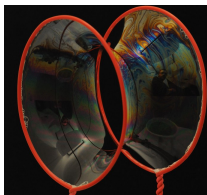


Beyond minimal surfaces: quantum strings for AdS/CFT

Lorenzo Bianchi

Universität Hamburg



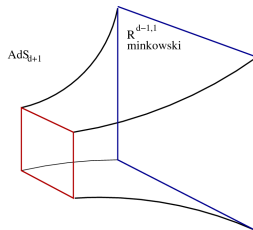
Hamburg
November 10th, 2015

AdS/CFT correspondence

$\mathcal{N} = 4$ Super Yang-Mills.
 $SU(N_c)$ gauge theory in 4d
 Supersymmetric and conformal

Type IIB superstring
 in $AdS_5 \times S^5$

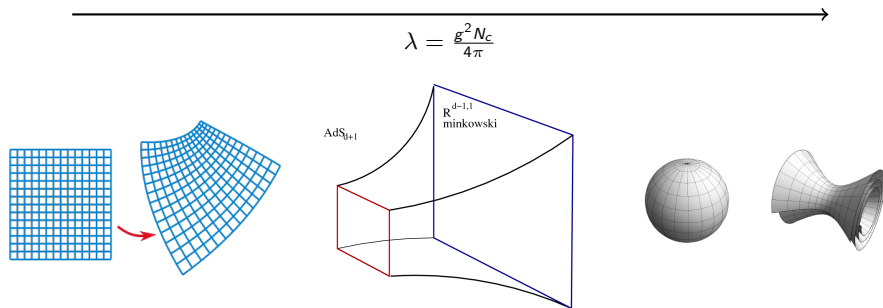
$$\lambda = \frac{g^2 N_c}{4\pi}$$



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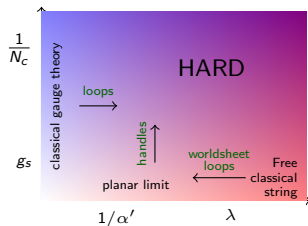
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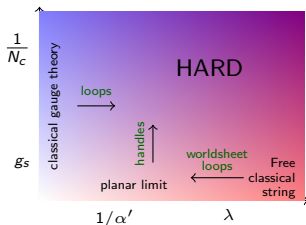
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 $(N_c \rightarrow \infty)$

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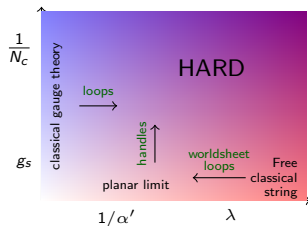
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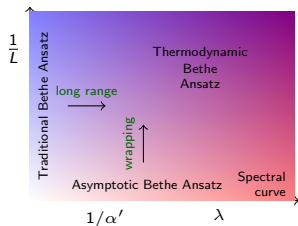
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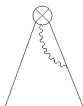
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An example: cusp anomalous dimension

Twist-two operators

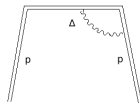
[Korchemsky, 1989; Korchemsky, Marchesini, 1993]



$$\Delta - S \sim 2 + f(\lambda) \log S$$

Wilson lines

[Polyakov, 1980; Korchemsky, Radyushkin, 1987]



$$\langle W \rangle \sim \left(\frac{L}{\epsilon}\right)^{f(\lambda) \log(p \cdot \Delta)}$$

Scattering amplitudes

[Magnea, Sterman, 1990; Bern, Dixon, Smirnov, 2005]

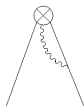


$$\log \frac{A}{A(0)} = \frac{2f(\lambda)}{\epsilon^2} + \mathcal{O}\left(\frac{1}{\epsilon}\right)$$

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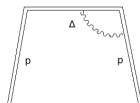
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In $\mathcal{N} = 4$ SYM [Beisert, Eden, Staudacher, 2006]

One can use the **Asymptotic Bethe Ansatz** to write down an **all-loop** integral equation (BES equation) for the cusp anomalous dimension $f(\lambda)$.

In $\mathcal{N} = 6$ Super Chern-Simons in 3d (ABJM) [Aharony, Bergman, Jafferis, Maldacena, 2008] [Gromov, Vieira, 2008]

$$f_{\text{ABJM}}(\lambda) = \frac{1}{2} f_{\mathcal{N}=4}(\lambda_{\text{YM}}) \Big|_{\frac{\sqrt{\lambda_{\text{YM}}}}{4\pi} \rightarrow h(\lambda)}$$

Perturbative procedure

How to compute the cusp anomalous dimension at strong coupling?

