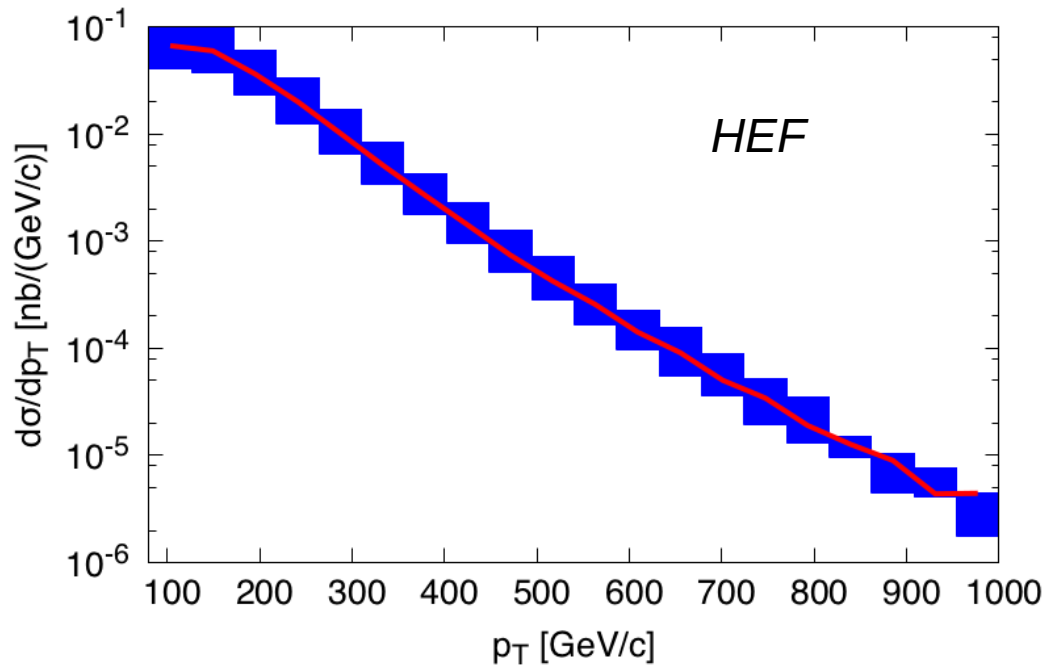
HEF and central jets

preliminary

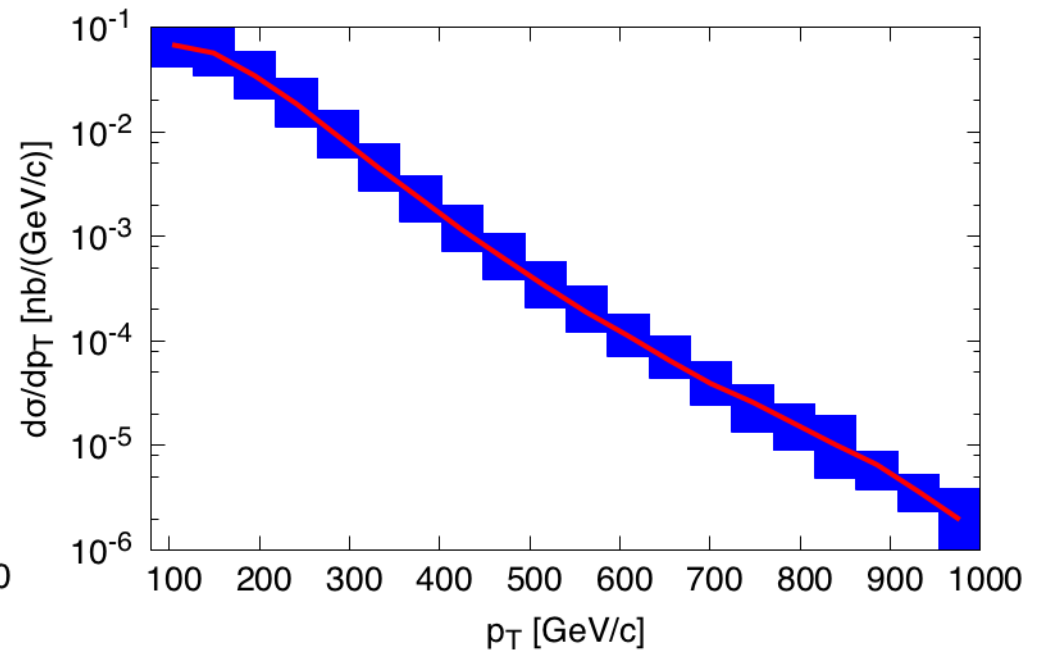
KK, Serino, van Hameren

pt spectra of leading jets

pp -> 4 jets + X, Leading jet, α_S from MSTW2008lo68cl



pp -> 4 jets + X, Leading jet, α_S from MSTW2008lo68cl

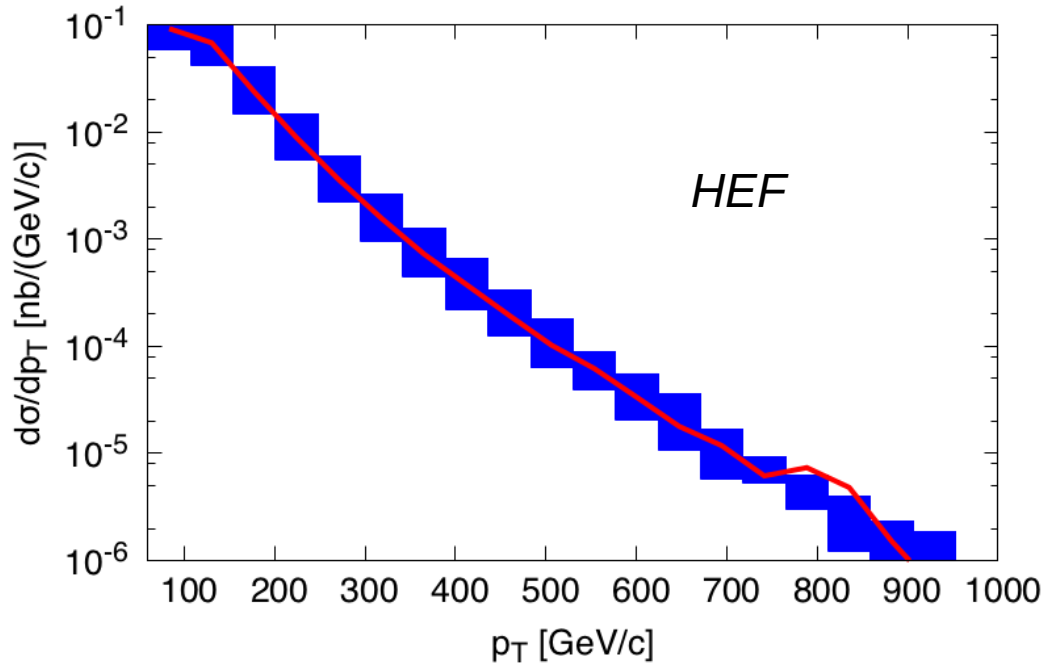


