

Measurement of the top quark mass in final states with three jets and one charged lepton at $\sqrt{s} = 8 \text{ TeV}$ with the ATLAS experiment

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Content

Motivation for the Analysis Analysis

Event Selection

Background Estimation

Neural Network

Template Method

Ensemble Tests and

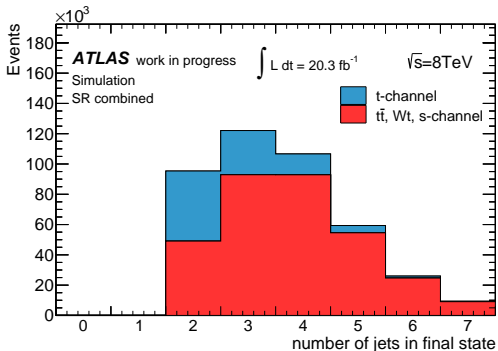
Estimation of Systematic Uncertainties

Optimization of the NN Cut Value

Conclusion

Motivation for the Analysis

- first mass measurement in data enhanced with single top quarks with two jets in the final state in 2014 (ATLAS-CONF-2014-055)
- now: modified phase space with **three jets** in the final state
- orthogonal phase space compared to any other selection in a top quark mass measurement done before



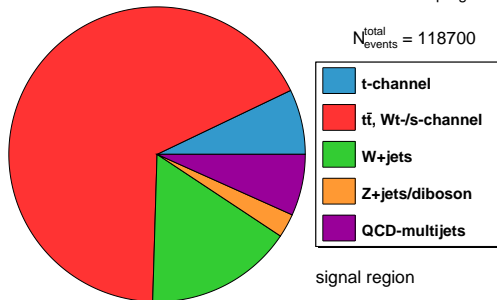
(preselection with ≥ 2 jets)

- more statistics for another top quark mass measurement
- composition of signal in 3 jet-bin:
 - ▶ 80 % Wt -, s-channel and $t\bar{t}$ processes
 - ▶ 20 % t -channel process

Event Selection

- exactly **3 jets** with $p_T > 30$ GeV: 1 b -tagged jet (MV1c50), 2 untagged jets (light jets)
 - exactly 1 charged lepton with $p_T > 25$ GeV
 - missing transverse momentum with $E_T^{\text{miss}} > 30$ GeV
 - triangular cut: $p_T(\ell) > 40 \text{ GeV} \cdot \left(1 - \frac{\pi - |\Delta\varphi(j_1, \ell)|}{\pi - 1}\right)$
 - transverse W -boson mass: $m_T(W) > 50$ GeV
- } cuts to reject QCD-multijet events

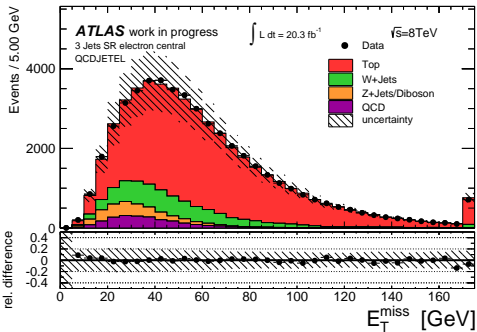
ATLAS work in progress



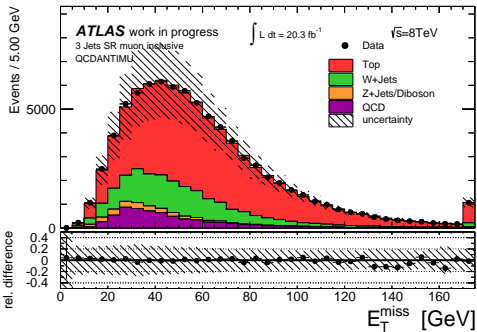
- selected sample is dominated by $t\bar{t}$ events

Background Estimation

- contribution of the QCD-multijet events is estimated by the use of two models:
 - ▶ electron channel: **jet-lepton model**; muon channel: **anti-muon model**
- data-driven determination of normalization for the QCD-models
- likelihood fit to the distribution of the missing transverse momentum

