

Antenna Subtraction with Hadronic Initial States

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Subtraction

$$\begin{aligned}d\sigma^{\text{NLO}} &= d\sigma_m^{\text{V+CT}} + d\sigma_{m+1}^{\text{R}} \\ &= (d\sigma_m^{\text{V+CT}} + d\sigma_{m+1}^{\text{R,S}}) + (d\sigma_{m+1}^{\text{R}} - d\sigma_{m+1}^{\text{R,S}})\end{aligned}$$

The subtraction term $d\sigma_{m+1}^{\text{R,S}}$

- reproduces the behaviour of the ME in all singular limits
- can be integrated analytically over the unresolved phase space
- does not introduce spurious singularities

Antenna subtraction

Many method available for NLO

- \mathcal{E} prescription [Frixione,Kunszt,Signer]
- Dipole formalism [Catani,Seymour]
- Antenna formalism [Kosower;Campbell,Cullen,Glover]
- ...

We use antenna functions to construct subtraction term

- + less complicated than dipoles (1 Antenna \simeq 2 Dipoles)
- + extension to NNLO is feasible
[Gehrmann-De Ridder,Gehrmann,Glover]
- + antenna functions are constructed from physical matrix elements
- + can be matched to parton shower [Giele,Kosower,Skands]
- no extension to initial state partons (so far)

Antenna functions

$$X_{gg\bar{q}} = \frac{\left[\text{Diagram 1} + \text{Diagram 2} \right]}{\left[\text{Diagram 3} \right]^2}$$

- contains by construction soft and collinear emission between two color connected hard radiators.
- are computed from ME of physical processes

Hadronic reaction

NLO corrections with hadronic initial state

$$d\sigma^{\text{NLO}} =$$

$$2\text{Re} \left(\left(\text{Virtual corrections} \right) + \left| \sum \text{Initial-final} + \sum \text{Initial-initial} \right|^2 \right)$$

- final-state radiation \rightarrow Virtual corrections
- initial-state radiation \rightarrow Virtual corrections+pdf factorisation

Hadronic reaction

$$\left| \Sigma \text{ (diagram 1)} + \Sigma \text{ (diagram 2)} \right|^2$$

$$\sim \underbrace{\left(\text{diagram 1} \text{ } \text{diagram 2} \right)}_{\text{fi nal-fi nal}} + \underbrace{\left(\text{diagram 3} \text{ } \text{diagram 4} \right)}_{\text{initial-fi nal}}$$

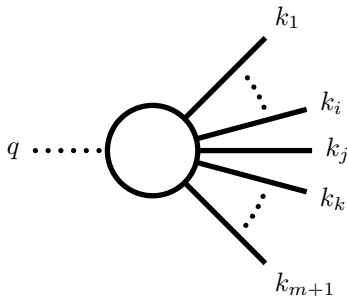
$$+ \underbrace{\left(\text{diagram 5} \text{ } \text{diagram 6} \right)}_{\text{initial-initial}}$$

The diagrams are Feynman diagrams for a hadronic reaction. Each diagram consists of a central circle with four external lines. In diagrams 1 and 2, one line is red and one is dotted. In diagrams 3 and 4, one line is red and one is dotted. In diagrams 5 and 6, one line is red and one is dotted. The diagrams are arranged in three rows, with the first row showing the sum of two diagrams squared, the second row showing the sum of two pairs of diagrams, and the third row showing a third pair of diagrams. The labels 'fi nal-fi nal', 'initial-fi nal', and 'initial-initial' are placed below their respective pairs of diagrams.

