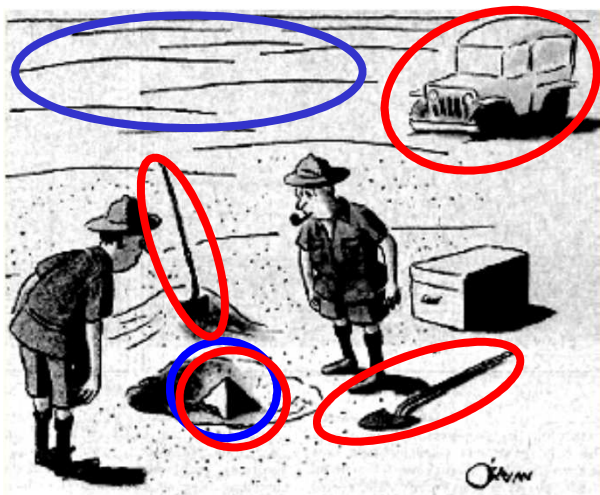


SUSY searches – Part I

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CPPM/IN2P3–Univ. de la Méditerranée (Marseille, FRANCE)

Terascale School on BSM Physics (DESY)



*“This could be the discovery of the century.
Depending, of course, on how far down it goes”*

Part I

Motivation, the tools to dig, EW SUSY

Part II

3rd gene and Strong SUSY, Beyond MSSM

Th 7-Apr	Fr 8-Apr
	Part II
Part I	

Who am I ?

□ Member of ATLAS since 1998

<http://www.cern.ch/pralavop/>

1998: CNRS Permanent Research position at CPPM (Marseille, FRANCE)

1998-2010: Building, qualification and commissioning of the ATLAS
Electromagnetic liquid argon calorimeter with cosmics and first data.
Co-editor of the paper summarizing the performance of the ATLAS
Liquid Argon Calorimeter [EPJ C70 \(2010\) 723](#).

2005-2010: Preparation of analysis in top and beyond the Standard Model
(W', Z' and Supersymmetry) physics. Prospects on top/W polarisation
[EPJ C44S2 \(2005\) 13](#).

2010: Co-editor of W/Z cross-section paper [JHEP 1012 \(2010\) 060](#).

2010-2012: Co-convenor of the ATLAS Supersymmetry Working group. 50 articles
submitted to journals (~ 20% of ATLAS articles by 2012).

2013: CNRS research director

2013-2015: 2 SUSY Searches ([JHEP 1406 \(2013\) 035](#), [JHEP 1405 \(2014\) 071](#)).

["Lessons for SUSY from the LHC"](#) in "SUSY after the Higgs discovery" book,
["Particle Physics and Cosmology"](#) in "100th Summer School Les Houches Proc."

2015-: Co-convenor of ATLAS Top Yukawa coupling measurement group

- **"Lessons for SUSY from the LHC after the first run"**, EPJ C74 (2014) 2801, 1404.7191
 - Lectures will be largely based on this – updated with latest results ...
- Bibliography page : <http://www.cern.ch/pralavop/phd.html>

❑ ATLAS / CMS SUSY results

[illegible]

67=58+9 papers for LHC Run-1 [Final ?]

1 paper and 7 CONF note for LHC Run-2

>150 papers !!!

Lecture Organization

Ia - Ingredients for a SUSY search @ LHC

“What drives the sensitivity to SUSY at LHC ?”

1. Teaser (1)
2. SUSY Framework @LHC (11)
3. Objects and Monte Carlo (3)
4. Discriminating variables (5)
5. Background Estimation (7)
6. Fit Results (7)
7. Result Interpretation (3)

~35 slides → 50'

Ib - SUSY ElectroWeak

“Exploring the bottom part of the SUSY spectrum”

1. Parameters of the EWK sector (2)
2. SUSY Higgses (8)
3. ElectroWeakinos (12)
4. Sleptons (1)
5. Summary (1)

~25 slides → 40'

IIa – 3rd gene. SUSY

“Best natural candidates”

1. Sbottom (7)
2. Stop (12)
3. Summary (1)

~20 slides → 30'

IIb – Strong SUSY

“Inclusive searches”

1. Gluinos (12)
2. 1/2nd generation squarks (1)
3. Summary (1)
4. Overall MSSM summary (2)

~15 slides → 25'

IIc – Beyond MSSM

“Last hopes (!) and conclusions”

1. Long lived (10)
2. RPV (8)
3. Other SUSY realization (3)
4. Overall conclusions (3)

~20 slides → 30'

Part Ia - Teaser

EW Neutral Current (Z^0)

Can we **discover** (at least) one super-partner or **invalidate** the ~~Natural version of MSSM~~?

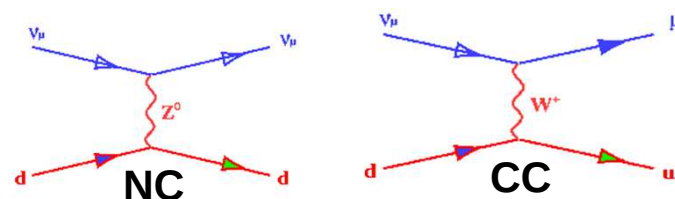
1- Need to be sure of your tools (background estimate, stat, ...)

The results that were obtained from this small CERN chamber (Block *et al.* 1964) included a limit on the ratio of "elastic" neutral-current (NC) events to charged-current (CC) events: $\sigma(\nu_\mu + p \rightarrow \nu_\mu + p) / \sigma(\nu_\mu + n \rightarrow \mu^- + p) \leq 0.03$ for events with proton momentum above 250 MeV/c. The accepted ratio is about 0.12. The CERN result was just wrong, following a book-keeping error (the actual 90% confidence limit was < 0.09). As is often the case in physics, the error was uncovered by a graduate student, Michel Paty from Strasbourg, who found a limit ≤ 0.20 and put this number in his thesis (Paty, 1965). The CERN group intended to publish an erratum, but (since there seemed to be little interest in our limit) we decided to wait for the results from a forthcoming propane run, since that would measure scattering from free as well as bound protons, and we could exploit the better kinematic constraints. However, the propane run was delayed by more than two years and the corrected limit for the ratio of $< 0.12 \pm 0.06$ was not published until 1970 (Cundy *et al.* 1970).

2- Need to have a well designed search strategy

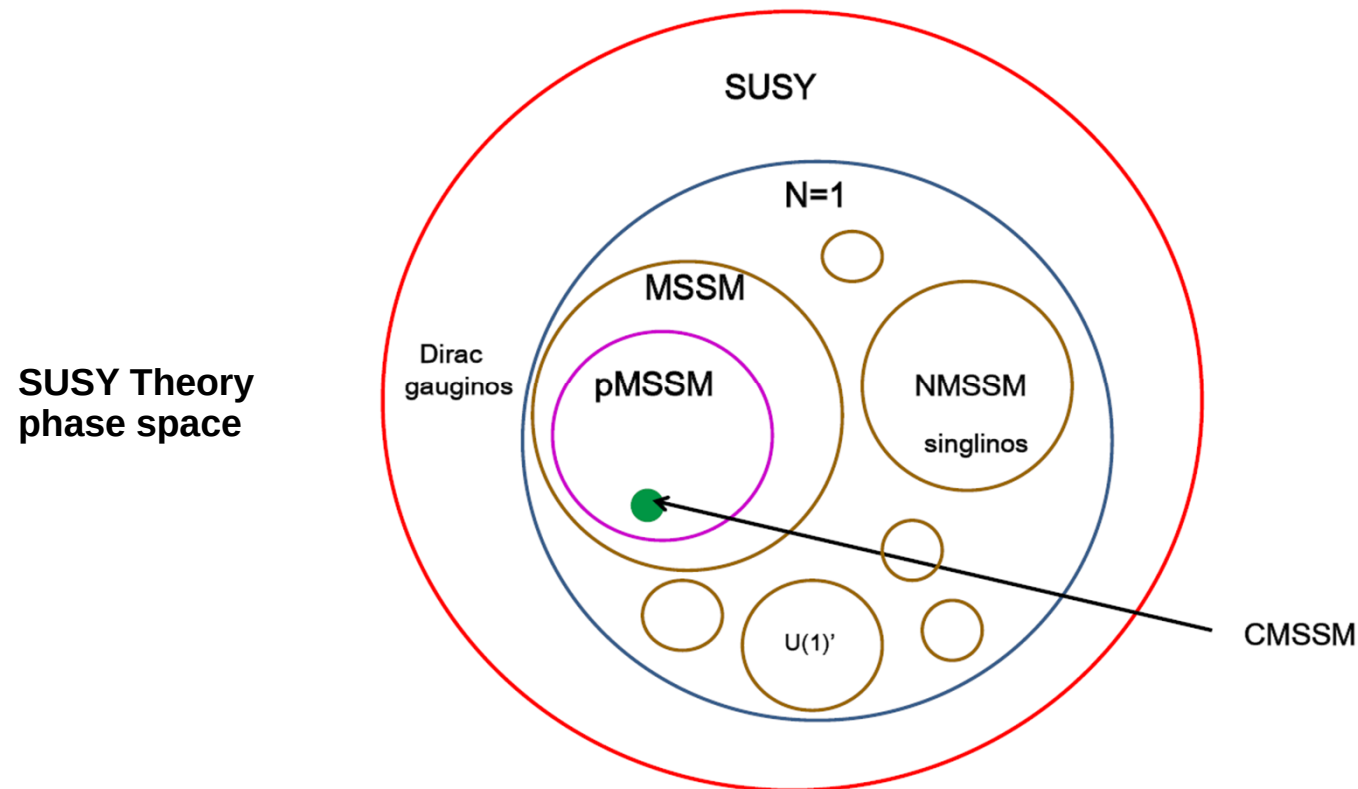
Gargamelle's priorities

1. W search
2. deep inelastic scattering, scaling
3. current algebra sum rules, CVC, PCAC
4. Diagonal Model
5. $\Delta S = 1$ processes, inverse hyperon decay, $\bar{\nu}_\mu + p \rightarrow \Lambda + \mu^+$
6. inverse muon decay, $\nu_\mu + e^- \rightarrow \mu^- + \nu_e$
7. electron-muon universality
8. neutral-current search
9. form factors in exclusive reactions
10. search for heavy leptons



That's why we will spend some time on understanding carefully all our tools today

SUSY Framework @LHC (1)



T. Rizzo (SLAC Summer Institute, 01-Aug-12)

Many SUSY realizations are possible ...
Mainly focus on MSSM

SUSY Framework @LHC (2)

A- MSSM, the larger SUSY framework we can test systematically at LHC ...

