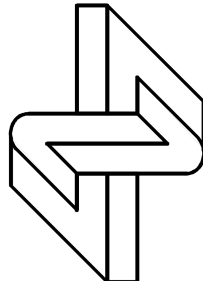


Reconstruction of the Higgs mass in $H \rightarrow \tau\tau$ events

Christian Veelken
NICPB Tallinn

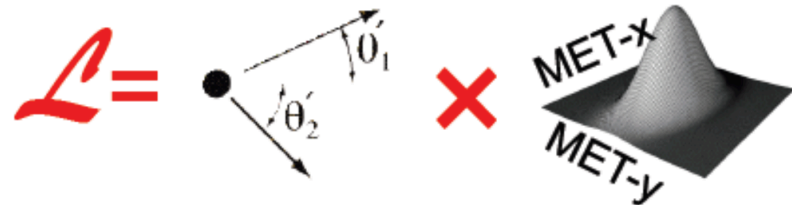


$\tau\tau$ Workshop DESY, November 16th 2015

Reminder: What is SVfit ?

Algorithm for reconstruction of Higgs mass in $H \rightarrow \tau\tau$ events, based on **likelihood method**, using as input:

- Measured e, μ, τ_h -jet momenta
- Reconstructed E_T^{miss} and event-by-event estimate of E_T^{miss} resolution (E_T^{miss} reconstructed by a multivariate regression technique)



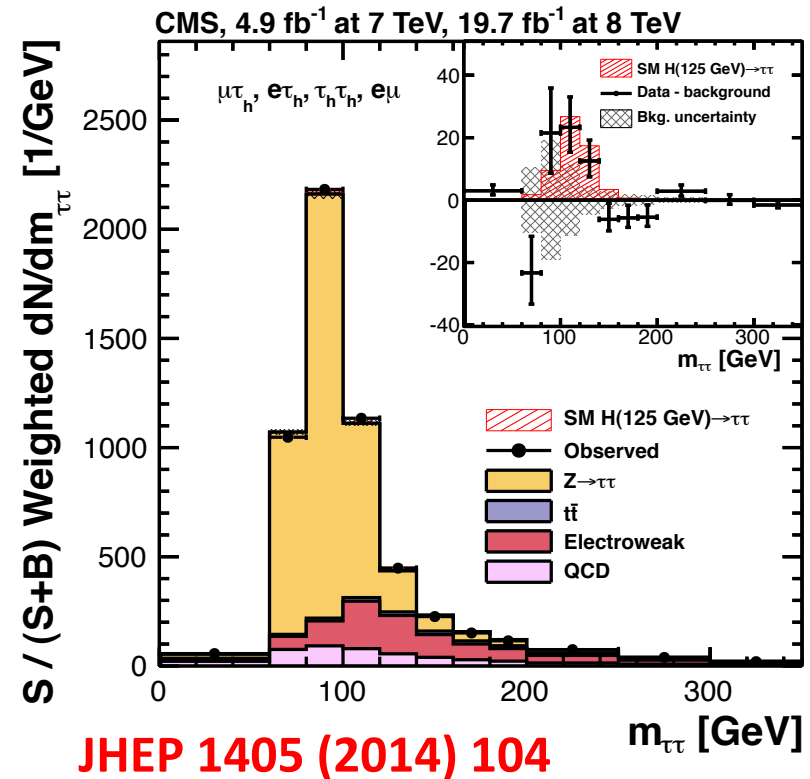
Focus of today's presentation:

- New version of SVfit rigorously based on matrix element (ME) method
- Comparison with Missing Mass Calculator (MMC)

Motivation for ME version of SVfit

$H \rightarrow \tau\tau$ signal shows-up as small bump on tail of large irreducible $Z/\gamma^* \rightarrow \tau\tau$ background

→ Sensitivity of the $H \rightarrow \tau\tau$ analysis crucially depends on good mass resolution to separate the signal from the background



Was “old” version of SVfit using an assumption that degrades the resolution ?

ME method provides clear description how to account for energy/ p_T resolution of τ_h -jet, which is non-negligible in CMS

Develop formalism to handle τ decays in ME method, which we can use for future applications in other analyses

Parametrization of τ Decays

Leptonic τ decays:

$\tau \rightarrow \ell \nu \nu$ decays parametrized by 3 variables

- θ^*
 - ϕ^*
 - $m_{\nu\nu}$
- } decay angles in τ restframe

Hadronic τ decays:

Treated as two-body decays into a hadronic system τ_h and a ν_τ

$\tau \rightarrow \tau_h \nu_\tau$ decays parametrized by 2 variables

- θ^*
- ϕ^*

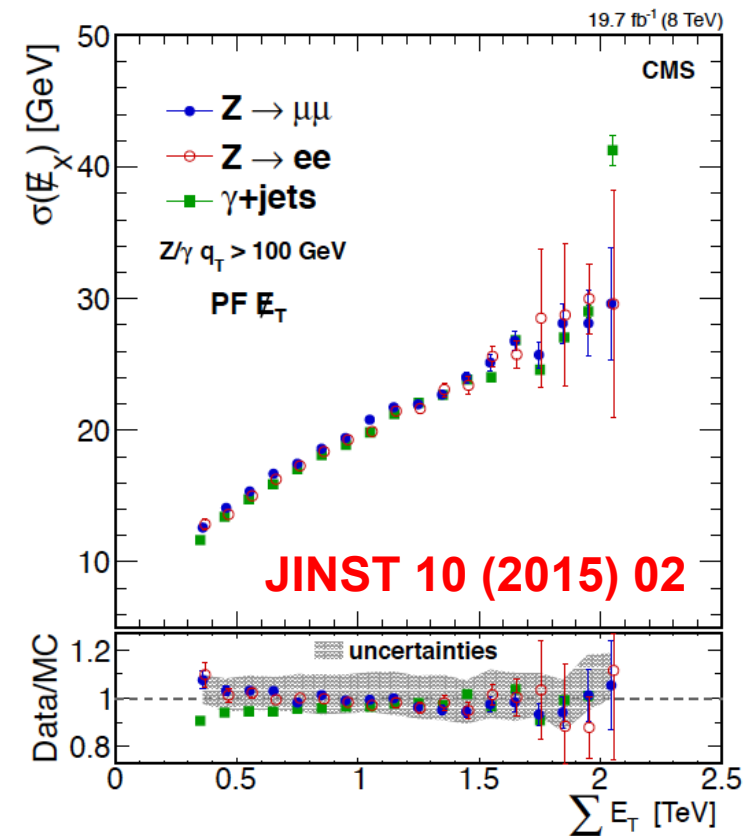
→ 4 unknown variables in $\tau\tau \rightarrow \tau_h \tau_h$, 5 in $\tau\tau \rightarrow \ell \tau_h$, 6 in $\tau\tau \rightarrow \ell\ell$ decays