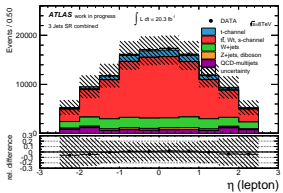
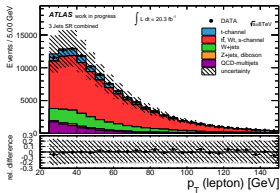
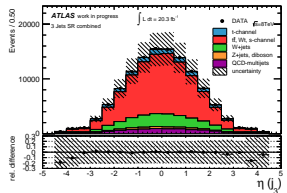
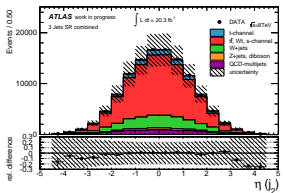
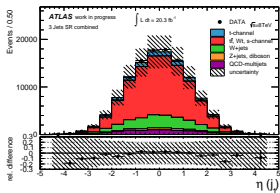
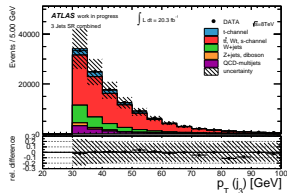
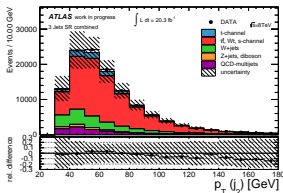
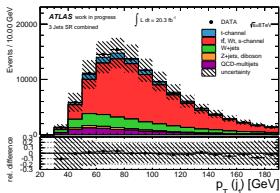


electron channel



muon channel

Control Plots



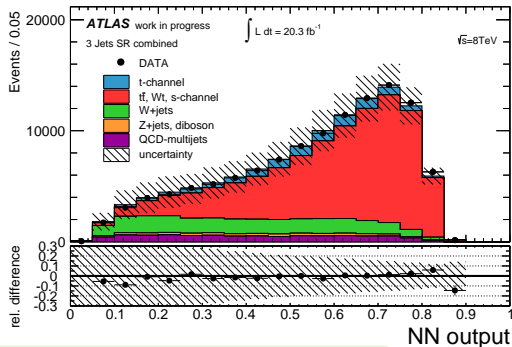
⇒ good agreement between Monte Carlo simulation and data in signal and control regions (not shown)

Neural Network

- use of a multivariate method as in the analysis with 2 jets (ATLAS-CONF-2014-055)
- training of all top processes ($t\bar{t}$, t -, Wt -, s -channel) versus W +jets, Z +jets and diboson processes
- eleven input variables, adopted from the t -channel cross section measurement (ATLAS-CONF-2012-132)

ATLAS work in progress

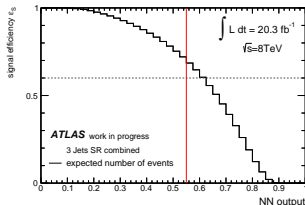
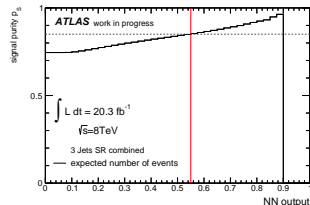
variable	ρ^{iter} [%]	σ^{iter}
$ \Delta\eta(\ell, b) $	30.26	124.41
$H_T(\ell, \text{jets}, E_T^{\text{miss}})$	18.22	74.92
$m(\ell, b)$	12.86	52.87
$m(j_1, j_2)$	10.63	43.70
$\Delta R(j_2, b)$	9.55	39.29
$\eta(\ell, E_T^{\text{miss}}, b)$	6.44	26.46
$m(j_1, j_3)$	6.51	26.75
$ \eta(j_2) $	1.62	6.65
$\sum_{i=1}^3 \eta(j_i)$	1.29	5.31
$ \eta(j_1) $	1.65	6.77
$ \Delta\eta(j_2, j_3) $	0.43	1.76



⇒ good separation between signal and background events

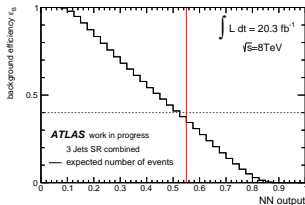
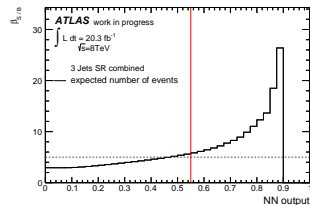
Choice of Cut Value

- selection of events exceeding a minimal neural network output value
- cut value fulfills minimal requirements on statistical quantities:
 - ▶ signal purity p_S and efficiency ε_S , ratio of signal to background events $\beta_{S/B}$ and background rejection ε_B



