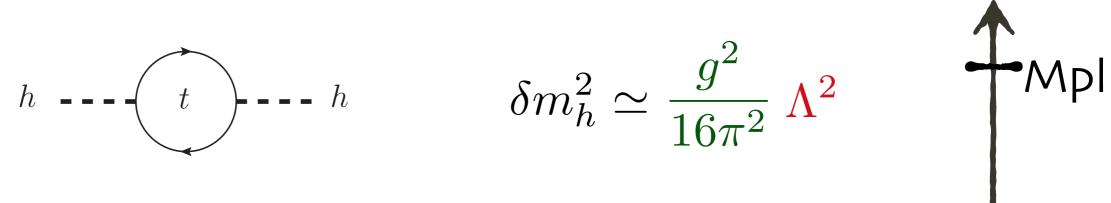
# Higgs as a Composite PNGB

Oleksii Matsedonskyi

expected ▶ Higgs mass is unstable under radiative corrections Higgs mass

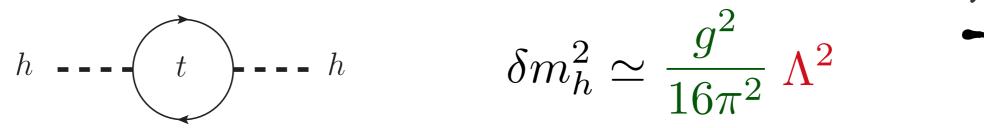


$$\delta m_h^2 \simeq \frac{g^2}{16\pi^2} \Lambda^2$$





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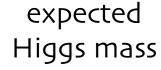


$$\Lambda \sim M_{Planck}$$

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$$\mathrm{tuning} \sim \frac{M_{Planck}^2}{m_h^2} \sim 10^{32}$$



▶ Higgs mass is unstable under radiative corrections



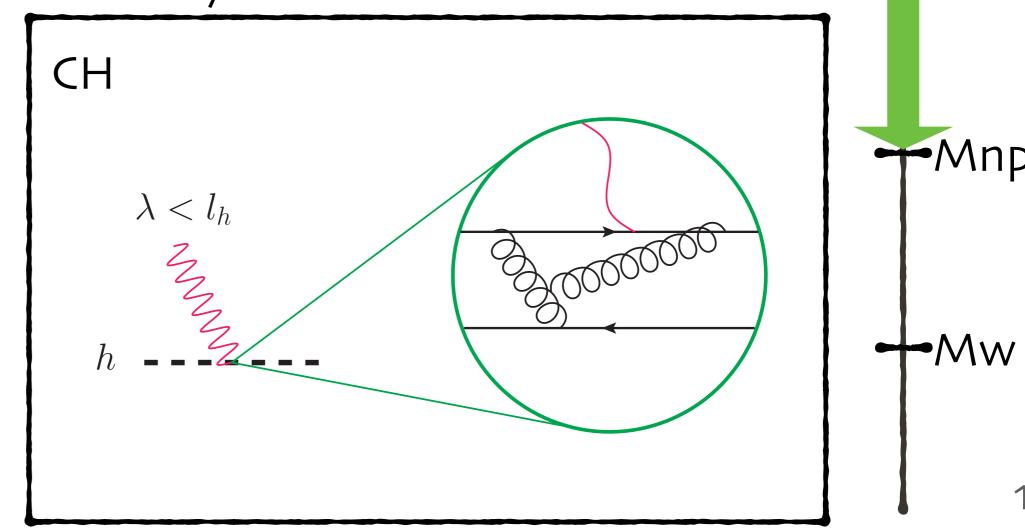


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UV can be screened by TeV scale

New Physics



▶ Higgs mass is unstable under radiative corrections



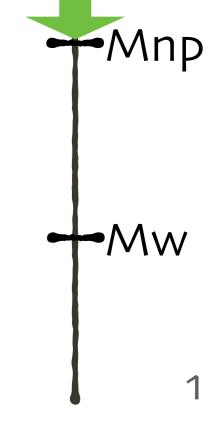


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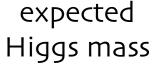


