

GROOMED JET OBSERVABLES AND CLUSTERING EFFECTS

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based on arXiv:1603.06375 with Chris Frye, Andrew Larkoski, Matthew Schwartz

DESY
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GOAL OF THIS TALK

Theoretical aspects of soft-drop groomed jet substructure:

- All-order factorization formula for soft-drop groomed observable
- NNLO analytic calculation that allows complete NNLL resummation

MOTIVATION

- precision comparison between experiment and data due to contamination from UE/PU

- Jet/Event Groomers: removing and mitigating contamination

filtering/mass drop: [Butterworth, Davison, Rubin, Salam 0802.2470](#)

trimming: [Krohn, Thaler, Wang 0912.1342](#)

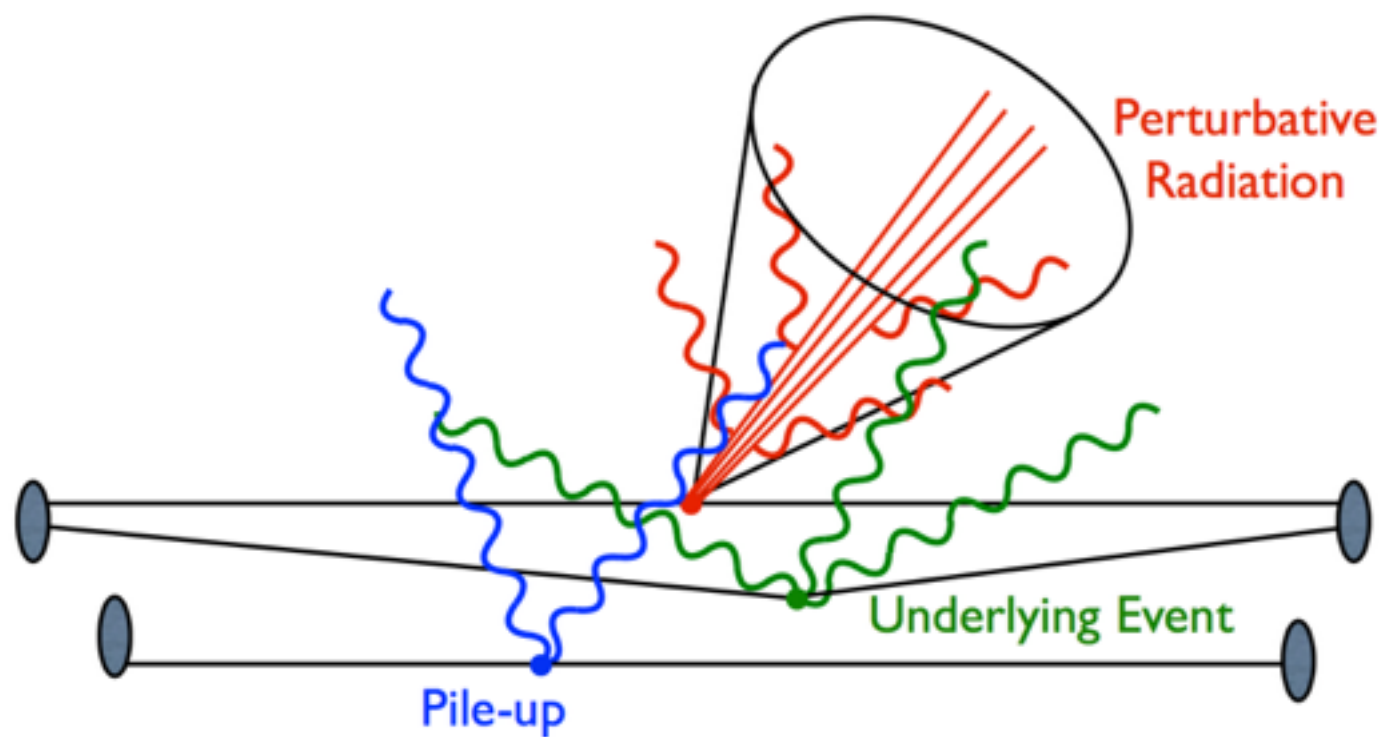
pruning: [Ellis, Vermilion, Walsh 0912.0033](#)

soft drop: [Larkoski, Marzani,](#)

[Soyez, Thaler 1402.2657](#)

modified mass drop: [Dasgupta,](#)

[Fregoso, Marzani, Salam 1307.0007](#)



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Challenges in resumming jet substructure at NLL and beyond :

- complication due to initial state radiation/ Non-global logs

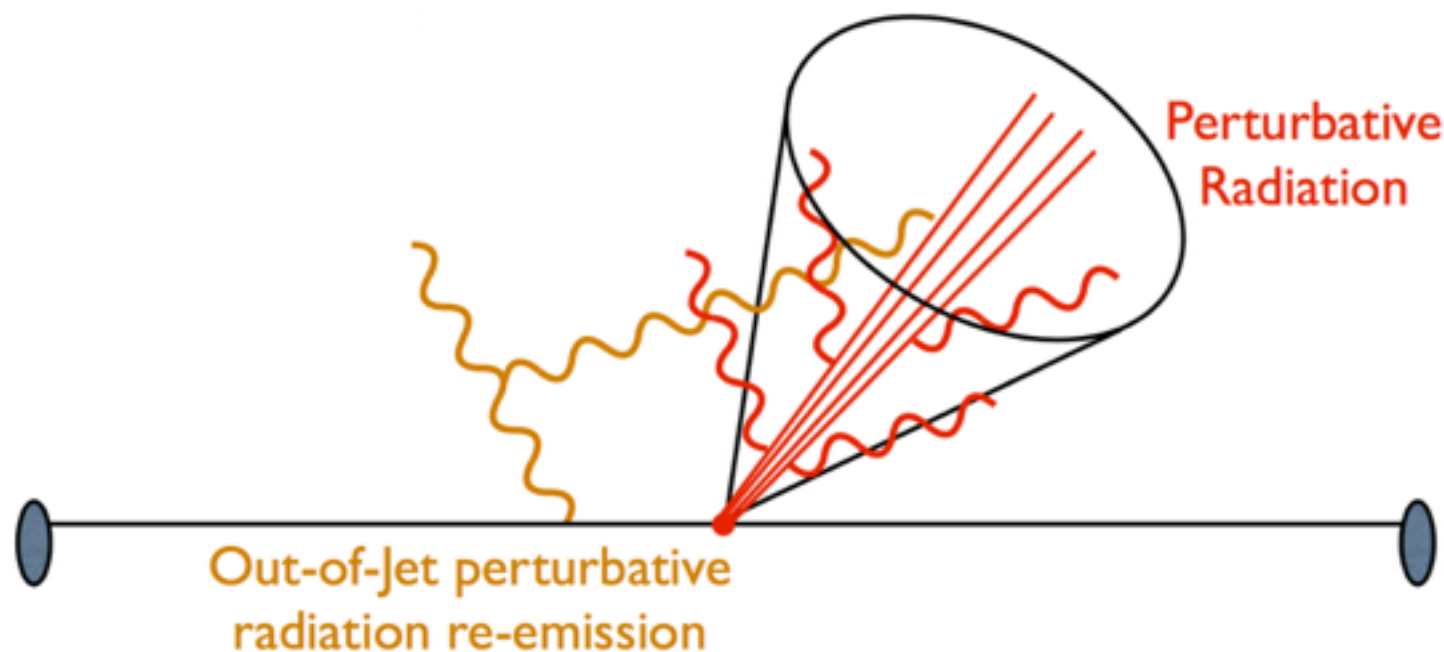
Removing perturbative soft radiation:

final state wide-angle radiation

initial state radiation

non-global radiation

Reduces process dependence
Eliminates non-global logarithms



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➤ Desired theoretical property for mMDT and soft drop:

No non-global logs at NLL, no dependence on jet radius.

[Dasgupta,Fregoso, Marzani, Salam 1307.0007][Larkoski, Marzani, Soyez,Thaler 1402.2657]

Going beyond

- Can we prove the absence of non-global logs and jet-radius dependence to all orders?
- What is the proper choice of grooming parameters/reclustering algorithm that allow a clean theoretical description?

OUTLINE

- define soft-drop algorithm
- Factorization at e^+e^-
- Obtaining two-loop soft non-cusp anomalous dimension
- study of clustering effects

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THE SOFT-DROP GROOMER

GROOMING PROCEDURE

- 1. Take a jet with radius $R \sim 1$, recluster with C/A.

$$d_{ij} = R_{ij} \equiv \sqrt{(y_i - y_j)^2 + (\phi_i - \phi_j)^2}$$

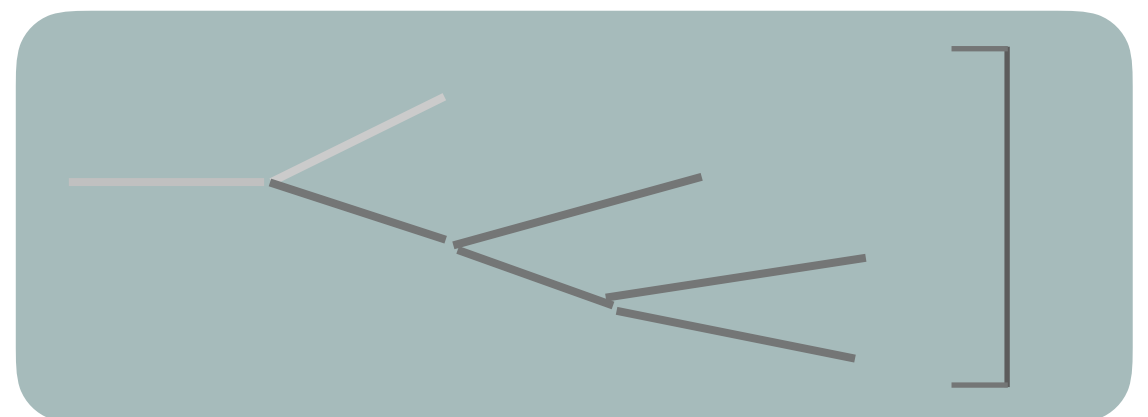
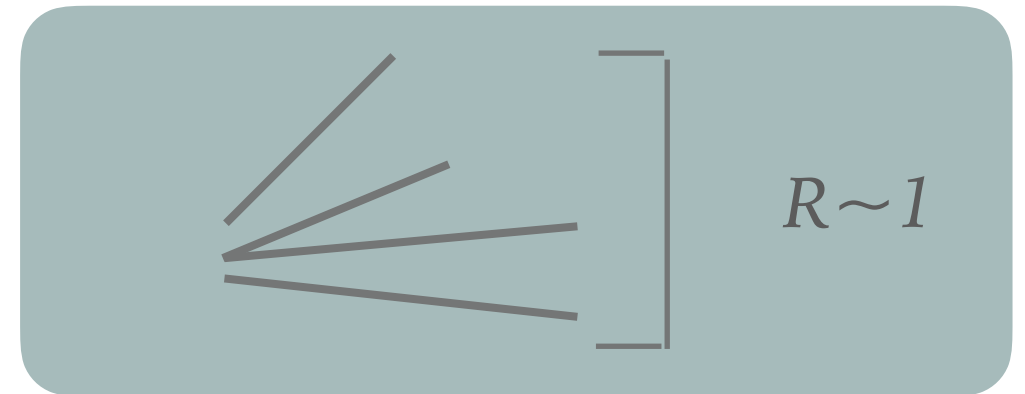
- 2 Undo the last step of reclustering, check if two branches (i,j) satisfy soft-drop condition. If not drop the softer branch.

$$\boxed{\frac{\min[p_{Ti}, p_{Tj}]}{p_{Ti} + p_{Tj}} > z_{\text{cut}} \left(\frac{R_{ij}}{R} \right)^\beta}$$

$$z_{\text{cut}} \sim 0.1$$

$$\beta = 0 \quad \text{check only energy fraction (mMDT)}$$

$$\beta = \infty \quad \text{no soft drop}$$

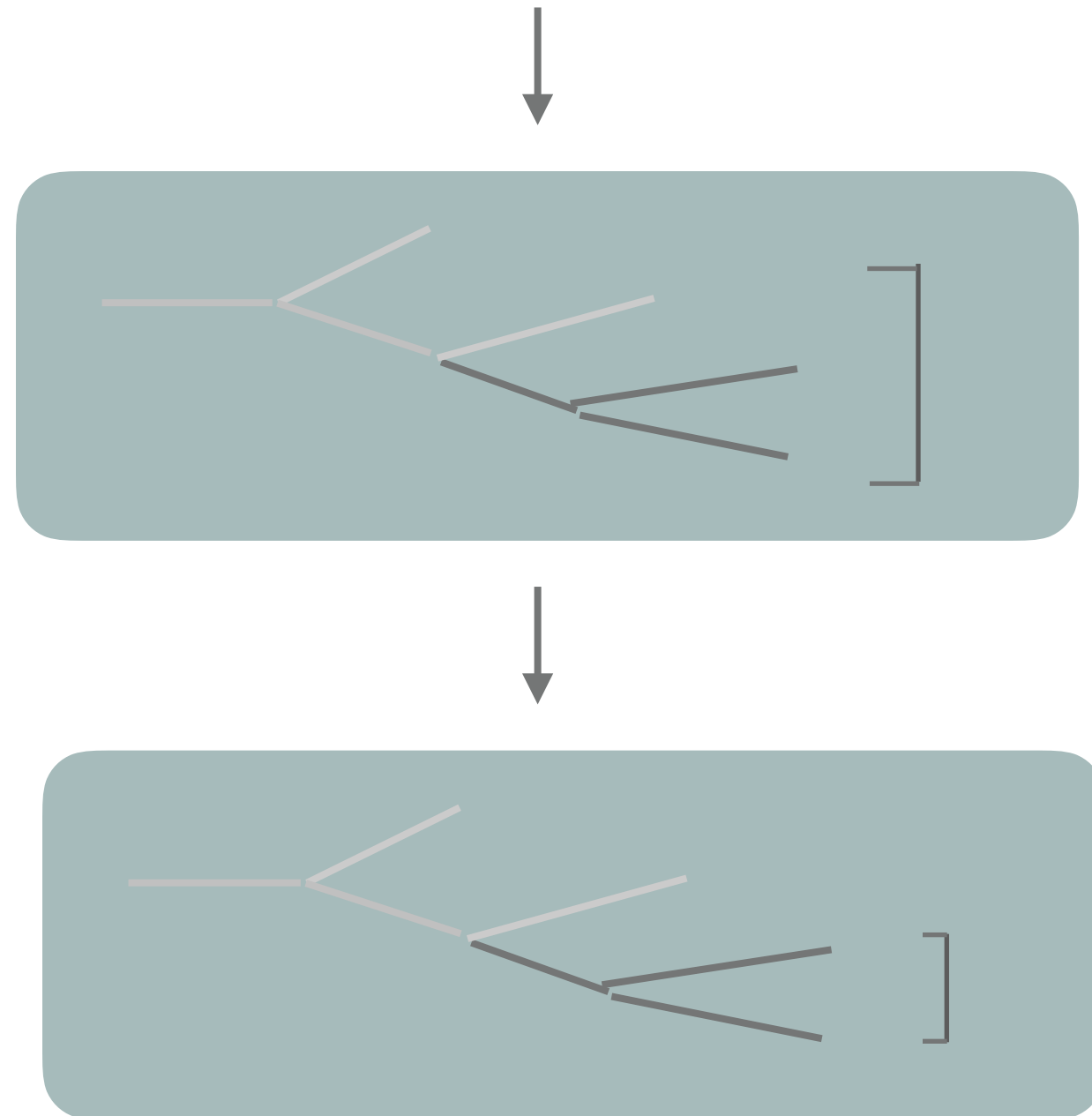


GROOMING PROCEDURE

- repeat step 2 on the harder branch, until the soft-drop condition is satisfied.
- measure energy correlation function of the groomed jet