Final-Final

We consider

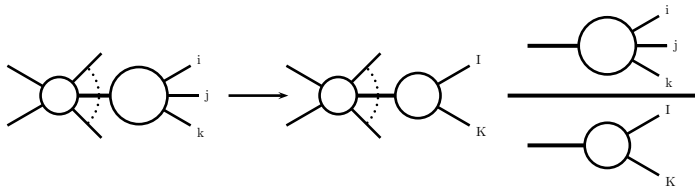
- No partons in the initial state
- Antenna function for hard radiator i, k in the final state
- j unresolved



Final state

$$\text{ME squared: } |\mathcal{M}_{m+1}(k_1, \dots, k_i, k_j, k_k, \dots, k_{m+1}; p_1, p_2)|^2$$

$$\text{Phasespace: } d\Phi_{m+1}(k_1, \dots, k_i, k_j, k_k, \dots, k_{m+1}; p_1, p_2)$$



ME squared, subtraction:

$$\sum_j X_{ijk}(k_i, k_j, k_k) |\mathcal{M}_m(k_1, \dots, K_I, K_K, \dots, k_{m+1}, p_1, p_2)|^2$$

Phasespace:

$$d\Phi_X(k_i, k_j, k_k) \times d\Phi_m(k_1, \dots, K_I, K_K, \dots, k_{m+1}, p_1, p_2)$$

Final-Final Subtraction term

Integrated antenna functions

$$\mathcal{X}_{ijk} = d\Phi_X(\mathbf{k}_i, \mathbf{k}_j, \mathbf{k}_k) X_{ijk}(\mathbf{k}_i, \mathbf{k}_j, \mathbf{k}_k) \quad \mathcal{A}_{qg\bar{q}} = -2\mathbf{I}_{q\bar{q}}^{(1)} + \mathcal{O}(\epsilon^0)$$

For combining with virtual corrections

$$d\sigma^{R,S} = \mathcal{N} \sum_{i,j,k} \mathcal{X}_{ijk} \int d\Phi_m(\dots, \mathbf{K}_I, \mathbf{K}_K, \dots) |\mathcal{M}_m(\dots, \mathbf{K}_I, \mathbf{K}_K, \dots)|^2 \\ \times \mathcal{J}_m^{(m)}(k_1, \dots, \mathbf{K}_I, \mathbf{K}_K, \dots, k_{m+1})$$

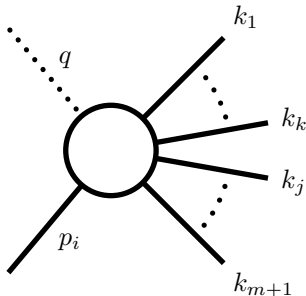
For combining with real corrections

$$d\sigma^{R,S} \\ = \mathcal{N} \int d\Phi_{m+1}(\dots, \mathbf{k}_i, \mathbf{k}_j, \mathbf{k}_l, \dots) \sum_{i,j,k} X_{ijk}(\mathbf{k}_i, \mathbf{k}_j, \mathbf{k}_k) |\mathcal{M}_m(\dots, \mathbf{K}_I, \mathbf{K}_K, \dots)|^2 \\ \times \mathcal{J}_m^{(m)}(k_1, \dots, \mathbf{K}_I, \mathbf{K}_K, \dots, k_{m+1})$$

Initial-Final

We consider

- One parton in the initial state
- Antenna function for one initial i and one final k hard radiator
- j unresolved



Phasespace Mapping

$$\left. \begin{array}{l} p \\ k_j \\ k_k \end{array} \right\} \rightarrow \left\{ \begin{array}{l} P = xp \\ K_K \end{array} \right.$$

$$x = \frac{S_{ij} + S_{ik} - S_{jk}}{S_{ij} + S_{ik}}$$

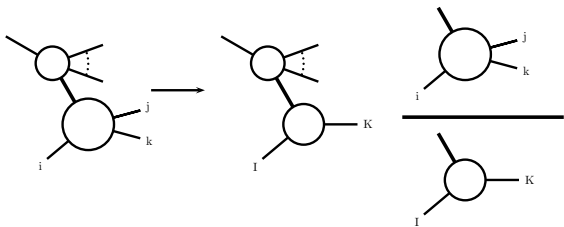
$$K_K = k_k + k_k - xp$$

Limits

$$\begin{array}{lll} xp \rightarrow p, & K_K \rightarrow k_k, & \text{when } j \text{ becomes soft} \\ xp \rightarrow p, & K_K \rightarrow k_j + k_k, & \text{when } j \text{ collinear with } k_k \\ xp \rightarrow p - k_j, & K_K \rightarrow k_k, & \text{when } j \text{ collinear with } p \end{array}$$