- For neutralino and charginos I'll use Greek or Capital Roman letter.
- LSP = Lightest Supersymmetric Particle. In this lecture LSP will be $\tilde{G}$ or $\tilde{\chi}_1^0$

MSSM a.k.a Weak Scale SUSY : $\longrightarrow$ MSSM=Minimal Supersymmetry Standard Model (i.e. based on parsimony principle)
29 sparticles + 4 Higgs undiscovered

	Names	Spin	P_R	Gauge Eigenstates	Mass Eigenstates
SUSY Higgses	Higgs bosons	0	+1	$H_u^0 H_d^0 H_u^+ H_d^-$	$h^0 H^0 A^0 H^\pm$
Strong SUSY	squarks	0	-1	$\tilde{u}_L \tilde{u}_R \tilde{d}_L \tilde{d}_R$	(same)
				$\tilde{s}_L \tilde{s}_R \tilde{c}_L \tilde{c}_R$	(same)
				$\tilde{t}_L \tilde{t}_R \tilde{b}_L \tilde{b}_R$	$\tilde{t}_1 \tilde{t}_2 \tilde{b}_1 \tilde{b}_2$
Weak SUSY	sleptons	0	-1	$\tilde{e}_L \tilde{e}_R \tilde{\nu}_e$	(same)
				$\tilde{\mu}_L \tilde{\mu}_R \tilde{\nu}_\mu$	(same)
				$\tilde{\tau}_L \tilde{\tau}_R \tilde{\nu}_\tau$	$\tilde{\tau}_1 \tilde{\tau}_2 \tilde{\nu}_\tau$
	neutralinos	1/2	-1	$\tilde{B}^0 \tilde{W}^0 \tilde{H}_u^0 \tilde{H}_d^0$	$\tilde{N}_1 \tilde{N}_2 \tilde{N}_3 \tilde{N}_4$
	charginos	1/2	-1	$\tilde{W}^\pm \tilde{H}_u^\pm \tilde{H}_d^\pm$	$\tilde{C}_1^\pm \tilde{C}_2^\pm$
	gluino	1/2	-1	$\tilde{g}$	(same)
	goldstino (gravitino)	1/2 (3/2)	-1	$\tilde{G}$	(same)

$\tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0$
 $\tilde{\chi}_1^{+/-}, \tilde{\chi}_2^{+/-}$

R-parity (P_R or R_P)= -1 SUSY, +1 SM

S. Martin, SUSY Primer, hep-ph/9709356

SUSY Framework @LHC (3)

See also Andreas Weiler

B- ... assuming ~no fine-tuning ...

- Range of MSSM particle masses can be predicted imposing ~no fine tuning ($\Delta \approx \delta m_h^2/m_h^2$)

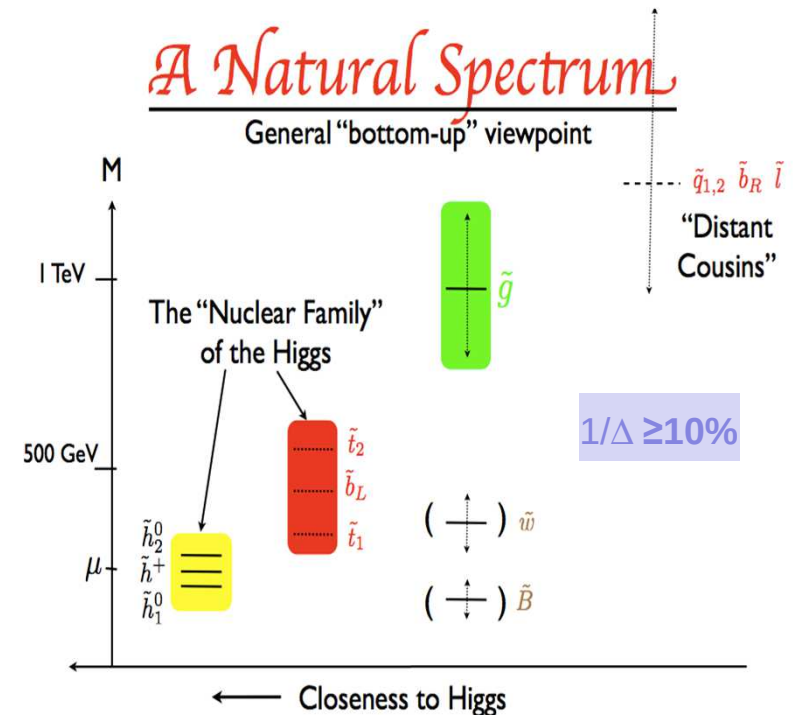
- Standard Model (SM) is an effective field theory valid up to Λ_{NP} M. Veltman, Acta Phys. Polon. B 12, 437 (1981)

$$(m_H^2)^2 = m_H^2 = (m_H^2)_0 - \frac{3G_F}{4\sqrt{2}\pi^2} (4m_t^2) \Lambda_{NP}^2 + \frac{3G_F}{4\sqrt{2}\pi^2} (2m_W^2 + m_Z^2 + m_h^2) \Lambda_{NP}^2 + O(m^2 \ln[\Lambda_{NP}/\mu])$$

- Some SUSY realizations are 'natural' candidate for New Physics (NP) (e.g. for the top loop)

$$m_H^2 = (m_H^2)_0 - \frac{3G_F}{4\sqrt{2}\pi^2} (4m_t^2) \Lambda_{NP}^2 + \frac{3G_F}{4\sqrt{2}\pi^2} (4m_{\tilde{t}}^2) \Lambda_{NP}^2 + O(m^2 \ln[\Lambda_{NP}/m_t])$$

Note: to have 4, two scalars are considered ($\tilde{t}_R$ and $\tilde{t}_L$)



L. Hall, BNL SUSY Workshop, Oct 2011 + 1110.6926

Note: no constraints on $M(H^0, H^\pm, A^0)$ and $M(\tilde{G})$

SUSY Framework @LHC (4)

C- ... and make three more assumptions

1. The nature of R-Parity (RP) :
 - a. RPC: SUSY particles are pair-produced at LHC and the lightest one is stable
 - b. RPV, including also long-lived if $\lambda < O(10^{-5})$.
2. The nature of the stable Lightest SUSY particle (LSP) :
 - a. χ_1^0 with $M_{\text{LSP}} > \text{GeV}$
 - b. $\tilde{G}$ with $M_{\text{LSP}} \ll \text{GeV}$
3. The value of $\Delta M = M_{\text{SUSY}}$ [highest particle produced at LHC] - M_{LSP}
 - a. Open spectra [$\Delta M > O(100) \text{ GeV}$]: high energetic objects, generic searches
 - b. Compressed spectra [$\Delta M < O(100) \text{ GeV}$]: dedicated searches (monojet ...) + long-lived

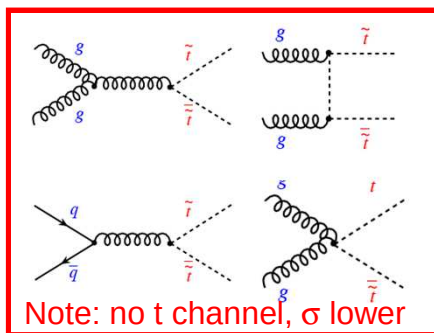
Vanilla framework used for **systematic** searches: **A, B, C1a, C2a, C3a***
Other combinations are searched for but not systematically.

* Default assumption in the following

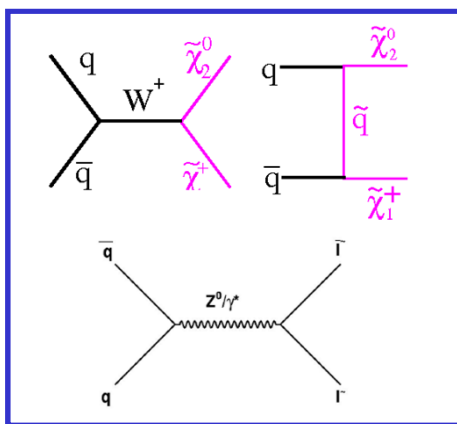
SUSY Framework @LHC (5)

❑ Sparticles are pair produced at LHC (R-Parity conserved)