The Problem

Unknown variables constrained by 2 observables only:

- $\sum p_x^v = E_x^{\text{miss}}$
- $\sum p_y^v = E_y^{\text{miss}}$

These are not even “hard” constraints, as E_T^{miss} is reconstructed with a resolution of 10-20 GeV



➔ Problem of reconstructing $m_{\tau\tau}$ is underconstrained

The Solution

Compute $m_{\tau\tau}$ using a likelihood approach:

$$\mathcal{L}(m_{\tau\tau}^{(i)}) = \int_{\Omega} d\vec{x} f(\vec{x}, \vec{y}) \delta \left(m_{\tau\tau}^{(i)} - m(\vec{x}, \vec{y}) \right)$$

x: unknown variables θ^*, ϕ^*, m_{vv}

y: measured observables $E_x^{\text{miss}}, E_y^{\text{miss}}$,
momenta of “visible” τ decay products (electrons, muons and τ_h)

The likelihood function \mathcal{L} is computed for a series of test mass hypotheses $m_{\tau\tau}^{(i)}$

The value of $m_{\tau\tau}^{(i)}$ that yields the maximal is taken as the best estimate $m_{\tau\tau}$
for the Higgs boson mass in a given event

The integral over dx has the following interpretation:

The value of $m_{\tau\tau}$ is computed by taking an “average” over all possible values of the unknown variables x that are compatible with the measured values y

The integral is computed numerically, using the VEGAS algorithm

The Matrix Element Method

Theory motivates a specific choice for the function f :

$$P(\mathbf{y}|a) = \frac{1}{\sigma(a)} \int \frac{f(x_a)f(x_b)}{2x_ax_bs} (2\pi)^4 \delta^4(x_a\mathbf{p}_a + x_b\mathbf{p}_b - \sum_i^n \mathbf{p}^{(i)})$$

$$|\mathcal{M}(\mathbf{p}, a)|^2 W(\mathbf{y}|\mathbf{p}) dx_a dx_b d\mathbf{p}$$

PDF x flux factor Energy and momentum conservation
Matrix element Transfer function

\mathbf{y} : Observables measured in the detector ($E_x^{\text{miss}}, E_y^{\text{miss}}$, momenta of e, μ, τ_h)

\mathbf{p} : “true” momenta of all particles in the final state

x_a, x_b : Bjorken scaling variables

a : unknown model parameter (here: “true” value of $m_{\tau\tau}$)

N.B.: Normalization factor $1/\sigma$ ensures that $P(\mathbf{y}|a)$ is a probability density,

i.e. $\int P(\mathbf{y}|a) d\mathbf{y} = 1 .$

Matrix Elements

Gluon fusion ME from literature/Madgraph

ME for τ decays

$$|\mathcal{M}(\mathbf{p}, a)|^2 = |\mathcal{M}_{\text{gg} \rightarrow \text{H} \rightarrow \tau\tau}(\mathbf{p}, a)|^2 \cdot |\mathcal{M}_{\tau}^{(1)}(\mathbf{p})|^2 \cdot |\mathcal{M}_{\tau}^{(2)}(\mathbf{p})|^2$$

Use narrow-width approximation for τ decays:

$$\rightarrow |\mathcal{M}_{\tau}^{(i)}|^2 = |\text{BW}_{\tau}|^2 \cdot |\mathcal{M}_{\text{decay}}^{(i)}|^2$$

$$\text{with } |\text{BW}_{\tau}|^2 = \frac{\pi}{m_{\tau} \Gamma_{\tau}} \delta(q_{\tau}^2 - m_{\tau}^2)$$

For $|\mathcal{M}_{\text{decay}}^{(i)}|^2$, we use ME from literature for leptonic τ decays:

$$|\mathcal{M}_{\tau \rightarrow \ell \bar{\nu}_{\ell} \nu_{\tau}}|^2 = 128 G_F^2 (\mathbf{p}_{\tau} \cdot \mathbf{p}_{\bar{\nu}_{\ell}}) (\mathbf{p}_{\ell} \cdot \mathbf{p}_{\nu_{\tau}})$$

Matrix Elements (cont' d)

For hadronic τ decays, we use the simplified model, assuming a two-body decay into τ_h and ν_τ and further assume the ME to be a constant:

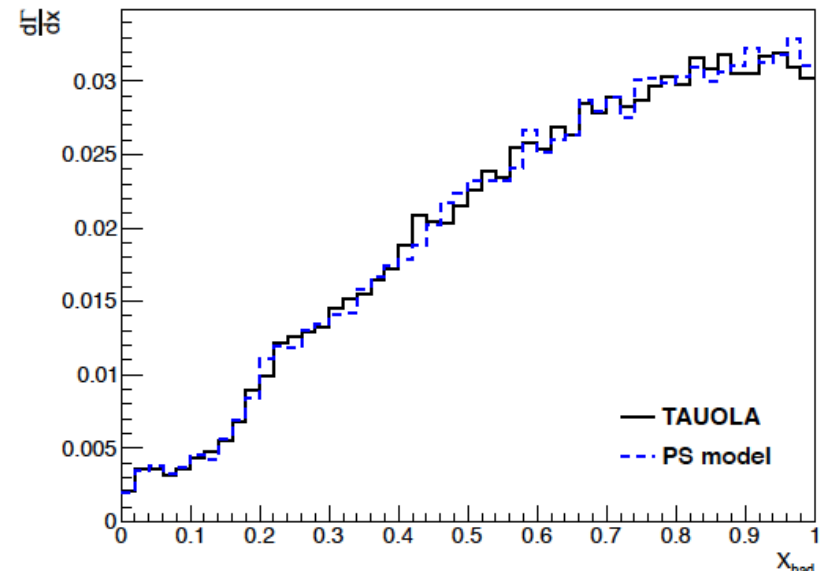
$$|\mathcal{M}_{\tau \rightarrow \tau_h \nu_\tau}^{\text{eff}}|^2 = \frac{16\pi m_\tau^3}{m_\tau^2 - m_{vis}^2} \cdot \frac{\hbar}{\Delta t} \cdot \mathcal{B}(\tau \rightarrow \tau_h \nu_\tau)$$