ATLAS work in progress

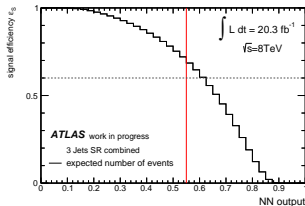
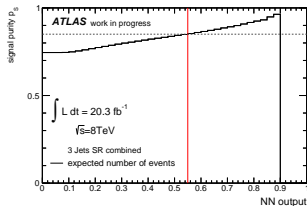
statistical quantity	arbitrary threshold
p_S	> 85 %
ε_S	> 60 %
$\beta_{S/B}$	> 5
ε_B	< 40 %



⇒ **NN > 0.55**

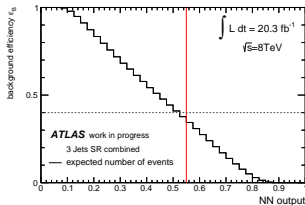
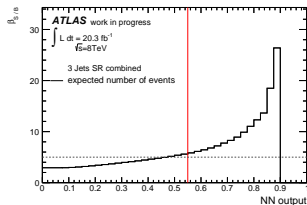
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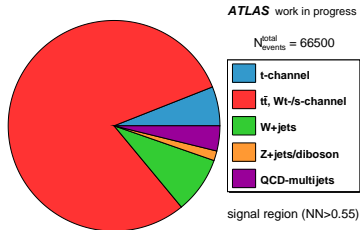
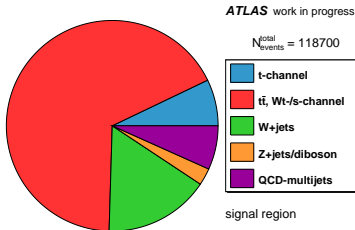
first guess:

$$\Rightarrow \boxed{NN > 0.55}$$

later on: cut variations to evaluate the minimal uncertainty on the measured mass

Event Yields

- decrease of total background fraction from 25 % to 14 %

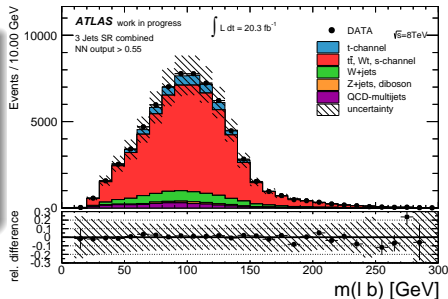
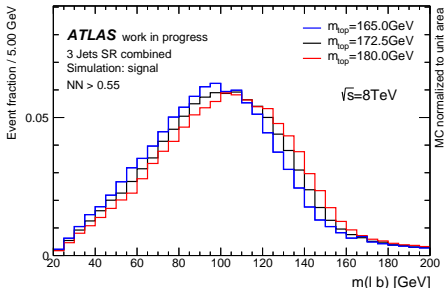


ATLAS work in progress

	signal region ($NN > 0.55$) process	combined channel total
expected signal events	t-channel	4000 ± 400
	s-channel	246 ± 25
	Wt-channel	3460 ± 350
	t \bar{t}	$49\,500 \pm 5000$
	W+jets	5800 ± 3500
expected background events	Z+jets/diboson	1000 ± 100
	QCD-multijets	2600 ± 1300
	total expected data	$66\,500 \pm 6500$
	data	67194
	bkgd. fraction r_{MC} [%]	14.0 ± 5.0

Template Method for Top Quark Mass Measurement

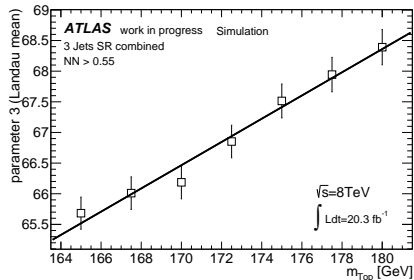
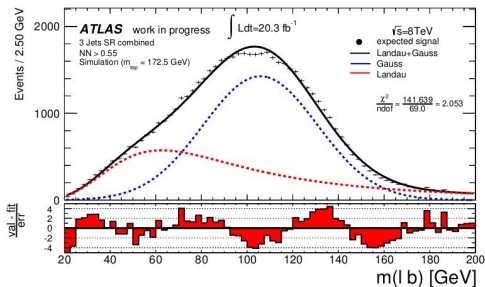
- $m(\ell b)$ -distribution depends on the top quark mass m_{top}
- used as estimator sensitive to m_{top} in the template method
- distribution must be effectively parametrized