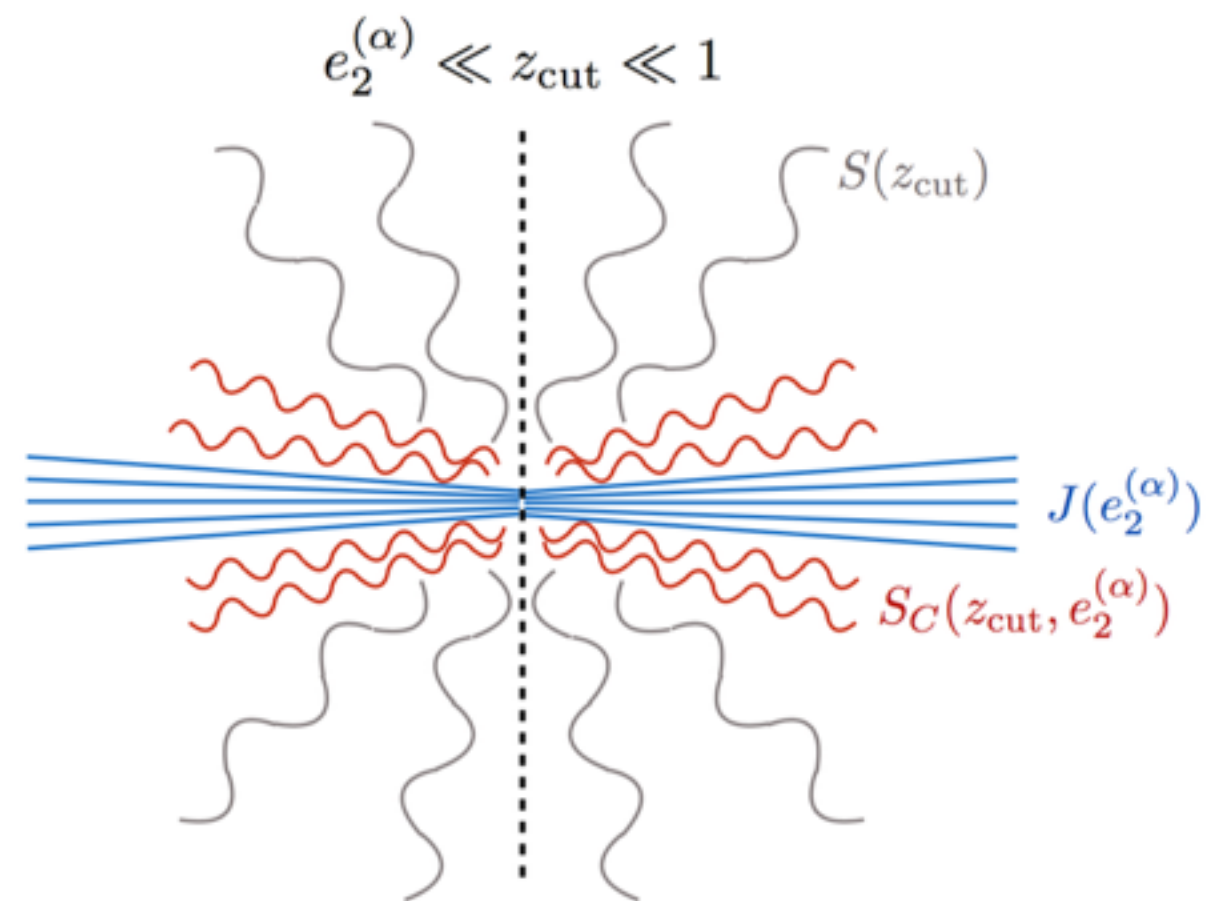
$$e_2^{(\alpha)} \Big|_{pp} = \frac{1}{p_{TJ}^2} \sum_{i < j \in J} p_{Ti} p_{Tj} R_{ij}^\alpha$$

groomed jet mass $e_2^{(2)} = \frac{m_g^2}{E_{Jg}^2}$

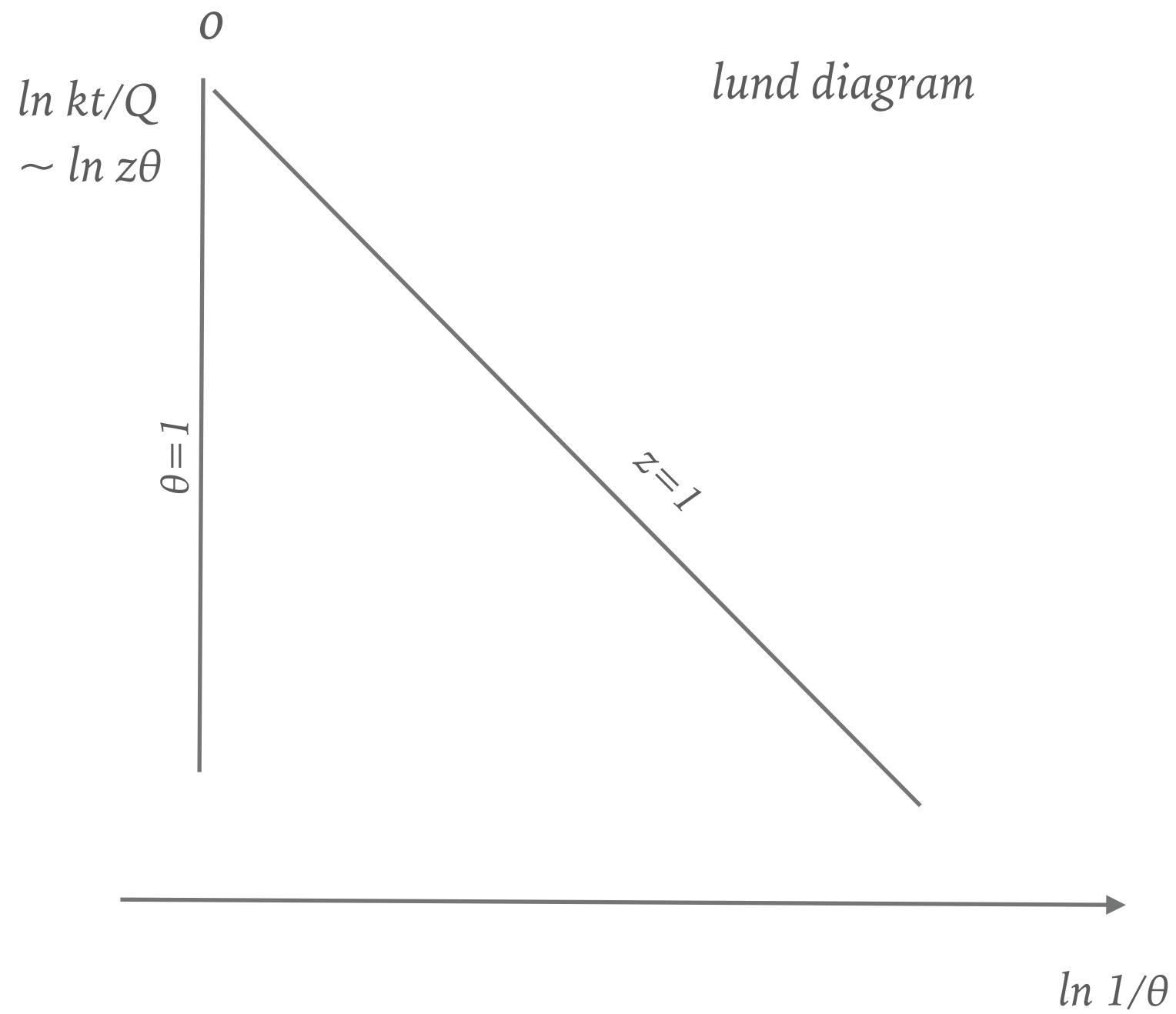


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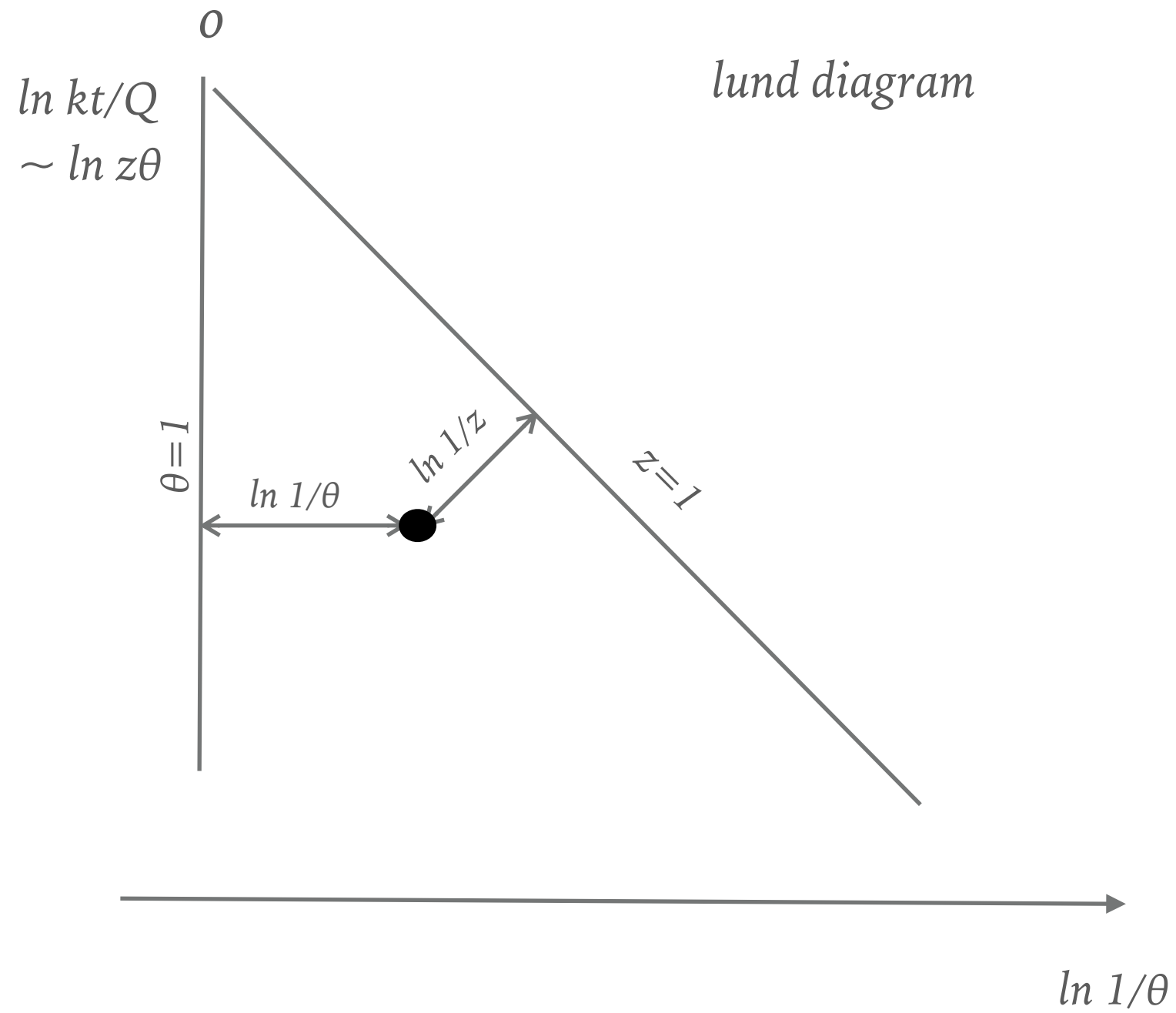
FACTORIZATION AT $E+E- \longrightarrow$ DIJETS