Factor chosen such that $\mathcal{B}(\tau \rightarrow \tau_h \nu) = 64.8\%$ is reproduced

Δt : τ lepton lifetime = $290 \cdot 10^{-15} \text{ s}$

→ Simplified model compares well with sum of all hadronic τ decay modes simulated by τ decay library TAUOLA

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Transfer Functions (TF)

The transfer functions $W(\mathbf{y}|\mathbf{p})$ represent the probability density for measuring observables \mathbf{y} in the detector, given that the “true” values of the momenta of all particles in the final state is \mathbf{p}

TF for visible τ Decay Products

For now, we assume that p_T , η and ϕ of electrons, muons and τ_h are measured with infinite precision

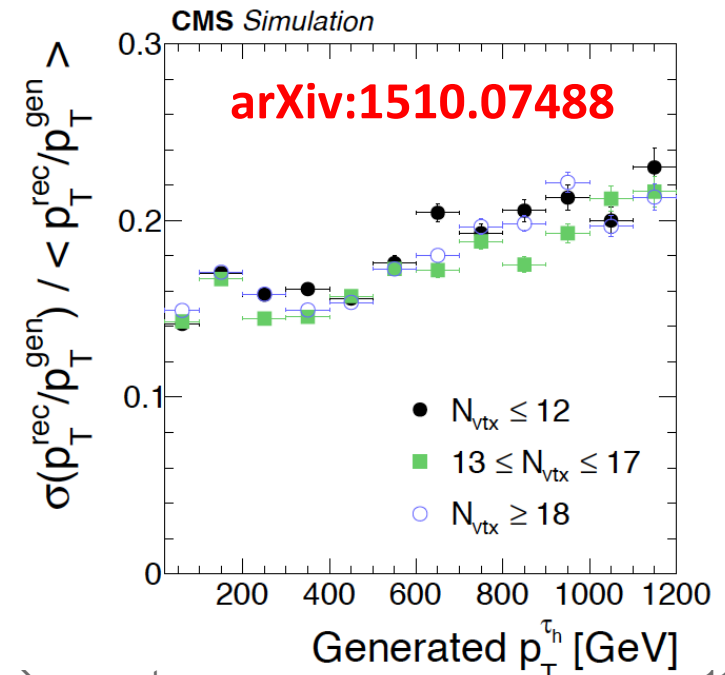
➔ The TF factorizes into a product of δ -functions:

$$W_\ell(p_T^{\text{rec}}, \eta^{\text{rec}}, \phi^{\text{rec}} | p_T^{\text{true}}, \eta^{\text{true}}, \phi^{\text{true}}) \\ = \frac{\sin^2 \theta^{\text{rec}}}{p_T^{\text{rec}}} \delta(p_T^{\text{rec}} - p_T^{\text{true}}) \delta(\theta^{\text{rec}} - \theta^{\text{true}}) \delta(\phi^{\text{rec}} - \phi^{\text{true}})$$

Experimental resolution on p_T of τ_h
actually not negligible

➔ Work ongoing to model resolution for τ_h
via transfer functions

Expected to improve $m_{\tau\tau}$ resolution by $O(10\%)$



TF for E_T^{miss}

$$W_\nu(E_x^{\text{miss}}, E_y^{\text{miss}} | \hat{E}_x^{\text{miss}}, \hat{E}_y^{\text{miss}}) = \frac{1}{2\pi \sqrt{|V|}} \exp \left(-\frac{1}{2} \begin{pmatrix} \Delta E_x^{\text{miss}} \\ \Delta E_y^{\text{miss}} \end{pmatrix}^T \cdot V^{-1} \cdot \begin{pmatrix} \Delta E_x^{\text{miss}} \\ \Delta E_y^{\text{miss}} \end{pmatrix} \right)$$

with:

