

Towards a world-sheet description of DFT

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This talk is **based on** ::

- *T-duality revisited* [arXiv:1310.4194]
- *On T-duality transformations for the three-sphere* [arXiv:1408.1715]
- and work in progress with I. Bakas & D. Lüst [arXiv:1602.xxxxx]

1. motivation
2. t-duality
3. example

1. motivation

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motivation :: dualities

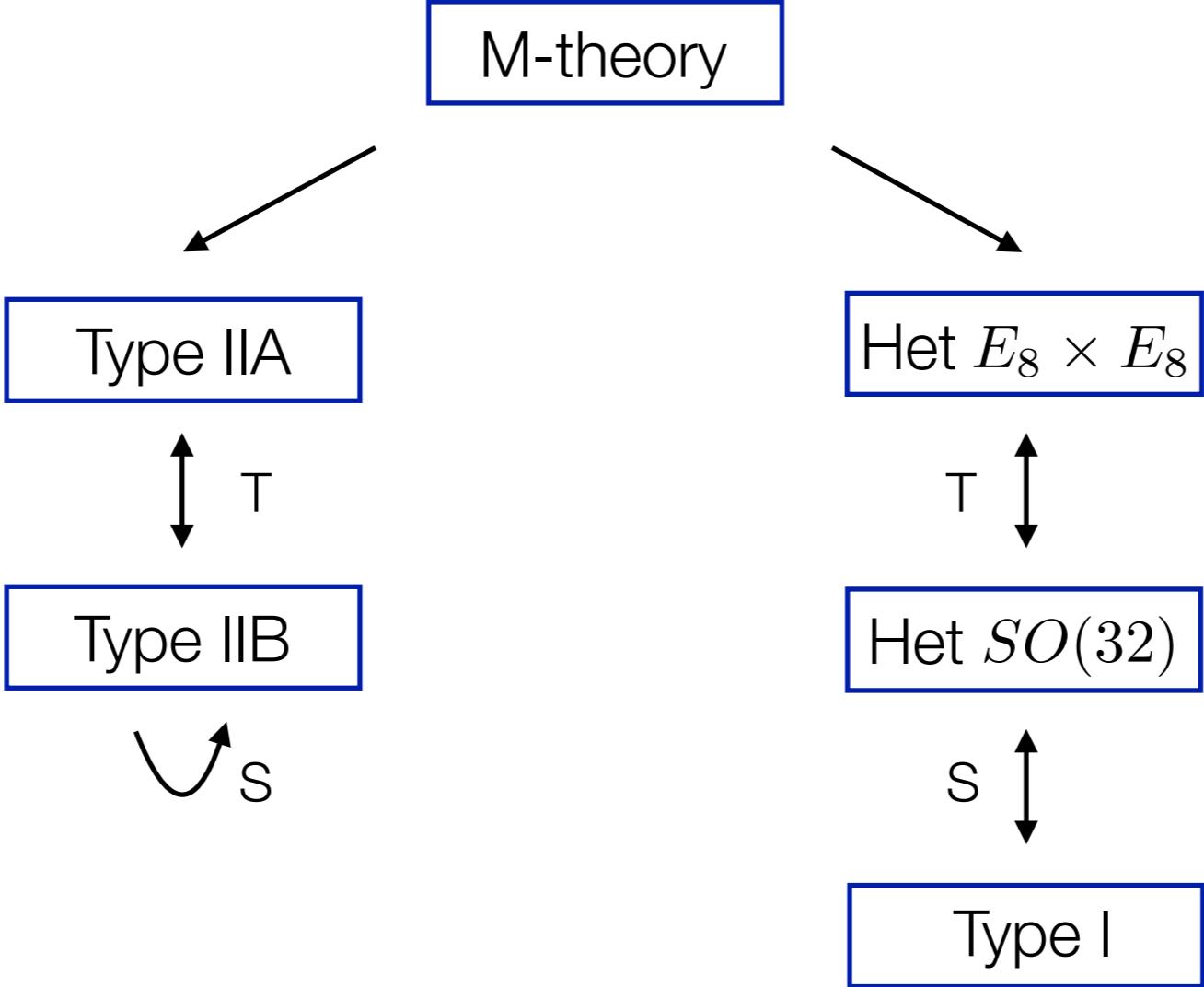
motivation :: dualities

Duality :: two **different theories** describe the **same physics**.

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Duality :: two **different theories** describe the **same physics**.

They play an important role in understanding string theory.



motivation :: main idea

Duality ::

Idea ::

Duality :: is a **symmetry** of the **equations of motion**.

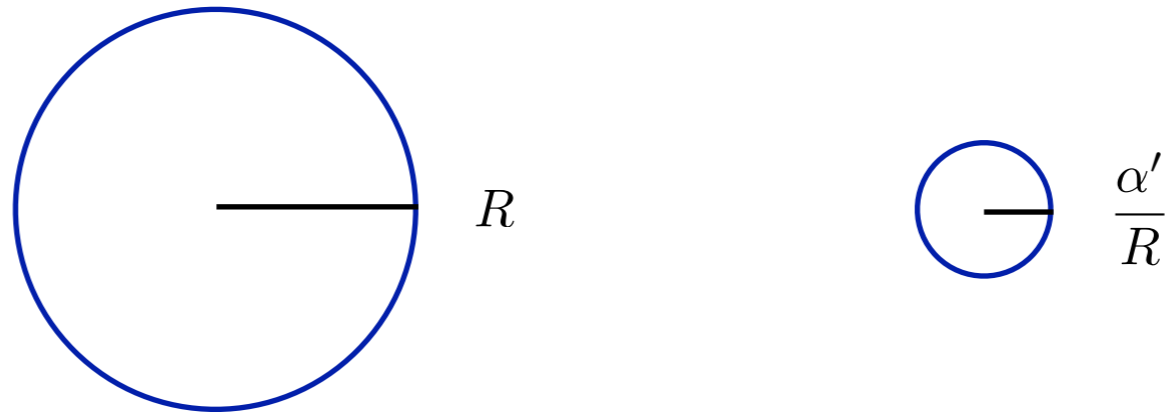
Idea ::

Duality :: is a **symmetry** of the **equations of motion**.

