

Integrability & Scattering Amplitudes as a Key to Gauge Theory

Till Bargheer



DESY Fellows Meeting, Nov 10, 2015

Collaborators: Niklas Beisert, Wellington Galleas, Nikolay Gromov, Song He, Yu-tin Huang, Florian Loebbert, Tristan McLoughlin, Carlo Meneghelli, Joe Minahan, Raul Pereira, Masahito Yamazaki, . . .

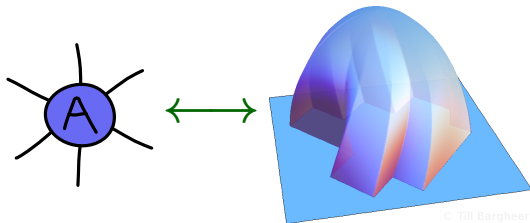
Motivation

Understand features of gauge theory

- At weak coupling
- At strong coupling
- Non-perturbatively

Two attack vectors:

- AdS/CFT and integrability
- Structures in scattering amplitudes

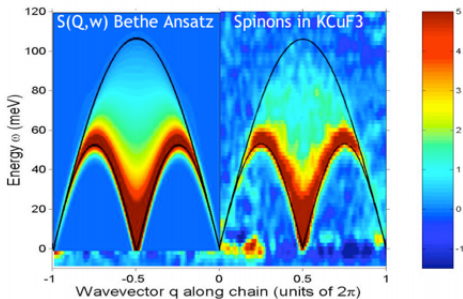


Integrability

A miraculous property of certain models *in two dimensions*

Occurs for very specific systems, but sometimes really models reality:

Results for KCuF3



Integrability in a Nutshell

In two dimensions:

If $Q_1 = \sum_j p_j$, $Q_2 = \sum_j p_j^2$ are conserved

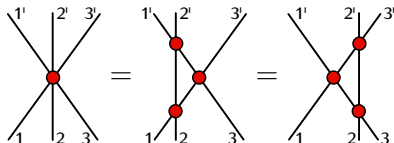
\Rightarrow two momenta conserved: $\{p_1, p_2\} = \{p'_1, p'_2\}$

Further conserved charge:

$Q_3 = \sum_j p_j^3 \Rightarrow \{p_1, p_2, p_3\} = \{p'_1, p'_2, p'_3\}$

Higher conserved charges \Rightarrow factorized scattering, $\{p_i\} = \{p'_i\}$

S-matrix obeys **Yang–Baxter equation**:



\Rightarrow In two dimensions, an interacting theory can be **integrable** (solvable)

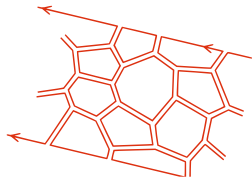
Gauge Theory

Simplest gauge theory: $\mathcal{N} = 4$ super Yang–Mills theory in $d = 4$

- Cancellations (susy), conformal symmetry, finiteness
- AdS/CFT \longrightarrow string theory at strong coupling
- Planar limit: Integrability \longrightarrow powerful techniques. Solvable?!

Why is planar $\mathcal{N} = 4$ SYM integrable?

- Planar limit: Leading piece in t'Hooft large N_c expansion
- Two-dimensionality in color space

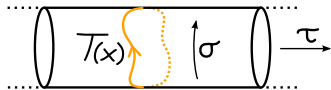


AdS/CFT: Dual to string theory in AdS

Worldsheet is two-dimensional

Family of flat connections

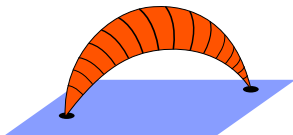
Higher charges: Holonomy $T(x) = \sum_j x^j Q_j$



\longrightarrow Integrable charges form an infinite-dimensional **Yangian algebra**

$\mathcal{N} = 4$ SYM: The Spectrum

AdS/CFT picture
for two-point functions:



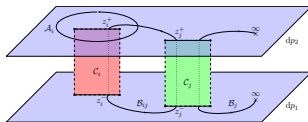
Spin chain of
single-trace operators:

$$\text{Tr}(ZZ\phi Z \dots Z\phi Z\phi) \sim$$
A diagram of a spin chain. It consists of a horizontal black line with a series of red arrows pointing up and blue arrows pointing down, alternating along the line. The line is enclosed within a black oval.