ME squared: $|\mathcal{M}(k_1, \dots, k_j, k_k, \dots, k_{m+1}; q, p)|^2$

Phasespace: $d\Phi_{m+1}(k_1, \dots, k_j, k_k, \dots, k_{m+1}; q, p)$



ME squared, subtraction:

$$X_{ijk}(k_j, k_k; p) |\mathcal{M}(k_1, \dots, K_K, \dots, k_{m+1}, q, xp)|^2$$

Phasespace:

$$d\Phi_X(k_j, k_k; q, p) \mathcal{J}(k_j, k_k, p) d\Phi_m(k_1, \dots, K_K, \dots, k_{m+1}, q, xp)$$

Subtraction Term for Initial-Final

Integrated antenna functions

$$\mathcal{X}_{i,jk}(x) = d\Phi_X(k_j, k_k; p, q) X_{i,jk}(k_j, k_k; p) \mathcal{J}(k_j, k_k; p)$$

$$\mathcal{A}_{q,qq}(x) = -2I_{q\bar{q}}^{(1)} \delta(1-x) + (Q^2)^{-\epsilon} \left[-\frac{1}{2\epsilon} P_{qq}(x) + \mathcal{O}(\epsilon^0) \right]$$

$d\sigma^{NLO,S}$

$$\begin{aligned} &= \mathcal{N} \sum_{ijk} \mathcal{X}_{i,jk}(x) \\ &\quad \times d\Phi_m(k_1, \dots, K_K, \dots, k_{m+1}; q, xp) \\ &\quad \times |\mathcal{M}_m(k_1, \dots, K_K, \dots, k_{m+1}; xp)|^2 J_m^{(m)}(k_1, \dots, K_K, \dots, k_{m+1}). \end{aligned}$$

Cancellation of the divergencies

Hadronic Cross section

$$\sigma_h = \mathcal{N} \int_0^1 \frac{d\xi}{\xi} f_i(\xi) \hat{\sigma}(\xi p_i)$$

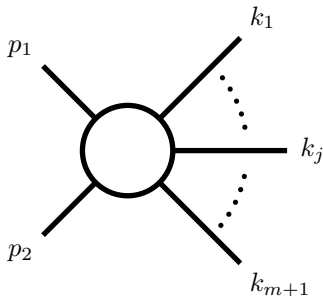
Mass factorization of the pdfs

$$f_i(\xi) = \sum_j \int_{\xi}^1 \frac{dx}{x} f_j\left(\frac{\xi}{x}\right) \left(\delta(1-x) \delta_{ij} + \frac{\alpha_s}{2\pi\epsilon} P_{ij}(x) \right)$$

Initial-Initial

We consider

- Two partons in the initial state
- Antenna function for initial p_1 and p_2 hard radiator
- j unresolved



Phasespace Mapping

$$\left. \begin{matrix} p_1 \\ p_2 \\ k_j \end{matrix} \right\} \rightarrow \begin{cases} P_1 = x_1 p_1 + (1 - x_2) p_2 - \frac{1 - \beta}{2} k_j + \alpha t \\ P_2 = (1 - x_1) p_1 + x_2 p_2 - \frac{1 + \beta}{2} k_j - \alpha t \end{cases}$$

where $t \perp (p_1, p_2, k_j)$, $t^\mu = \frac{\epsilon^{\mu\nu\rho\sigma} p_1^\nu p_2^\rho k_j^\sigma}{2p_1 p_2}$, $t^2 = -\frac{p_1 \cdot k_j p_2 \cdot k_j}{2p_1 \cdot p_2}$