Idea :: make duality into a **symmetry** of an **action**.

1. motivation
2. t-duality
3. example

T-duality :: the physics of string theory **compactified** on two **circles**



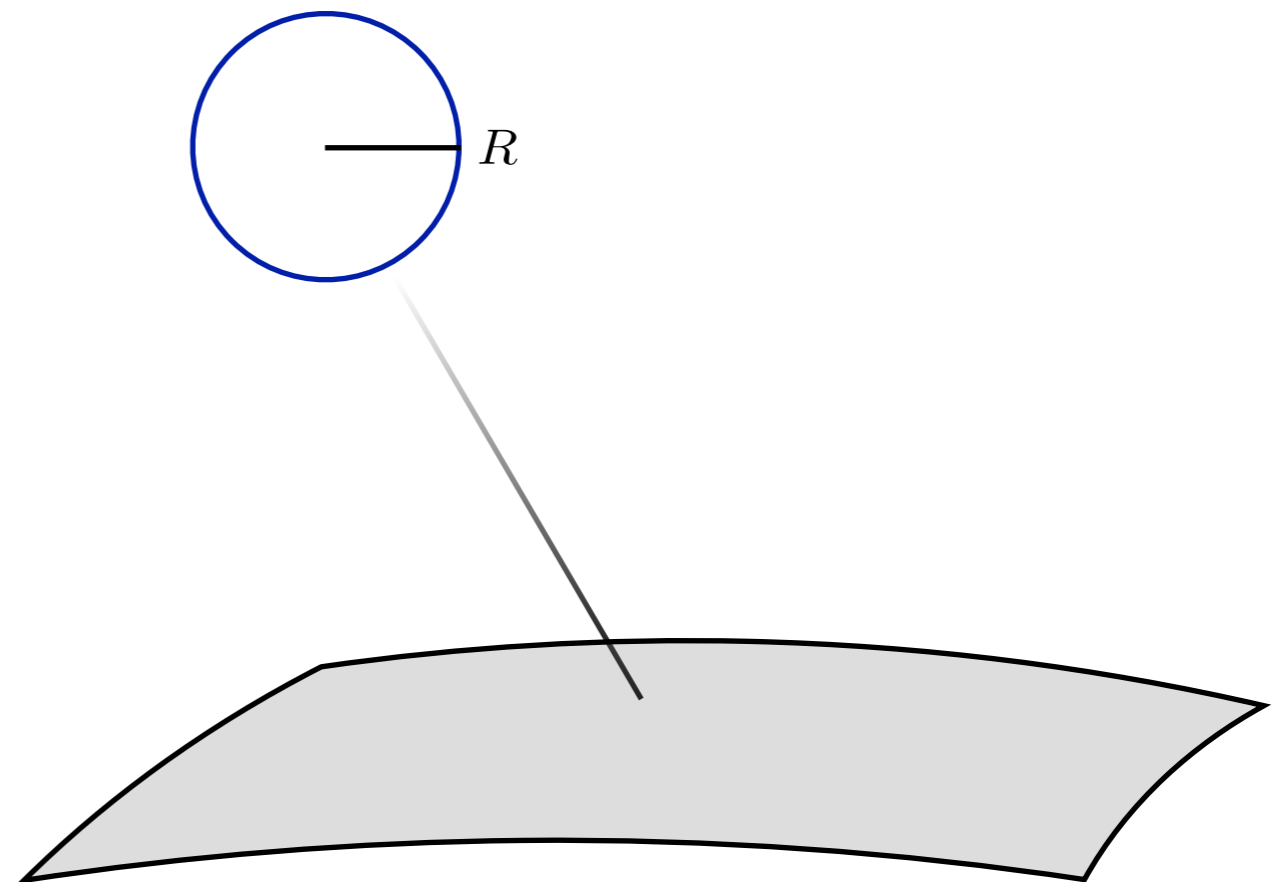
is indistinguishable.