UV can be screened by TeV scale New Physics

$$\delta m_h^2 \simeq \frac{1}{l_h^2} \simeq m_{
ho}^2$$



▶ Higgs mass is unstable under radiative corrections





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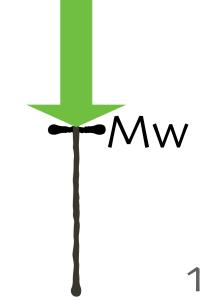
UV can be screened by TeV scale New Physics

$$\delta m_h^2 \simeq \frac{1}{l_h^2} \simeq m_{\rho}^2$$

Higgs realized as a Goldstone boson can be naturally light



$$\delta m_h^2 \simeq \frac{g_{\mathcal{G}}^2}{16\pi^2} \ m_{\rho}^2$$



### composite sector

 $SO(5) \rightarrow SO(4)$  at a scale f >> Mz

(analogous to  $SU(2)L \times SU(2)R \rightarrow SU(2)V$  in QCD)

H

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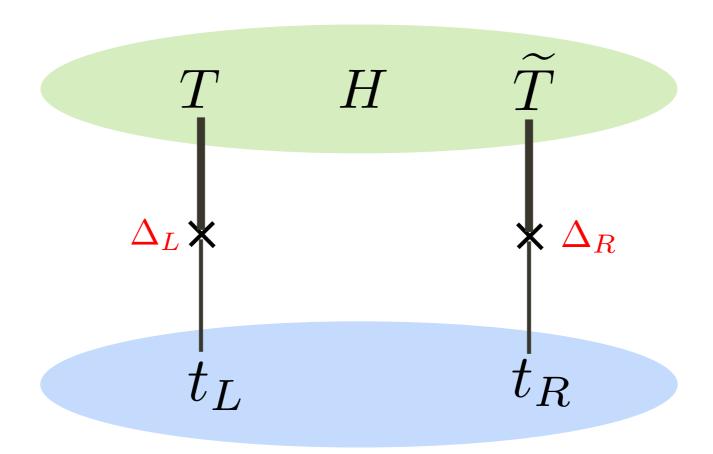
H

 $t_L$   $t_R$ 

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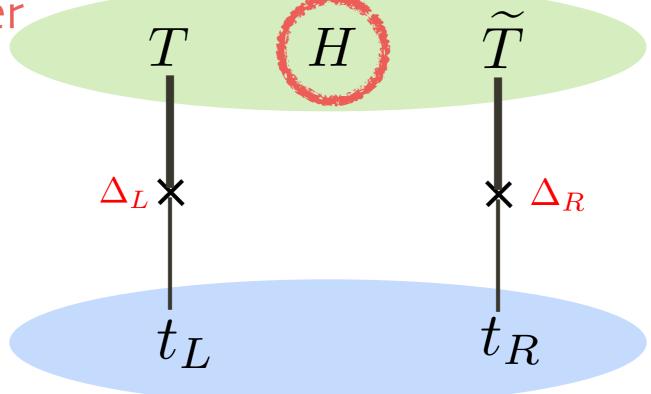
SM fermions become partially composite 
$$t_L^{``}=\cos\phi_L\,t_L+\sin\phi_L\,T_L \ \sin\phi_L\simeq rac{\Delta_L}{M_T}$$

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Higgs self-inter

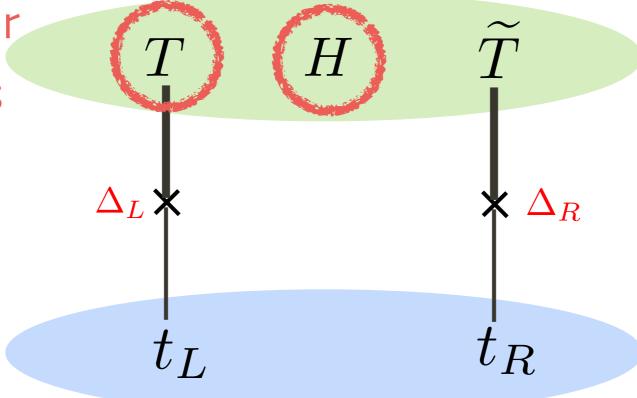


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- Higgs self-inter
- direct searches