Step 1: The action [Metsaev, Tseytlin, 1998]

$$S = T \int d\tau d\sigma \gamma^{\alpha\beta} \partial_\alpha X^a \partial_\beta X^b G_{ab}(X) + \text{fermions}$$

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Step 2: The classical solution [Uvarov, 2009; LB, M. Bianchi, A. Bres, V. Forini, E. Vescovi, 2014]

- Poincaré coordinates on $AdS_4 \times CP^3$

$$ds^2 = \frac{dx^\mu dx_\mu + dw^2}{w^2} + ds_{CP^3}^2 \quad \mu = 0, 1, 2$$

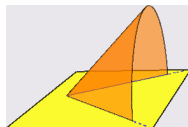
- Minimal surface

$$x^+ = \tau \quad x^- = -\frac{1}{2\sigma} \quad x^+ x^- = -\frac{1}{2} w^2$$

- AdS light-cone gauge [Metsaev, Thorn, Tseytlin, 2000]

$$x^+ = x^0 + x^2 = \tau$$

$$\langle W_{cusp} \rangle = Z_{string} = e^{-\frac{1}{2} f(\lambda) V}$$



Cusp anomaly at two loops [LB, M. Bianchi, A. Bres, V. Forini, E. Vescovi, 2014]

Step 3: Fluctuations around the classical solution - Asymptotic spectrum

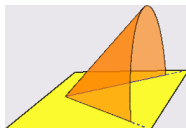
- **Bosons:** 1 mode $m^2 = 1$; 1 mode $m^2 = 1/2$; 6 modes $m^2 = 0$.
- **Fermions:** 6 modes $m^2 = \frac{1}{4}$; 2 modes $m^2 = 0$.

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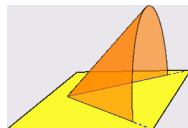
Expand the action up to quartic order in fluctuations and compute all **connected 2-loop vacuum Feynman diagrams**.

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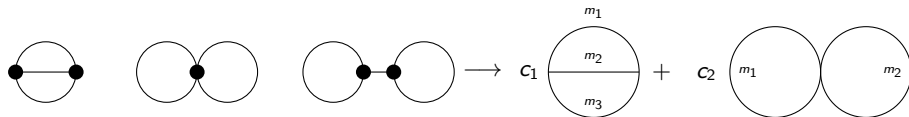
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Expand the action up to quartic order in fluctuations and compute all **connected 2-loop vacuum Feynman diagrams**.

The two-loop result **gives strong support** to a recent conjecture for $h(\lambda)$ [Gromov, Sizov, 2014].



Quantum dispersion relation [LB, M. Bianchi, 2015]

Asymptotic spectrum

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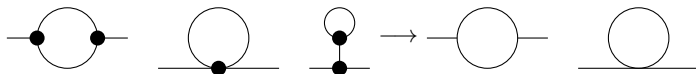
Non-relativistic theory $\Rightarrow$ Corrections to the dispersion relation

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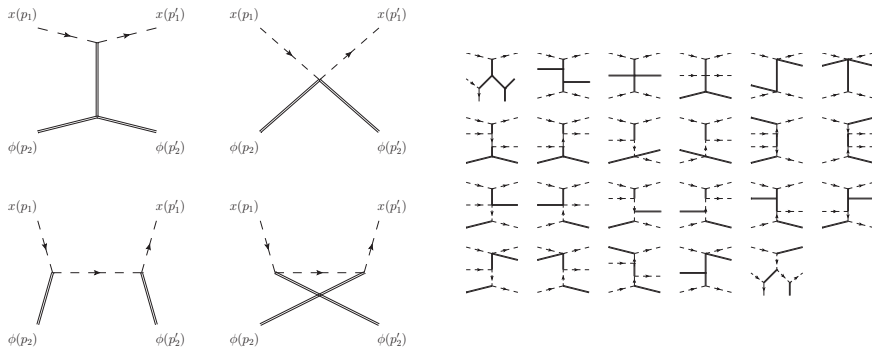


We found agreement with the integrability predictions up to some subtleties.

S-matrix [LB, M.S. Bianchi 2015]

$$\mathbb{S} |\Phi_A(p)\Phi_B(p')\rangle = |\Phi_C(p)\Phi_D(p')\rangle S_{AB}^{CD}(p, p')$$

Tree-level



One-loop S-matrix by unitarity [LB, Hoare, Forini, 2013; Englund, McKeown, Roiban, 2013.]

Standard unitarity in 4d [Bern, Dixon, Dunbar, Kosower, 1994]

$$\mathcal{A}|_{\text{cut}} = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots$$

Glue together the two amplitudes and **uplift** the integral with

$$i\pi\delta^+(p^2 - m^2) \rightarrow \frac{1}{p^2 - m^2 - i\epsilon}$$

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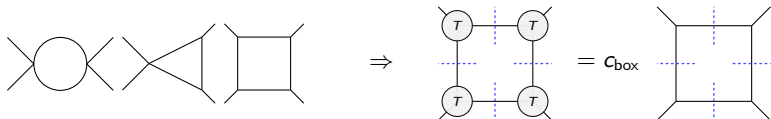
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Generalized unitarity in 4d [Bern, Dixon, Kosower, 1998; Britto, Cachazo, Feng, 2004]

$$\mathcal{A}^L = \sum_i c_i \mathcal{I}_i^{(L)} \rightarrow \text{Known basis of L-loop scalar integrals}$$

For L=1



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Conclusion

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- Perturbative calculations are essential to give a solid foundation and inspiration to any integrability-based construction, and thus to guarantee its **predictivity**.
- It is often possible to improve current perturbative techniques.
- We developed a technique to perform **perturbative computations** in $1+1$ dimensions at one loop and applied it also to **off-shell quantities**.
- Higher loop computations often reveal subtleties regarding **regularization**, with impact on the quantum integrability of the theory.