t-duality :: doubled geometry and dft

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circle fibration

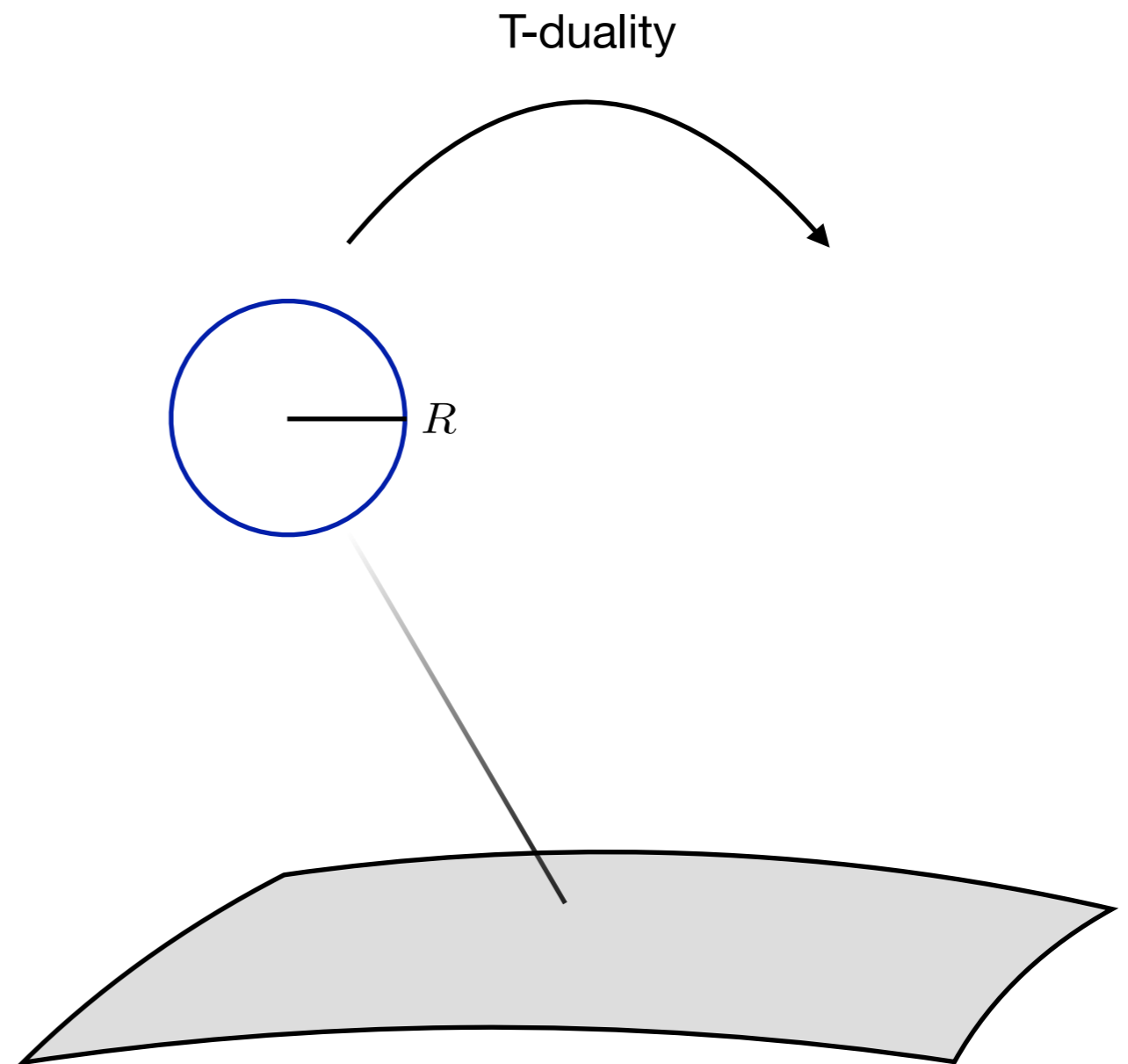
base manifold



t-duality :: doubled geometry and dft

circle fibration