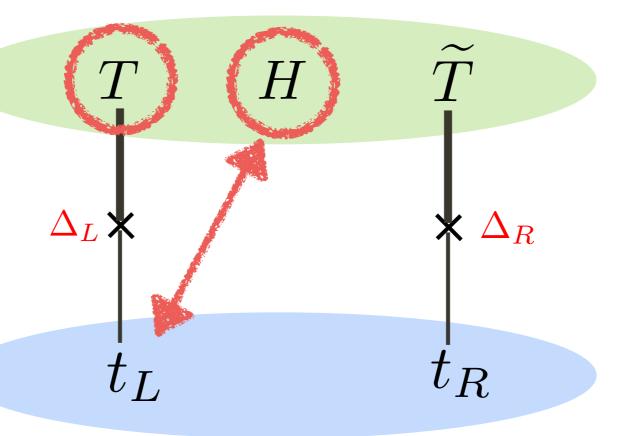


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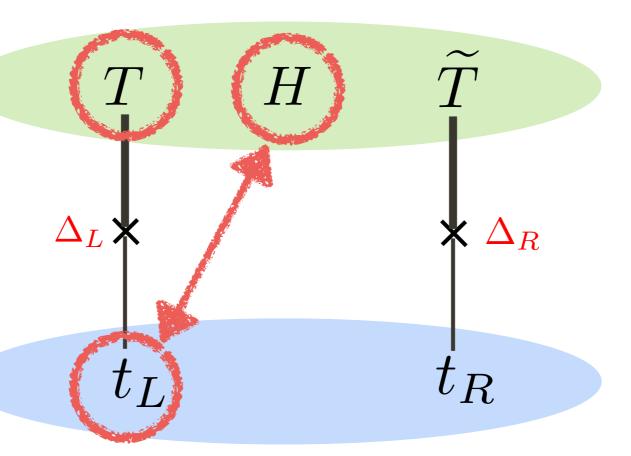


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- Higgs self-inter
- direct searches
- Higgs couplings
- **◆** EWPT
- ◆ flavour



elementary sector

SM without Higgs

## Mass Spectrum

- f sets all the mass scales of the theory
- ullet masses are proportional to f and the strength of coupling to it:

NP: 
$$m_{
ho} \sim g_{
ho} f$$

SM:  $m_{\rm SM} \sim g_{\rm SM} \, v_{\rm SM}$ 

$$v_{\rm SM} \ll f$$

extra scale separation from tuning

$$\xi = \left(\frac{v}{f}\right)^2$$

possible since Gsm can be embedded into SO(4)