$$\Delta E_x^{\text{miss}} = E_x^{\text{miss}} - \hat{E}_x^{\text{miss}}$$

$$\Delta E_y^{\text{miss}} = E_y^{\text{miss}} - \hat{E}_y^{\text{miss}}$$

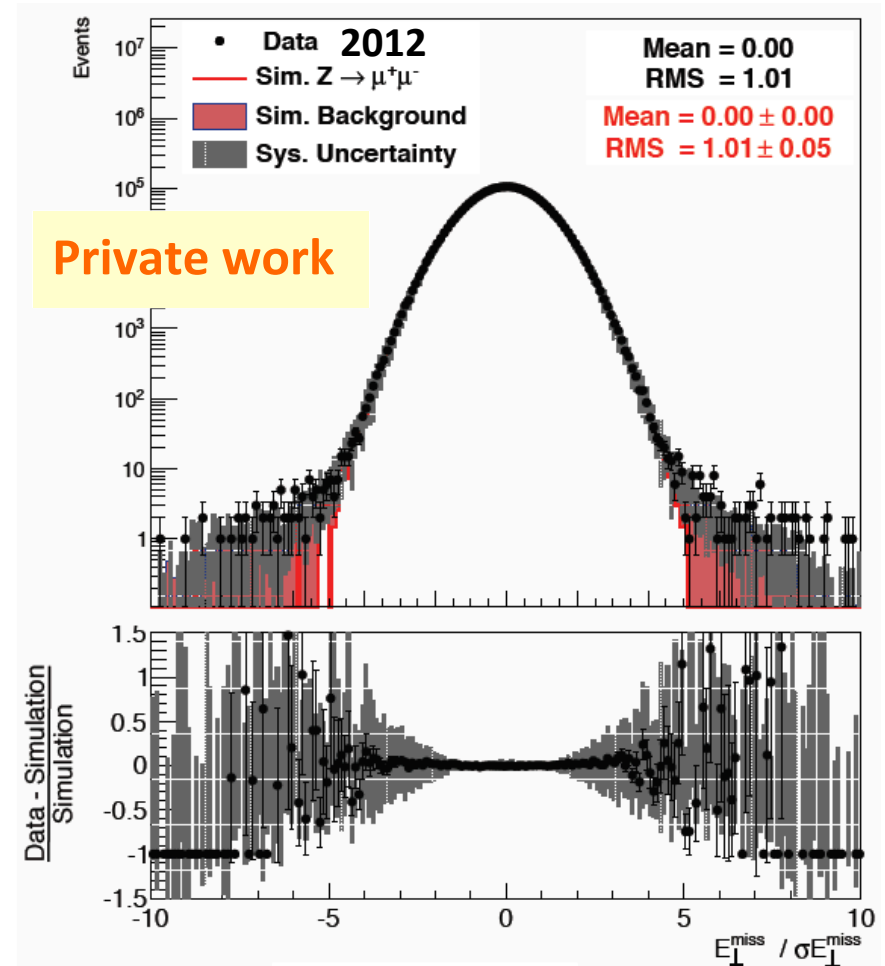
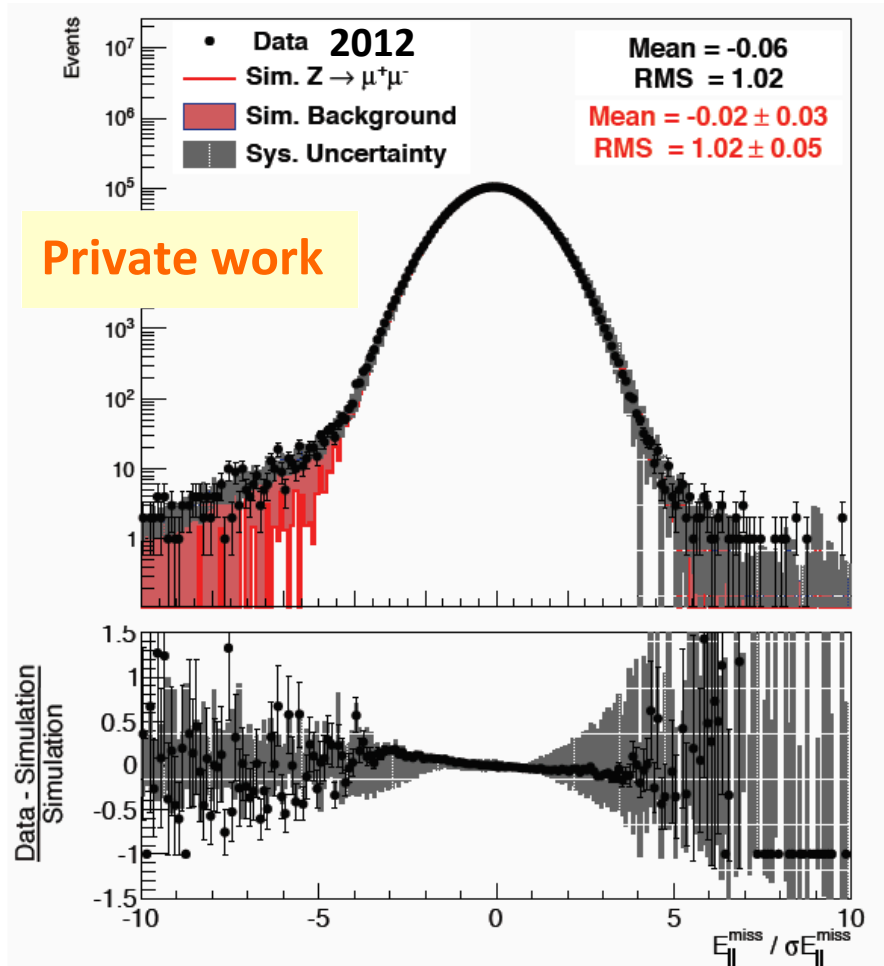
$E_x^{\text{miss}}, E_y^{\text{miss}}$: x and y components of reconstructed E_T^{miss}

$\hat{E}_x^{\text{miss}}, \hat{E}_y^{\text{miss}}$: x and y components of the sum of momenta of all neutrinos produced in the τ decays

E_T^{miss} resolution is given by covariance matrix $V = \begin{pmatrix} \sigma_x^2 & \rho \sigma_x \sigma_y \\ \rho \sigma_x \sigma_y & \sigma_y^2 \end{pmatrix}$

computed on event-by-event basis using resolution functions
obtained from the MC simulation **JME-10-009**

Validation of TF for E_T^{miss}



Distribution of pulls approximately Gaussian

Data and MC simulation agree within systematic uncertainties

Computation of Phase-space Integral

$$d\mathbf{p} = \begin{cases} d\Phi_{\tau_h \nu_\tau}^{(1)} d\Phi_{\tau_h \nu_\tau}^{(2)} & \text{if } \tau\tau \rightarrow \tau_h \nu_\tau \tau_h \nu_\tau, & (12\text{-dim}) \\ d\Phi_{\ell \bar{\nu}_\ell \nu_\tau}^{(1)} d\Phi_{\tau_h \nu_\tau}^{(2)} & \text{if } \tau\tau \rightarrow \tau_h \nu_\tau \ell \bar{\nu}_\ell \nu_\tau, & (15\text{-dim}) \\ d\Phi_{\ell \bar{\nu}_\ell \nu_\tau}^{(1)} d\Phi_{\ell \bar{\nu}_\ell \nu_\tau}^{(2)} & \text{if } \tau\tau \rightarrow \ell \bar{\nu}_\ell \nu_\tau \ell \bar{\nu}_\ell \nu_\tau, & (18\text{-dim}) \end{cases}$$

where

$$\begin{aligned} d\Phi_{\tau_h \nu_\tau}^{(i)} &= d^3\mathbf{p}_{vis}^{(i)} d^3\mathbf{p}_\nu^{(i)} \\ d\Phi_{\ell \bar{\nu}_\ell \nu_\tau}^{(i)} &= d^3\mathbf{p}_{vis}^{(i)} d^3\mathbf{p}_\nu^{(i)} d^3\mathbf{p}_{\bar{\nu}}^{(i)}. \end{aligned}$$