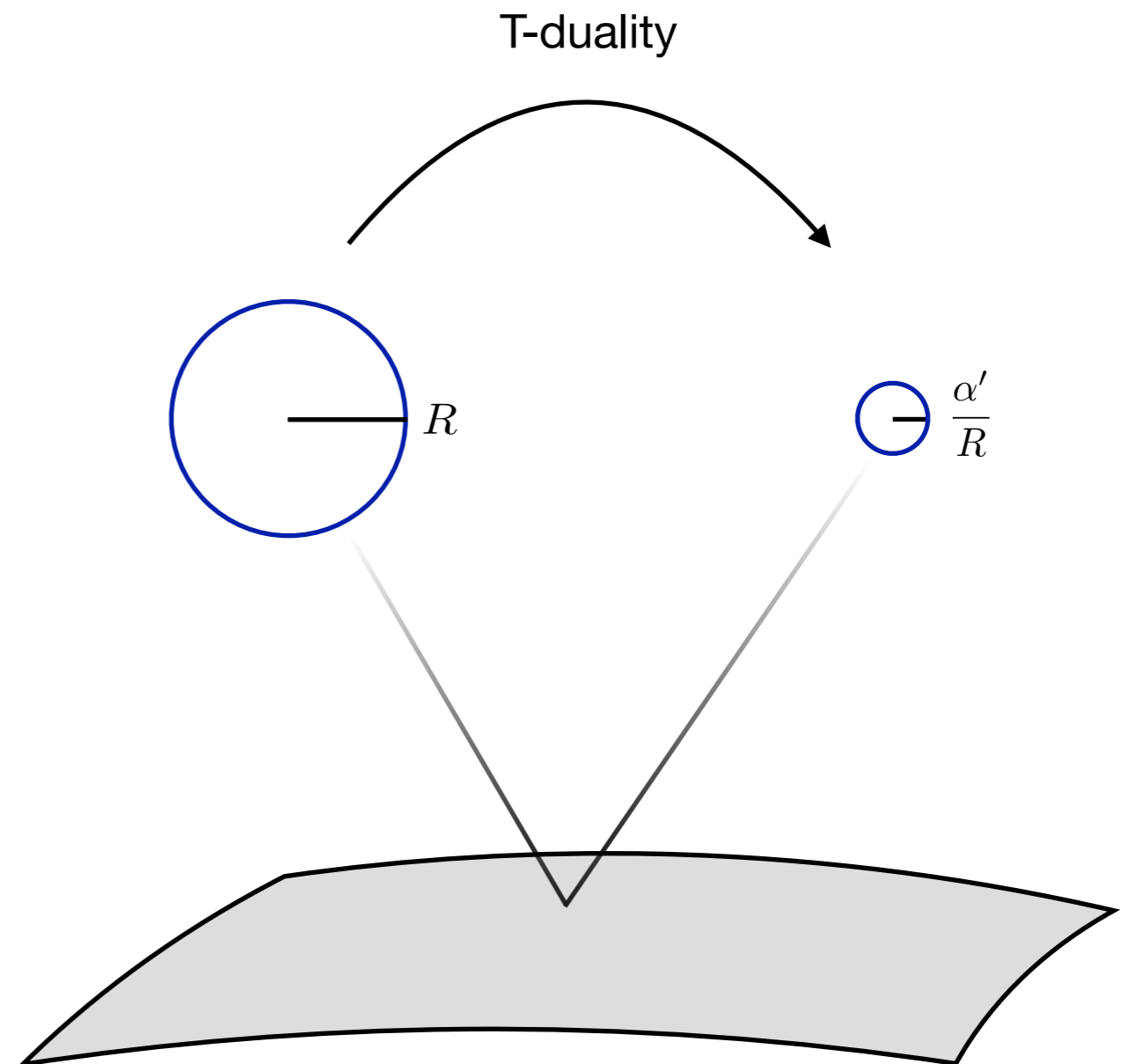
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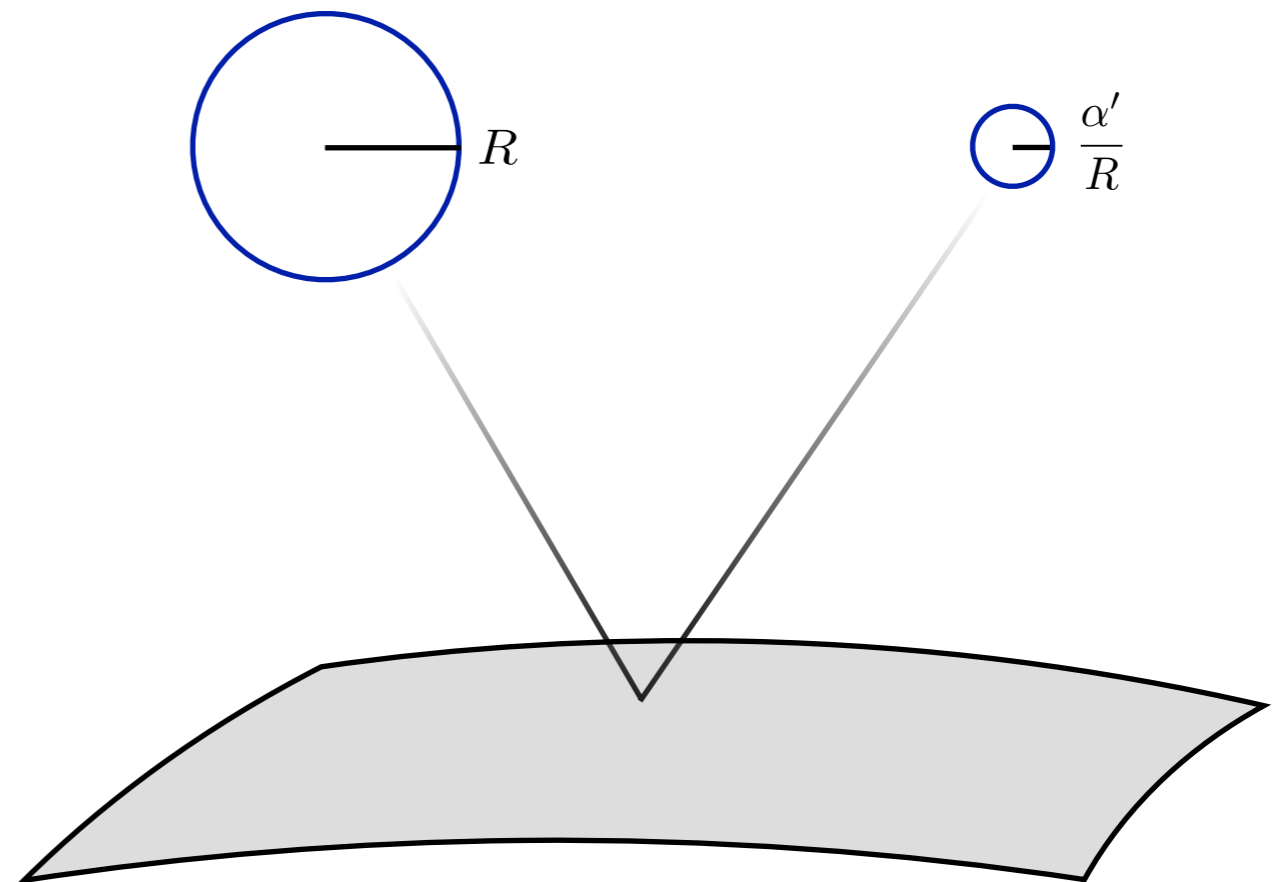
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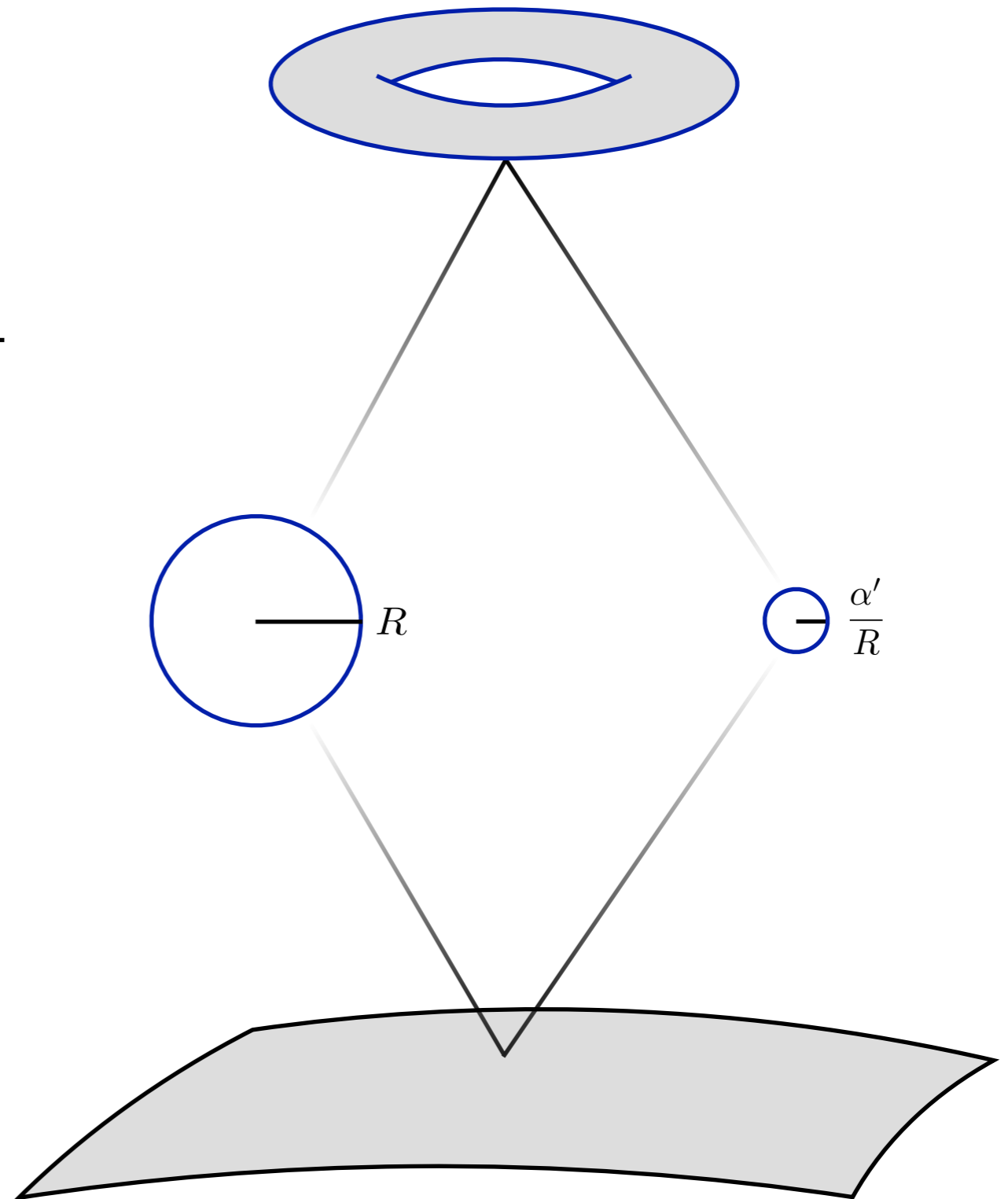