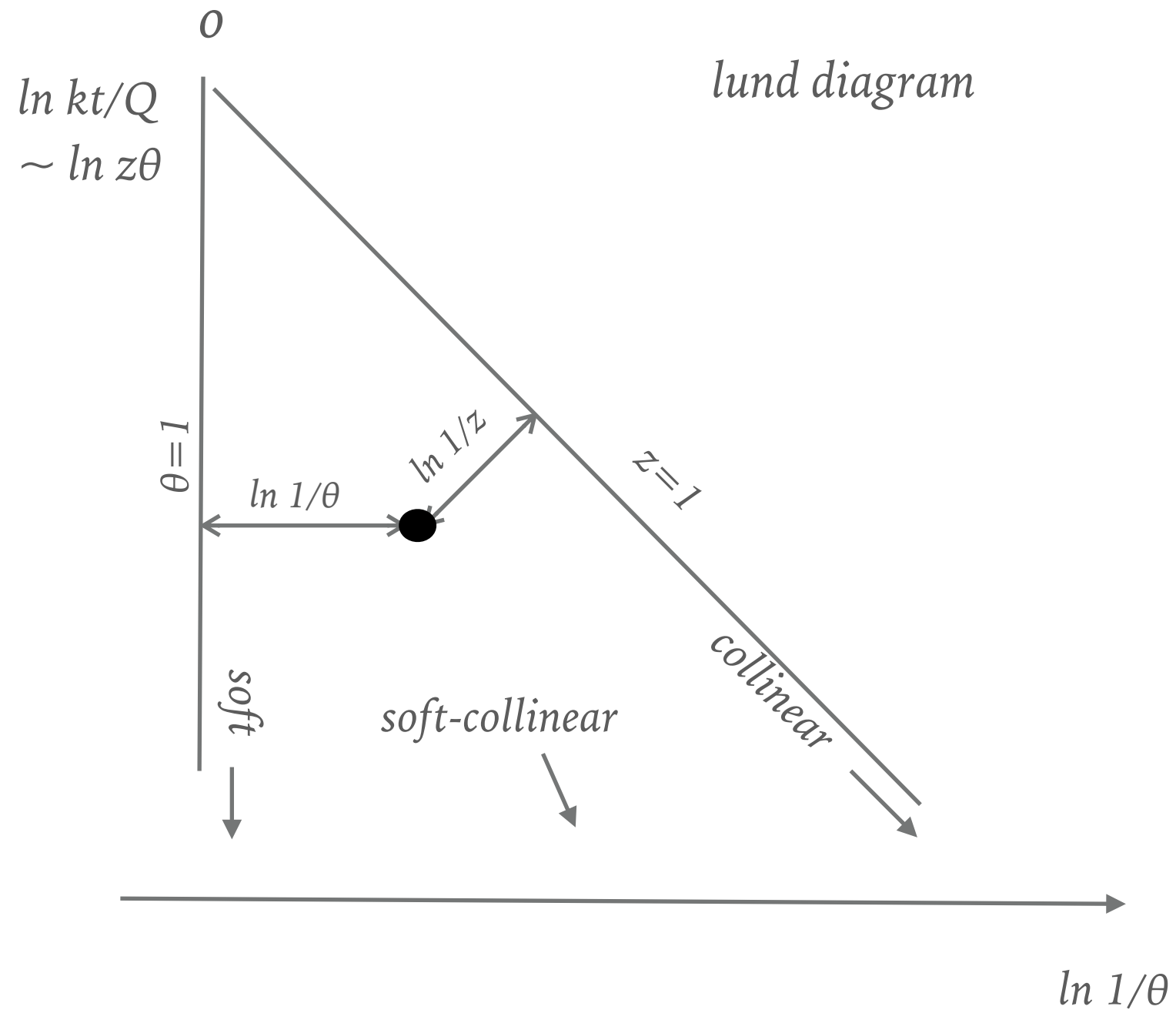
KINEMATIC REGIONS



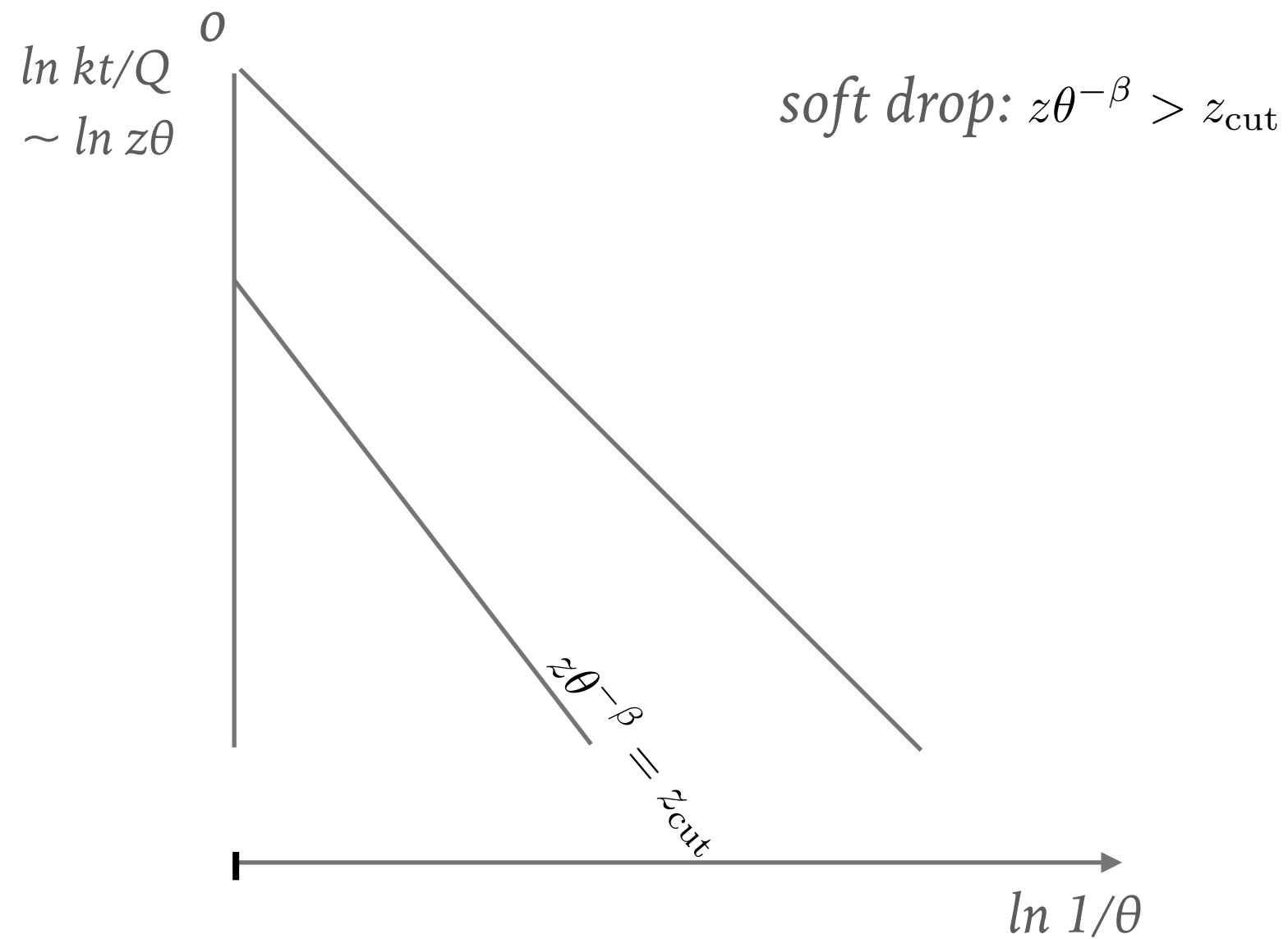
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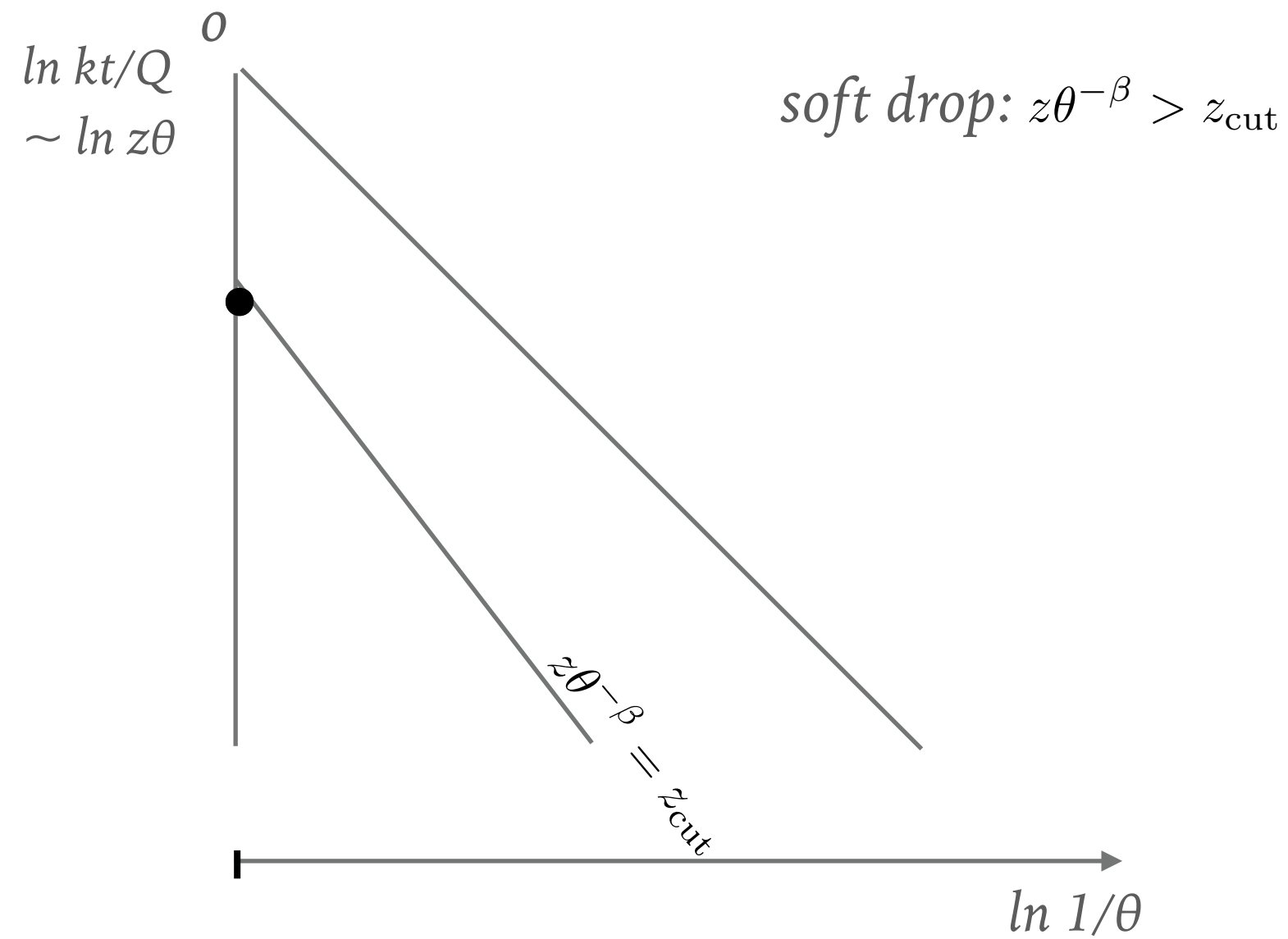
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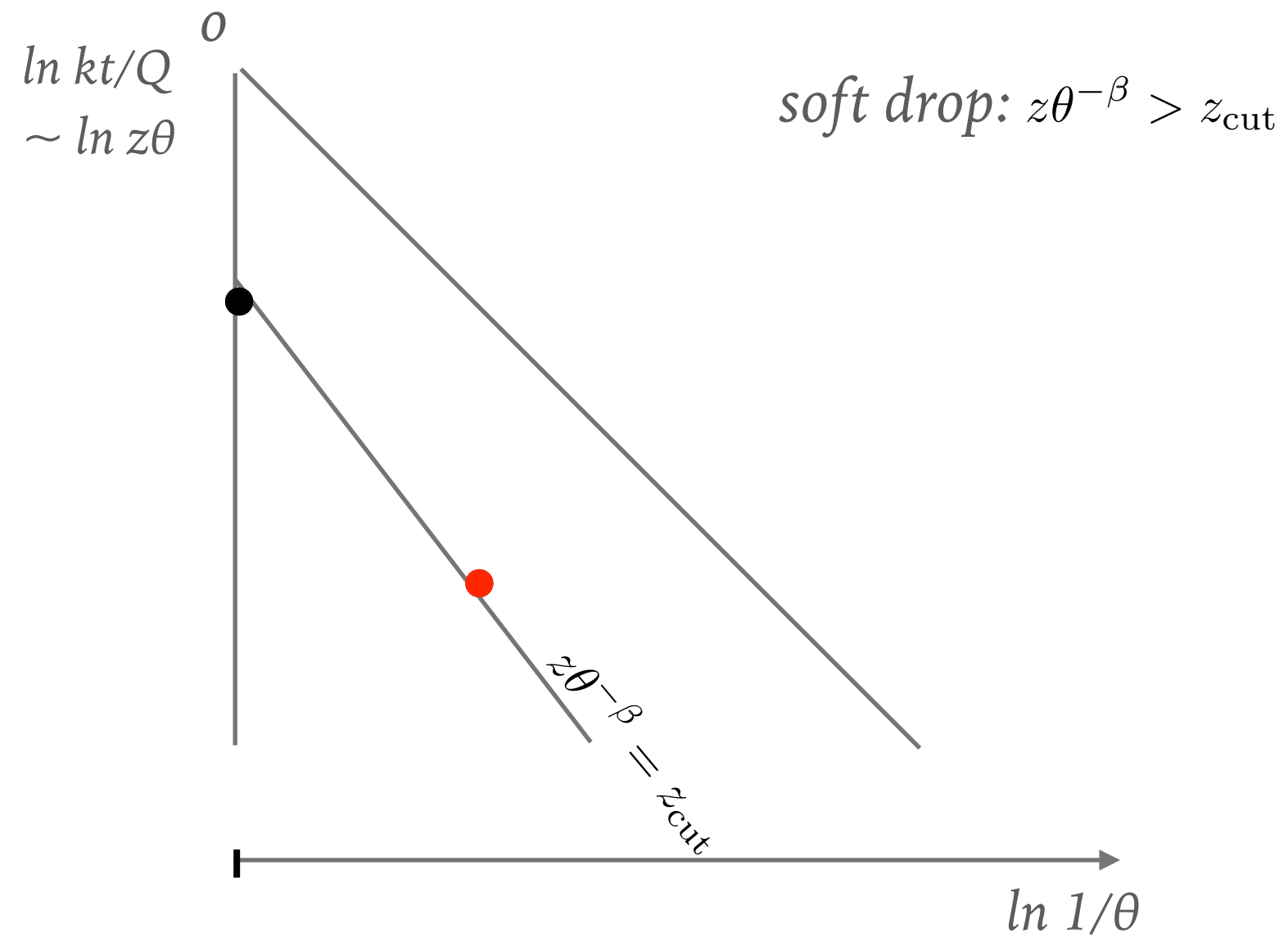
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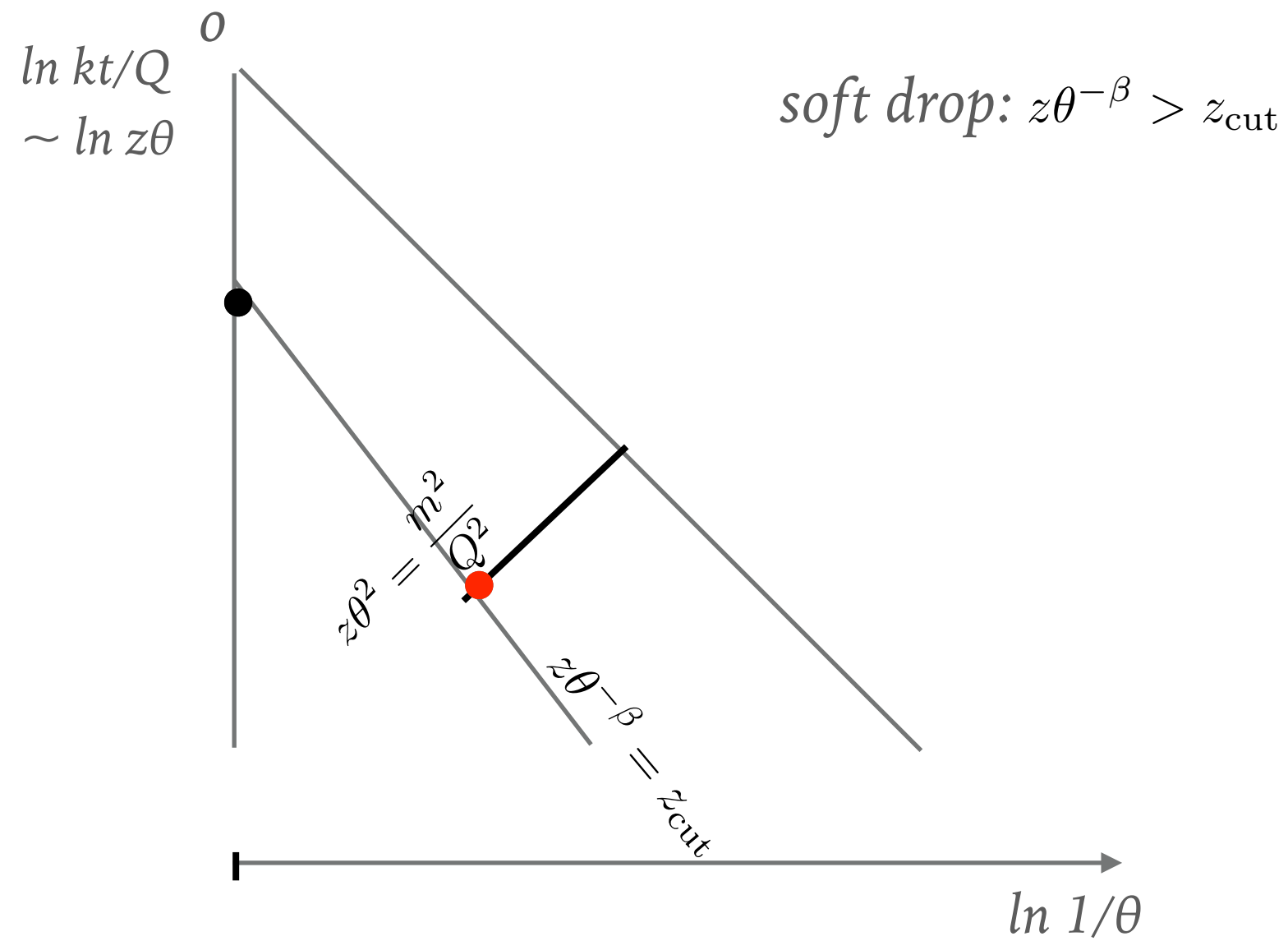
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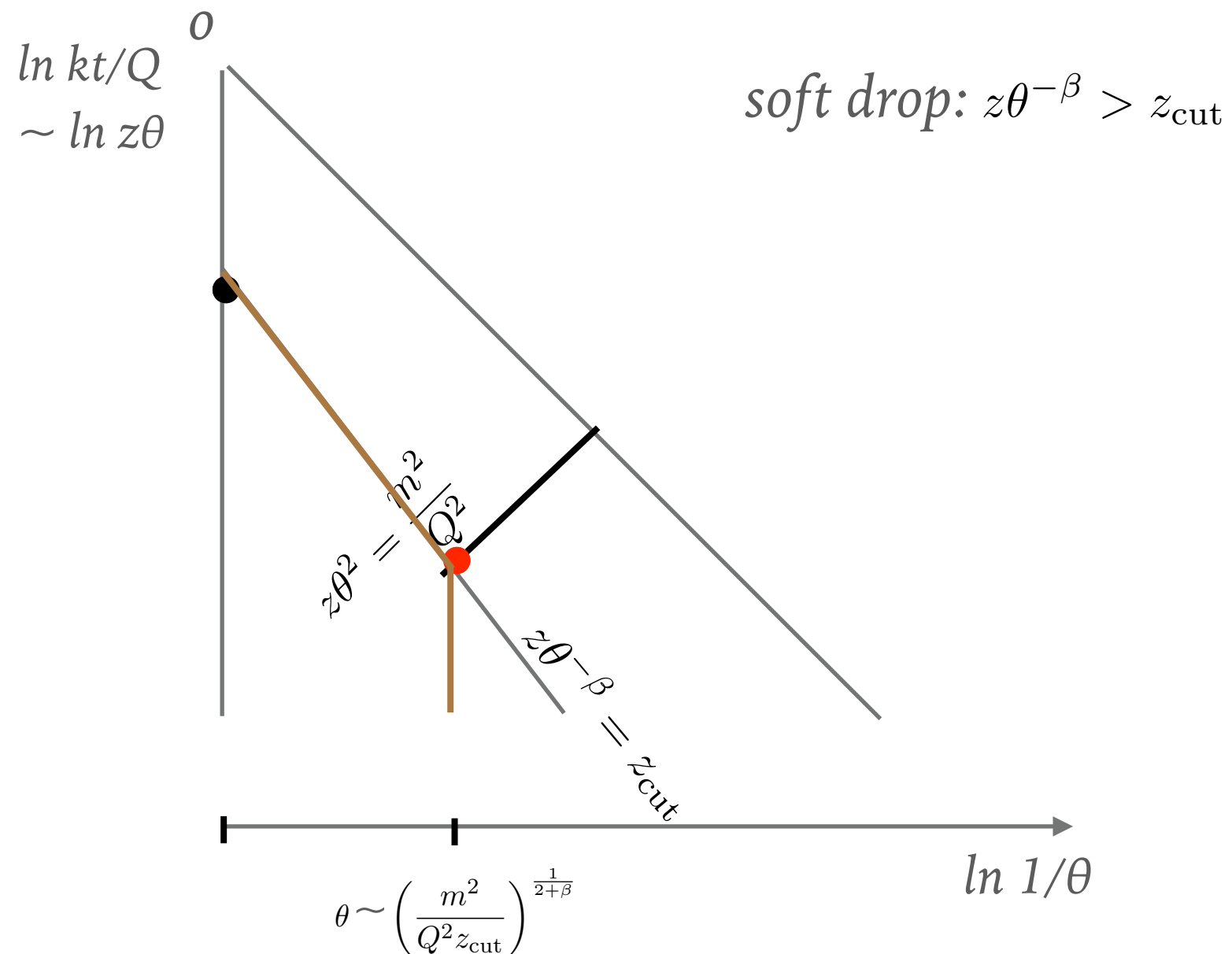
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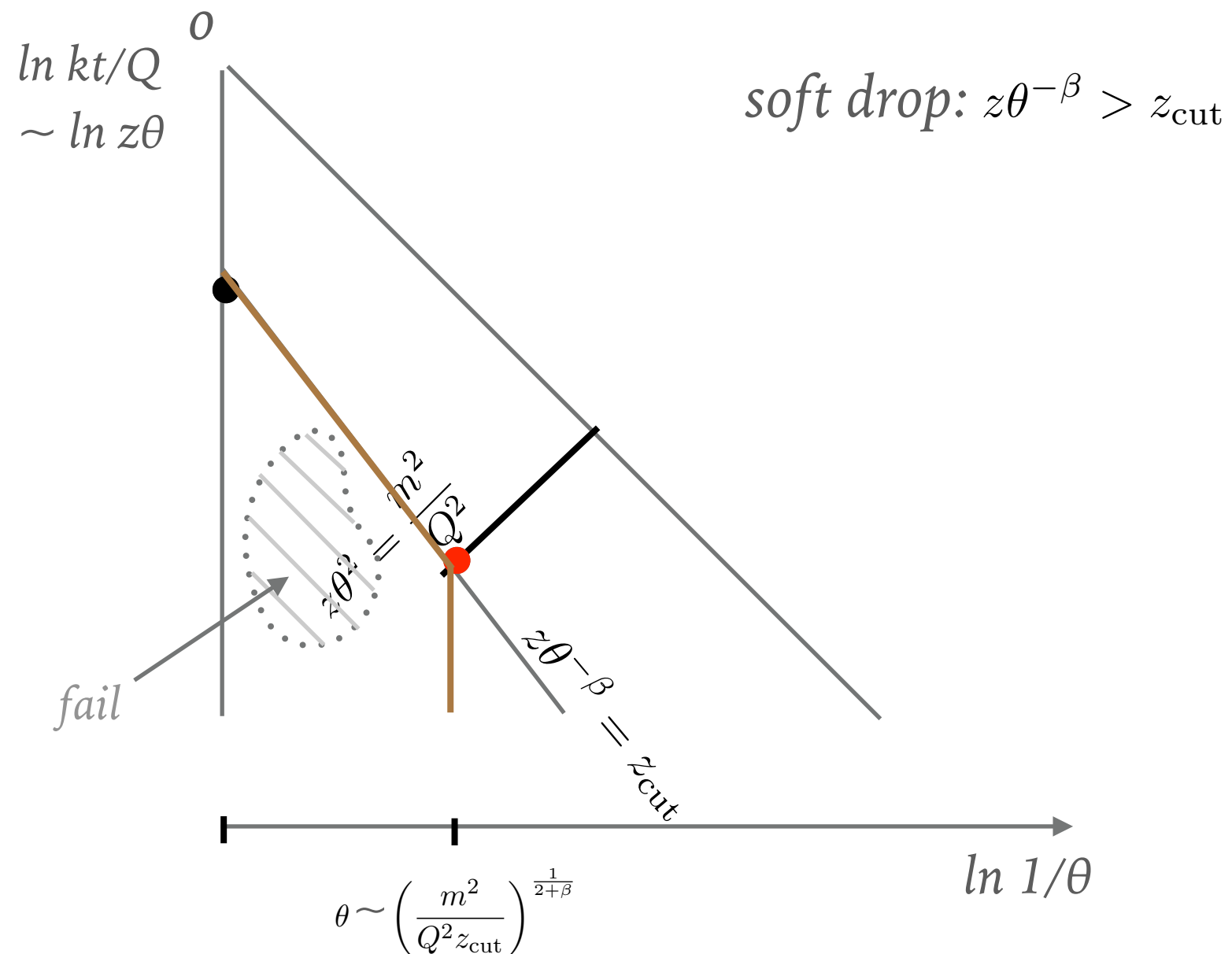
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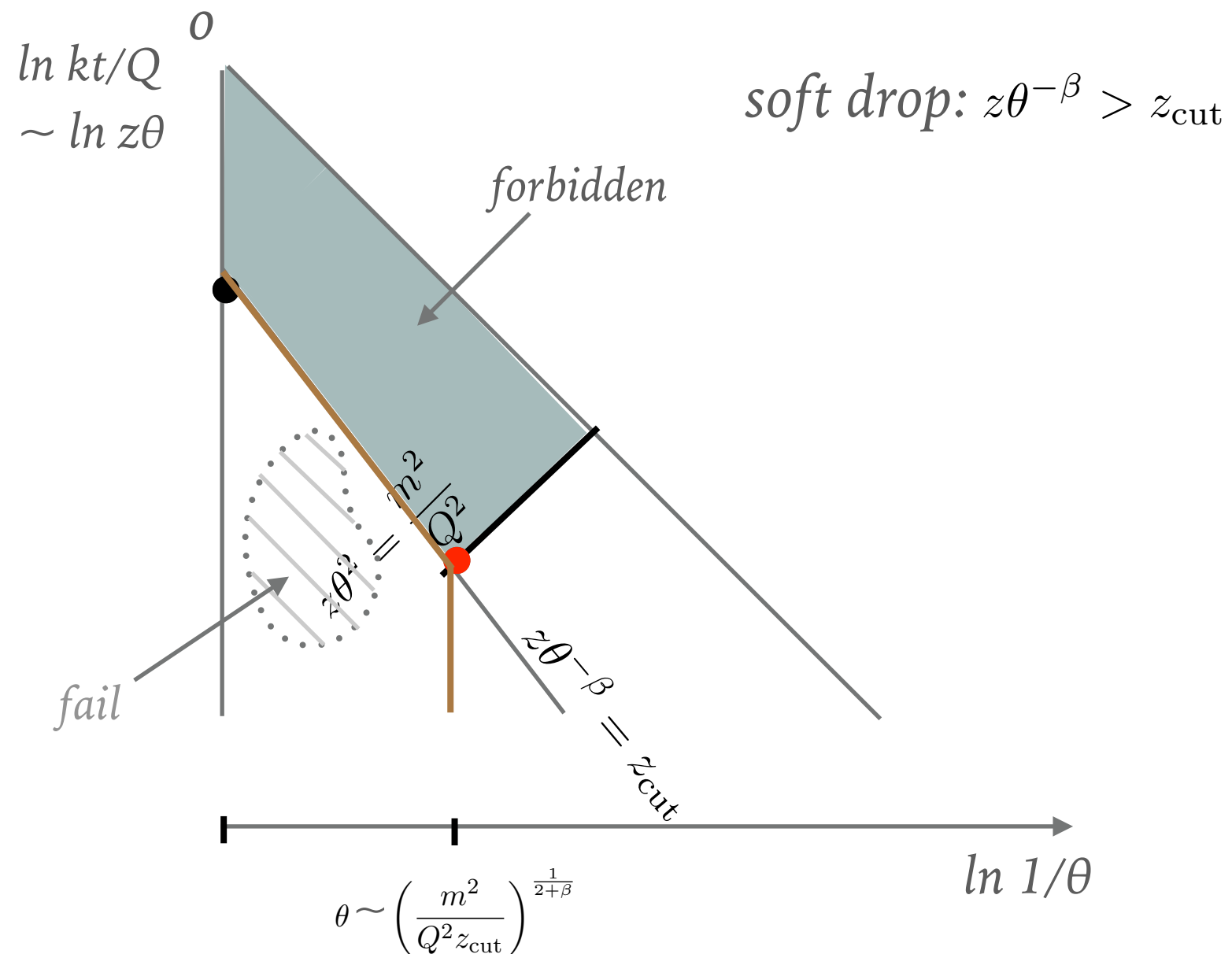
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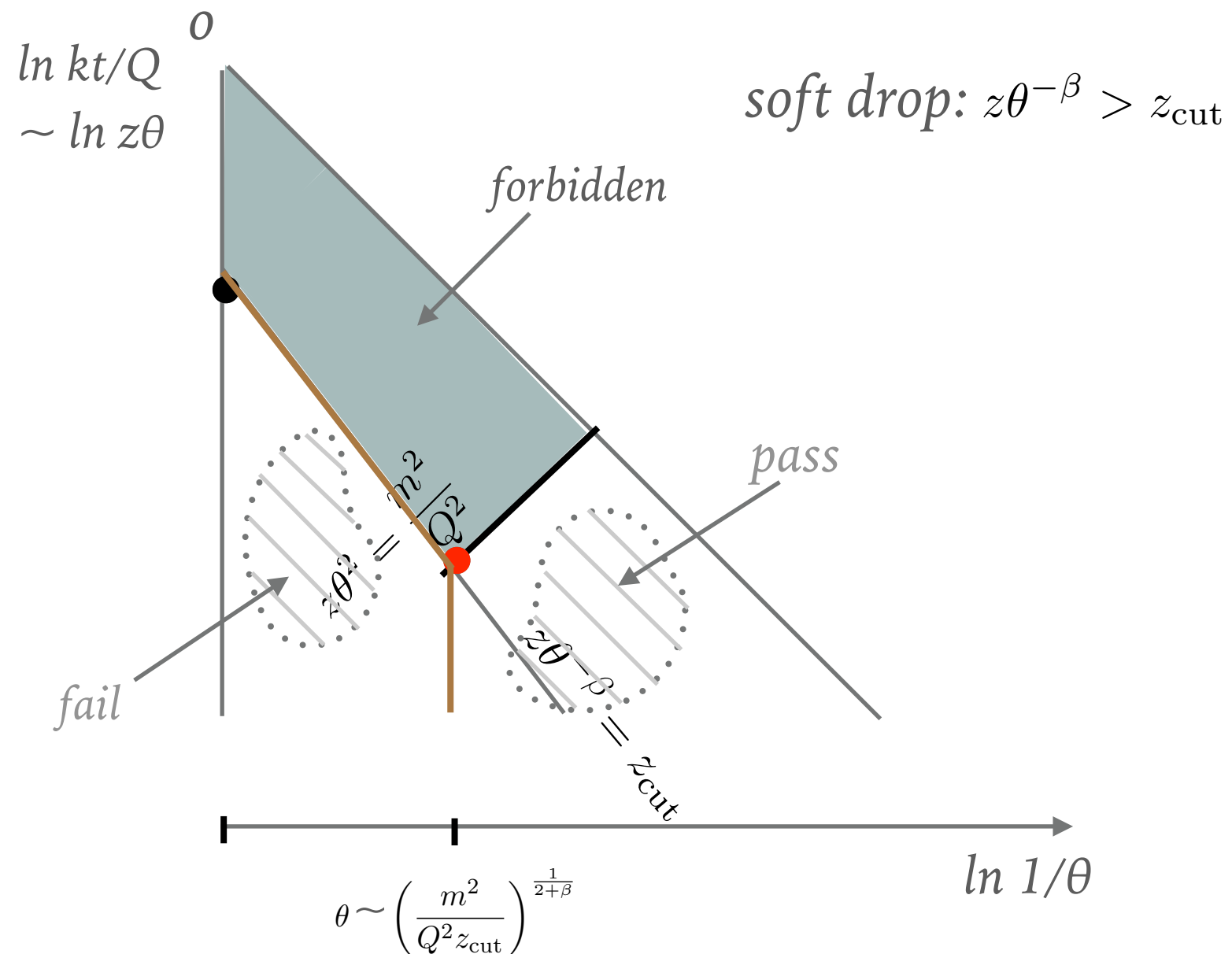
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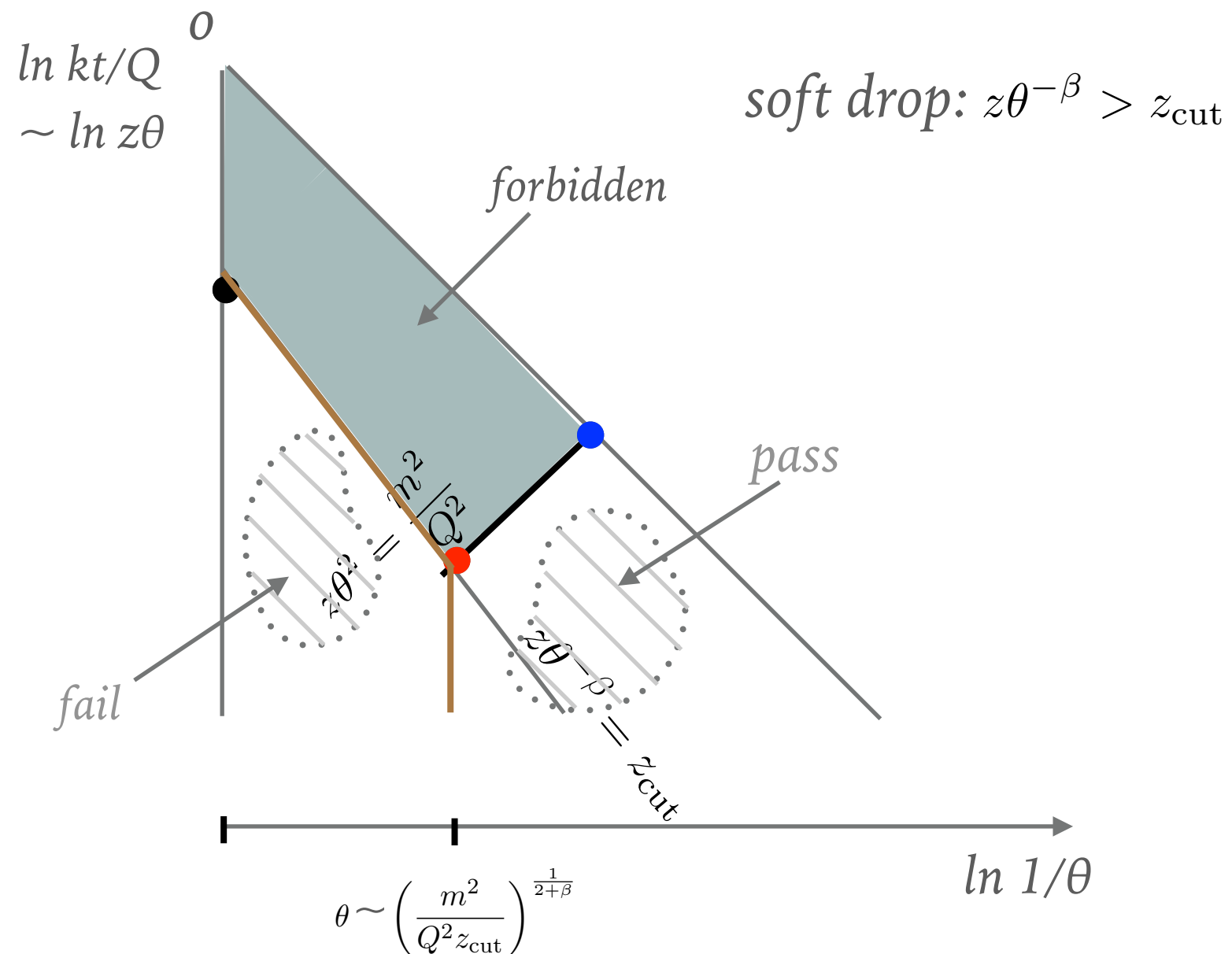
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HIERARCHY OF SCALES