currently

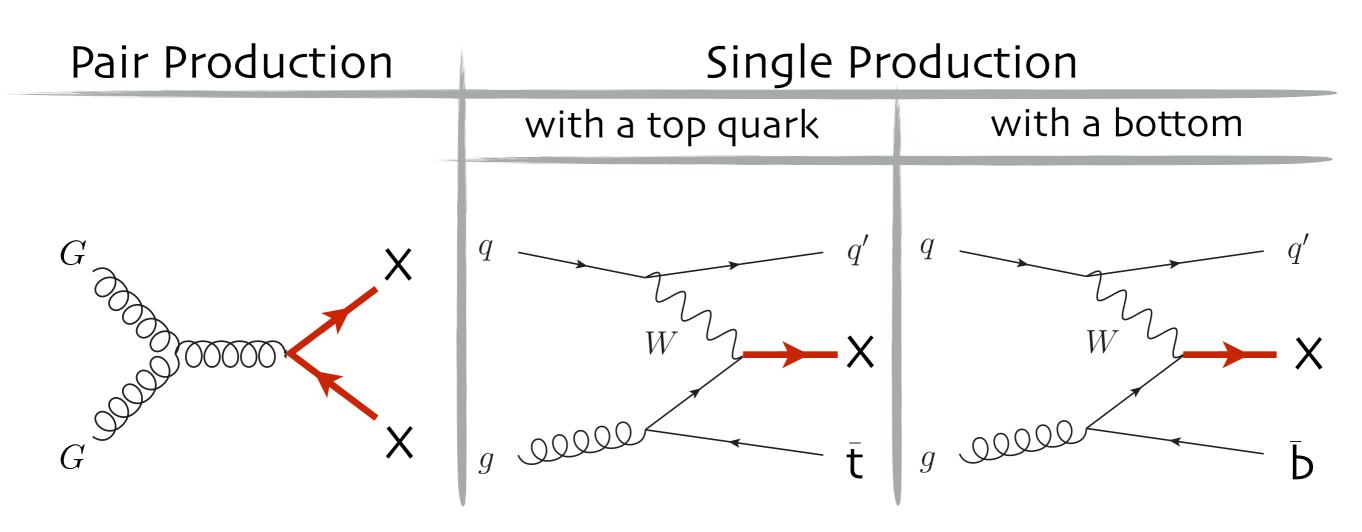
$$\xi \lesssim$$
 0.1

$$f \gtrsim$$
 800  $GeV$ 

## Direct Production

Simplest realizations of CH with a moderate tuning (~10%) require a presence of 1-1.7 TeV composite fermions

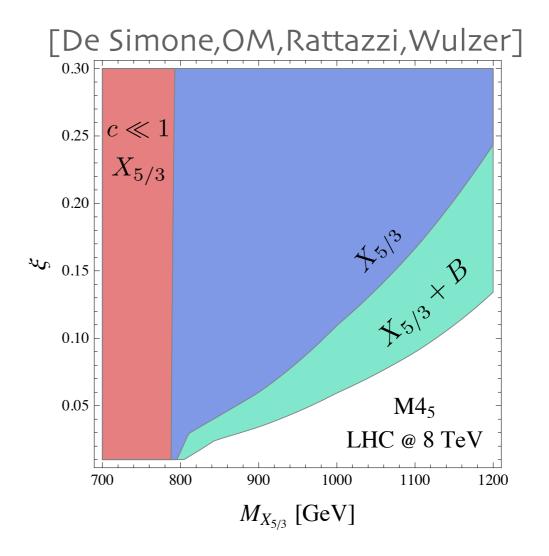
[OM,Panico,Wulzer]



[Contino, Servant; Mrazek, Wulzer]

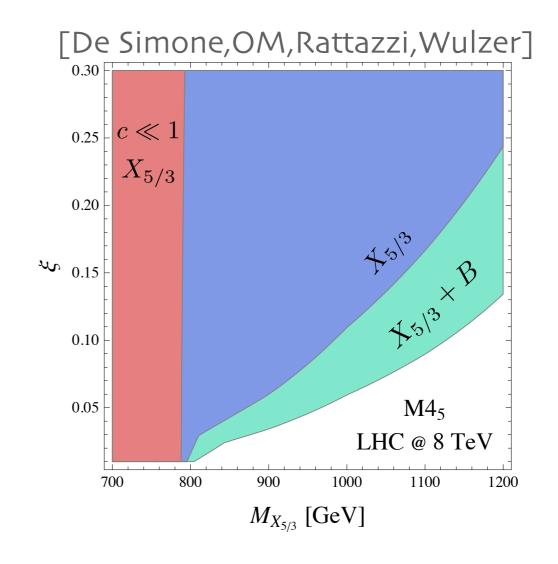
## Direct Production

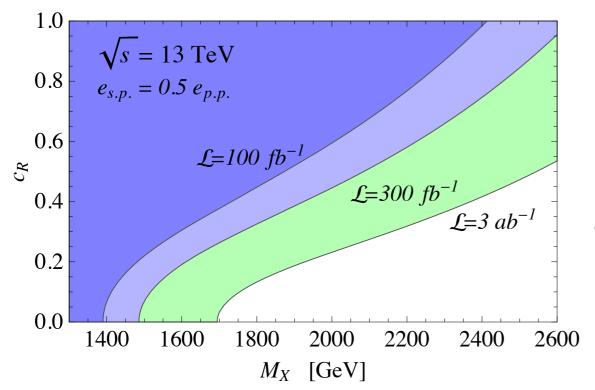
Direct searches can bound the degree of CH models naturalness