t-duality :: doubled geometry and dft

doubled geometry

circle fibration

base manifold



t-duality :: doubled geometry and dft

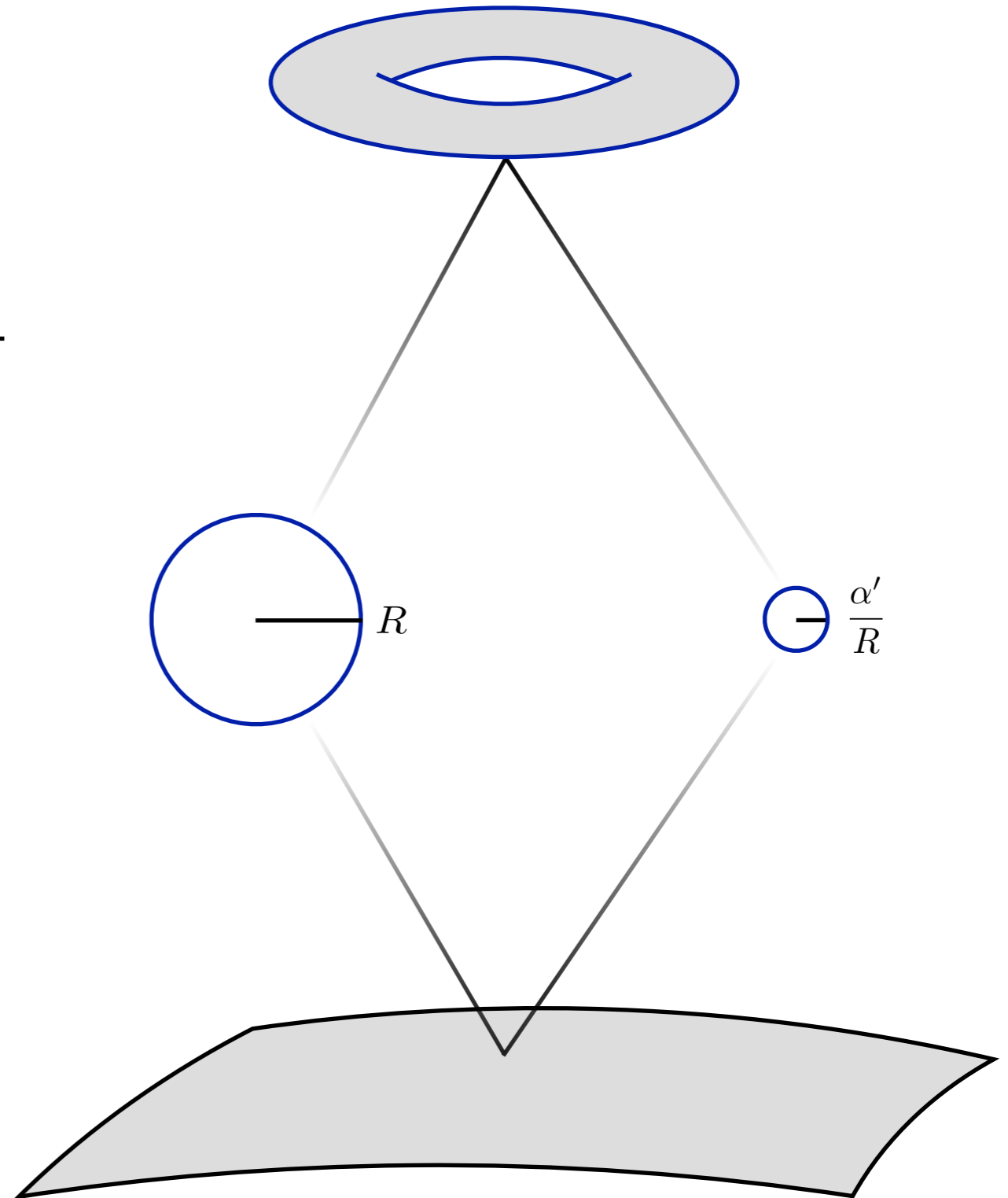
doubled geometry

symmetry

circle fibration

duality

base manifold



t-duality :: doubled geometry and dft

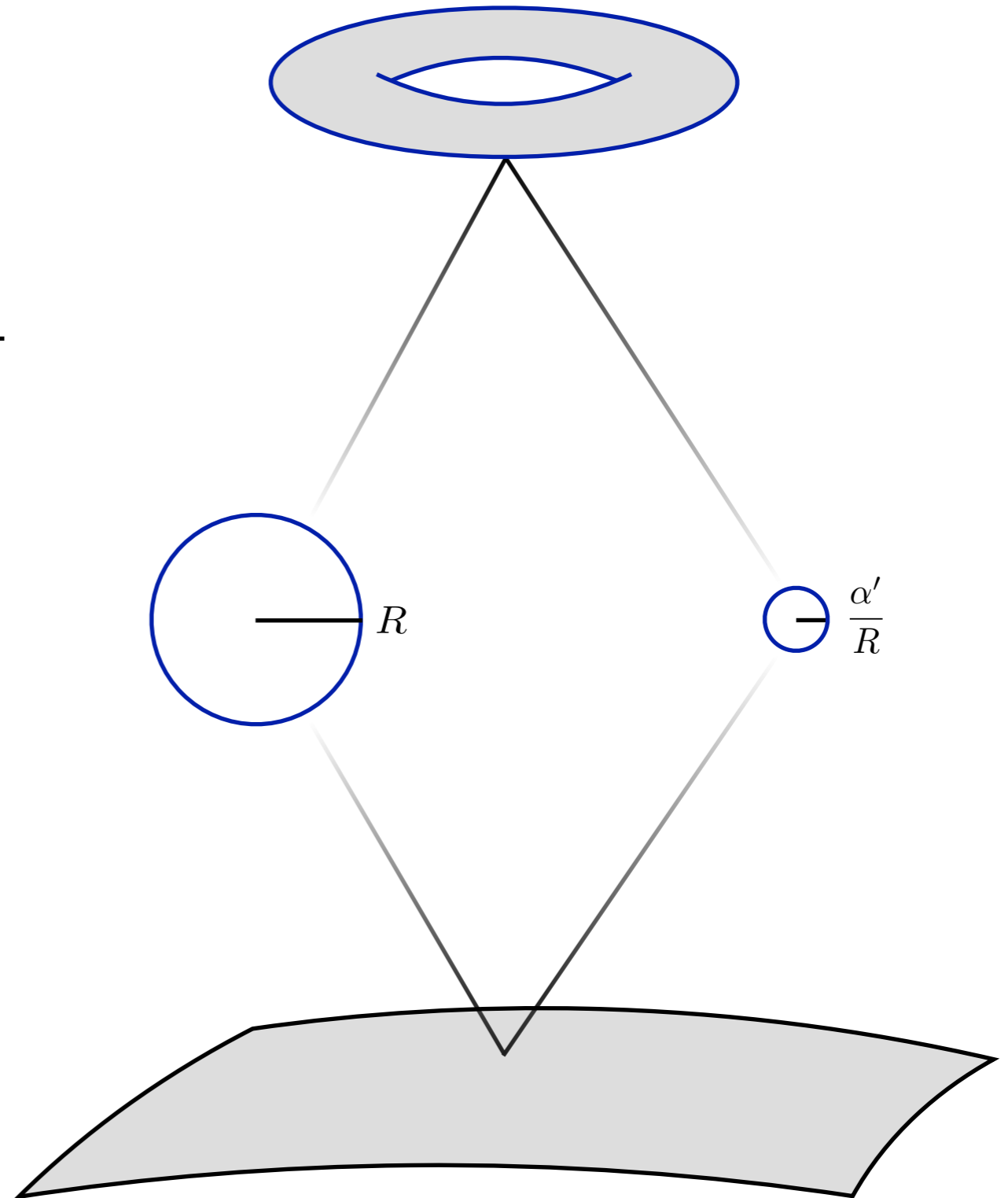
doubled geometry

symmetry

circle fibration

duality

base manifold



t-duality :: reduction via gauging

The reduction

doubled geometry



physical geometry

can be realized via a **gauging procedure**.

Main idea ::

1. Identify a global **symmetry** (circle) in the action.
2. **Gauge** this symmetry.
3. **Integrate-out** the gauge field.

t-duality :: reduction — step I

Consider the following sigma-model **action** for the **doubled geometry**

$$\mathcal{S} = -\frac{1}{4\pi\alpha'} \int_{\partial\Sigma} \left[G_{ij} dX^i \wedge \star dX^j + \alpha' R \Phi \star 1 \right] - \frac{i}{2\pi\alpha'} \int_{\Sigma} \frac{1}{3!} H_{ijk} dX^i \wedge dX^j \wedge dX^k .$$

This action is **invariant** under global transformations $\delta_{\epsilon} X^i = \epsilon k^i(X)$ if

$$\mathcal{L}_k G = 0 ,$$

$$\iota_k H = d\mathbf{v} ,$$

$$\mathcal{L}_k \Phi = 0 .$$



isometry (circle)



one-form \mathbf{v}

The **gauged action** takes the following form

$$\begin{aligned}\widehat{\mathcal{S}} = & -\frac{1}{4\pi\alpha'} \int_{\partial\Sigma} \left[G_{ij} (dX^i + k^i A) \wedge \star (dX^j + k^j A) + 2i\mathbf{v} \wedge A + \alpha' R \Phi \star 1 \right] \\ & - \frac{i}{2\pi\alpha'} \int_{\Sigma} \frac{1}{3!} H_{ijk} dX^i \wedge dX^j \wedge dX^k .\end{aligned}$$