Focus on regime $m^2/Q^2 \ll z_{\text{cut}} \ll 1$

► $m^2/Q^2 \ll z_{\text{cut}}$

Wide angle emissions are not allowed to pass soft-drop

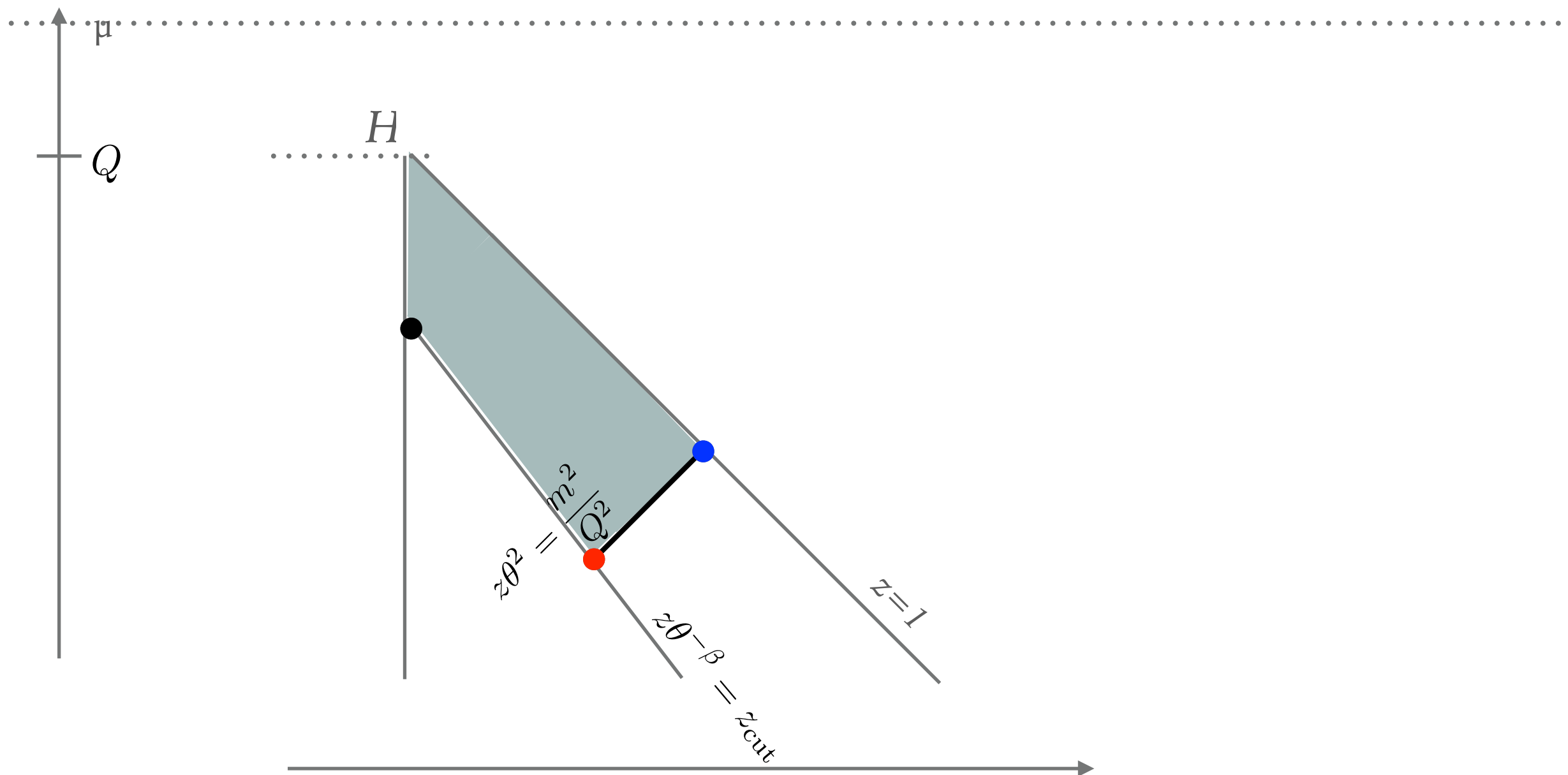
● After soft drop, all remaining particles in the jet must be collinear!

► $z_{\text{cut}} \ll 1$

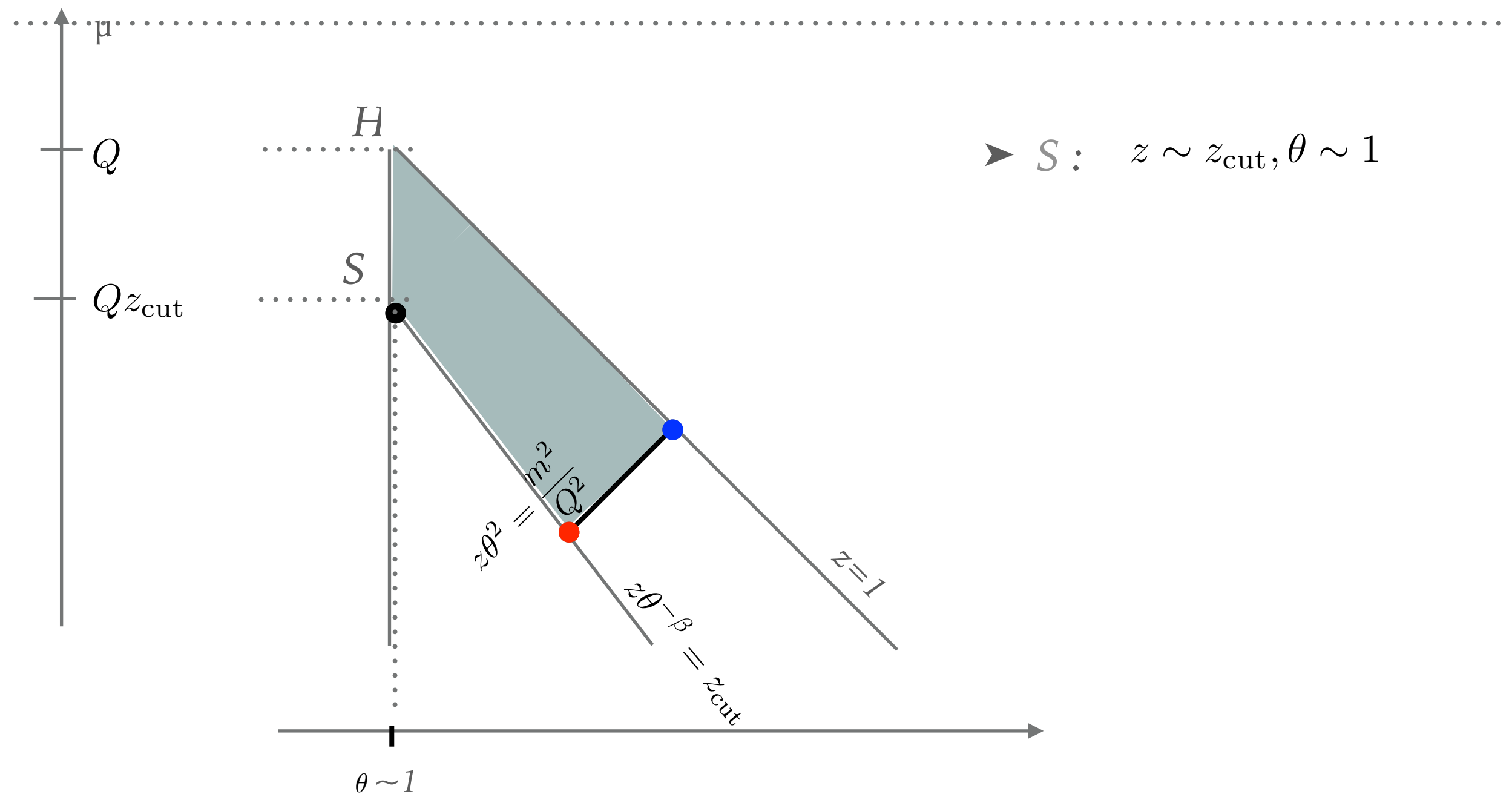
All hard-collinear particles survive soft-drop grooming.