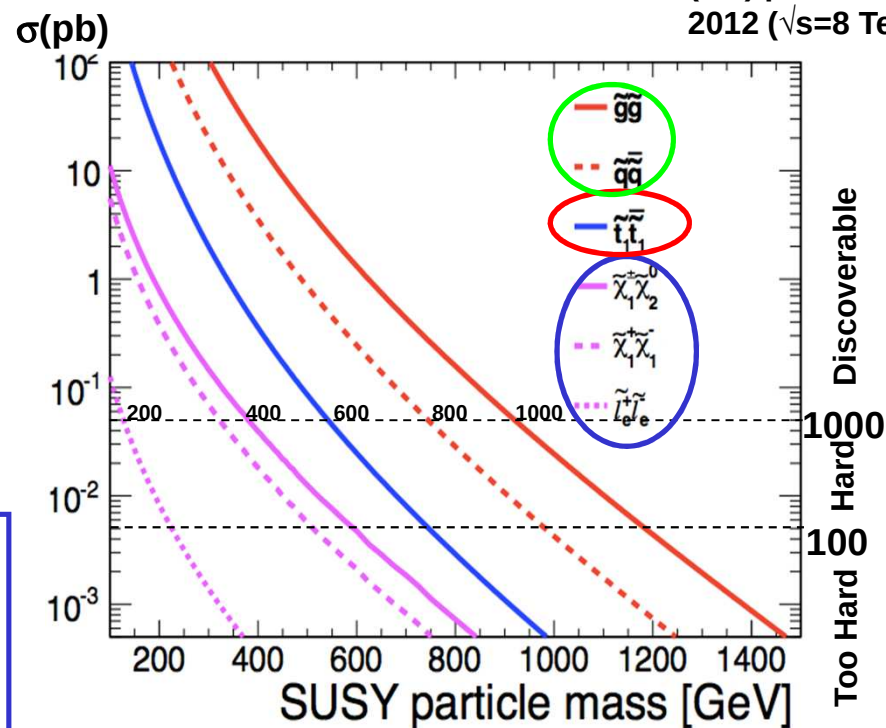
SUSY 3rd gene



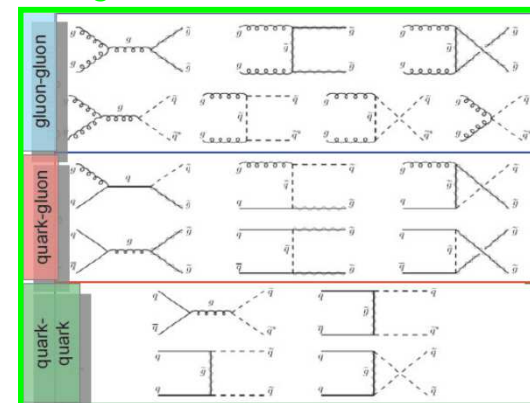
Weak SUSY



$\sigma(\text{pb})$

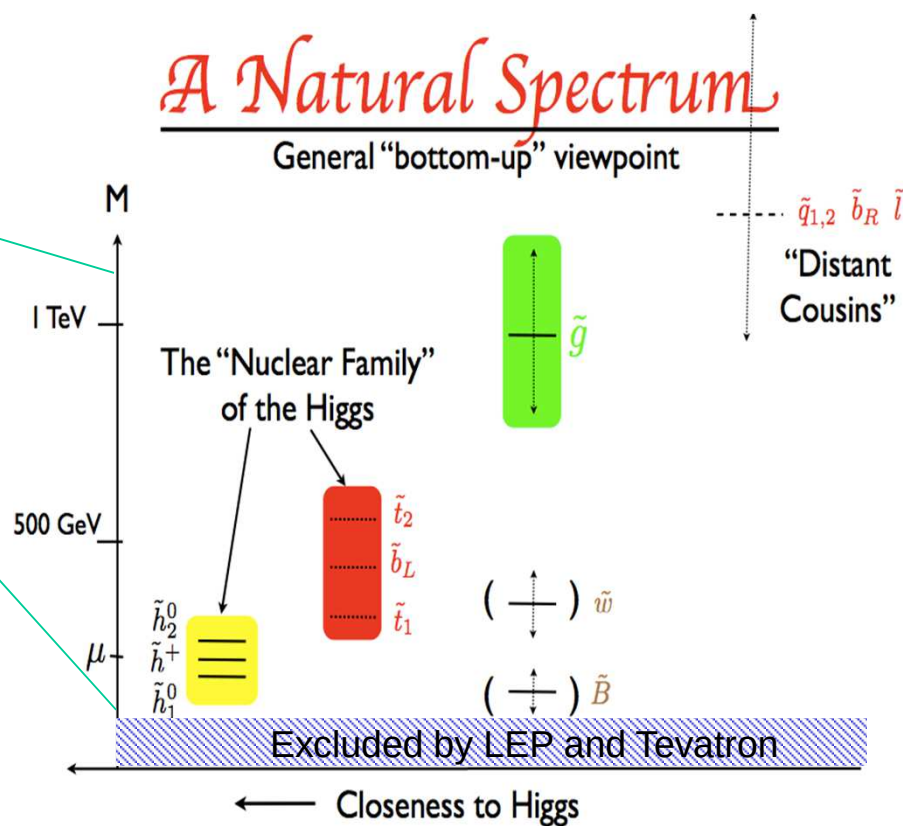
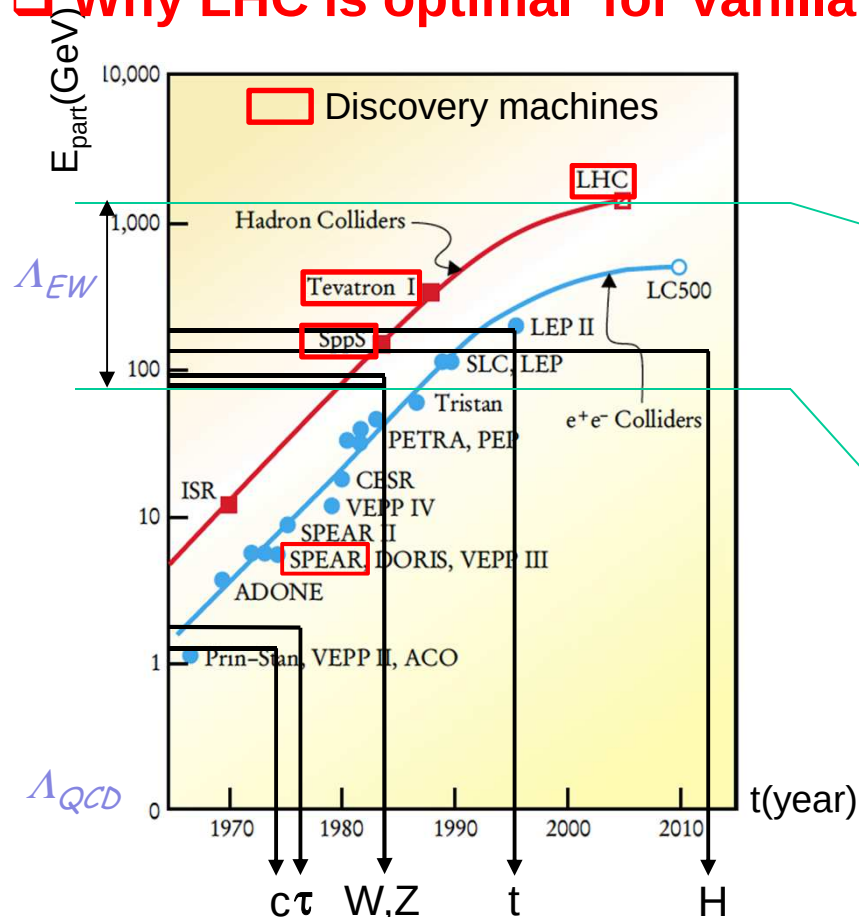


Strong SUSY



SUSY Framework @LHC (6)

Why LHC is optimal for Vanilla SUSY ?

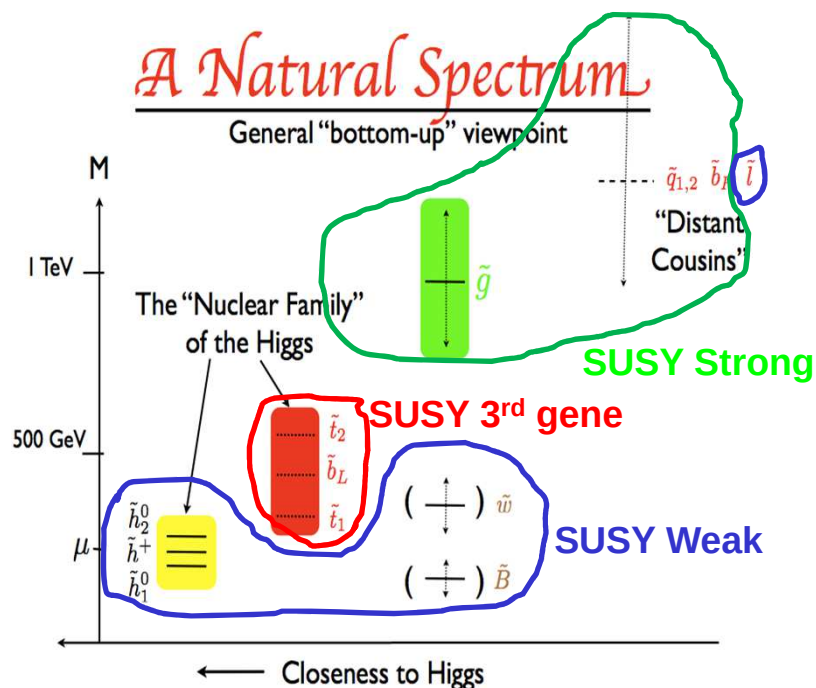


LHC* can cover the **best motivated SUSY phase space for the first time !!**

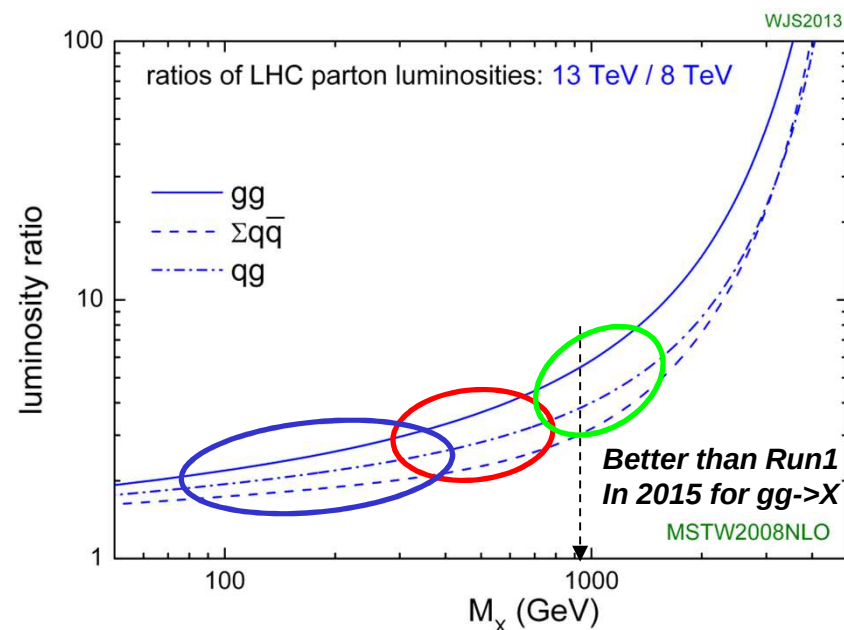
*The largest and the most complex machine ever constructed by humans (From Nobel Prize 2013 Press release)

SUSY Framework @LHC (7)

□ SUSY searches at 8 and 13 TeV



- SUSY Weak: Lecture Ib
- SUSY 3rd gene.: Lecture IIa
- SUSY Strong: Lecture IIb



- SUSY Weak: No results at 13 TeV
- SUSY 3rd gene.: ~Almost no results at 13 TeV
- SUSY Strong: Many results at 13 TeV

SUSY Framework @LHC (8)

□ A very predictive theory !