This action is **invariant** under local symmetry transformations provided that

$$\iota_k \mathbf{v} = k^m \mathbf{v}_m = 0 .$$

Hull, Spence - 1989 & 1991

For the gauged and ungauged theories to be equivalent, impose the **constraint**

$$0 = F = dA .$$

t-duality :: reduction — step III

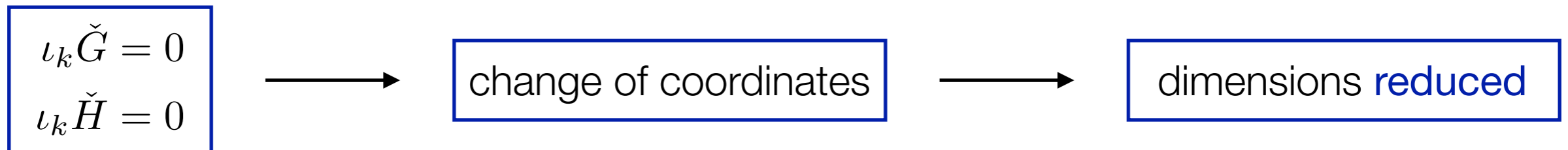
The **equation of motion** for the **gauge field** has the solution

$$|k|^2 A = -k^i G_{ij} dX^j - i \star v .$$

Integrating-out the gauge field results in an action specified by

$$\check{G} = G - \frac{1}{|k|^2} k \wedge \star k + \frac{1}{|k|^2} v \wedge \star v , \quad \check{H} = H + d \left(\frac{1}{|k|^2} k \wedge v \right) .$$

Note that k is a **null-eigenvector** for \check{G} and \check{H} ::



t-duality :: reduction — step III

The equation of motion for the gauge field has the solution

$$\mathcal{T}^i_j = \begin{pmatrix} k^1 & 0 \\ k^2 & \\ \vdots & \\ k^D & \mathbb{1} \end{pmatrix} \quad \check{G}_{ij} = (\mathcal{T}^T \check{G} \mathcal{T})_{ij} = \begin{pmatrix} 0 & 0 \\ 0 & \check{G}_{ab} \end{pmatrix}$$
$$\check{\mathcal{H}}_{ijk} = \check{H}_{lmn} \mathcal{T}^l_i \mathcal{T}^m_j \mathcal{T}^n_k$$
$$\check{\mathcal{H}}_{1jk} = 0$$

Integrating

Note that

$$\begin{aligned} \iota_k \check{G} &= 0 \\ \iota_k \check{H} &= 0 \end{aligned}$$



change of coordinates



dimensions reduced

t-duality :: reduction — step III

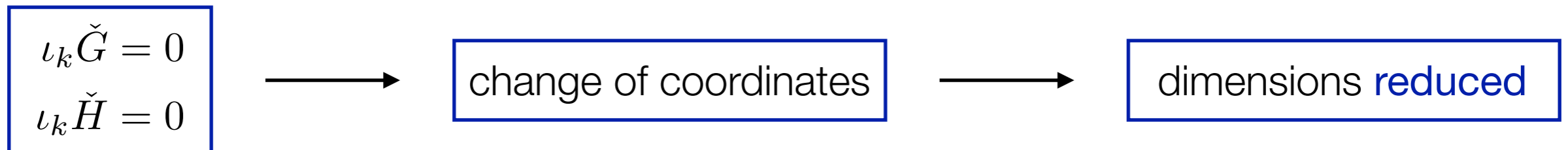
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- Summary ::
- through a **gauging procedure** on the **world-sheet**,
 - the **target-space** dimensions can be **reduced**.

- Constraints ::
- global symmetry $\mathcal{L}_k G = 0, \quad \iota_k H = d\mathbf{v}, \quad \mathcal{L}_k \Phi = 0,$
 - gauging $\iota_k \mathbf{v} = 0 = \iota_v \mathbf{k},$
 - vanishing F $\mathcal{L}_v G = 0, \quad \iota_v H = d\mathbf{k}, \quad \mathcal{L}_v \Phi = 0.$

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- Outlook ::
- T-duality corresponds to interchanging $k \longleftrightarrow v$.

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example :: su(2) wzw

Consider the **SU(2) WZW** model (three-sphere with H -flux) ::

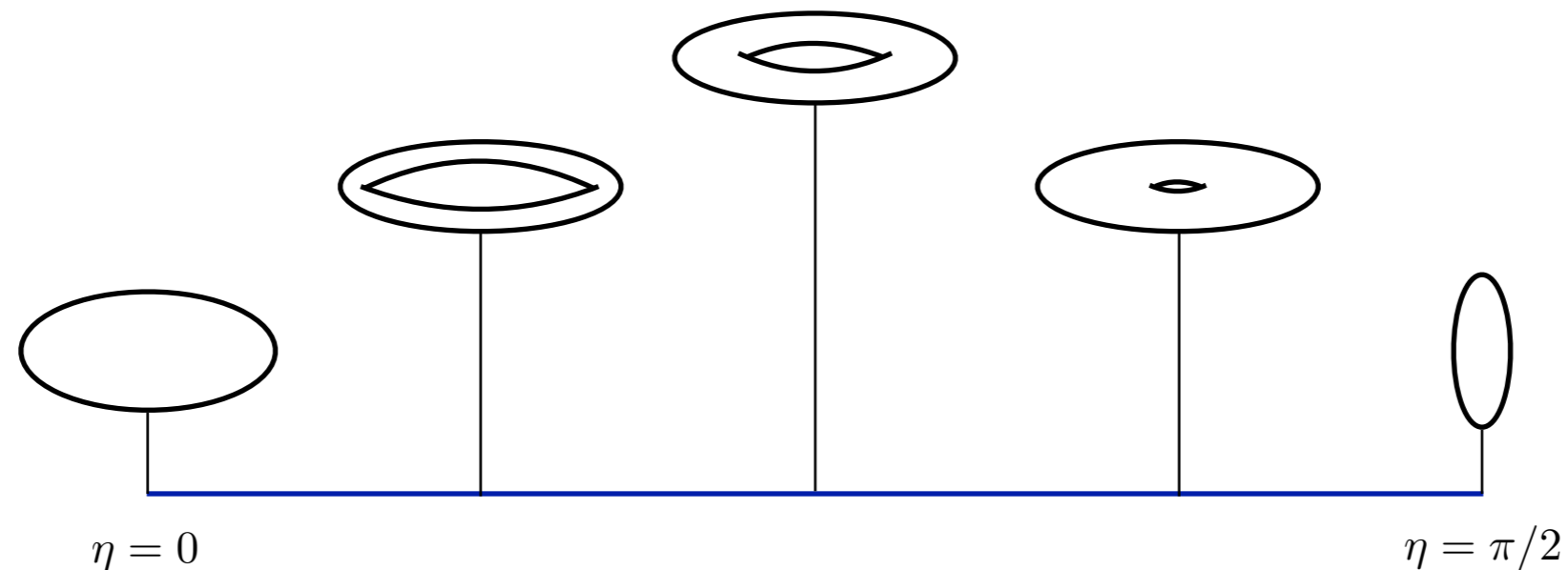
$$ds^2 = R^2 \left(\sin^2 \eta d\zeta_1^2 + \cos^2 \eta d\zeta_2^2 + d\eta^2 \right),$$