- two templates are generated separately for the signal and background contribution
- 7 discrete mass points are simulated in a range from 165.0 GeV to 180.0 GeV
 - ▶ generating a **mass dependent** signal template

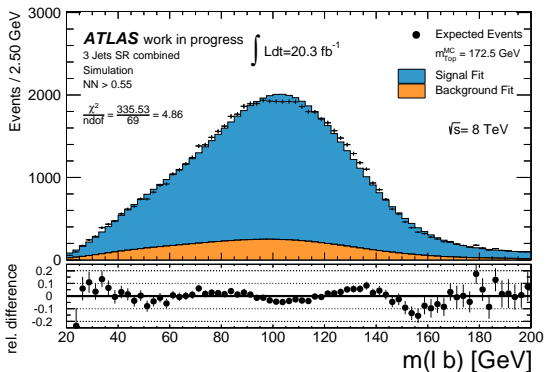
Parametrization and Calibration Curves

- signal and background distributions can be parametrized by the same function
 - ▶ sum of Landau and Gaussian function
- parameters are linearly interpolated in dependence of the top quark mass
 - ▶ mass dependent calibration curves



- final templates are given by the probability density functions $P_{\text{signal}}(m(\text{lb})|m_{\text{top}})$ and $P_{\text{bkgd}}(m(\text{lb}))$

Template Fit



Final template fit has three parameters

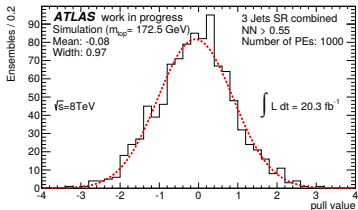
1. Top quark mass: m_{top}
2. Background fraction: f_{back}
3. Overall normalization: N

Estimated templates are used as the input to a binned maximum likelihood fit to the data with:

$$\mathcal{L}(m_{\text{top}}, N, f_{\text{back}}) = \prod_{\text{bin}} \text{Poisson}_{\lambda_{\text{bin}}} \left(m(lb)_{\text{bin}}^{\text{data}} \right) \cdot G(f_{\text{back}}, r_{\text{MC}}, \sigma_{r_{\text{MC}}})$$

$$\lambda_{\text{bin}} = N \cdot [(1 - f_{\text{back}}) P_{\text{signal}}(m(lb)_{\text{bin}} | m_{\text{top}}) + f_{\text{back}} P_{\text{bkgd}}(m(lb)_{\text{bin}})]$$

Statistical Validation and Estimation of Systematic Uncertainties



- method is tested by generating sets of pseudodata
- each mass point is validated
- constructing pull distributions: $\text{pull} = \frac{\langle m_{\text{top}}^{\text{fit}} \rangle - m_{\text{top}}^{\text{MC}}}{\sigma_{\text{top}}^{\text{fit}}}$
 - ▶ pull distribution: $\mu \stackrel{!}{=} 0$ and $\sigma \stackrel{!}{=} 1$

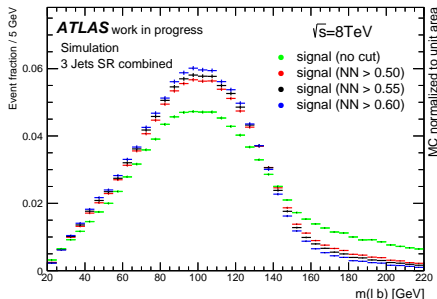
- various sources of systematic uncertainties influence the measurement
 - ▶ object energy scale/resolution and efficiencies
 - ▶ modeling uncertainties of signal processes
 - ▶ modeling uncertainties of background processes

dominant systematic sources:

- jet energy scale
- electron energy scale
- $t\bar{t}$ ISR/FSR
- $t\bar{t}$ MC generator
- QCD-multijet normalization

Influence of the NN Cut Value on the Total Uncertainty

- $m(\ell b)$ -distribution depends on the cut value selected
- four cut scenarios are assumed
 - ▶ no cut, $NN > 0.50$, $NN > 0.55$, $NN > 0.60$
- for each scenario own templates are constructed
- statistical validations show no deviations



ATLAS work in progress

threshold value	stat. unc. Δm_{top}	syst. unc. Δm_{top}	tot. unc. Δm_{top}
no cut	<i>a</i>	<i>b</i>	
$NN > 0.50$	+15 %	-15 %	
$NN > 0.55$	+27 %	-20 %	
$NN > 0.60$	+42 %	-25 %	