- *on both sides unintegrated pdfs.*
- *all channels take into account i.e. off-shell quarks and off-shell gluons*
- *gluon density obtained using the KMR prescription.*
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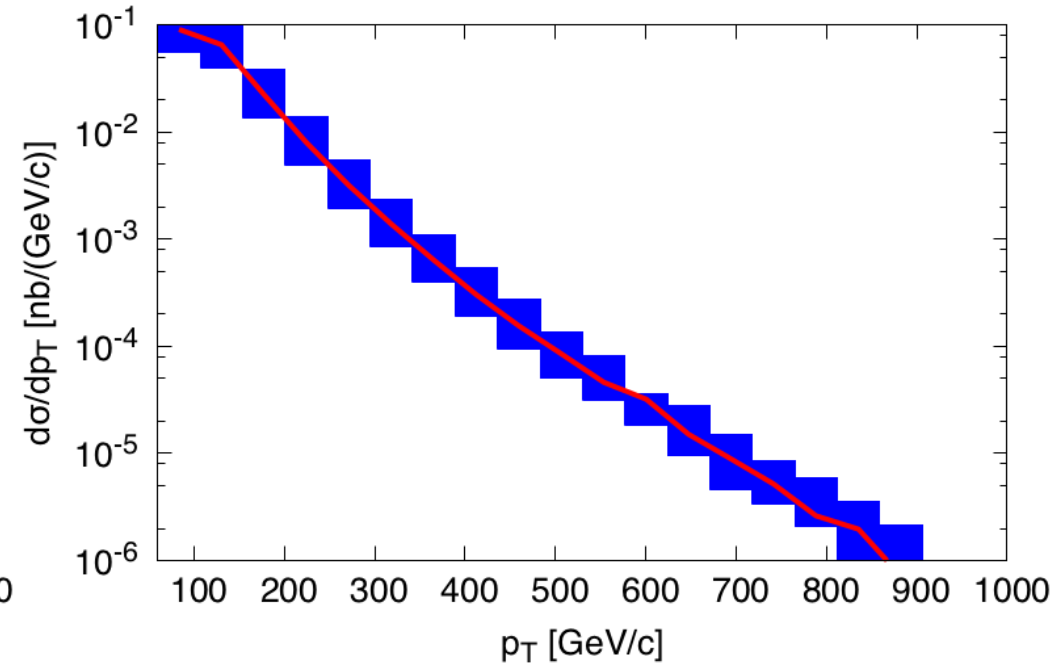
HEF and central jets