● effects of jet energy loss and change of jet flavor during grooming is power-suppressed by z_{cut} . [*Dasgupta, Fregoso, Marzani, Salam 1307.0007*]

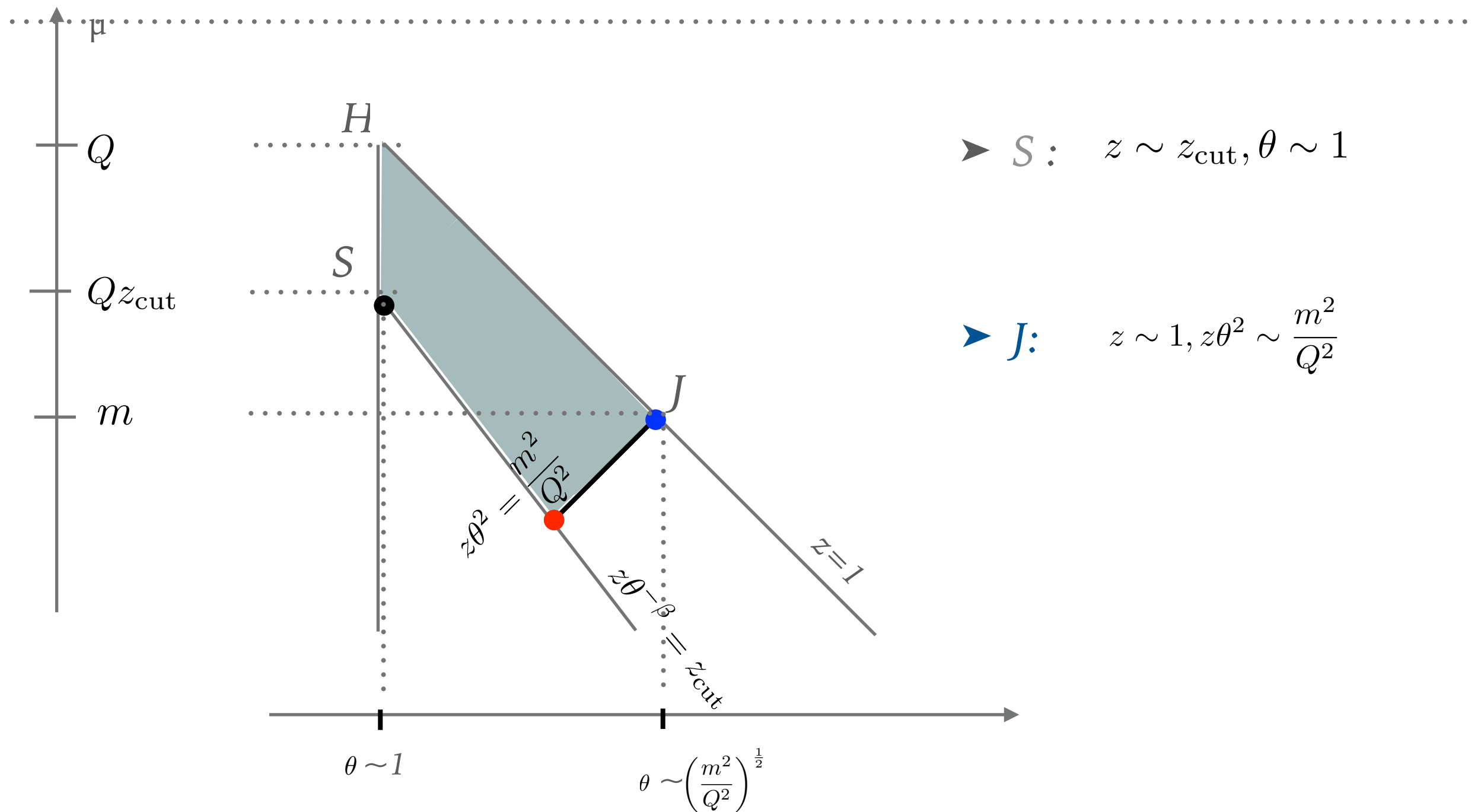
FINDING RELEVANT MODES



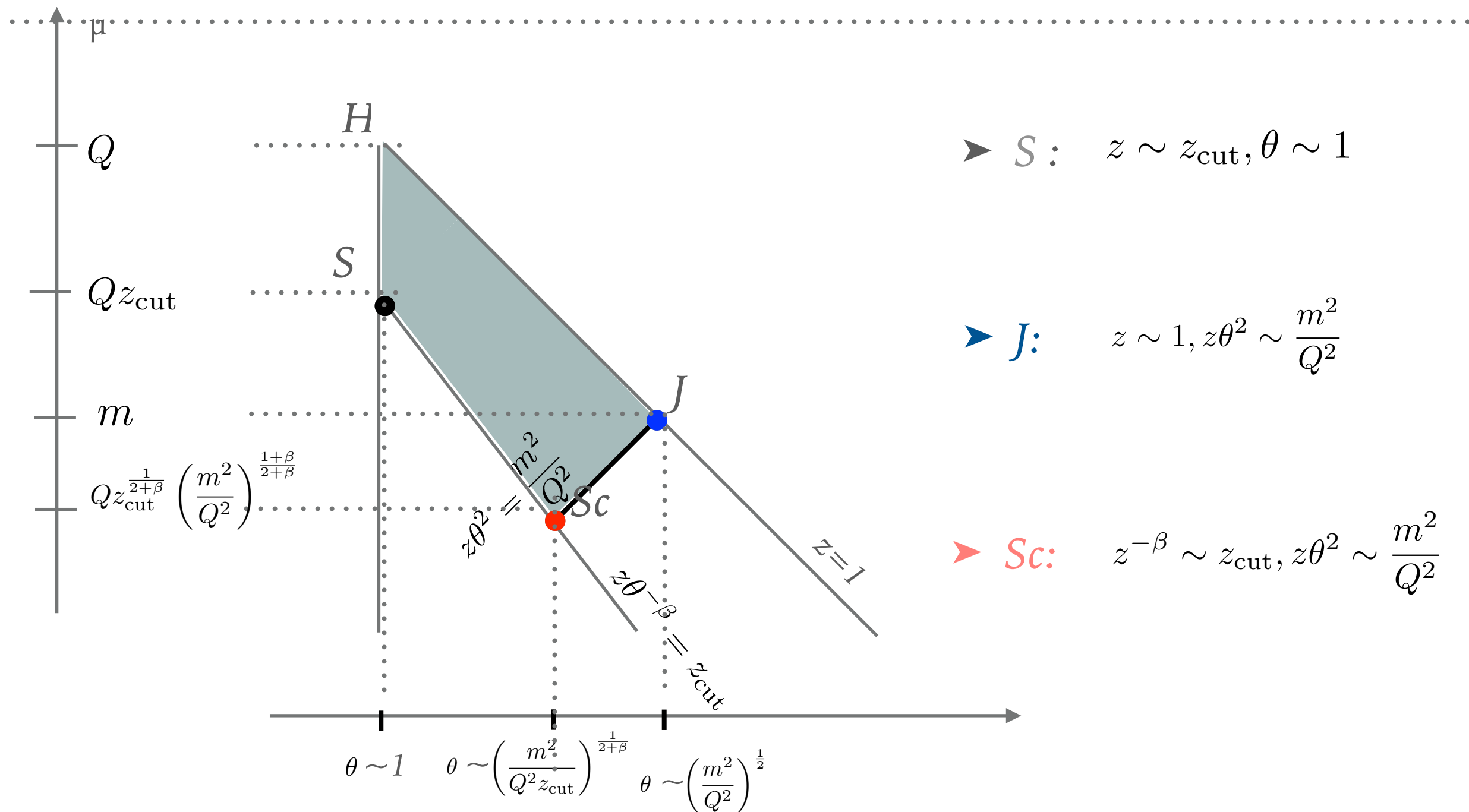
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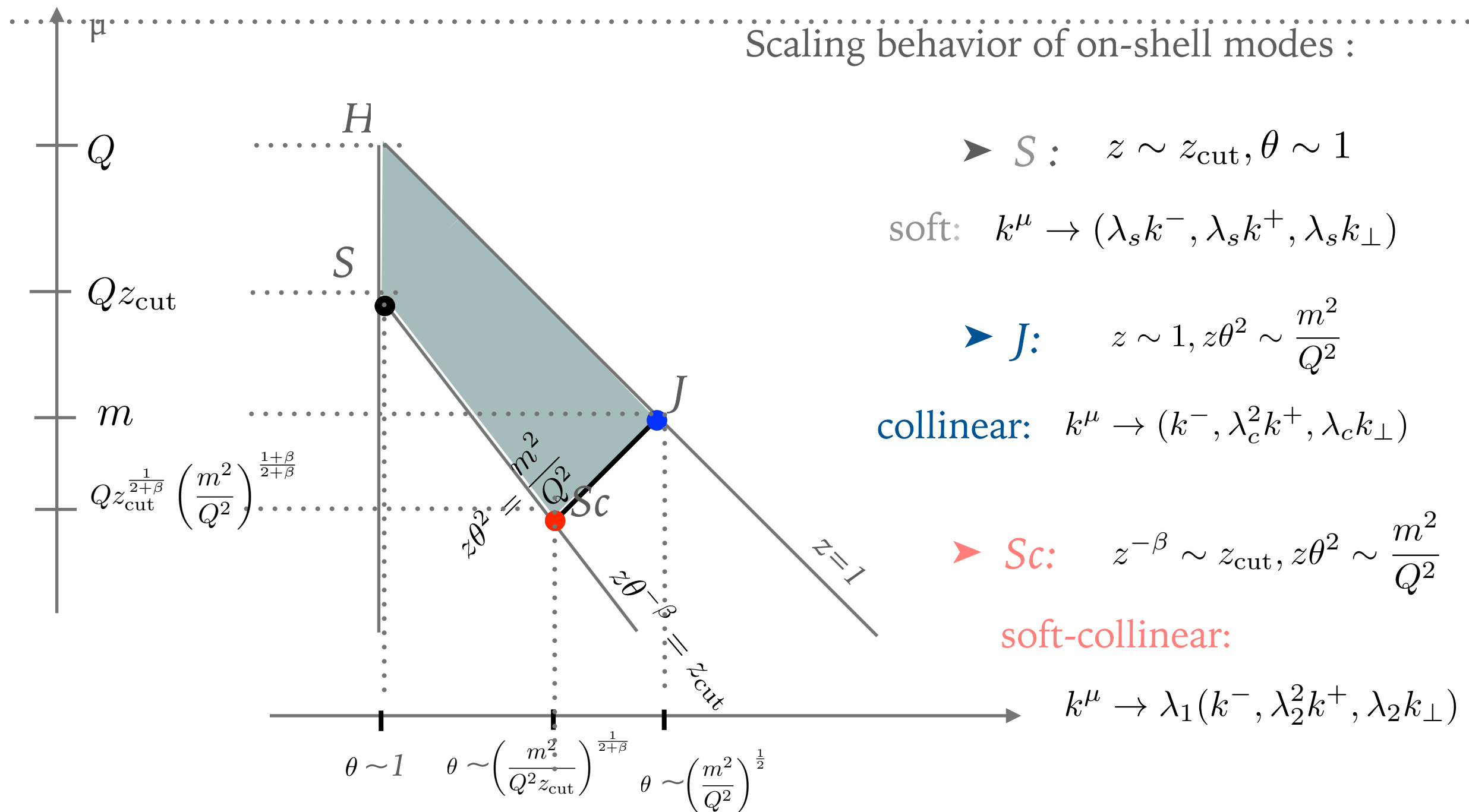
FINDING RELEVANT MODES



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FINDING RELEVANT MODES



➤ Factorization theorem

only collinear objects contribute to the groomed mass

$$\frac{d^2\sigma}{de_{2,L}^{(\alpha)} de_{2,R}^{(\alpha)}} = H(Q^2) S(z_{\text{cut}}) \left[S_{C,L}(z_{\text{cut}} e_{2,L}^{(\alpha)}) \otimes J_L(e_{2,L}^{(\alpha)}) \right] \left[S_{C,R}(z_{\text{cut}} e_{2,R}^{(\alpha)}) \otimes J_R(e_{2,R}^{(\alpha)}) \right]$$

fail soft-drop

pass soft-drop near the boundary

➤ Consequences of factorization

- Collinear universality allows to predict the shape of the distribution at the LHC using results obtained at e^+e^-
- the shape is unaffected by non-global logarithms and is independent of the jet radius.
- (e^+e^-) single scale dependence: the absence of clustering logs/non-global logs

CLUSTERING EFFECT

Check single scale dependence:

- The jet function is inclusive
- Can check either the soft function or soft-collinear function depend on one single scale

$$\Theta_{SD} : C/A \text{ clustering} + \text{soft-drop veto}$$



Strict factorization holds

- Θ_{SD} acts on soft/collinear/soft-collinear final state independently
- For C/A, clustering sequence is invariant under soft/collinear/soft-collinear scaling.

CLUSTERING EFFECT

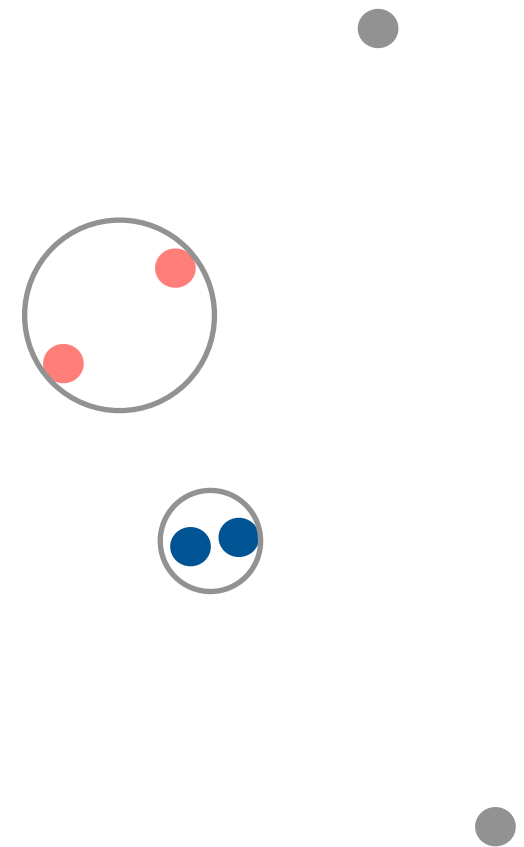
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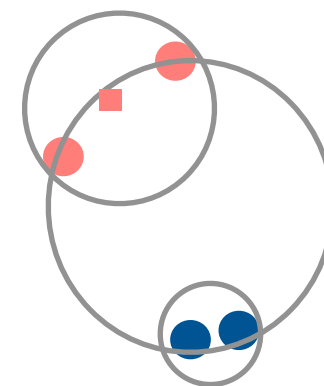
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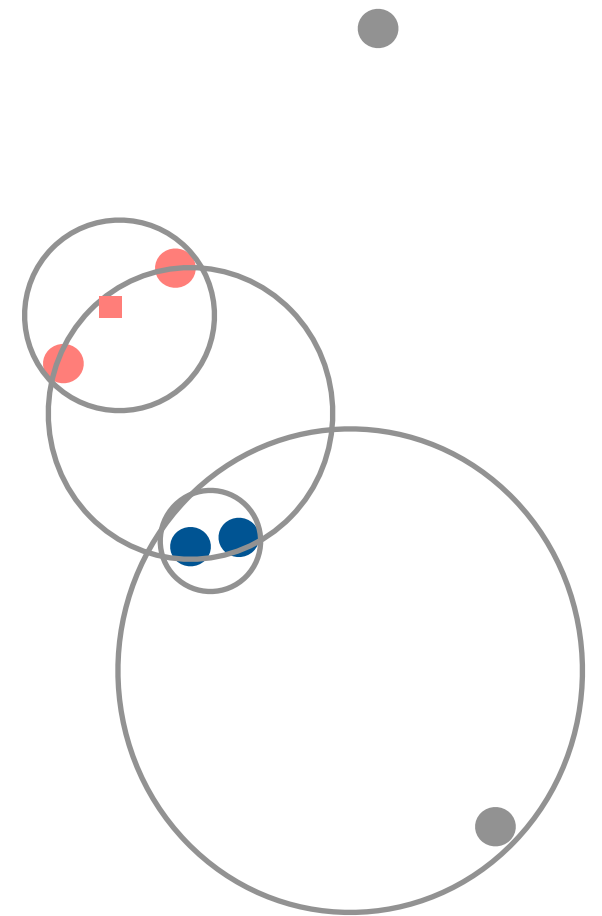
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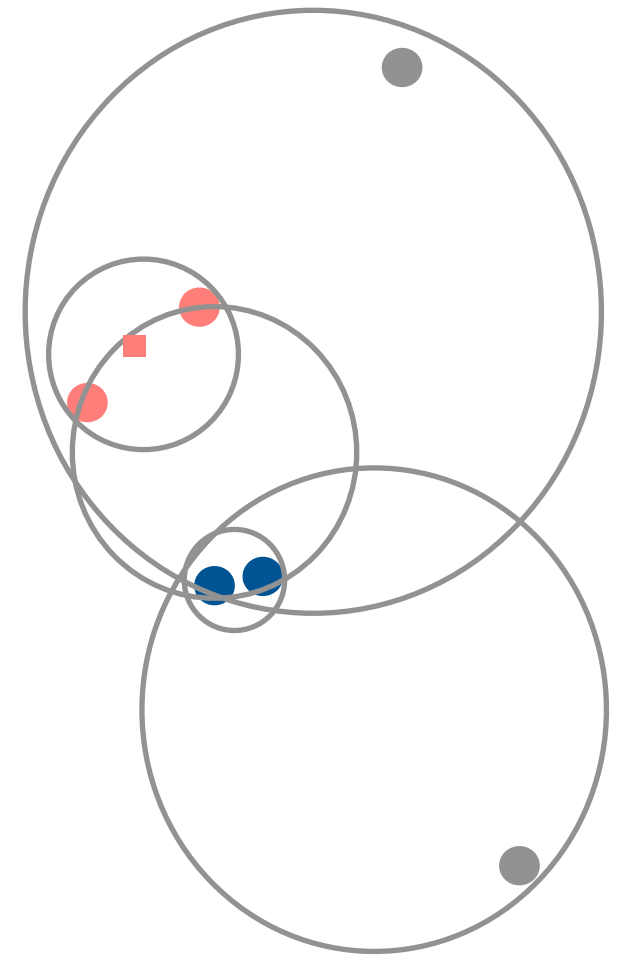
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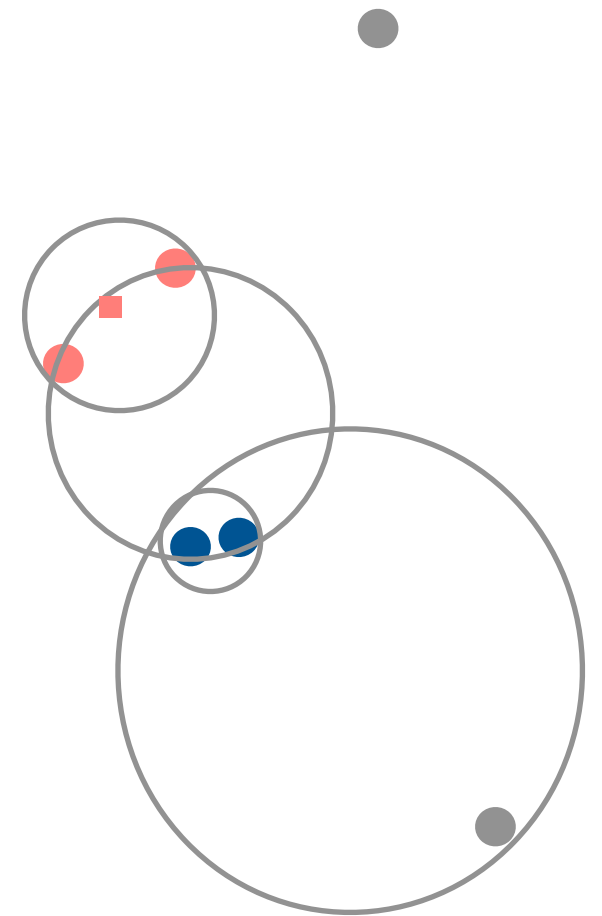
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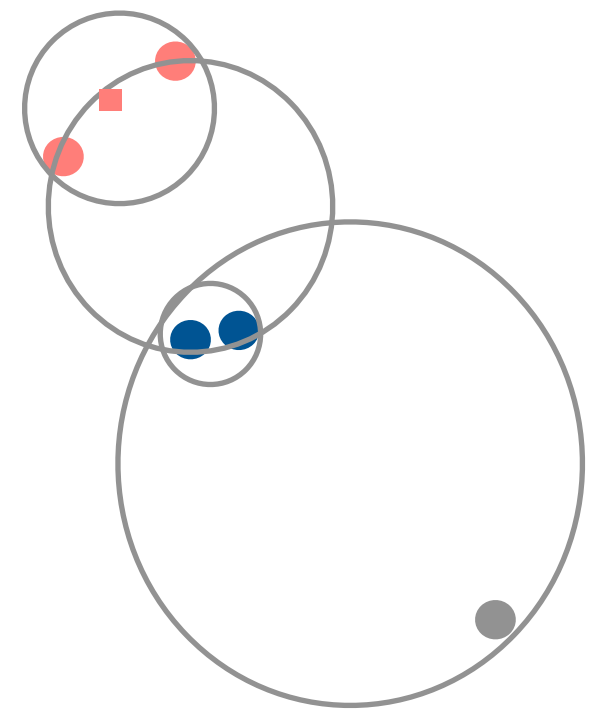
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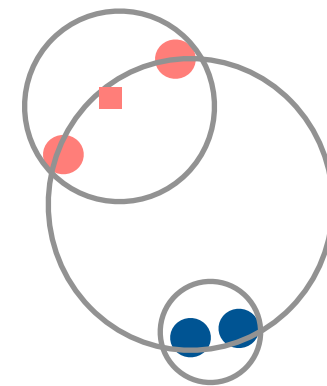