Product of phase-space element and τ decay ME simplified analytically before numeric integration:

$$|\mathcal{M}_\tau^{(i)}|^2 d\Phi_{\tau_h \nu_\tau}^{(i)} = \frac{\pi}{m_\tau \Gamma_\tau} f_h \left(\mathbf{p}^{\text{vis}(i)}, m^{\text{vis}(i)}, \mathbf{p}^{\text{inv}(i)} \right) \frac{d^3\mathbf{p}^{\text{vis}}}{2E_{\text{vis}}} dz d\phi_{\text{inv}}$$

$$|\mathcal{M}_\tau^{(i)}|^2 d\Phi_{\ell \bar{\nu}_\ell \nu_\tau}^{(i)} = \frac{\pi}{m_\tau \Gamma_\tau} f_\ell \left(\mathbf{p}^{\text{vis}(i)}, m^{\text{vis}(i)}, \mathbf{p}^{\text{inv}(i)} \right) \frac{d^3\mathbf{p}^{\text{vis}}}{2E_{\text{vis}}} dz dm_{\text{inv}}^2 d\phi_{\text{inv}}$$

Phase-space Integration (cont' d)

$$f_h(\mathbf{p}^{\text{vis}}, m_{vis}, \mathbf{p}^{\text{inv}}) = \frac{|\mathcal{M}_{\tau \rightarrow \tau_h \nu_\tau}^{\text{eff}}|^2}{512\pi^6 |\mathbf{p}^{\text{vis}}| z^2}$$

$$f_\ell(\mathbf{p}^{\text{vis}}, m_{vis}, \mathbf{p}^{\text{inv}}) = \frac{|\mathcal{M}_{\tau \rightarrow \ell \bar{\nu}_\ell \nu_\tau}^{\text{eff}}|^2}{512\pi^6} \cdot \frac{I_{inv}}{|\mathbf{p}^{\text{vis}}| z^2}$$

with

$$I_{inv} = \frac{G_F^2}{\pi^2} m_{inv}^2 \left(2 E_\tau E_\ell - \frac{2}{3} \sqrt{E_\tau^2 - m_\tau^2} \sqrt{E_\ell^2 - m_\ell^2} \right)$$

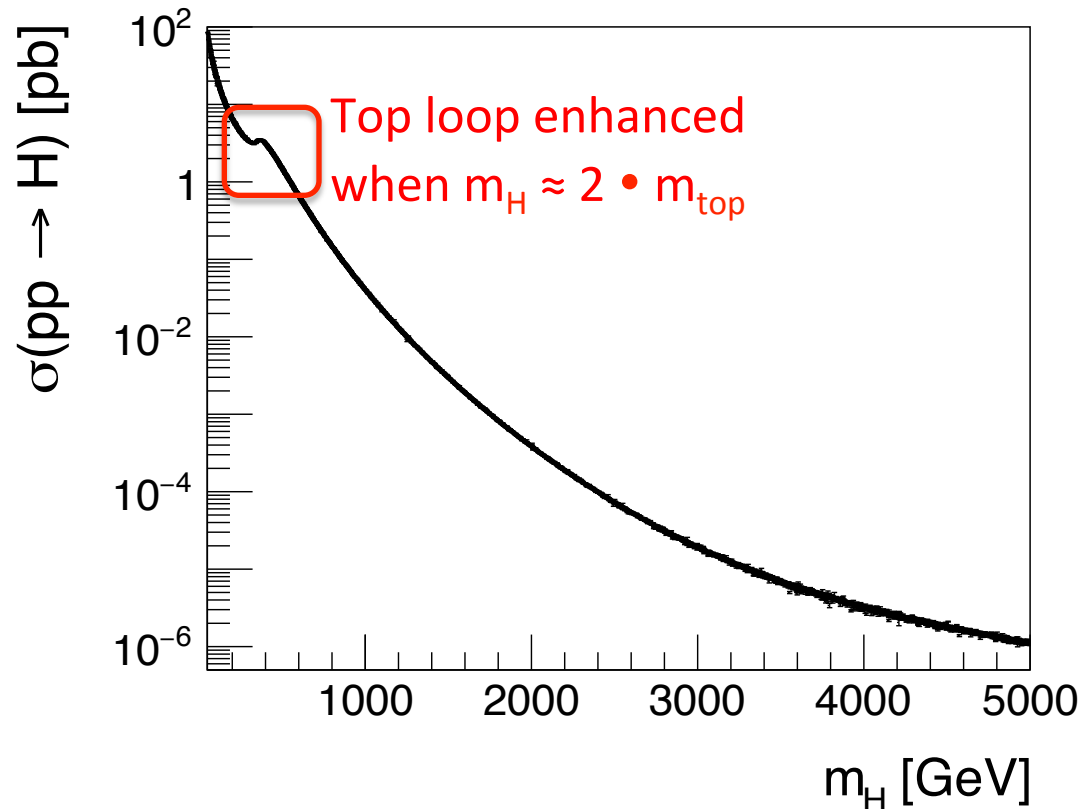
The τ lepton momenta can be computed as function of \mathbf{p}^{vis} and the integration variables z and ϕ_{inv} respectively z , ϕ_{inv} and m_{inv}^2

This then allows to evaluate the squared modulus of the ME

$|\mathcal{M}_{pp \rightarrow H \rightarrow \tau\tau}(\mathbf{p}, m_H)|^2$ and the PDF in the same way as if the τ leptons were stable particles

Normalization

Cross-section $\sigma = 1/\text{Normalization factor}$ is steeply falling as function of $m_{\tau\tau}^{\text{test}} = m_H$



- ➔ Normalization factor $1/\sigma$ increases probability to reconstruct events in high mass tail compared to “standard” SVfit
- ➔ Add artificial “penalty” term of form $k \cdot \text{Log}(m_{\tau\tau})$ to reduce high mass tail