Large volume:

Weak coupling: Bethe ansatz

Strong coupling: Spectral curve



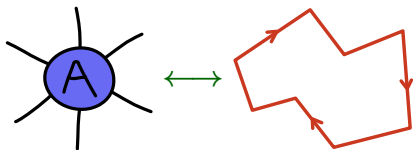
Finite-size corrections: Wrapping

Finite size, exact:

Thermodynamic Bethe Ansatz, FinLIE, Quantum spectral curve

⇒ Anomalous dimensions can now be computed (numerically)
for any operator at finite coupling

Scattering Amplitudes



Zoo of interesting structures

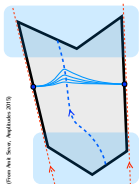
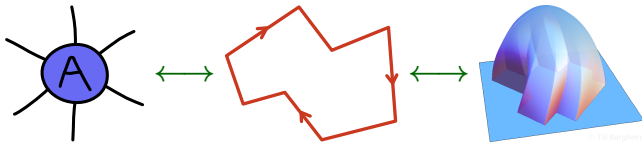
- Duality to polygon Wilson loops
- Dual superconformal symmetry
- Twistors / momentum twistors
- Grassmannian integral / on-shell recursion for trees & integrands
- Differential equations for loop integrals
- Amplitude bootstrap program

Relation to integrability only gradually uncovered

\Rightarrow **Room for fruitful interplay**

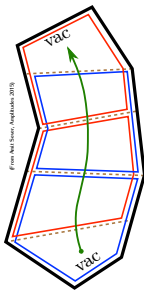
Integrability for Amplitudes

What is the map to 2d integrable system for scattering amplitudes?



Key: Look at Wilson loop as a flux tube

Decomposition: pentagon transition functions



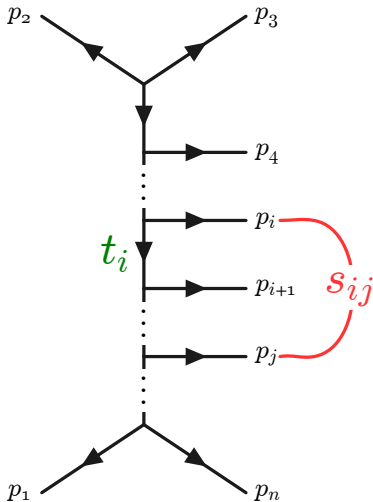
Flux tube states \leftrightarrow large spin operators

$$\mathcal{O} = \text{Tr}(ZDD \dots DD \overset{p_1}{\overrightarrow{F}} DD \dots DD \overset{p_2}{\overrightarrow{F}} DD \dots DDZ)$$

Exact spectrum $E(p_i)$ and scattering phases $S(p_1, p_2)$ from integrability

Proposal for pentagon transitions: $P \sim \sqrt{S(p, q)/S(p, q')}$

Physics: Regge Limit & BFKL



- Transverse momenta \ll rapidities
- Hierarchy of rapidities:

$$s = s_{3,n} \gg s_{3,n-1}, s_{4,n} \gg \dots \\ \dots \gg s_{i,i+2} \gg s_i \gg -t_i$$

- Large logarithms: $\log(s_i) \gg 1$
- Double expansion in $\log(s_i)$ and coupling g_{YM}
- Very generic:
Applies to general gauge theory
- Useful in practical computations

Physics: Regge Limit & BFKL

Resum perturbation theory, “Reggeization”

- Large-log approximation, but all orders in g_{YM}
- Schematically,

$$R_n \xrightarrow{\text{MRL}} \sum (\text{Regge poles}) + \sum (\text{Regge cuts})$$

- Sum over $\mathfrak{sl}(2)$ representations (ν, n) :

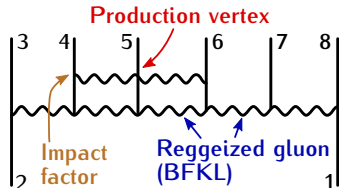
$$(\text{Regge cut}) \sim \sum_{n=-\infty}^{\infty} \int d\nu \Phi_{\nu,n}^* s_i^{\omega(\nu,n)} \Phi_{\nu,n}$$

Ingredients:

- Energies $\omega(\nu, n)$
- Impact factors $\Phi_{\nu,n}$
- Production vertices (at higher points)

Energies $\omega(\nu, n)$ are eigenvalues of

BFKL Hamiltonian of $\mathfrak{sl}(2)$ spin chain \rightarrow **integrable!**



Apply Integrability!

Goal: **Extract** BFKL data (eigenvalues, impact factors, production vertices) from newly discovered structures in amplitudes;
make use of integrability!

- Six points: understood
(analytic structure, BFKL data from Wilson loop OPE)
- A lot of structure beyond six points that needs to be understood
 - ▶ Analytic structure in different regions
 - ▶ BFKL data

Plan:

- *Step 1:* **Extract MRL data** from known amplitude data
(from bootstrap: two loops, any n ; three loops, $n = 7$)
- *Step 2:* **Apply Wilson loop OPE** to obtain all-order or even exact BFKL data beyond $n = 6$
- (*Step 3:* Data can be used for practical computations)