- Once mass spectrum known theoretically computable decay rate
- Mix of on-shell (2 body decay) and off-shell (3-body decay)

MSSM: 29 sparticles + 4 Higgs undiscovered

Main decay channels in MSSM

Names	Spin	P_R	Gauge Eigenstates	Mass Eigenstates
Higgs bosons	0	+1	$H_u^0 H_d^0 H_u^+ H_d^-$	$h^0 H^0 A^0 H^\pm$
squarks	0	-1	$\tilde{u}_L \tilde{u}_R \tilde{d}_L \tilde{d}_R$ $\tilde{s}_L \tilde{s}_R \tilde{c}_L \tilde{c}_R$ $\tilde{t}_L \tilde{t}_R \tilde{b}_L \tilde{b}_R$	(same) (same) $\tilde{t}_1 \tilde{t}_2 \tilde{b}_1 \tilde{b}_2$
sleptons	0	-1	$\tilde{e}_L \tilde{e}_R \tilde{\nu}_e$ $\tilde{\mu}_L \tilde{\mu}_R \tilde{\nu}_\mu$ $\tilde{\tau}_L \tilde{\tau}_R \tilde{\nu}_\tau$	(same) (same) $\tilde{\tau}_1 \tilde{\tau}_2 \tilde{\nu}_\tau$
neutralinos	1/2	-1	$\tilde{B}^0 \tilde{W}^0 \tilde{H}_u^0 \tilde{H}_d^0$ (Bino) (Wino) (Higgsino)	$\tilde{N}_1 \tilde{N}_2 \tilde{N}_3 \tilde{N}_4$
charginos	1/2	-1	$\tilde{W}^\pm \tilde{H}_u^\pm \tilde{H}_d^\pm$ (Wino) (Higgsino)	$\tilde{C}_1^\pm \tilde{C}_2^\pm$
gluino	1/2	-1	$\tilde{g}$	(same)
goldstino (gravitino)	1/2 (3/2)	-1	$\tilde{G}$	(same)

$h \rightarrow bb, WW, \tau\tau; H^0 \rightarrow hh, WW, tt, bb; A^0 \rightarrow tt, bb; H^{\pm} \rightarrow \tau\nu, tb$

$\tilde{q} \rightarrow q\tilde{g}, q\tilde{\chi}_1^0, q'\tilde{\chi}_1^{+/-}, q'W^{(*)}\tilde{\chi}_1^0$ $\left\{ \begin{array}{l} \tilde{q}_L \rightarrow q\tilde{\chi}_{1(2)}^0, q'\tilde{\chi}_1^{+/-} (\tilde{\chi}_2^0 \text{wino}) \\ \tilde{q}_R \rightarrow q\tilde{\chi}_1^0 (\tilde{\chi}_1^0 \text{bino}) \end{array} \right.$
 $\tilde{g} \rightarrow q\bar{q}, q\bar{q}\tilde{\chi}_1^0, q\bar{q}'\tilde{\chi}_1^{+/-}, g\tilde{\chi}_1^0$
STRONG, 3rd gene

$\tilde{l} \rightarrow l\tilde{\chi}_{1(2)}^0, \nu\tilde{\chi}_1^{+/-}$ $\left\{ \begin{array}{l} \tilde{l}_L \rightarrow l\tilde{\chi}_{1(2)}^0, \nu\tilde{\chi}_1^{+/-} (\tilde{\chi}_2^0 \text{wino}) \\ \tilde{l}_R \rightarrow l\tilde{\chi}_1^0 (\tilde{\chi}_1^0 \text{bino}) \end{array} \right.$
 $\tilde{\nu} \rightarrow \nu\tilde{\chi}_{1(2)}^0, l\tilde{\chi}_1^{+/-}$
 $\tilde{\chi}_2^0 \rightarrow W^{(*)}\tilde{\chi}_1^{+/-}, Z^{(*)}\tilde{\chi}_1^0, \tilde{l}l, \tilde{\nu}\nu, \tilde{q}q$
 $\tilde{\chi}_{1(2)}^{+/-} \rightarrow W^{(*)}\tilde{\chi}_1^0, Z^{(*)}\tilde{\chi}_1^{+/-}, l\bar{\nu}, \nu\bar{l}, q\bar{q}'$
Weak

➔ Huge combinatorics: Need to look at lots of final states !

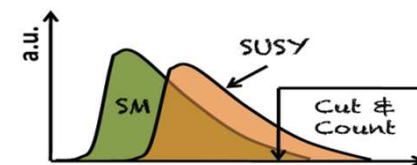
SUSY Framework @LHC (9)

❑ Main characteristics of Vanilla SUSY events at LHC

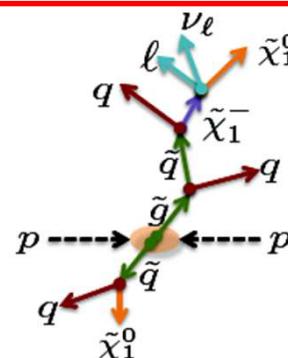
- SUSY appears as excess in tails

MSSM: 29 sparticles + 4 Higgs undiscovered

Names	Spin	P_R	Gauge Eigenstates	Mass Eigenstates
Higgs bosons	0	+1	$H_u^0 \ H_d^0 \ H_u^+ \ H_d^-$	$h^0 \ H^0 \ A^0 \ H^\pm$
squarks	0	-1	$\tilde{u}_L \ \tilde{u}_R \ \tilde{d}_L \ \tilde{d}_R$ $\tilde{s}_L \ \tilde{s}_R \ \tilde{c}_L \ \tilde{c}_R$ $\tilde{t}_L \ \tilde{t}_R \ \tilde{b}_L \ \tilde{b}_R$	(same) (same) $\tilde{t}_1 \ \tilde{t}_2 \ \tilde{b}_1 \ \tilde{b}_2$
sleptons	0	-1	$\tilde{e}_L \ \tilde{e}_R \ \tilde{\nu}_e$ $\tilde{\mu}_L \ \tilde{\mu}_R \ \tilde{\nu}_\mu$ $\tilde{\tau}_L \ \tilde{\tau}_R \ \tilde{\nu}_\tau$	(same) (same) $\tilde{\tau}_1 \ \tilde{\tau}_2 \ \tilde{\nu}_\tau$
neutralinos	1/2	-1	$\tilde{B}^0 \ \tilde{W}^0 \ \tilde{H}_u^0 \ \tilde{H}_d^0$	$\tilde{N}_1 \ \tilde{N}_2 \ \tilde{N}_3 \ \tilde{N}_4$
charginos	1/2	-1	$\tilde{W}^\pm \ \tilde{H}_u^\pm \ \tilde{H}_d^\pm$	$\tilde{C}_1^\pm \ \tilde{C}_2^\pm$
gluino	1/2	-1	$\tilde{g}$	(same)
goldstino (gravitino)	1/2 (3/2)	-1	$\tilde{G}$	(same)



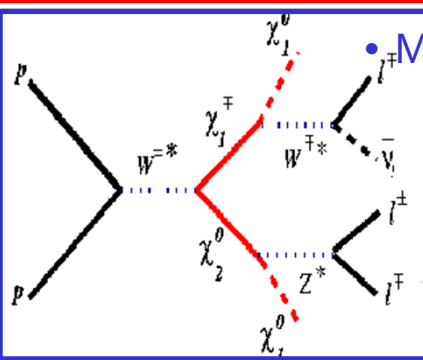
@ LHC



- Missing Energy=MET ($\tilde{\chi}_1^0$)
 - “Jetty” (some from b)
 - Eventually leptons
 - Eventually γ