Artificial Regularization Term

Rather than obtaining $m_{\tau\tau}$ via maximization of $P(\mathbf{y} | m_{\tau\tau}^{(i)})$,
we may choose to maximize

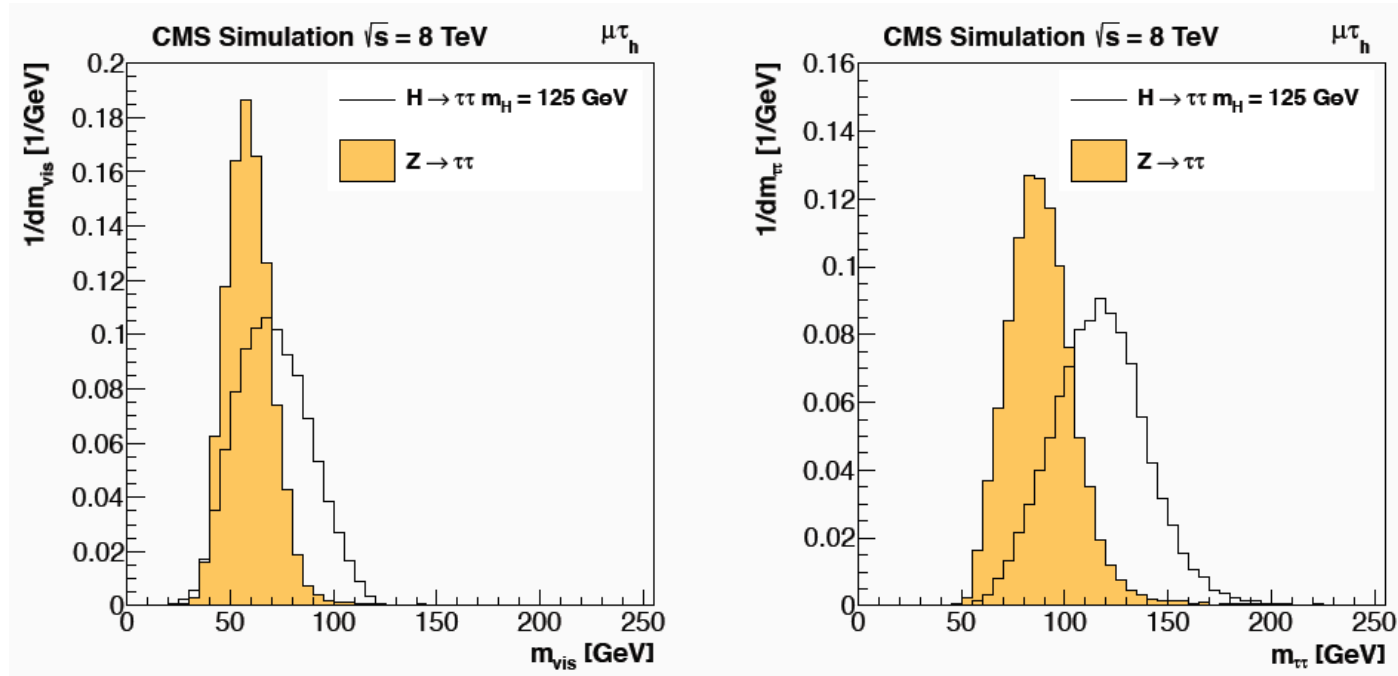
$$\text{Log}(P(\mathbf{y} | m_{\tau\tau}^{(i)})) - k \cdot \text{Log}(m_{\tau\tau}^{(i)})$$

instead.

The parameter k is found by optimizing the $m_{\tau\tau}$ resolution for a sample of events:
We find that a value $k = 5$ significantly reduces the high $m_{\tau\tau}$ tail for the irreducible
 $Z/\gamma^* \rightarrow$ background, while causing a small bias on the $m_{\tau\tau}$ distribution
reconstructed in Higgs $\rightarrow \tau\tau$ signal events

N.B.: Adding the artificial term $k \cdot \text{Log}(m_{\tau\tau})$ is known as penalized likelihood
method in the literature

Mass Resolution



J.Phys.Conf.Ser. 513 (2014) 022035

SVfit algorithm achieves resolution of typically 15-20% relative to true value of $m_{\tau\tau}$

This resolution significantly improves separation between Higgs $\rightarrow \tau\tau$ signal and irreducible $Z/\gamma^* \rightarrow \tau\tau$ background compared to alternative mass observables

Reducible backgrounds (W+jets, QCD, $t\bar{t}$) approximately flat in $m_{\tau\tau}$

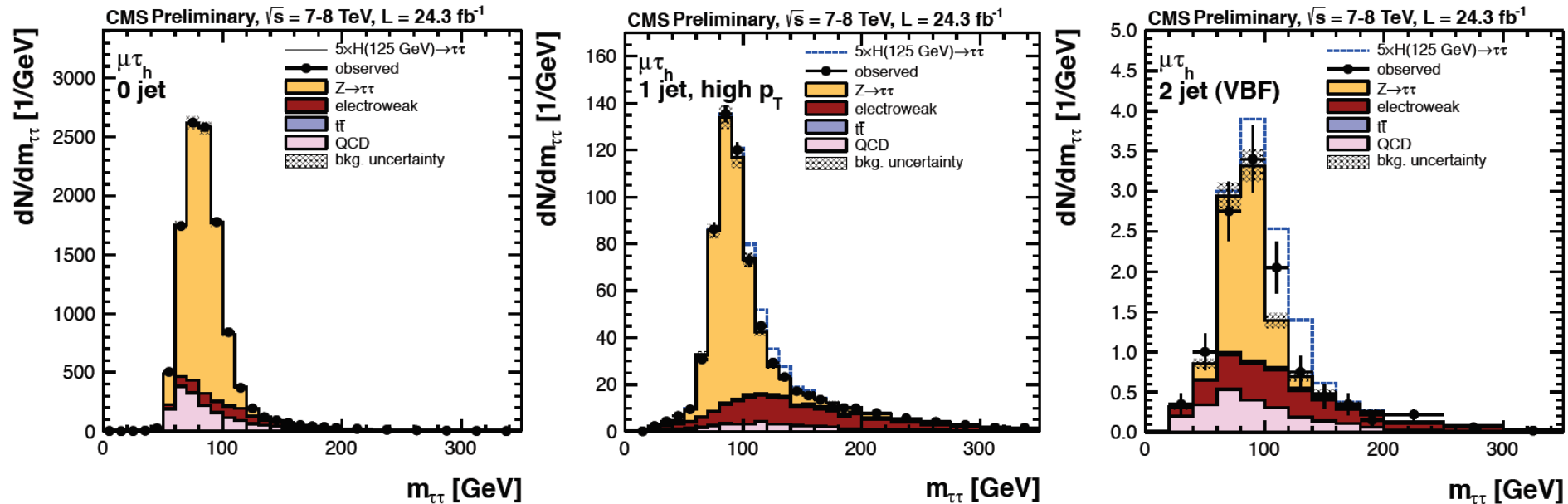
➔ Effective background contribution decreases