$$\beta = \frac{s_{2j} - s_{1j}}{s_{1j} + s_{2j}} \quad \alpha \text{ free}$$

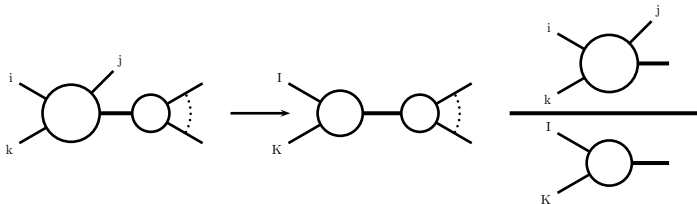
$P_1 \rightarrow p_1$, $P_2 \rightarrow p_2$, when k becomes soft

$P_1 \rightarrow p_1 - k_j$, $P_2 \rightarrow p_2$, when j collinear with p_1

$P_1 \rightarrow p_1$, $P_2 \rightarrow p_2 - k_j$, when j collinear with p_2

ME squared: $|\mathcal{M}_{m+1}(k_1, \dots, k_{j-1}, k_j, k_{j+1}, \dots, k_{m+1}; p_i, p_j)|^2$

Phasespace: $d\Phi_{m+1}(k_1, \dots, k_{j-1}, k_j, k_{j+1}, \dots, k_{m+1}; p_1, p_2)$



ME squared, subtraction:

$$\sum_k X_{12,k}(k_j; p_1, p_2) |\mathcal{M}_m(k_1, \dots, k_{j-1}, k_{j+1}, \dots, k_{m+1}, P_1, P_2)|^2$$

Phasespace:

$$d\Phi_X(k_j; p_1, p_2) \mathcal{J}(k_j, p_1, p_2) d\Phi_m(k_1, \dots, k_{j-1}, k_{j+1}, \dots, k_{m+1}; P_1, P_2)$$

Integrated Antenna function

Phase space integration freezes the value of x_1, x_2

$$\mathcal{A}_{q\bar{q},g} = -2\delta(1-x_1)\delta(1-x_2)I_{q\bar{q},g}^{(1)} - s^{-\epsilon} \frac{P_{qq}(x_1)\delta(1-x_2) + P_{qq}(x_2)\delta(1-x_1)}{2\epsilon} + \mathcal{O}(\epsilon)$$

Final Form of The Subtraction term

$$\begin{aligned}
 d\sigma_{NLO}^{S,(ii)} &= \mathcal{N} d\Phi_{m+1}(k_1, \dots, k_{k-1}, k_k, k_{k+1}, \dots, k_{m+1}; p_1, p_2) \\
 &\quad \times \sum_k X_{12,k}^0 |\mathcal{M}_m(k_1, \dots, k_{k-1}, k_{k+1}, \dots, k_{m+1}; P_1, P_2)|^2 \\
 &\quad J_m^{(m)}(k_1, \dots, k_{k-1}, k_{k+1}, \dots, k_{m+1}) \\
 &= \mathcal{N} \sum_k \mathcal{X}_{12,k}(x_1, x_2) d\Phi_m(k_1, \dots, k_{k-1}, k_{k+1}, k_{m+1}; P_1, P_2) \\
 &\quad \times |\mathcal{M}_m(k_1, \dots, k_{k-1}, k_{k+1}, \dots, k_{m+1}; P_1, P_2)|^2 \\
 &\quad J_m^{(m)}(k_1, \dots, k_{k-1}, k_{k+1}, \dots, k_{m+1})
 \end{aligned}$$

Conclusion and Outlook

- formulated antenna subtraction for processes with initial state partons (at NLO)
- method can be extended to NNLO (work in progress)
- possible applications
 - matching to antenna-based parton shower
 - NNLO $t\bar{t}$ production
 - NNLO $pp \rightarrow 2j$ production
 - NNLO DIS $(2 + 1)j$ production