 $\varepsilon \approx 1\%$

STRONG, 3rd gene



- Missing Energy = MET ($\tilde{\chi}_1^0$)
- Jet veto
- 2 or 3 leptons

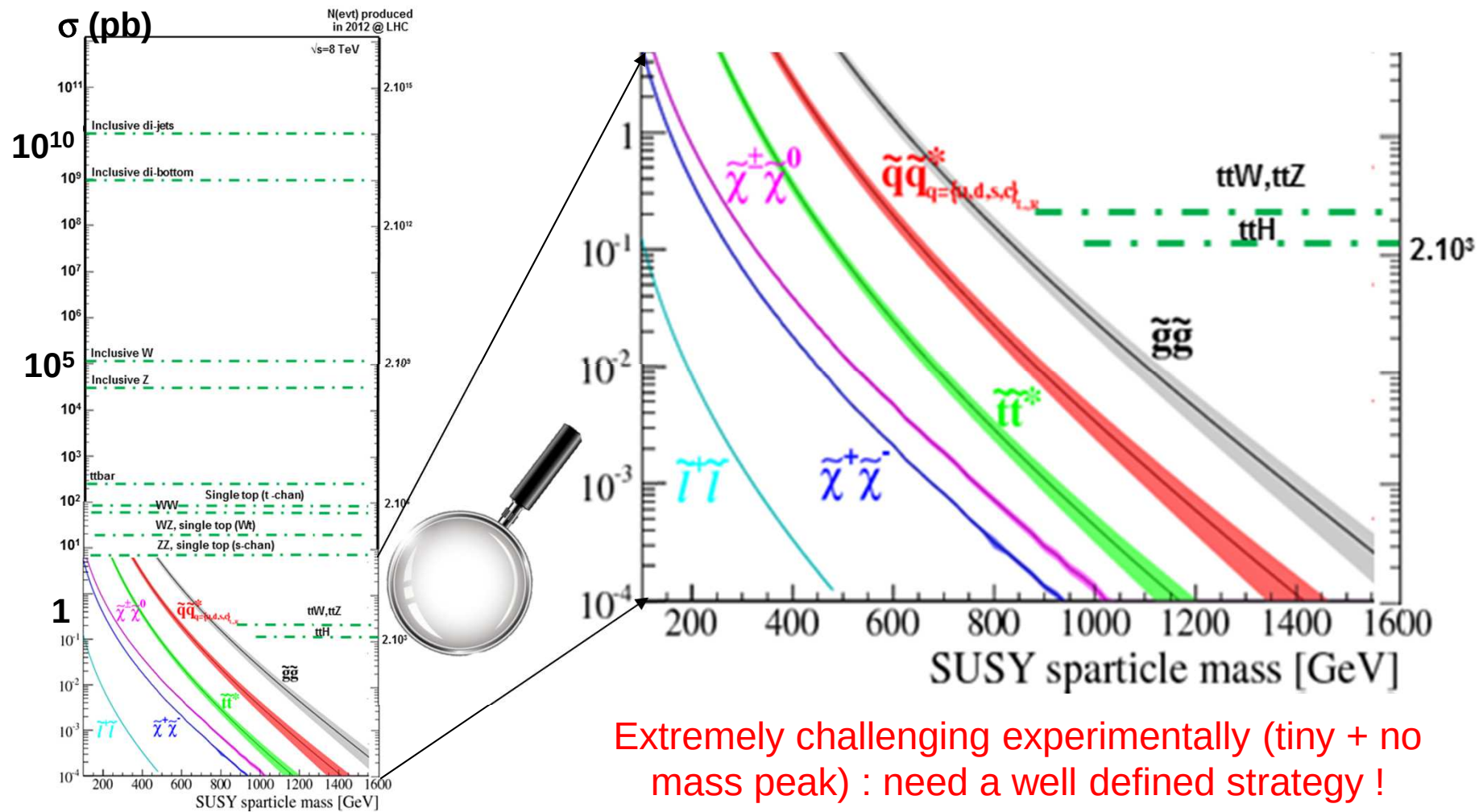
 $\varepsilon \approx 10\%$

Electro-Weak (EWK)

➔ Critical to master well the reconstruction of all objects (esp. MET) !

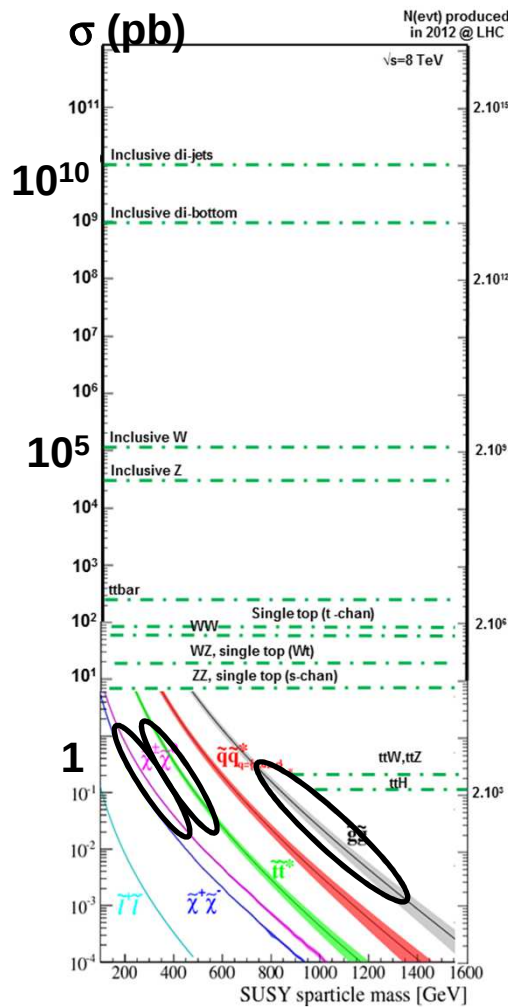
SUSY Framework @LHC (10)

❑ But SUSY is tiny compared to SM processes !

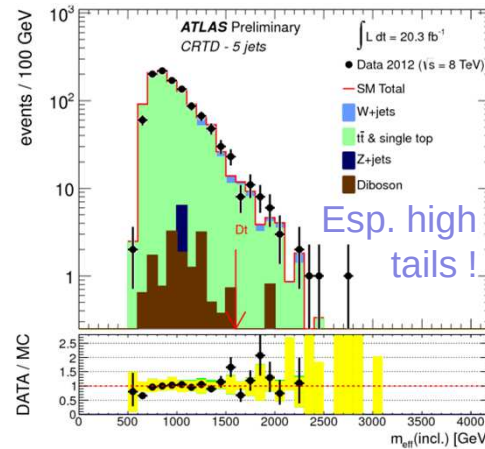


SUSY Framework @LHC (11)

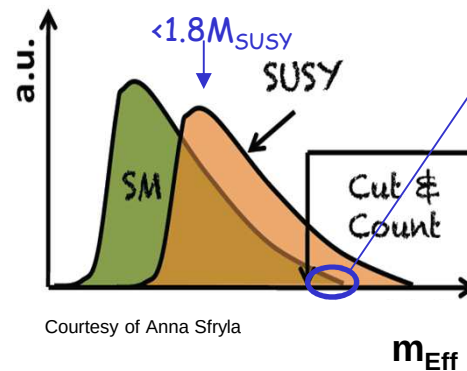
□ Strategy



1 Control objects and Monte Carlo



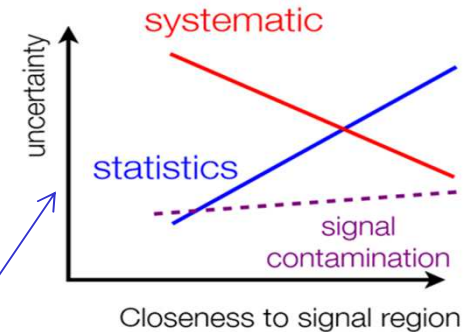
2 Powerful discriminants



Courtesy of Anna Sfryla

3 Background estimation

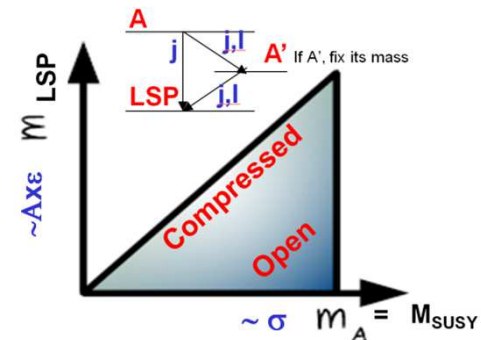
- 1- Multijets: jet smearing method
- 2- W, Z, t, VV : control regions



3- $t\bar{t}V$: Monte Carlo

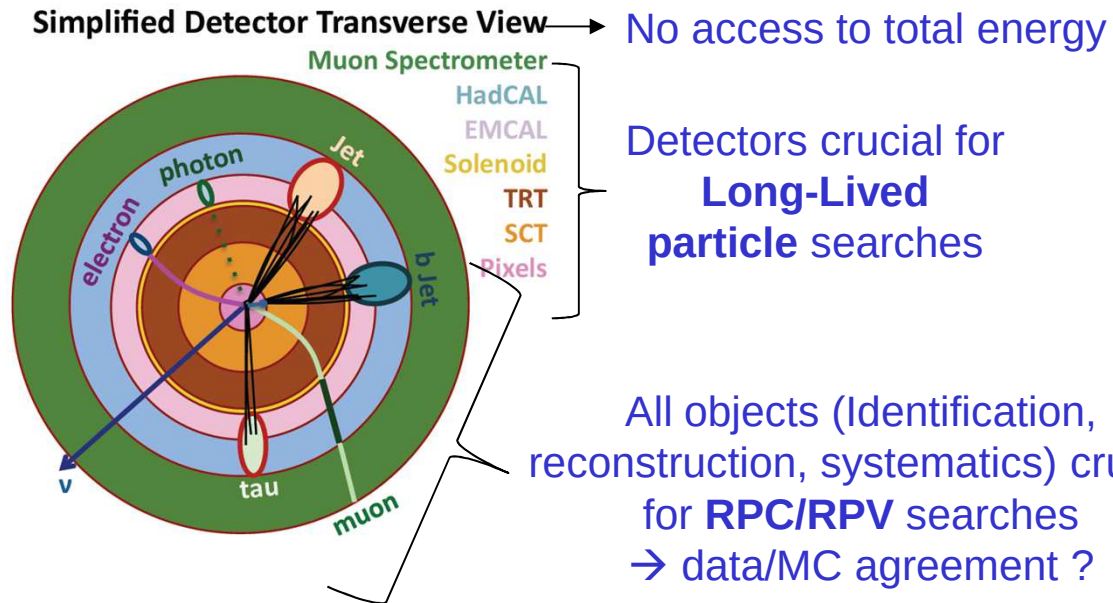
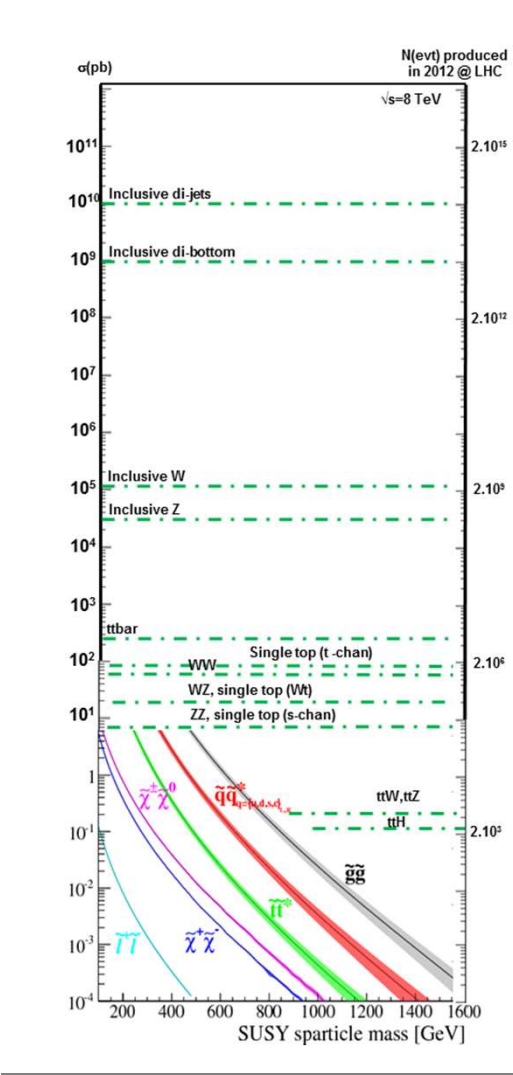
4 Interpretation if no excess