- total uncertainty decreases with increasing cut value
 - ▶ more studies necessary to evaluate the optimal cut value

Summary and Conclusion

Summary

- selection of events in a phase space that has never been used in a top quark mass measurement before
- full analysis has been done
- studies on different cut values on the neural network output distribution have been performed

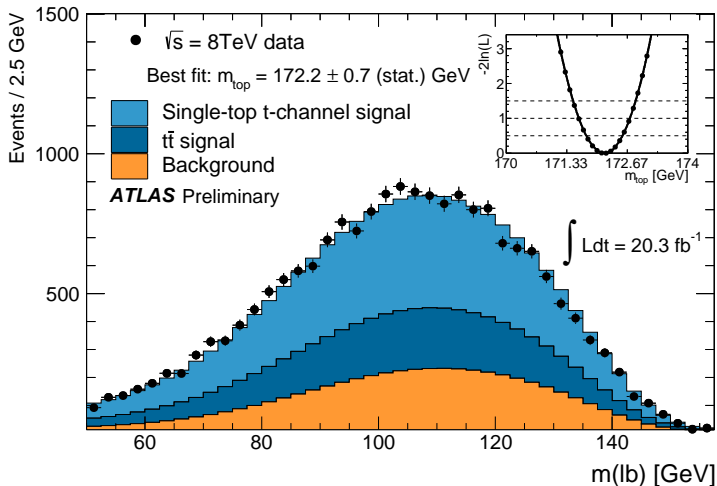
Conclusion

- investigation of this phase space for a top quark mass measurement has good prospects

Backup

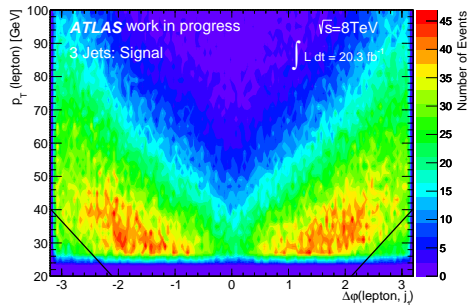
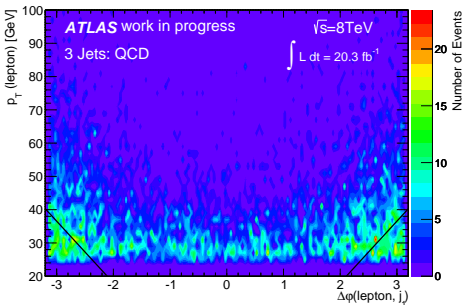
Result of the Analysis with 2 Jets in the Final State

(ATLAS-CONF-2014-055)



$$m_{\text{top}} = 172.2 \pm 0.7(\text{stat.}) \pm 2.0(\text{syst.}) \text{ GeV}$$

Visualization of the Triangular Cut



$$p_T(\ell) > 40 \text{ GeV} \cdot \left(1 - \frac{\pi - |\Delta\varphi(j_1, \ell)|}{\pi - 1}\right)$$

- cut is used to suppress QCD-multijet events
- they mostly arise from dijet events that show a different kinematic signature compared to signal events

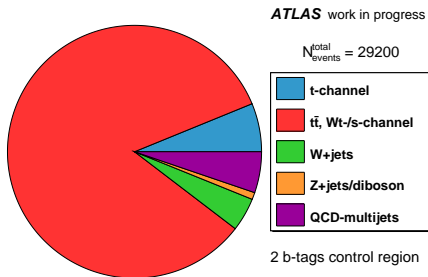
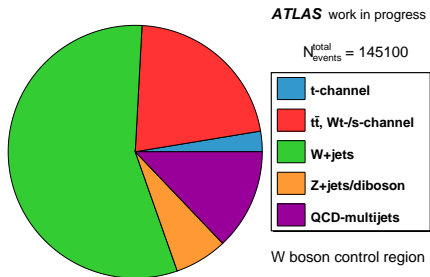
Control Regions