if Higgs signal is concentrated in narrow mass window

Performance in CMS Higgs $\rightarrow \tau\tau$ Analysis

Events analyzed in different event categories:

JHEP 1405 (2014) 104



Most sensitive categories are 1 jet and VBF, in which Higgs is typically boosted:

➔ SVfit resolution improves compared to 0 jet category in which visible τ decay products are typically “back-to-back” in transverse plane

Overall sensitivity of CMS SM Higgs $\rightarrow \tau\tau$ analysis improves by $\approx 40\%$ when using $m_{\tau\tau}$ reconstructed by SVfit algorithm compared to m_{vis}

ATLAS Higgs boson ML challenge

ATLAS setup a kaggle contest for improving the sensitivity of their SM Higgs $\rightarrow \tau\tau$ analysis

<https://www.kaggle.com/c/higgs-boson>

ATLAS provided Ntuples with MMC* mass, the momenta and energies of the visible tau decay products plus E_T^{miss}

- \rightarrow Information sufficient to compute SVfit mass for events in ATLAS Ntuples
- \rightarrow Compare SVfit vs. MMC mass resolution for SM Higgs $\rightarrow \tau\tau$ events using the event categories of our SM Higgs \rightarrow analysis **HIG-13-004**

Note:

- ATLAS does not (yet) estimate the E_T^{miss} resolution on an event-by-event basis
 - \rightarrow I estimated ATLAS' E_T^{miss} resolution based on public performance numbers
- The ATLAS τ_h reconstruction assumes the τ_h to be massless
 - \rightarrow I extended SVfit to take an average over the distribution of true τ_h mass

*Missing Mass Calculator **arXiv:1012.4686**



Completed • \$13,000 • 1,785 teams

Higgs Boson Machine Learning Challenge

Mon 12 May 2014 – Mon 15 Sep 2014 (4 months ago)

Dashboard

Private Leaderboard - Higgs Boson Machine Learning Challenge

This competition has completed. This leaderboard reflects the final standings.

See someone using multiple accounts?

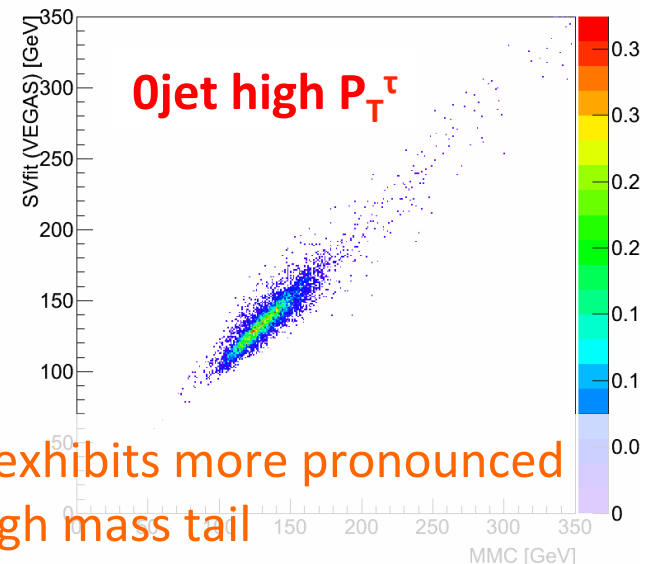
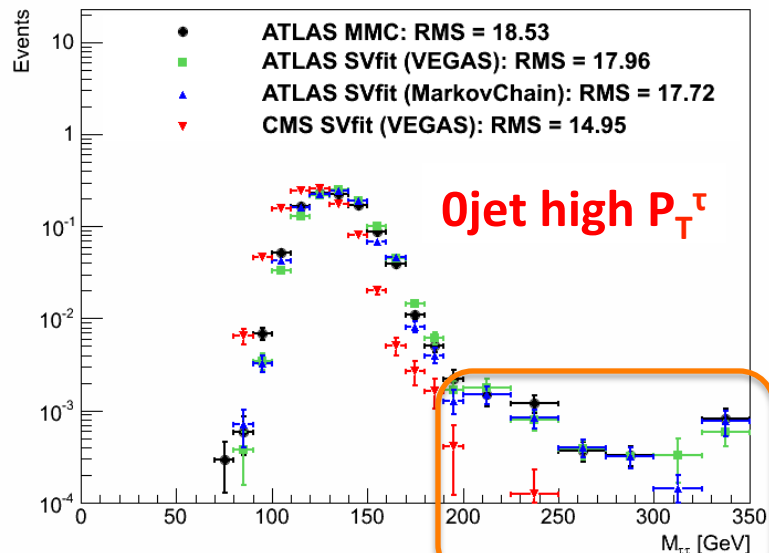
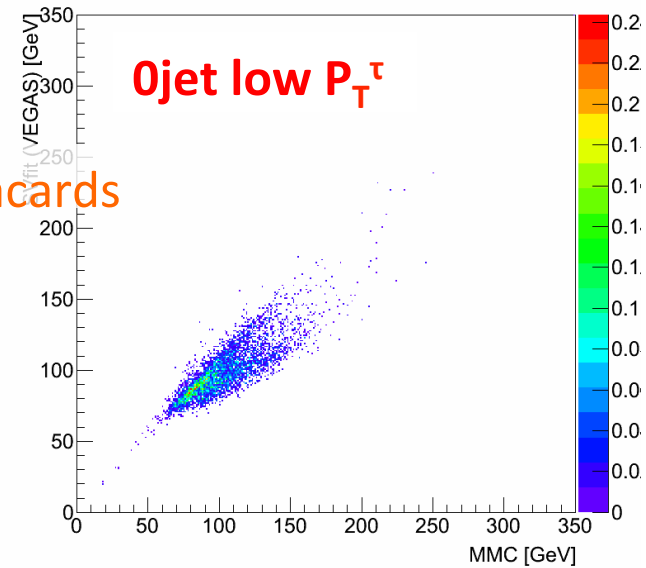
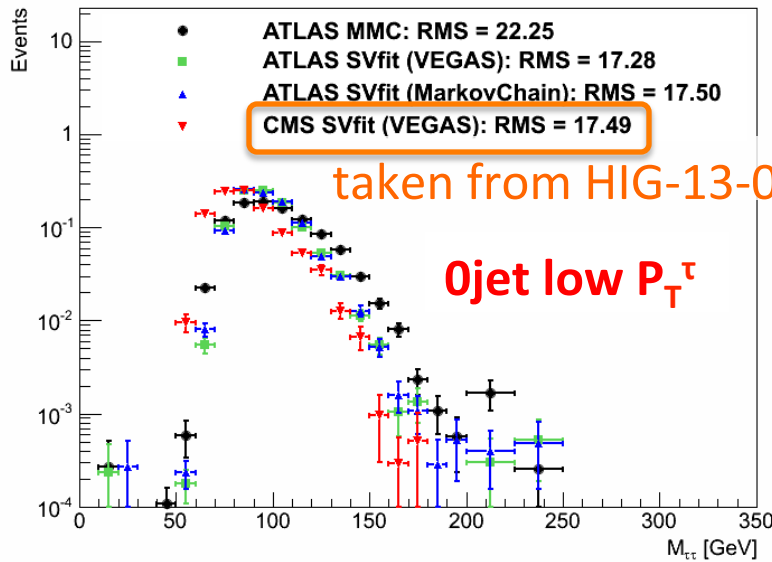
[Let us know.](#)