- 1- Constraint model, e.g. MSUGRA
- 2- Simplified/topological models

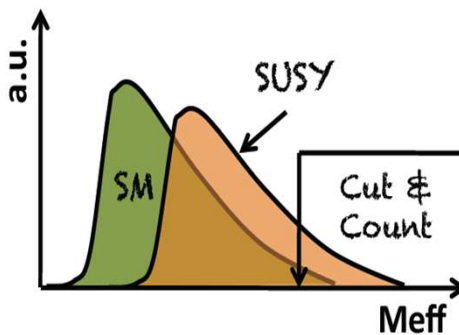


3- Pheno. models, eg. pMSSM

Objects and Monte-Carlo



All objects (Identification, reconstruction, systematics) crucial for **RPC/RPV** searches
→ data/MC agreement ?

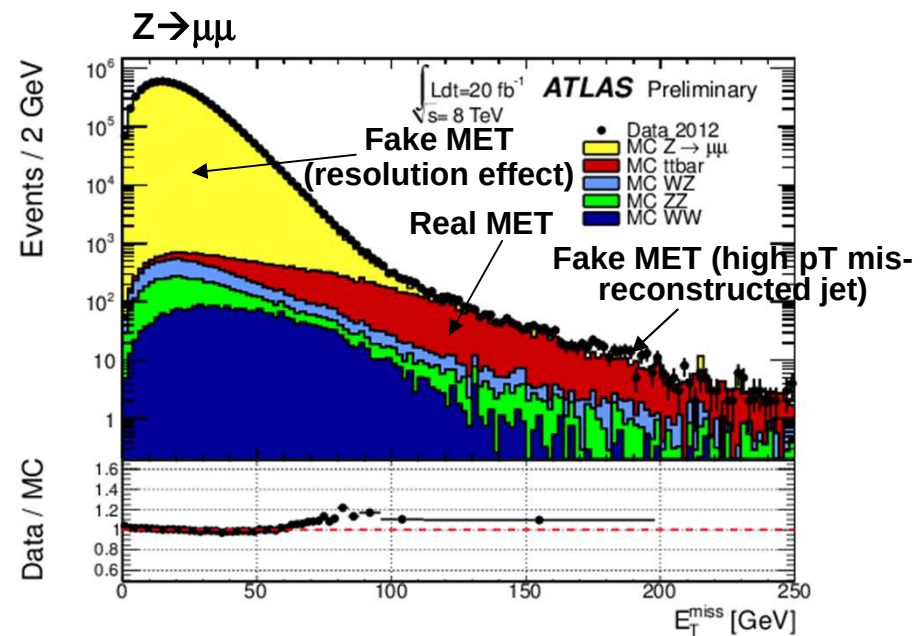
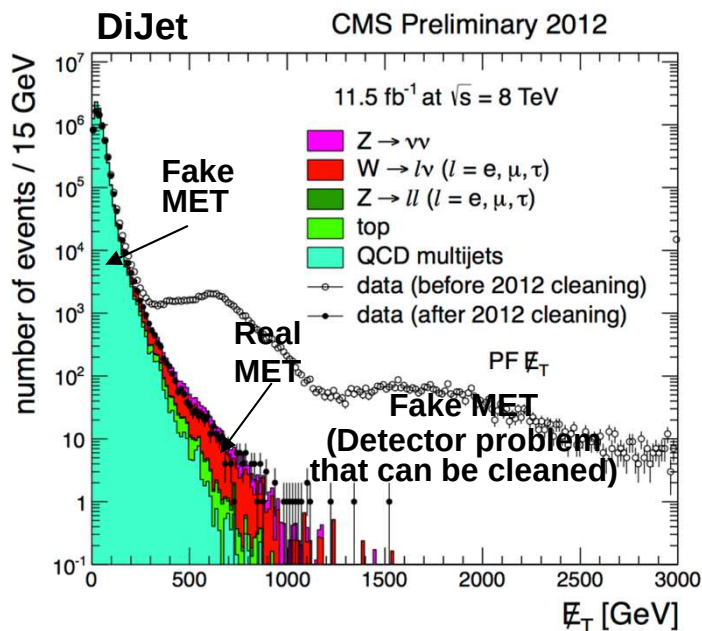


Very challenging to describe $t\bar{t}W$ +jets in the tails
→ Huge effort from the theory community

Objects and Monte-Carlo (1)

❑ Transverse Missing Energy crucial for all SUSY searches

- Energy conservation (transverse plane) : $\vec{MET} = \vec{E}_T^{\text{non-int}} = -\sum \vec{E}_T(\text{calo}) - \sum \vec{E}_T(\text{muon})$
 ✓ If calorimeter not enough use also tracker (Particle Flow a la CMS)
- Real MET : Presence of a neutral weakly interacting particle in the event (i.e. ν)
- Fake MET : Mismeasurement + detector malfunctions, poorly instrumented regions



➔ Agreement data – Monte Carlo key for SUSY searches (systematics not dominant)

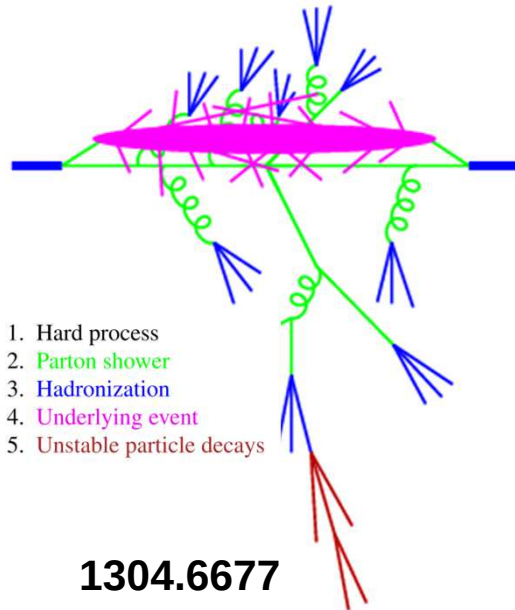
Objects and Monte-Carlo (2)

□ Challenging to model SM processes with high jet multiplicities

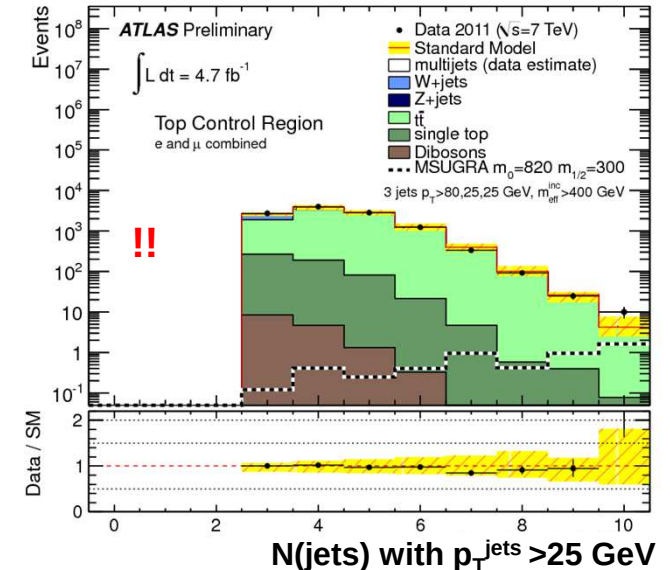
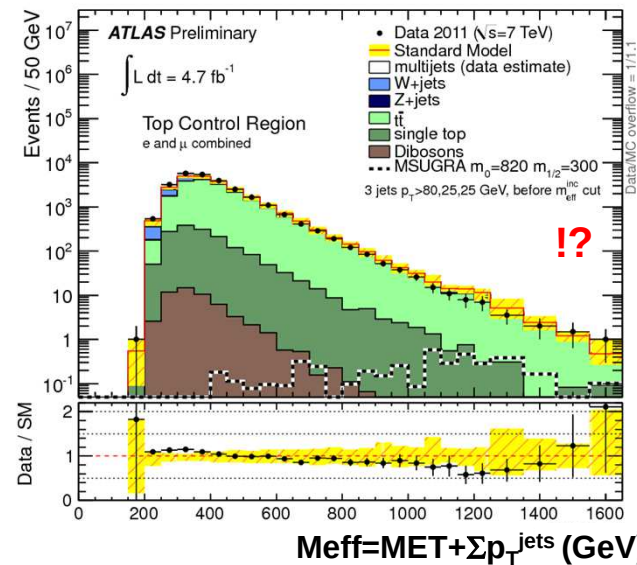
- Parton Shower (PS) : PYTHIA, HERWIG
- Matrix Element (ME) + PS : MADGRAPH, ALPGEN

Note: SHERPA, HERWIG++, NLO+PS (MC@NLO, POWHEG) also used

→ 'Best' to describe large-angle emissions beyond the hardest jet (jets well separated)