W-boson control region

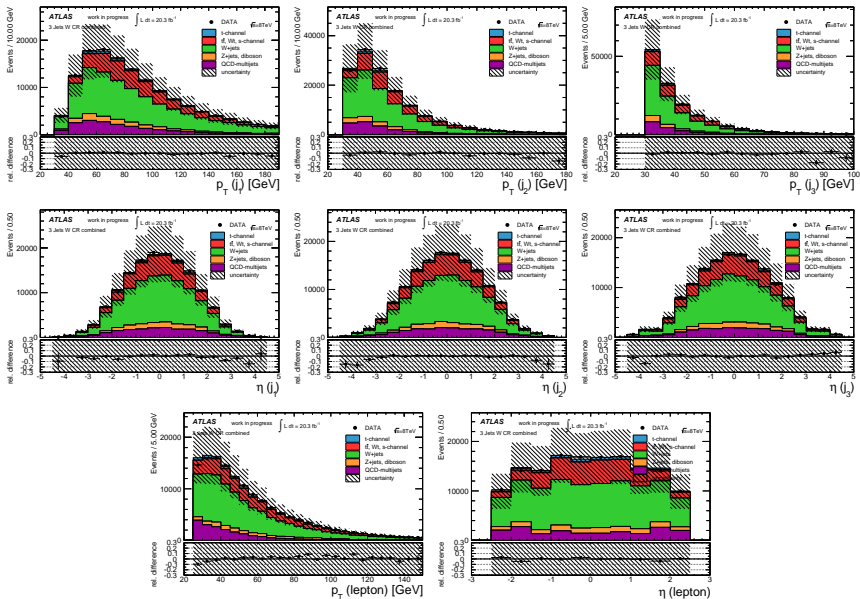
- looser b -tagging efficiency
- signal region is excluded

2 b -tags control region

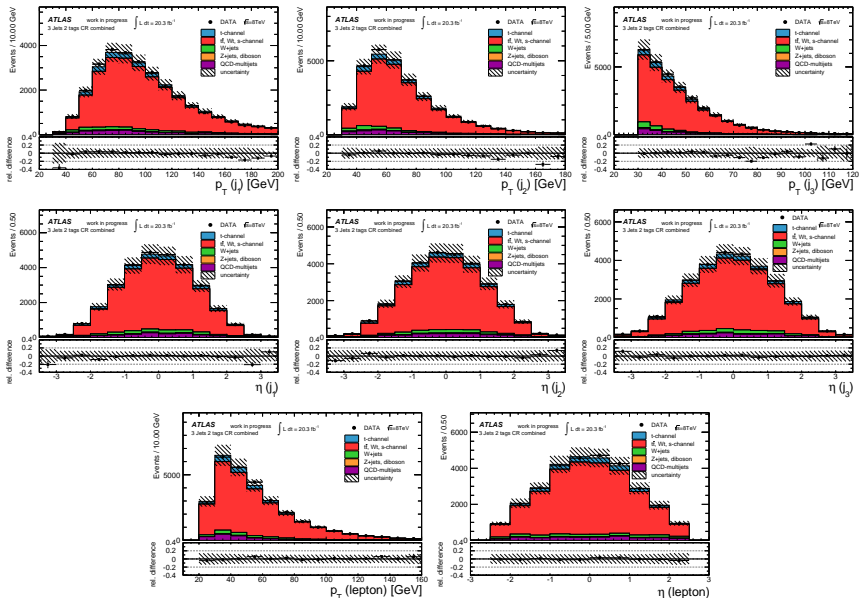
- second b -tagged jet is required
- no overlap to signal region



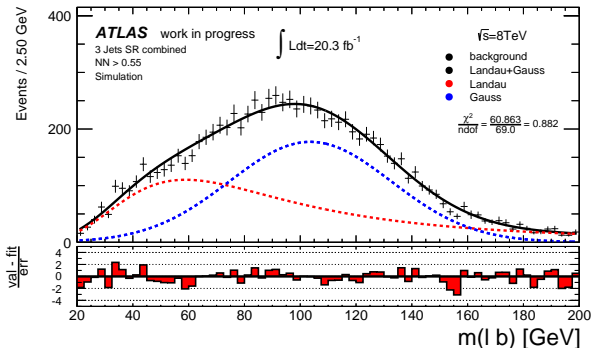
Control Plots in the W-Boson Control Regions



Control Plots in the 2 b -tags Control Regions



Parametrization of the Background Contribution to the $m(\ell b)$ -Distribution



- the same effective parametrization as for the signal contribution:

$$f(p_1 \dots p_6, x = m(\ell b)) = p_1 \cdot [p_2 \cdot L(x, p_3, p_4) + (1 - p_2) \cdot G(x, p_5, p_6)]$$