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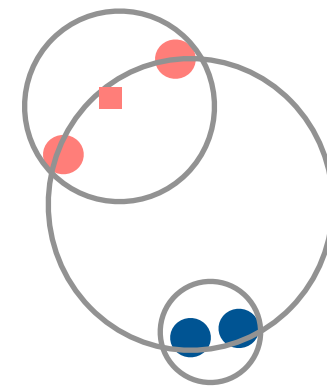
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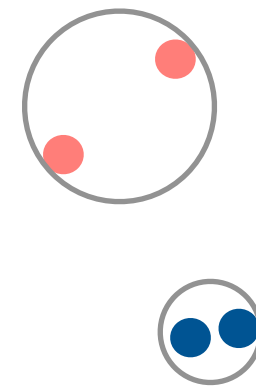
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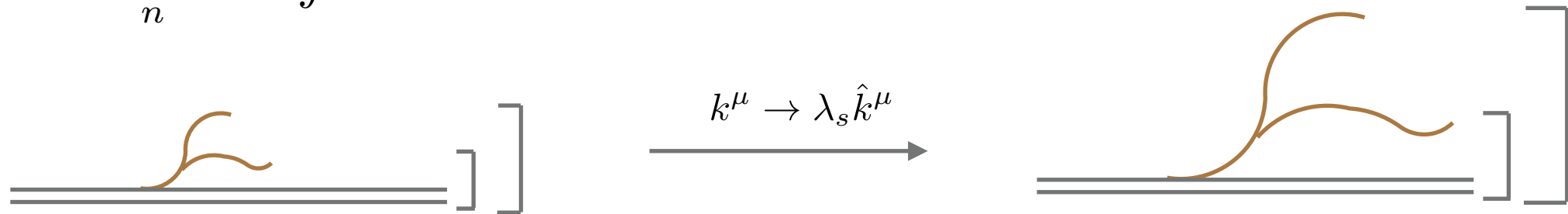
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$$S(z_{\text{cut}}) = \sum_n \mu^{2n\epsilon} \int d\Pi_n |\mathcal{M}_n|^2 \Theta_{\text{SD}}$$



C/A clustering

$$\Theta(\theta_{ij} - \min[\theta_i, \theta_j])$$



$$\Theta(\hat{\theta}_{ij} - \min[\hat{\theta}_i, \hat{\theta}_j])$$

soft drop

$$\Theta(z_{\text{cut}} \theta^\beta - \sum_i z_i)$$



$$\Theta(\theta^\beta - \sum_i z_i)$$

SCET matrix element

$$d\Pi_n |\mathcal{M}_n|^2$$



$$z_{\text{cut}}^{-2n\epsilon} d\Pi_n |\mathcal{M}_n|^2$$

$$S(z_{\text{cut}}) = \sum_n \mu^{2n\epsilon} (z_{\text{cut}})^{-2n\epsilon} \int d\Pi_n |\mathcal{M}_n|^2 \Theta_{\text{SD}}^{z_{\text{cut}}=1} \quad \text{depends on a single soft scale } Q z_{\text{cut}}.$$

By similar argument,

$$S_C(z_{\text{cut}} m^2) = \sum_n \mu^{2n\epsilon} \left(z_{\text{cut}}^{\frac{1}{2+\beta}} (m^2)^{\frac{1+\beta}{2+\beta}} \right)^{-2n\epsilon} \frac{Q^2}{m^2} \int d\Pi_n |\mathcal{M}_n|^2 \Theta_{SD}^{z_{\text{cut}}=1} \delta_{\frac{m^2}{Q^2}=1}$$

ACHIEVING NNLL ACCURACY

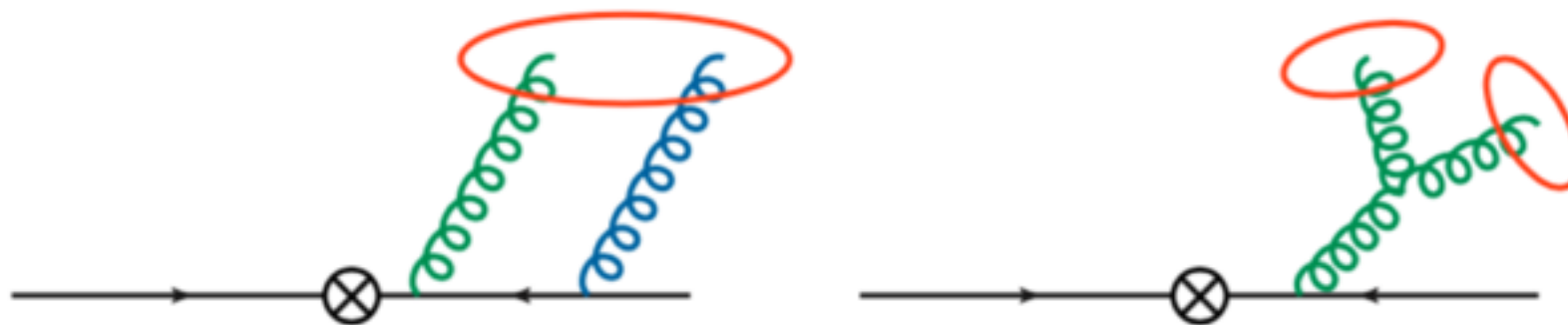
- missing ingredient: two-loop non cusp anomalous dimension of the soft-collinear function.

$$\frac{d^2\sigma}{de_{2,L}^{(\alpha)} de_{2,R}^{(\alpha)}} = H(Q^2) S(z_{\text{cut}}) \left[S_{C,L}(z_{\text{cut}} e_{2,L}^{(\alpha)}) \otimes J_L(e_{2,L}^{(\alpha)}) \right] \left[S_{C,R}(z_{\text{cut}} e_{2,R}^{(\alpha)}) \otimes J_R(e_{2,R}^{(\alpha)}) \right]$$

$$0 = \gamma_H + \gamma_S + 2\gamma_J + 2\gamma_{S_C}$$

- calculate the soft anomalous dimension from two-loop hemisphere soft function at e^+e^- :
 - two wilson-line fixed angle matrix element
 - simple phase-space constraint
 - for $\beta = 0$, related to calculations from literature [[Manteuffel, Schabinger, Zhu, 1309.3560](#)];
 - for $\beta = 1$, we extracted the anomalous dimension from EVENT2

GROOMING CALCULATION AT NNLO



TWO-LOOP HEMISPHERE SOFT FUNCTION

can be related to the two-loop soft-function with a global energy veto

[Manteuffel, Schabinger, Zhu, 1309.3560]

$$S_{\text{veto}} = \int d\Pi_2 |\mathcal{M}(k_1, k_2)|^2 \Theta\left(\frac{Q}{2} z_{\text{cut}} - k_1^0 - k_2^0\right)$$

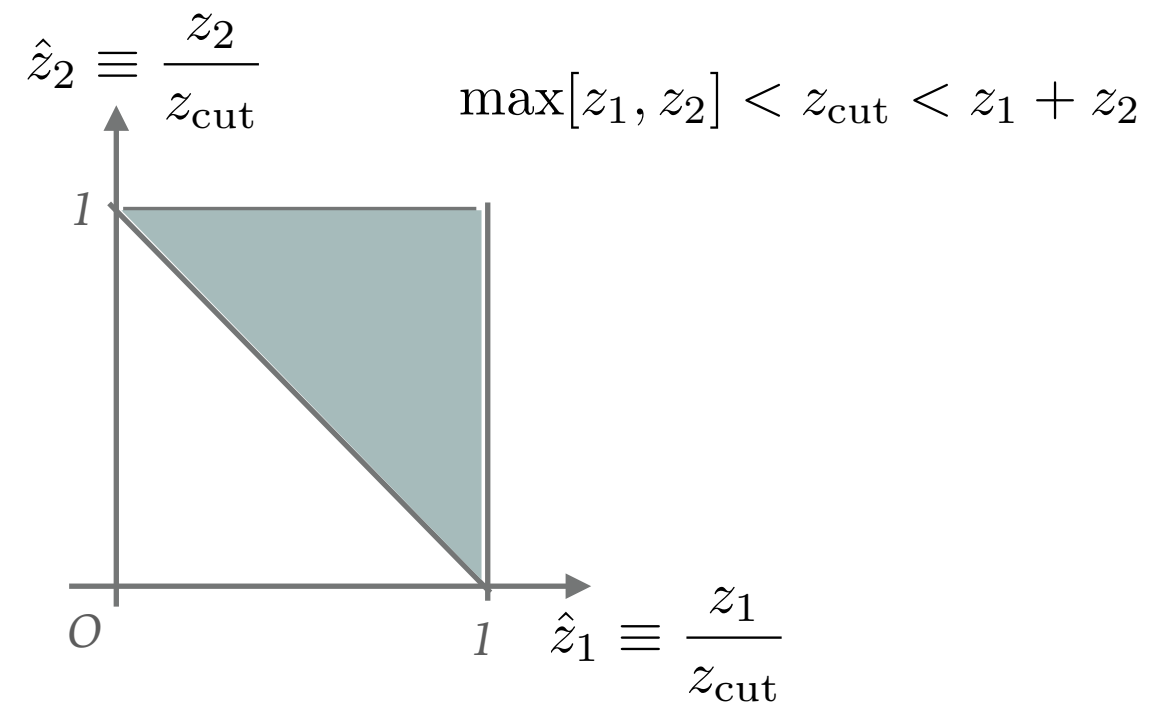
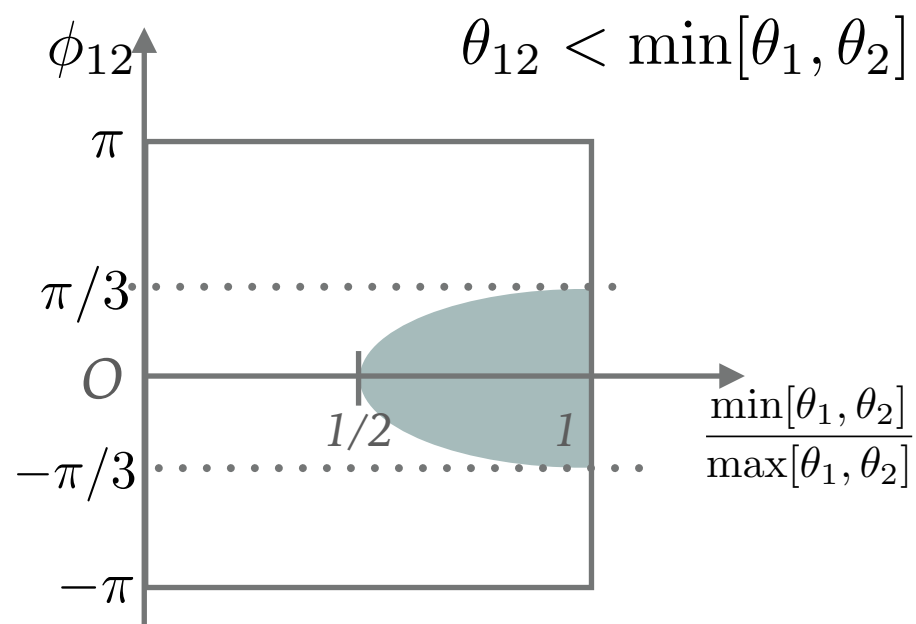
Two-loop non-cusp anomalous dimension

$$\boxed{\gamma_{\text{SD}} = \gamma_{\text{veto}} + \gamma_{\text{Clust}}^{\text{alg.}}}$$

- Abelian contribution: purely comes from clustering effect that violates abelian exponentiation
- Non-abelian contribution:
only need to consider the difference between soft drop and energy veto phase-space constraints, which is a highly restricted phase-space region.

INDEPENDENT EMISSIONS

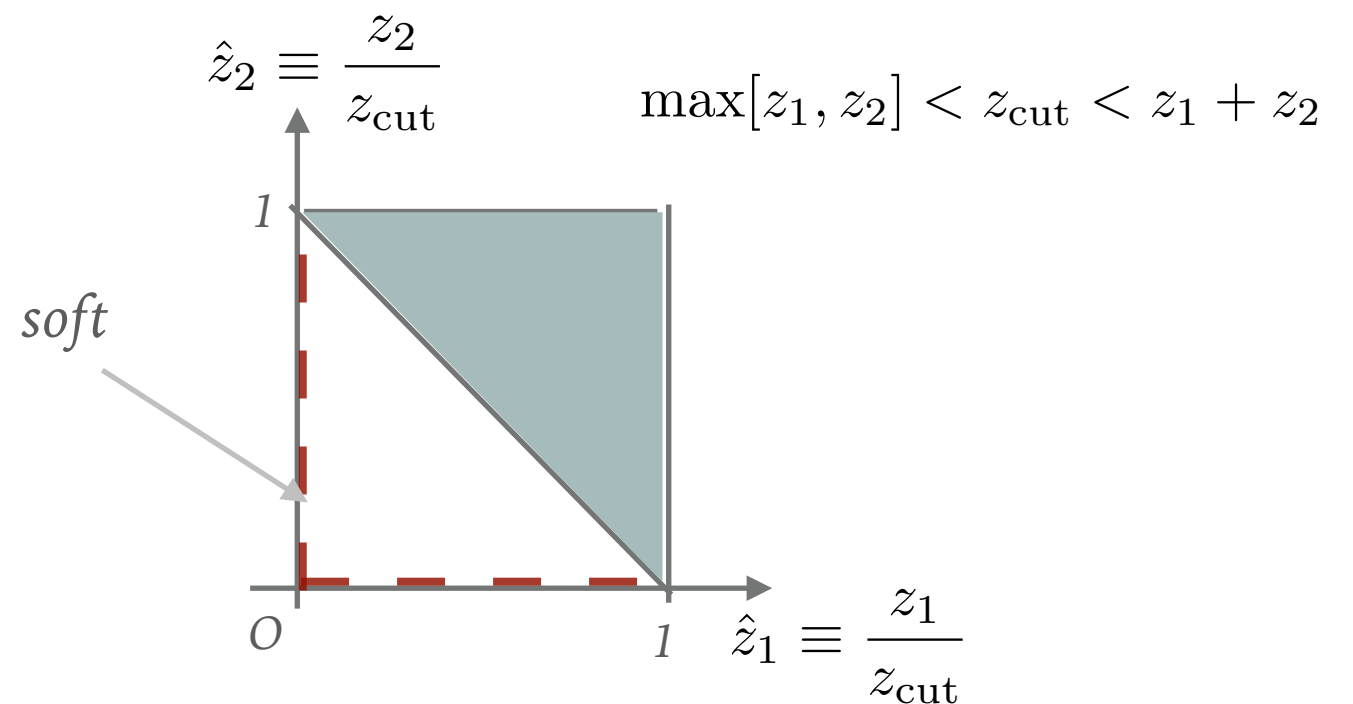
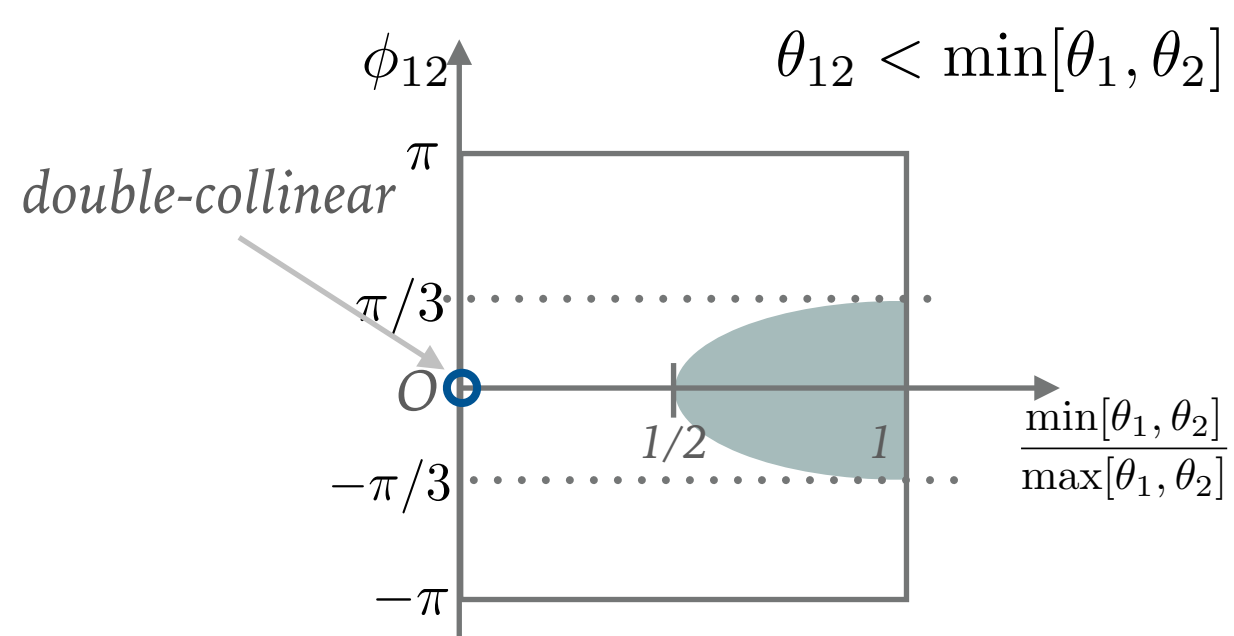
$$\begin{aligned}
 S(z_{\text{cut}})|_{A, \alpha_s^2} &= \frac{1}{2!} [S^{1\text{-loop}}(z_{\text{cut}})]^2 + \left[\text{diagram with two emissions and a clustering correction bracket} \right] \text{clustering correction} \\
 &+ \frac{1}{2!} \int [d^d k_1]_+ [d^d k_2]_+ |\mathcal{M}(k_1)|^2 |\mathcal{M}(k_2)|^2 \left[\Theta_{\text{SD}} - \Theta \left(z_{\text{cut}} \frac{Q}{2} - k_1^0 \right) \Theta \left(z_{\text{cut}} \frac{Q}{2} - k_2^0 \right) \right]
 \end{aligned}$$



$$\gamma_{C/A}^{A, \alpha_s^2} = \left(\frac{\alpha_s}{4\pi} \right)^2 34.01 C_F^2$$