SUSY -1lepton Top Control Region



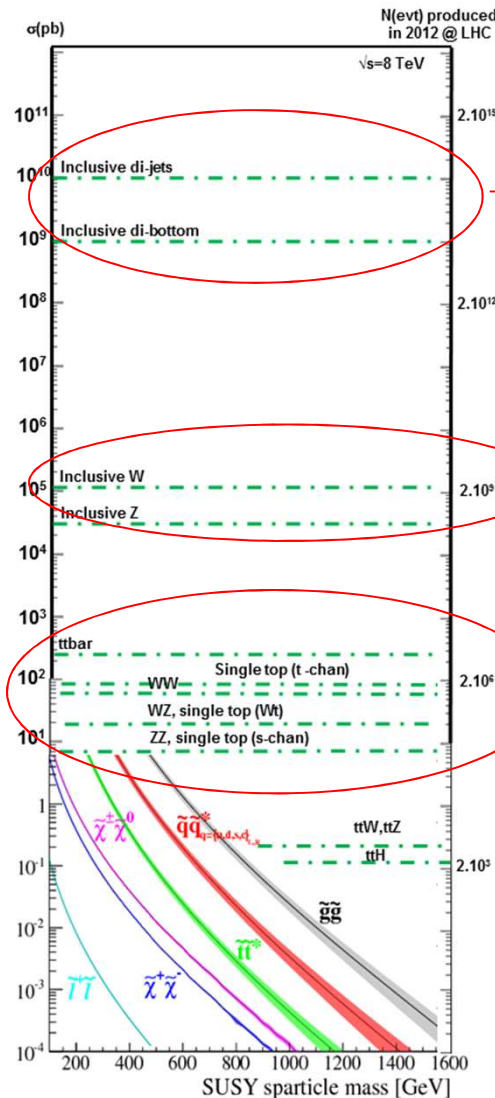
Objects and Monte-Carlo (3)

□ They drive the sensitivity of the searches !

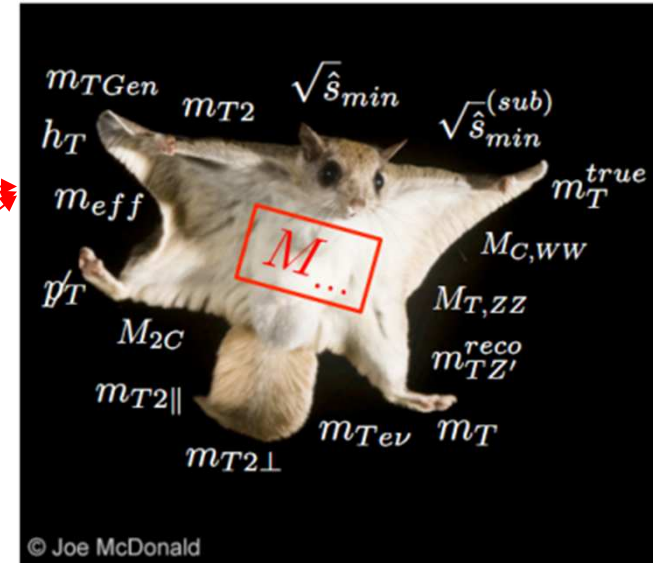
	Systematics	SM background estimate	SUSY Signal <1 TeV	Comments
Experimental	Pile-up			Negligible or Small
	Trigger			Small
	Jet Energy scale (JES)	Dominant.	Dominant	Generally dominates exp
	Jet Energy resolution (JER)	Less than JES (apart Z+jets)	Less than JES	
	b/ τ -tagging	Could dominate	Could dominate	Take over for ≥ 2 b/ τ
	Lepton/ γ energy scale			Small (even for multilep.) except τ
	Lepton/ γ energy resolution			negligible
Theory	Scale, PDF uncertainties	Not for data-driven methods	~20% for NLO+NLL	Depend on many parameters
	Generators+Showering	Poor man's method	N.A.	
	ISR/FSR	Generally important for ttbar	Up to 30% for Compressed spectra	
	MC stat			Depend on grid computing !
	Total (indicative)	~10-100%	~20-50%	

Fully correlated between signal & backgrd

Discriminant variables



Excellent providers
 of jets+MET+(lepton)
 signature

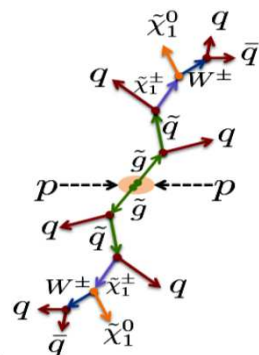


- 1) Need to suppress QCD / W,Z / top by $\sim 10^{10} / 10^5 / 10^2$
- 2) Estimate small remaining quantities
- 3) Fit results
- 4) Interpret the results if no excess

Discriminant variables (1)

□ First need hard kinematic cuts

- To reduce “difficult” background (Fake MET/ lepton, pile-up)

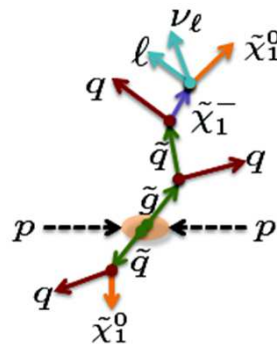


SM Background

1. $Z(\rightarrow \nu\nu) + \text{jets}$
2. $t\bar{t} + W + \text{jets}$
3. QCD

0lepton+jets+MET:

- ✓ Jet+MET trigger
- ✓ Ask several high p_T jets
- ✓ High MET cuts to kill QCD

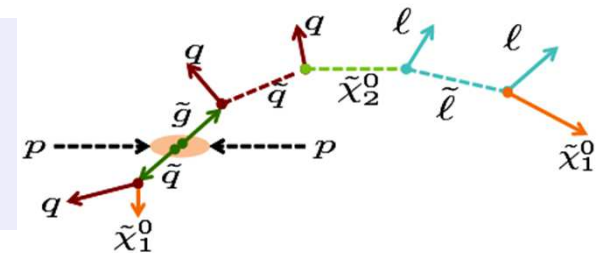


SM Background

1. $t\bar{t}$
2. $W + \text{jets}$
- (3. Fake lepton)

1lepton+jets+MET:

- ✓ Lepton trigger
- ✓ Ask several high p_T jets
- ✓ Lower MET cuts than 0lep
- ✓ $m_T(W) > m(W)$



≥2lepton+jets+MET:

- ✓ Dilepton trigger
- ✓ MET and/or high p_T jets
- ✓ 2l Opp. sign: Z or non Z
- ✓ 2l Same sign
- ✓ 3, 4leptons

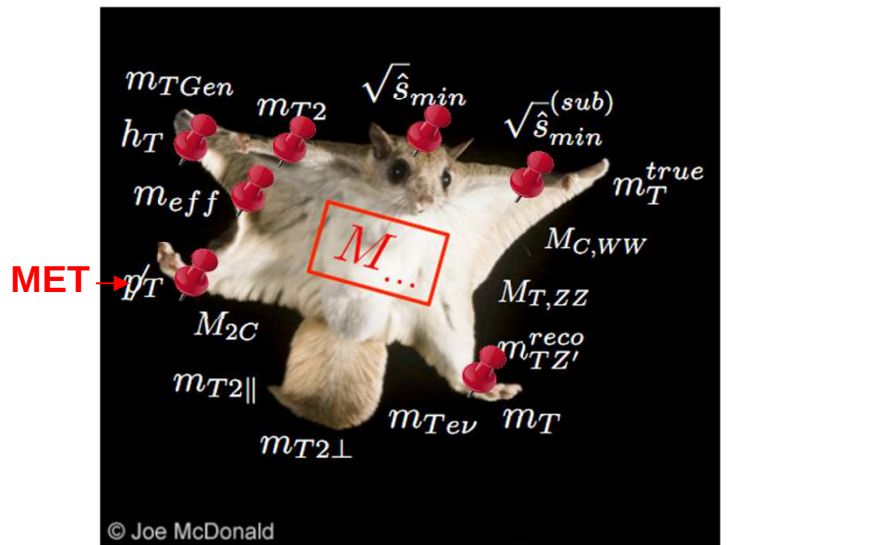
Discriminant variables (2)

❑ Then add powerful discriminant variables

- LHC: unknown momentum along the beams
- SUSY: Sparticles pair produced + Presence of invisible (massive) particles

Assume knowledge of SUSY decay chain

→ Transverse mass-like variables



© Joe McDonald

PRD 84 (2011) 095031 =used in SUSY analyses

Other approaches w/o this assumption:

- Reconstruction of 2 megajets: Razor, α_T
- QCD killers: $\Delta\phi(\text{jets}, \text{MET})$
- QCD+EWK killers: b-jets
- $t\bar{t}$ killers : 2lepton Same sign, 3 bjets, 3leptons, >6 jets
- Z veto