#	Δrank	Team Name <small>‡ model uploaded * in the money</small>	Score <small>?</small>	Entries	Last Submission UTC (Best – Last Submission)
1	↑1	Gábor Melis ‡ *	3.80581	110	Sun, 14 Sep 2014 09:10:04 (-0h)
2	↑1	Tim Salimans ‡ *	3.78913	57	Mon, 15 Sep 2014 23:49:02 (-40.6d)
3	↑1	nhlx5haze ‡ *	3.78682	254	Mon, 15 Sep 2014 16:50:01 (-76.3d)
4	↑38	ChoKo Team <small>👤</small>	3.77526	216	Mon, 15 Sep 2014 15:21:36 (-42.1h)
5	↑35	cheng chen	3.77384	21	Mon, 15 Sep 2014 23:29:29 (-0h)
6	↑16	quantify	3.77086	8	Mon, 15 Sep 2014 16:12:48 (-7.3h)
7	↑1	Stanislav Semenov & Co (HSE Yandex)	3.76211	68	Mon, 15 Sep 2014 20:19:03
8	↓7	Luboš Motl's team <small>👤</small> <ul style="list-style-type: none">• Luboš Motl• Christian Veelken	3.76050	589	Mon, 15 Sep 2014 08:38:49 (-1.6h)
9	↑8	Roberto-UCIIM	3.75864	292	Mon, 15 Sep 2014 23:44:42 (-44d)
10	↑2	Davut & Josef <small>👤</small>	3.75838	161	Mon, 15 Sep 2014 23:24:32 (-4.5d)

➔ SVfit finished 8th place out of 1785 teams!

Comparison SVfit vs. MMC

(on $gg \rightarrow \text{Higgs} + \text{VBF Higgs signal events}$)



MMC exhibits more pronounced
high mass tail

Conclusions: SVfit vs. MMC

- Difference between MMC and SVfit is small in boosted events, in which $m_{\tau\tau}$ is well constrained by E_T^{miss}
- MMC likelihood model seems not optimal for 0jet events, in which both taus are typically “back-to-back”

Extensions of SVfit Algorithm

Combination of likelihood approach and matrix element method is very universal, can be used for estimating any quantity of interest

Recent developments:

- Estimation of Higgs boson p_T , η and ϕ in addition to $m_{\tau\tau}$
- Reconstruction of Higgs boson mass in search for lepton flavor violating (LFV) Higgs $\rightarrow \mu\tau \rightarrow \mu\tau_h$ and Higgs $\rightarrow e\tau \rightarrow e\tau_h$ decays
- Reconstruction of di-Higgs mass m_{hh} in search for $X \rightarrow hh \rightarrow bb\tau\tau$ decays of high mass resonances X
- Reconstruction of “transverse” Higgs mass

$$m_{T\tau\tau} = \sqrt{(E_T^{\tau_1} + E_T^{\tau_2})^2 - ((p_x^{\tau_1} + p_x^{\tau_2})^2 + (p_y^{\tau_1} + p_y^{\tau_2})^2)}$$

in CMS MSSM Higgs $\rightarrow \tau\tau$ analysis

➔ In most cases SVfit algorithm yields a significant improvement in analysis sensitivity

Summary

- SVfit algorithm allows to reconstruct Higgs boson mass in Higgs $\rightarrow \tau\tau$ events with a resolution of typically 15-20% relative to the true value of $m_{\tau\tau}$
- This resolution significantly improves the separation of the Higgs $\rightarrow \tau\tau$ signal from the irreducible $Z/\gamma^* \rightarrow \tau\tau$ background, improving the sensitivity of CMS SM Higgs $\rightarrow \tau\tau$ analysis by $\approx 40\%$
- The approach is very universal and has been customized for applications to LFV Higgs decays, the search for $X \rightarrow hh$ decays and for the CMS MSSM Higgs $\rightarrow \tau\tau$ analysis
- I expect the formalism to handle τ lepton decays in the ME method, developed for the SVfit algorithm, to be useful for future applications of the ME method to analyses with τ leptons in the final state

 **Backup** 