INDEPENDENT EMISSIONS

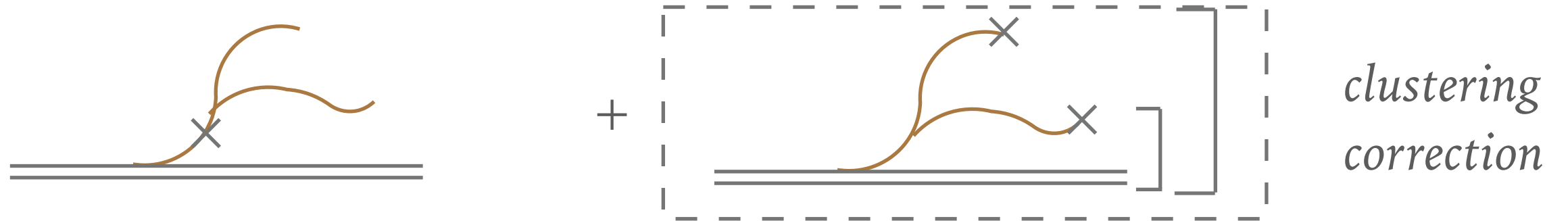
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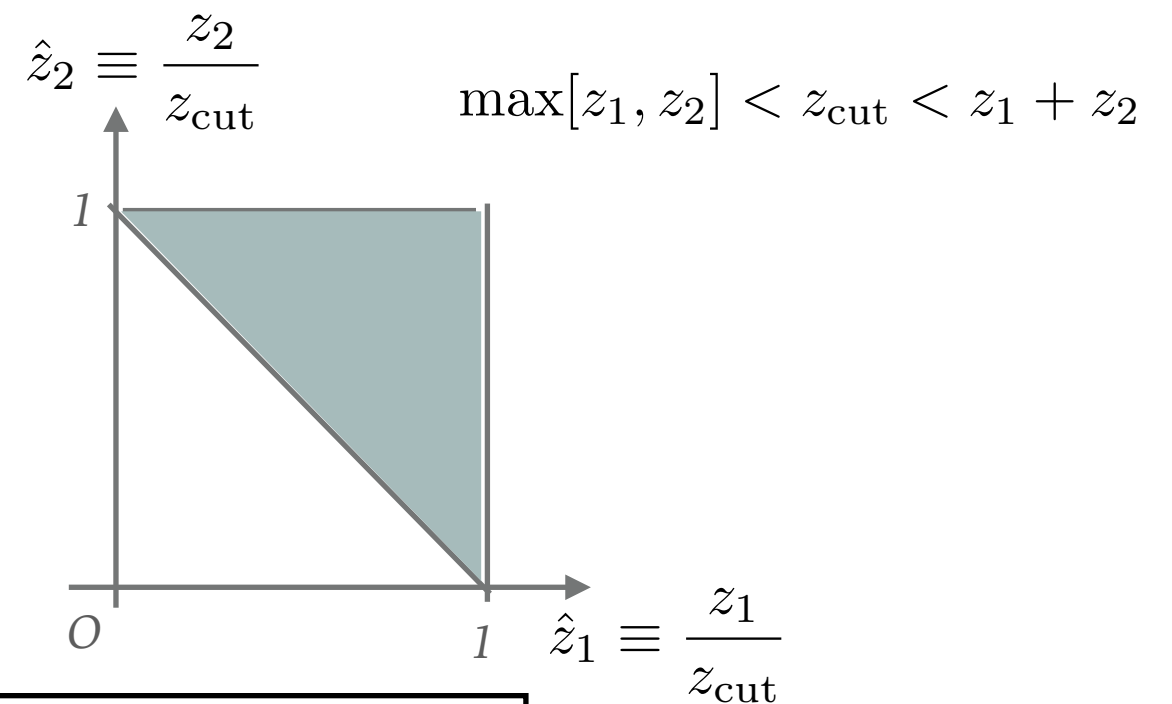
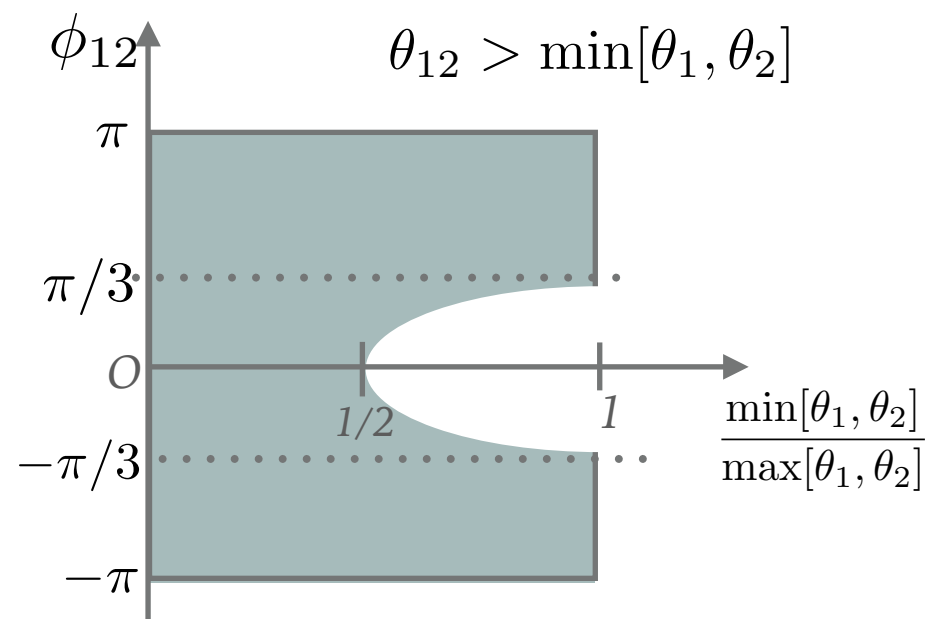
→

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CORRELATED EMISSIONS

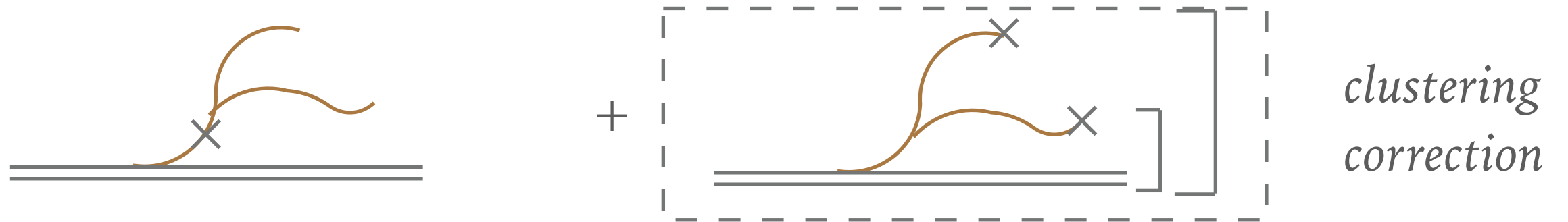


$$S(z_{\text{cut}})|_{\text{n-A}, \alpha_s^2} = S_{\text{veto}}|_{\text{n-A}, \alpha_s^2} + \int [d^d k_1]_+ [d^d k_2]_+ |\mathcal{M}_{\text{n-A}}(k_1, k_2)|^2 [\Theta_{\text{SD}} - \Theta_{\text{veto}}]$$

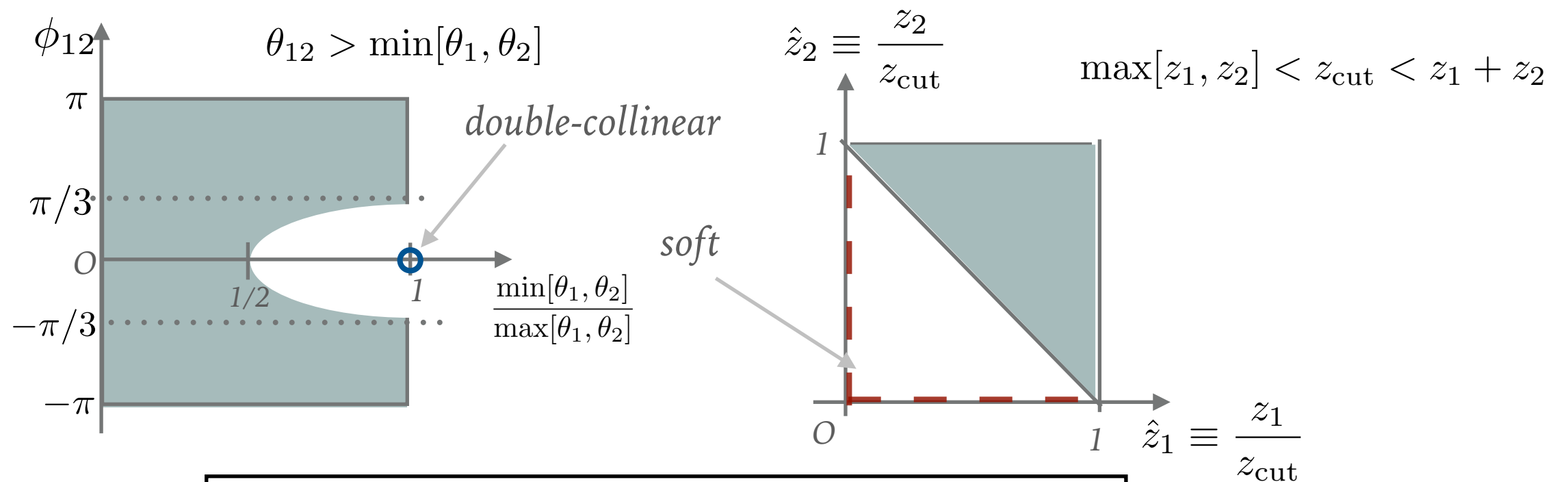


$$\gamma_{\text{C/A}}^{\text{n-A}, \alpha_s^2} = \left(\frac{\alpha_s}{4\pi}\right)^2 C_F [-9.31 C_A - 14.04 n_f T_R]$$

CORRELATED EMISSIONS



$$S(z_{\text{cut}})|_{\text{n-A}, \alpha_s^2} = S_{\text{veto}}|_{\text{n-A}, \alpha_s^2} + \int [d^d k_1]_+ [d^d k_2]_+ |\mathcal{M}_{\text{n-A}}(k_1, k_2)|^2 [\Theta_{\text{SD}} - \Theta_{\text{veto}}]$$

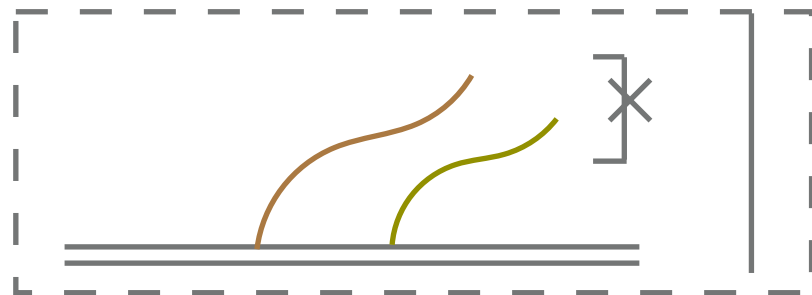


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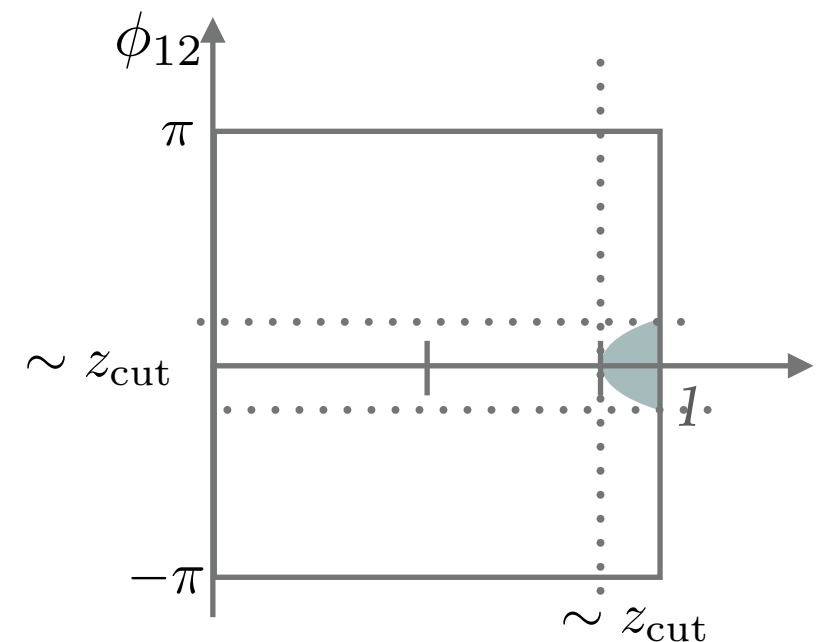
RECLUSTERING WITH ANTI-KT ALGORITHM

► Clustering metric

$$d_{ij} = \min[p_{T,i}^{-2}, p_{T,j}^{-2}] R_{ij}^2$$



*clustering
correction*



cluster soft emissions i,j before clustering them with the jet only if

$$\min[\hat{z}_i^{-2}, \hat{z}_j^{-2}] \theta_{ij}^2 < z_{\text{cut}}^2 \min[\theta_i^2, \theta_j^2]$$

The probability of clustering independent emissions together is suppressed by z_{cut}^2 .

► No abelian clustering correction at leading power.

.....