$$H = 2 R^2 \sin \eta \cos \eta d\zeta_1 \wedge d\zeta_2 \wedge d\eta,$$

$$\Phi = \text{const.}$$

$$\zeta_{1,2} \in [0, 2\pi),$$

$$\eta \in [0, \pi/2].$$



This geometry has **two** directions of **isometry**.

example :: reduced theories

The **reduction constraints** are satisfied by **precisely two** Killing vectors ::

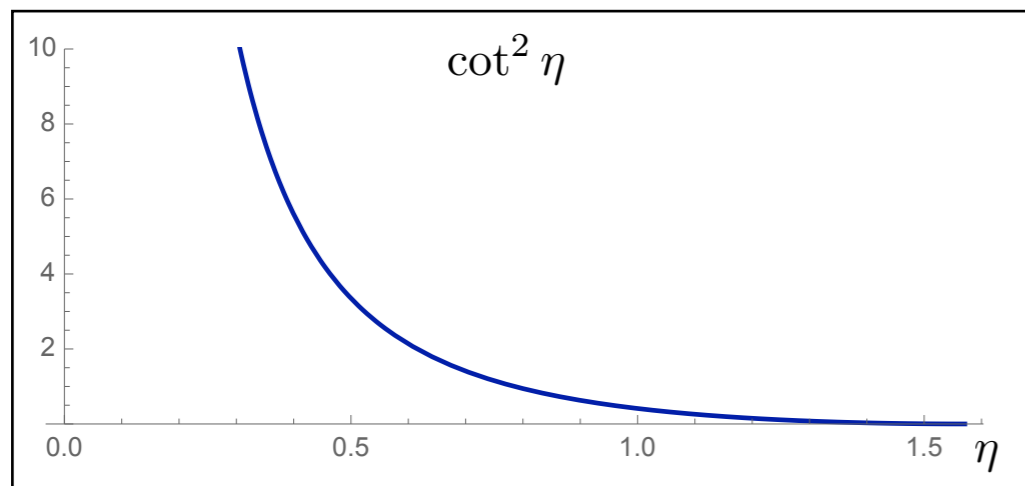
choice 1 — $k = \partial_{\zeta_1}$



$$\check{d}s^2 = R^2 (\cot^2 \eta d\zeta^2 + d\eta^2),$$

$$\check{H} = 0,$$

$$\check{\Phi} = \Phi - \log(\sin \eta),$$



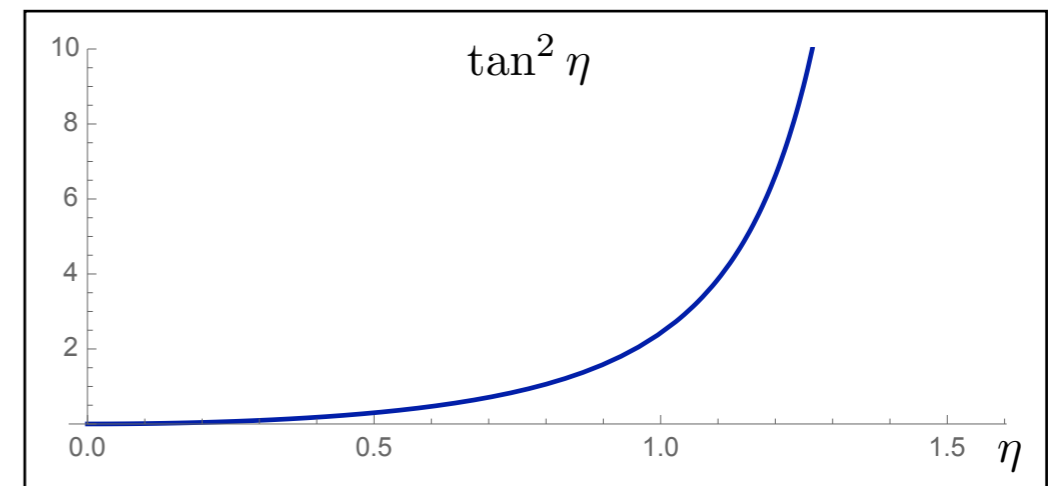
choice 2 — $k = \partial_{\zeta_2}$



$$\check{d}s^2 = R^2 (\tan^2 \eta d\zeta^2 + d\eta^2),$$

$$\check{H} = 0,$$

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$$\text{choice 2} - k = \partial_{\zeta_2}$$



$$\check{d}s^2 = R^2 (\tan^2 \eta d\zeta^2 + d\eta^2),$$

$$\check{H} = 0,$$

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The reduced theories are **conformal**, and **T-dual** to each other.

Summary ::

- the **SU(2) WZW** model ($D=3$) allows for **two** reductions,
- leading to two **conformal** models in $D=2$,
- which are **T-dual** to each other.

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More results ::

- generalization to **multiple** (non-abelian) gaugings,
- T-duality is a **symmetry** of the doubled theory,
- the **Buscher** rules are reproduced,
- results on **conformality** of reduced theories,
- WZW models with arbitrary group,
- ...