pt spectra of second jet

pp → 4 jets + X, 2nd leading jet, α_S from MSTW2008lo68cl



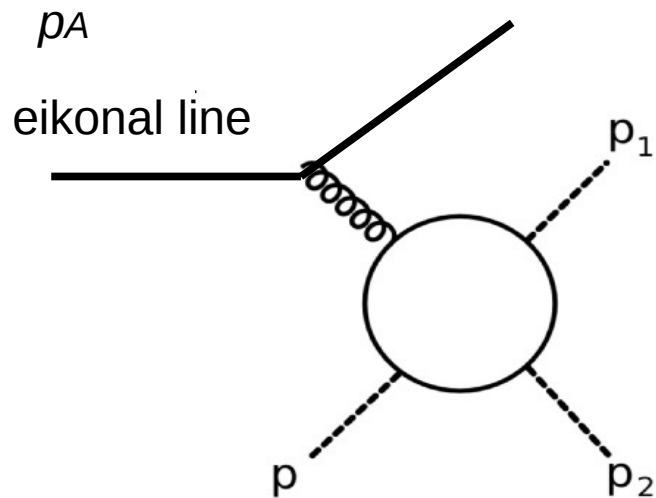
pp → 4 jets + X, 2nd leading jet, α_S from MSTW2008lo68cl



- *on both sides unintegrated pdfs.*
- *all channels take into account i.e. off-shell quarks and off-shell gluons*
- *gluon density obtained using the KMR prescription.*
- *the underlying collinear pdf is at NLO accuracy*

BACK UP

Example of matrix element for hybrid factorization

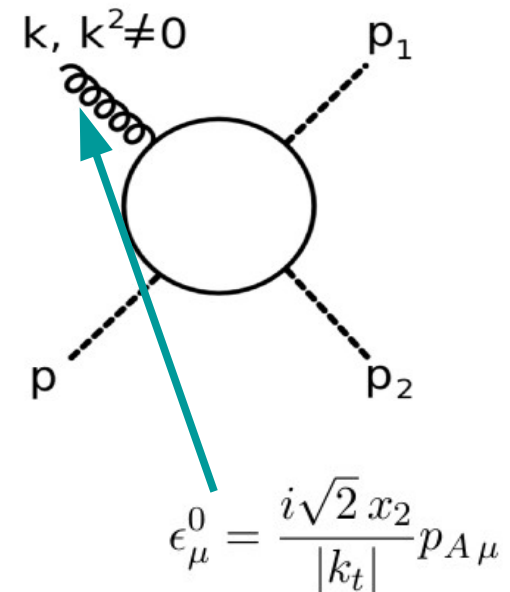


Decomposition
of propagator

$$g_{\mu\nu} = g_{\mu\nu\perp} + \frac{p_{A\mu}p_\nu + p_{A\nu}p_\mu}{s}$$

$$k = x_2 p_A + k_T$$

$$x_2 = k_\mu p^\mu / p_{A\nu} p^\nu$$



$$\epsilon_\mu^0 = \frac{i\sqrt{2} x_2}{|k_t|} p_{A\mu}$$

Polarization sum
for onshell gluons

$$\sum_{\lambda=\pm} \epsilon_\mu^\lambda \epsilon_\nu^{\lambda*} = g_{\mu\nu} - \frac{p_{A\mu}q_\nu + q_\mu p_{A\nu}}{q^\rho p_{A\rho}}$$

Polarization of off-shell gluon

Effective action based approach
Lipatov 95, Lipatov, Vyazovsky 2000

Gauge link based derivation
Kotko'14