## Direct Production

Direct searches can bound the degree of CH models naturalness





Growth of single production at 13TeV requires a universal parametrization of experimental analyses output

$$\sigma_{\text{sing}}(X\overline{t}) = \left(c_R^2 + c_L^2\right)\sigma_{Wt}(m_X) + 2c_R c_L \left(\frac{m_t}{m_X + m_t}\right)\sigma'_{Wt}(m_X)$$

[OM,Panico,Wulzer]

# Higgs Couplings Deformation

Higgs as NGb generically induces non-renormalizable interactions

$$H \to f \exp i\frac{H}{f} = f + iH - \frac{H^2}{2f} - \frac{iH^3}{6f^2} + \cdots$$

• e.g. HVV coupling 
$$g^2 v_{\mathrm{SM}} h W^2 o g^2 \left(1 - \frac{1}{2} \xi + \cdots \right) v_{\mathrm{SM}} h W^2$$

- one expects NP in higher order operators  $\frac{g_{
  ho}^2\,v^2}{m_{
  ho}^2}\sim \xi$
- this is generic to strongly coupled resonances, but the Goldstone symmetry imposes additional constraints on the deviations

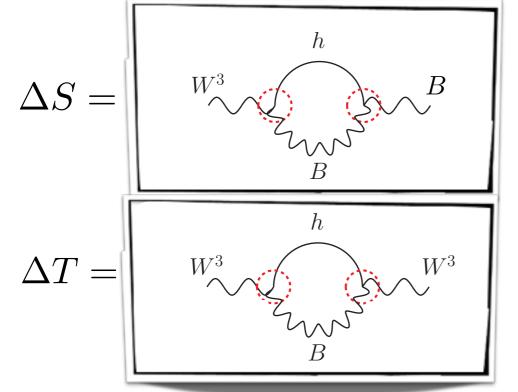
$$g_{\rm SM} \to g_{\rm SM}(1+c\,\xi)$$

## **EWPT**

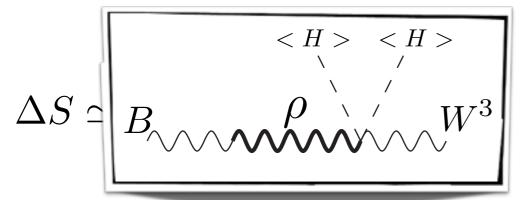
strong constraints on NP scenarios by LEP measurements:

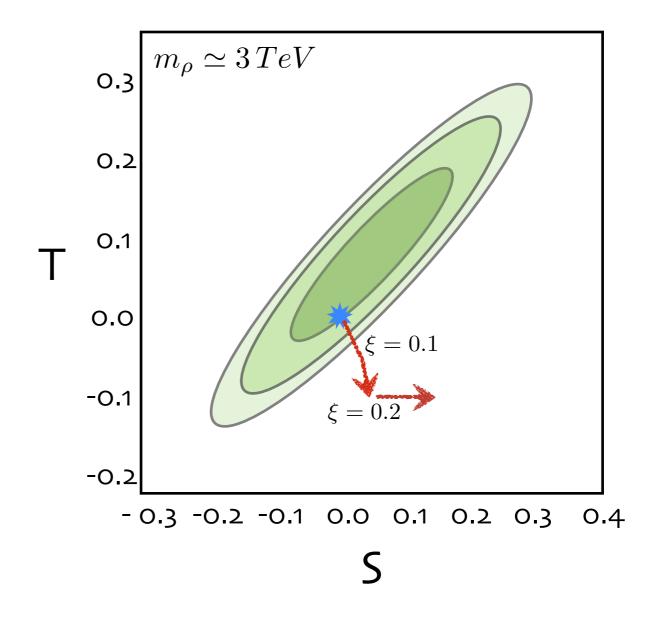
$$S \sim \Pi'_{BW_3}(0)$$
  
 $T \sim \Pi_{W^{\pm}}(0) - \Pi_{W_3}(0)$ 

generic for CH models



vector resonances



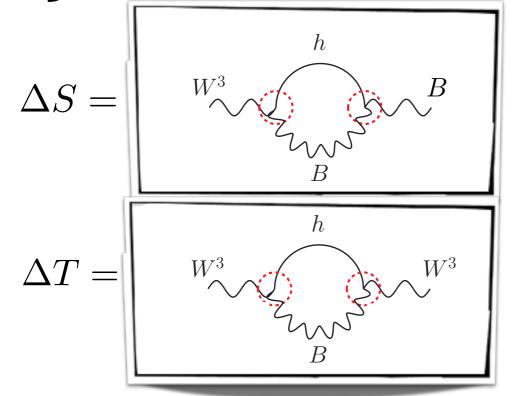


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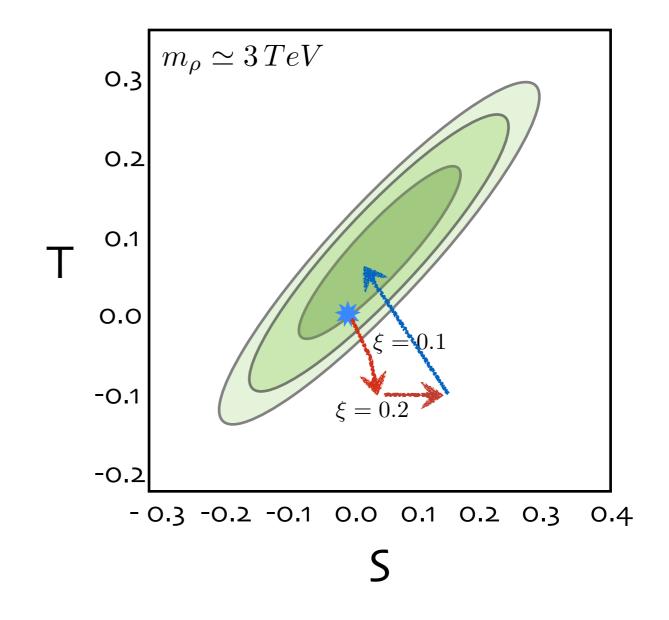
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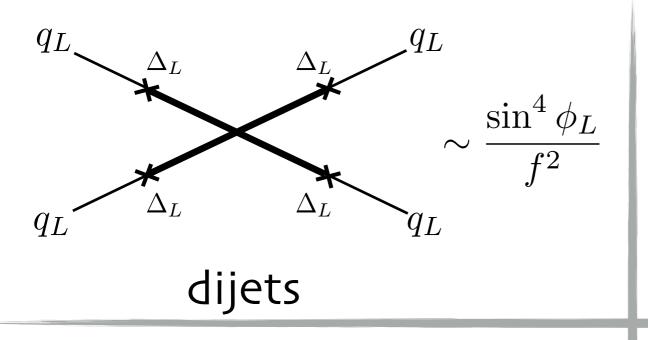
$$\Delta S = \begin{bmatrix} \langle H \rangle & \langle H \rangle \\ \rho & \langle M \rangle \end{bmatrix}$$

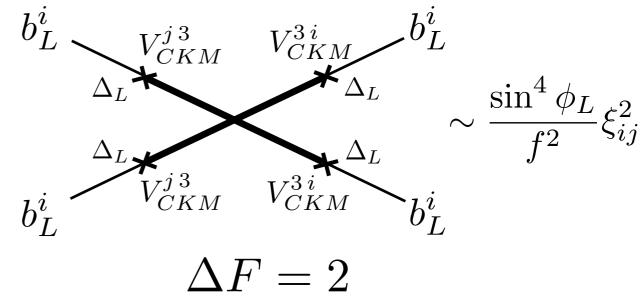


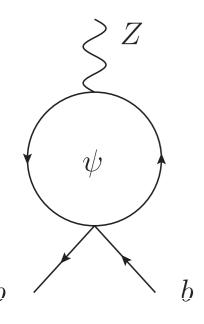
$$\Delta S \simeq \pm \xi \log \left( \frac{m_{\rho}^2}{m_4^2} \right)$$