The non-abelian clustering logs can be obtained by computing the correction to S_{veto}

$$\gamma_{\text{ak}_T} = -8 \left(\frac{\alpha_s}{4\pi} \right)^2 C_F \left\{ \left[\left(\frac{131}{9} - \frac{4}{3}\pi^2 - \frac{44}{3}\log 2 \right) C_A + \left(-\frac{46}{9} + \frac{16}{3}\log 2 \right) n_f T_R \right] \log z_{\text{cut}} \right. \\ \left. + \left(-\frac{269}{6} + \frac{7}{2}\zeta_3 + \frac{274}{9}\log 2 + \frac{11\pi^2}{9} + \frac{44}{3}\log^2 2 \right) C_A \right. \\ \left. + \left(\frac{53}{3} - \frac{4\pi^2}{9} - \frac{116}{9}\log 2 - \frac{16}{3}\log^2 2 \right) n_f T_R \right\}$$

*identical to the leading
clustering log in jet-
veto calculation*

*[Tackmann, Walsh, Zuberi
1206.4312][Banfi, Salam,
Zanderighi 1203.5773]*

$$C_2(R) = 2C_A \left[\left(1 - \frac{8\pi^2}{3} \right) C_A + \left(\frac{23}{3} - 8\ln 2 \right) \beta_0 \right] \ln R^2 \\ + 15.62C_A^2 - 9.17C_A\beta_0 + C_2^{\text{Rsub}}(R)$$

RELATING TO JET VETO CALCULATION

- Clustering metric at the LHC with anti-kT algorithm: cluster two soft emissions into one jet if $d_{ij}^{\text{eff}} < 1$

$$d_{ij}^{\text{eff},pp} \equiv \min[\hat{z}_i^{-2}, \hat{z}_j^{-2}] \frac{R_{ij}^2}{R^2}$$

boost-invariant geometrical separation phase space

- Effective clustering metric at e+e-: combine two soft emissions (i,j) before combining them with hard jet core if $d_{ij}^{\text{eff}} < 1$

$$d_{ij}^{\text{eff},e^+e^-} \equiv \min[\hat{z}_i^{-2}, \hat{z}_j^{-2}] \frac{\theta_{ij}^2}{z_{\text{cut}}^2 \min[\theta_i^2, \theta_j^2]}$$

boost-invariant angle along the jet axis

- In large rapidity region (where $\log R$ arises), set $R=z_{\text{cut}}$,

$$d_{ij}^{\text{eff},pp} \xrightarrow[y_i, y_j \rightarrow \infty]{i \parallel j} d_{ij}^{\text{eff},e^+e^-}$$

SUMMARY

- grooming parameter and measuring regime: $z_{\text{cut}} \sim 0.1 \quad m^2/Q^2 \ll z_{\text{cut}}$
- preferred reclustering algorithm: C/A

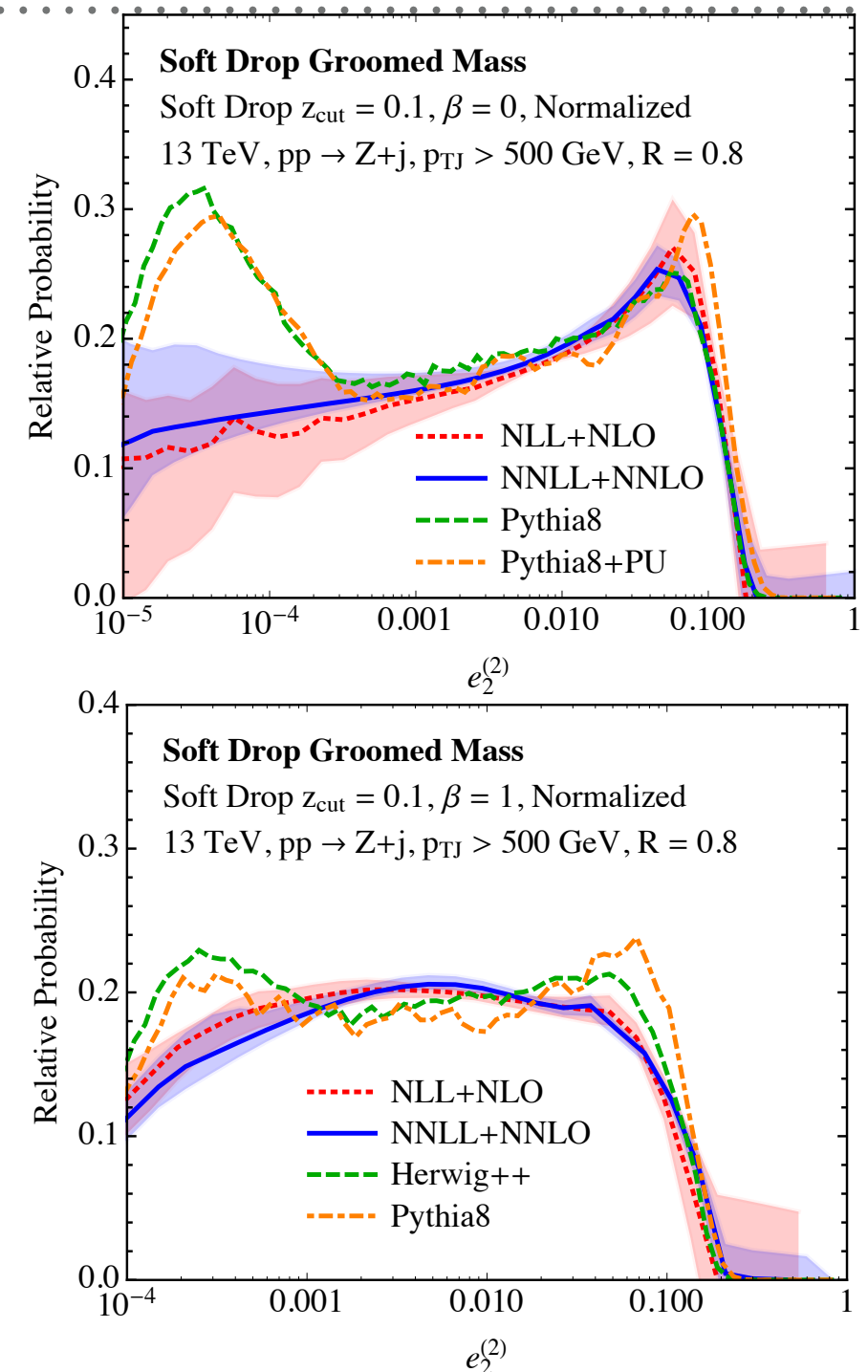
With the above choices, we obtained :

- All-order factorization at leading power of pile-up insensitive, process independent jet observable.
- First calculation done at NNLL accuracy with no non-global logs of jet substructure observable

SUMMARY

What can be done in the future:

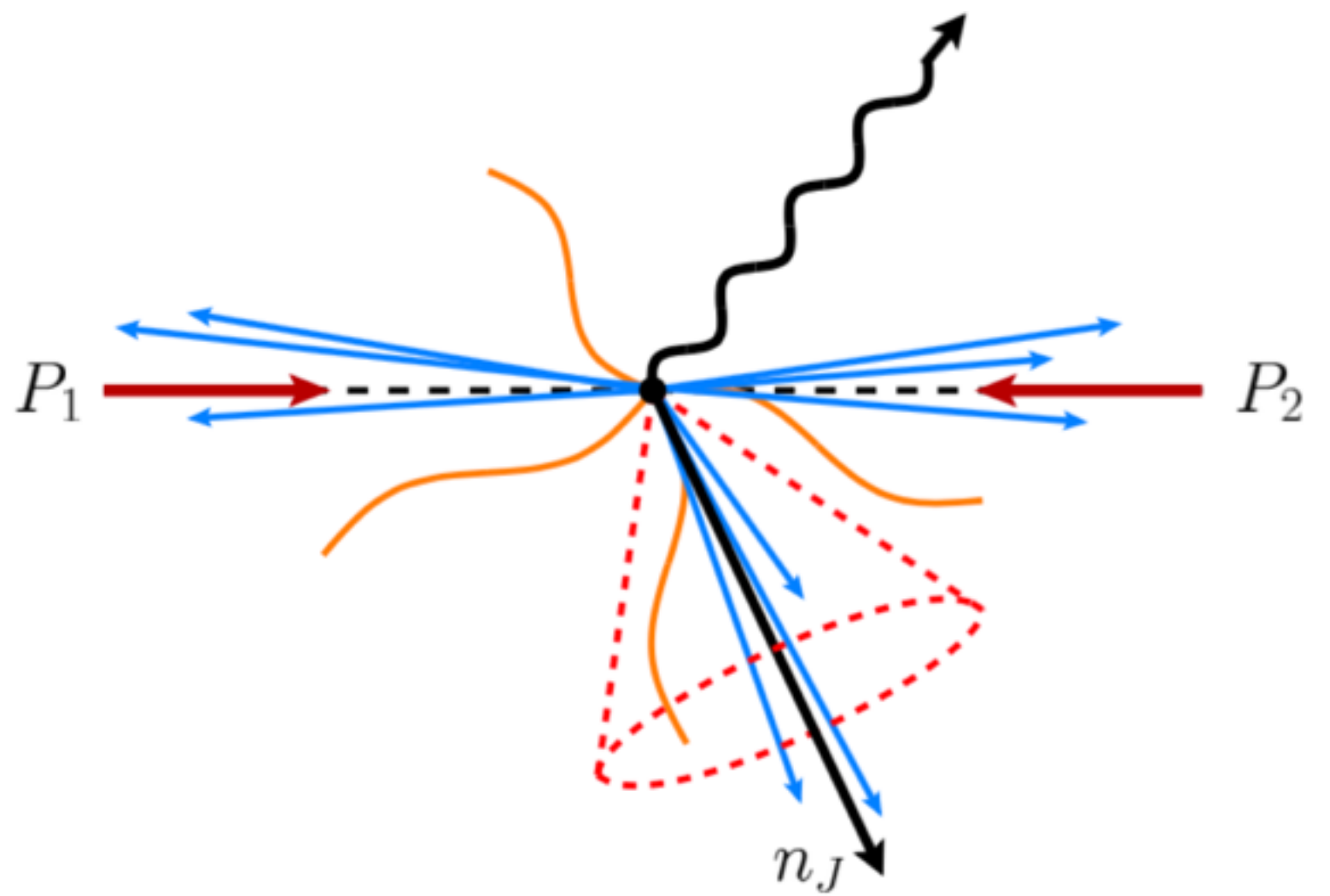
- sub-leading power factorization theorem needed to predict Z_{cut} power-corrections
- Soft-drop observables that extract more information on jet substructure: D2 etc ;
- precision grooming calculation:
going to N3LO only requires two-Wilson line soft function calculation



prediction at pp collision

(to be continued...)

THANKS!



COMPLICATIONS AT HADRON COLLIDERS

Factorization formula (in general , at leading z_{cut})

$$\frac{d\sigma_{\text{resum}}}{de_2^{(\alpha)}} = \sum_{k=q,\bar{q},g} D_k(p_T^{\text{min}}, \eta_{\text{max}}, z_{\text{cut}}, R) S_{C,k}(z_{\text{cut}} e_2^{(\alpha)}) \otimes J_k(e_2^{(\alpha)})$$

CLUSTERING LOGS (ANTI-KT ALGORITHM)

- study the clustering effect by looking at a different algorithm: anti-kT.
- Two soft emissions are clustered together before clustering with hard-collinear emission if

$$\min[z_i^{-1}, z_j^{-1}]\theta_{ij} < \min[\theta_i, \theta_j]$$

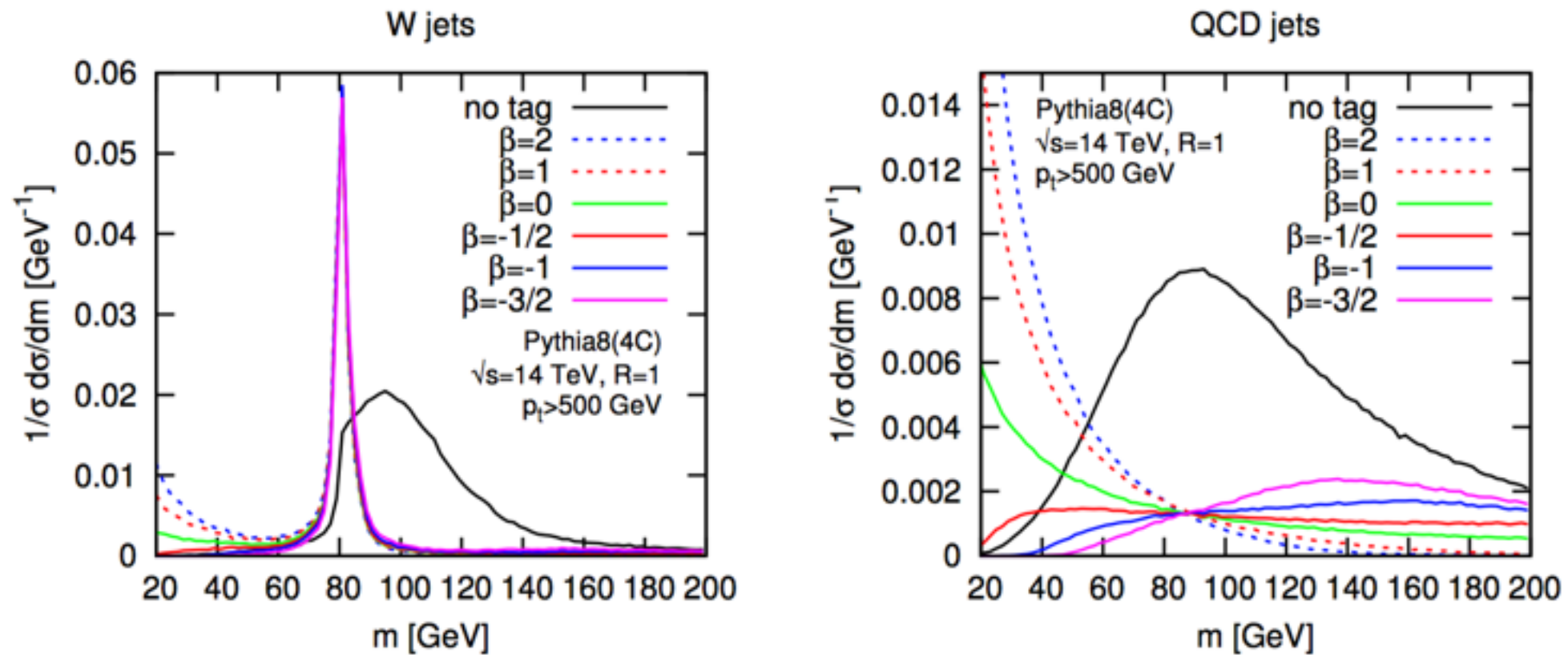
This inequality is not preserved after rescaling the soft momenta.

$$\begin{array}{ccc}
 \begin{array}{c} \text{Diagram of two emissions clustered together} \end{array} & \xrightarrow{k^\mu \rightarrow \lambda_s \hat{k}^\mu} & \begin{array}{c} \text{Diagram of two emissions clustered together} \end{array} \\
 \Theta(\min[z_i^p, z_j^p]\theta_{ij} - \min[\theta_i, \theta_j]) & \longrightarrow & \Theta(\min[\hat{z}_i^p, \hat{z}_j^p]\hat{\theta}_{ij} - z_{\text{cut}}^{|p|}\min[\hat{\theta}_i, \hat{\theta}_j]) \\
 \Theta_{\text{SD}}(\{k_i\}, z_{\text{cut}}) & \not\longrightarrow & \Theta_{\text{SD}}(\{\hat{k}_i\}, 1)
 \end{array}$$

- If two independent emissions are clustered into the same jet then the rate is suppressed by the area of the jet which is $O(z_{\text{cut}})$. At leading power, this implies the absence of abelian clustering effect.

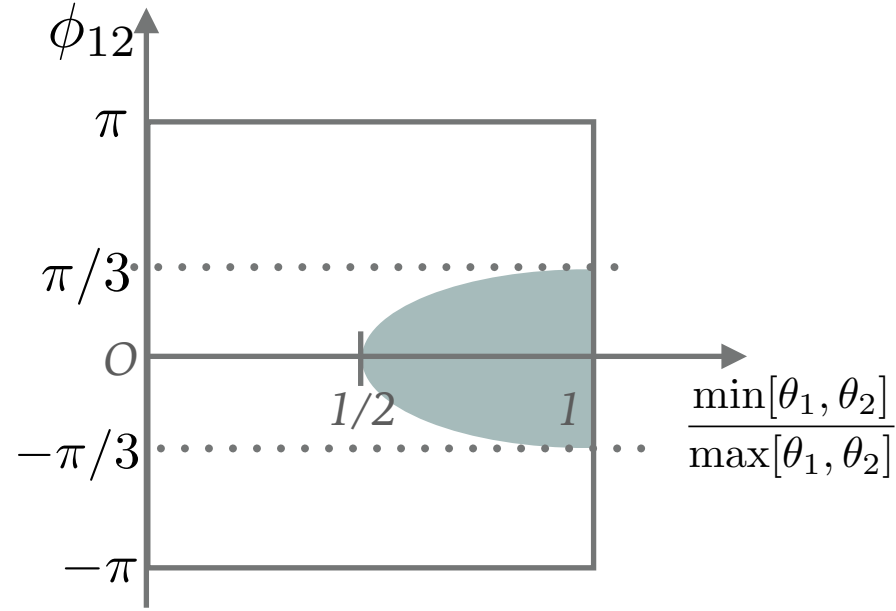
.....

► Experimental control of signal over background

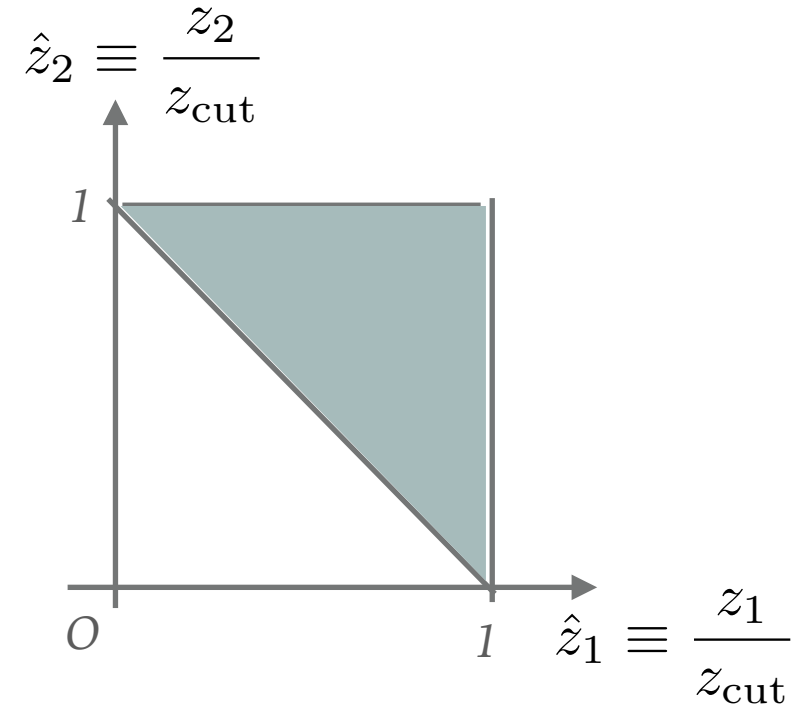


Performance of soft drop as a boosted W tagger : mass distribution of signal (left) and background (right) jets before and after soft drop

[Larkoski, Marzani, Soyez, Thaler 1402.2657]



$$\Theta_{\text{veto}} = \Theta(\Lambda - k_1^0 - k_2^0)$$



$$\begin{aligned} \Theta_{\text{SD}} - \Theta_{\text{veto}} &= \left\{ \Theta(\eta_1 \eta_2) [1 - \Theta(\theta_{1J} - \theta_{12}) \Theta(\theta_{2J} - \theta_{12})] + \Theta(-\eta_1 \eta_2) \right\} \\ &\quad \times \Theta\left(z_{\text{cut}} \frac{Q}{2} - k_1^0\right) \Theta\left(z_{\text{cut}} \frac{Q}{2} - k_2^0\right) \Theta\left(k_1^0 + k_2^0 - z_{\text{cut}} \frac{Q}{2}\right) \\ &= -\Theta(\eta_1 \eta_2) \Theta(\theta_{1J} - \theta_{12}) \Theta(\theta_{2J} - \theta_{12}) \\ &\quad \times \Theta\left(z_{\text{cut}} \frac{Q}{2} - k_1^0\right) \Theta\left(z_{\text{cut}} \frac{Q}{2} - k_2^0\right) \Theta\left(k_1^0 + k_2^0 - z_{\text{cut}} \frac{Q}{2}\right) \end{aligned}$$