[Grojean,OM,Panico]

## Flavour

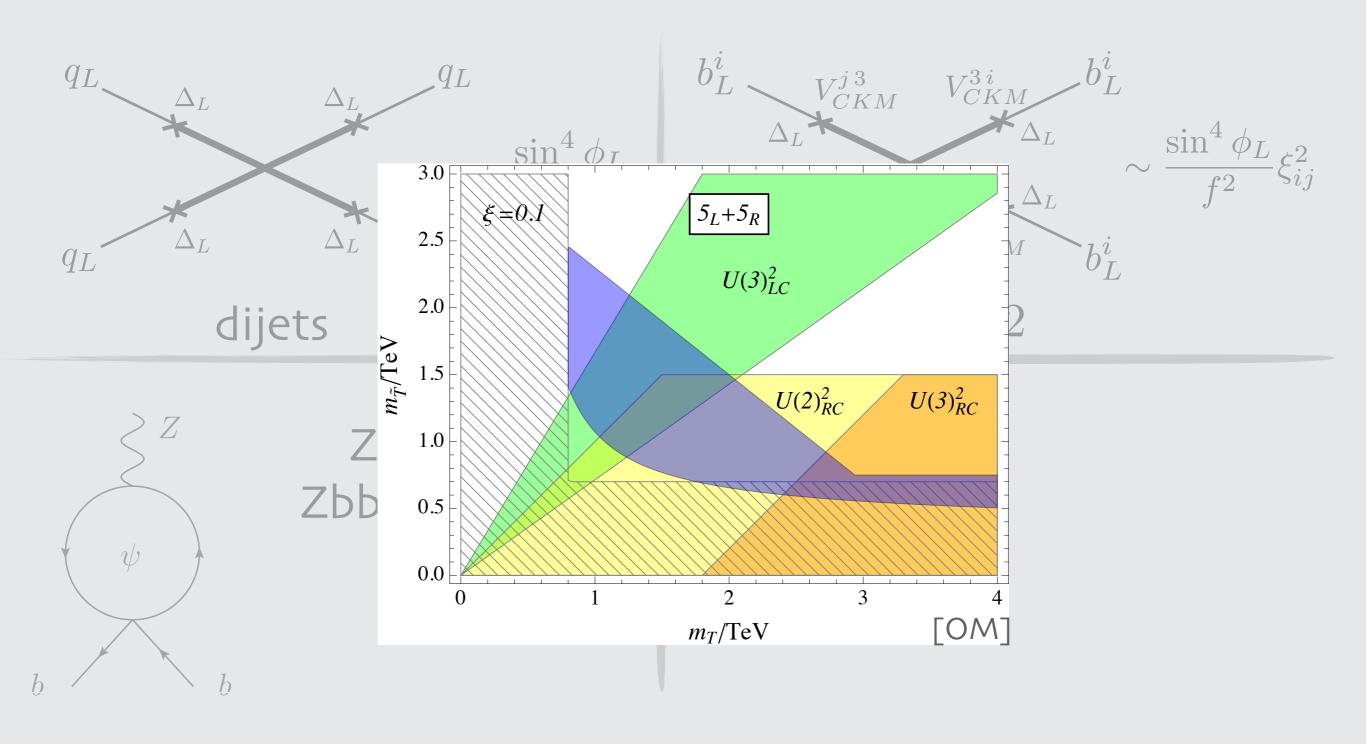






Z width, Zbb coupling

## Flavour



## Summary

CH can address both large and (with ~10% tuning in the simplest models) little hierarchy problems

The simplest CH realizations require 1-2TeV composite resonances

Besides that, deviations from the SM can show up in many other observables