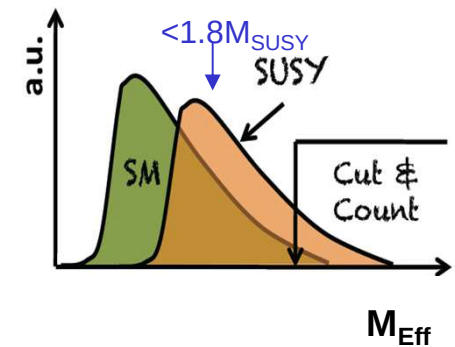
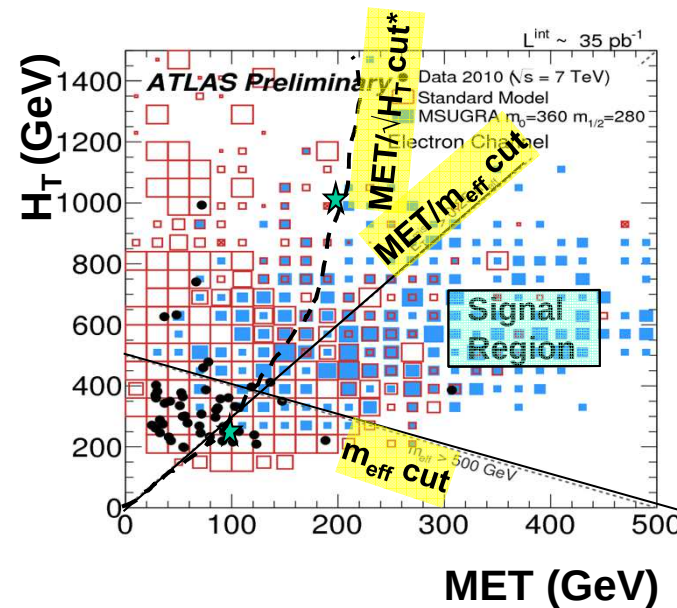
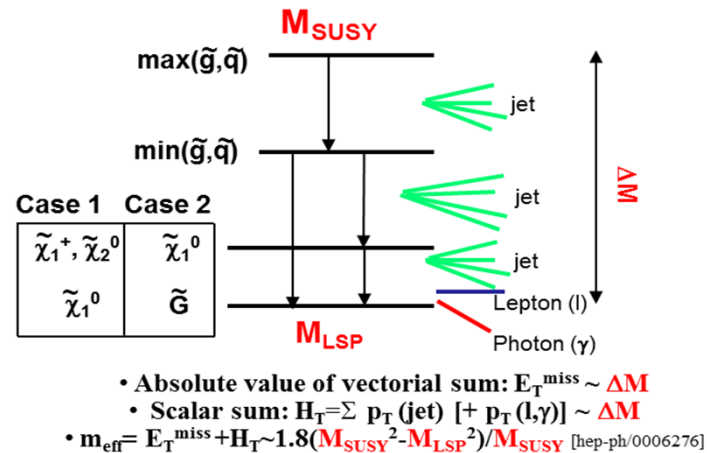
→ Optimal choice of variable(s)/method(s) is analysis dependent

Discriminant variables (3)

Exemple1: Effective mass (M_{Eff})

used w or w/o leptons in final states

- Profit from the correlation between H_T and MET in SUSY (absent in SM)

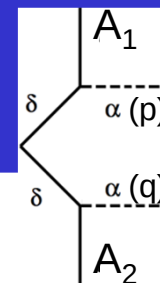


arXiv:hep-ph/9610554

➔ Hard M_{Eff} and MET cuts: signal efficiency $\sim 0.1-10 \%$, high purity for signal

* Useful for lower MET SUSY signal (here cut $\sim 6 \text{ GeV}^{1/2}$)

Discriminant variables (4)

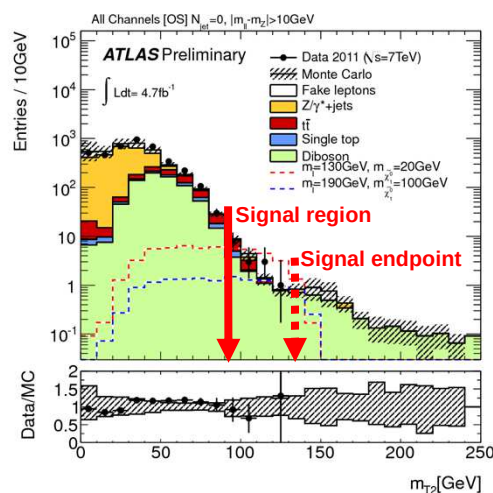


Example 2: M_{T2}, m_{CT}

- Generate end-point at different position than SM because of massive LSP*

$$M_{T2}^2 = \min_{\vec{p}_T + \vec{q}_T = \text{MET}} [\max\{M_T^2(A_1, p), M_T^2(A_2, q)\}]$$

- Min.: most 'consistent' missing momentum sharing between invisibles
- Max.: Better of the 2 lower bounds



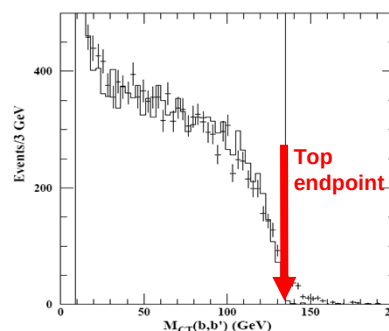
For $\tilde{l} \rightarrow l \tilde{\chi}$ endpoint
 $[M(\tilde{l})^2 - M(\tilde{\chi})^2] / M(\tilde{l})$
 $\sim 130 \text{ GeV}$ for $m(l) \sim 0$

JP G29 (2003) 2343

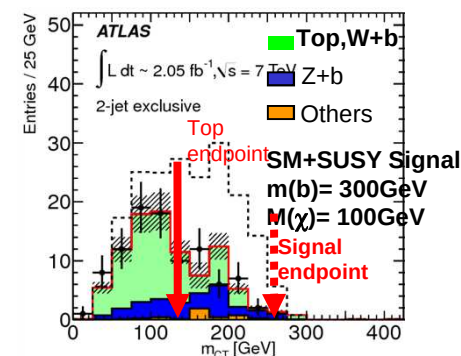
$$M_{CT}^2 = [E_T(A_1) + E_T(A_2)]^2 - [\vec{p}_T(A_1) - \vec{p}_T(A_2)]^2$$

For $t\bar{t}$ endpoint
 $[M(t)^2 - M(W)^2] / M(t)$
 $\sim 135 \text{ GeV}, m(b, \nu) \sim 0$

For $\tilde{b} \rightarrow b \tilde{\chi}$ endpoint
 $[M(\tilde{b})^2 - M(\tilde{\chi})^2] / M(\tilde{b})$
 $\sim 260 \text{ GeV}, m(b) \sim 0$



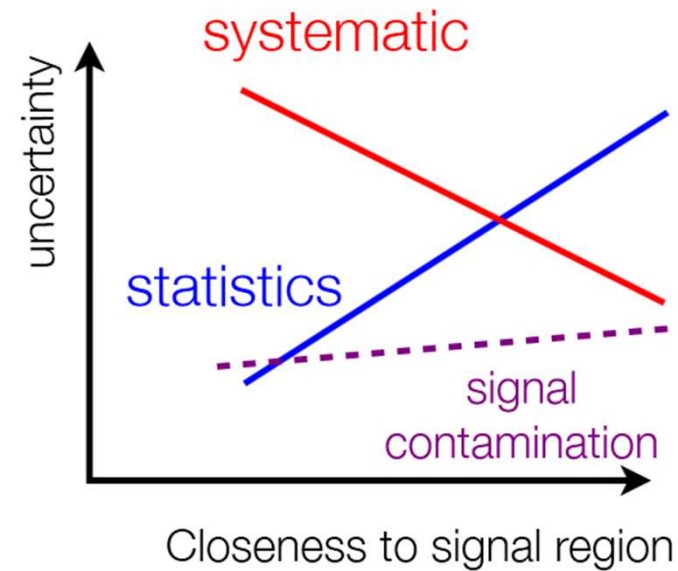
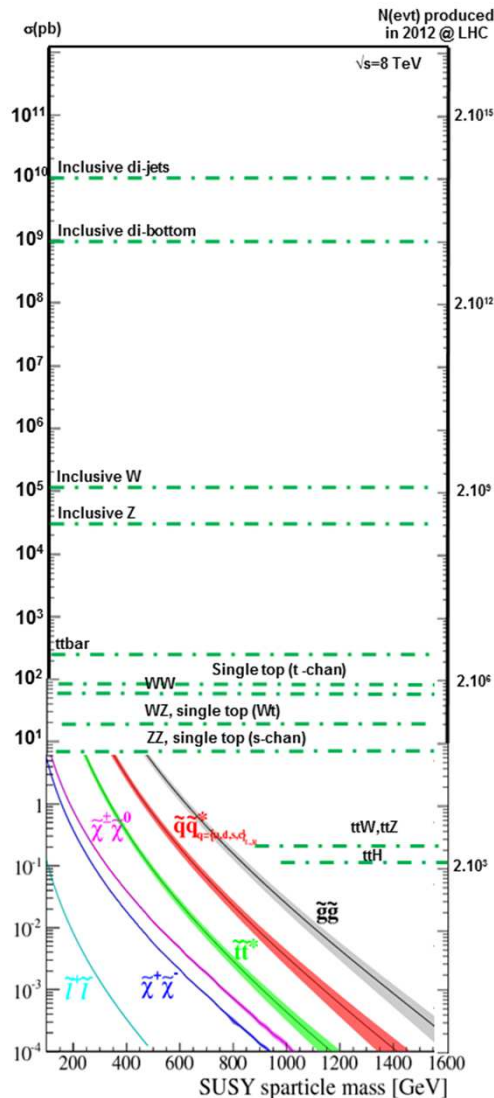
JHEP 0804 (2008) 034, JHEP 1003 (2010) 030



→ Powerfull to reject SM background but need to assume value of endpoint to cut !

*Originally designed to measure SUSY masses

Background estimation



- 1) Need to suppress QCD / W,Z / top by $\sim 10^{10} / 10^5 / 10^2$
- 2) Estimate small remaining quantities
- 3) Fit results
- 4) Interpret the results if no excess

Background estimation (1)

□ Different strategies for different background

- Note: out of the box Data-Monte Carlo agreement is generally very good at LHC

Pure MC

Methods : none !

Pros: Easy, helpful to start and design Signal Regions

Cons: Suffer from large syst and/or statistical errors

Targets: Well suited for small backgrounds

Semi data-driven

Methods : i) isolate a pure background sample, ii) normalise MC iii) assume MC shape to transfer it to Signal Region

Pros: Main systematics cancel in the transfer factor

Cons: full study of possible theory systematics

Targets: Main irreducible background (top, W/Z+jets)

Fully data-driven

Methods : a lot !

Pros: i) Don't rely on potential failures in simulation, ii) Suited for large σ

Cons: Rely strongly on simplifying assumptions → systematics

Targets: Fake MET (QCD, Z+jets), fake leptons, long-lived particle (high pT muons with mis-measured β)

→ Precision in background determination drives the SUSY sensitivity

Background estimation (2)

ATLAS-CONF-2012-033

□ Take the example of 0lepton + ≥ 4 jets + MET channel

- Production : $\tilde{g}\tilde{g}$
- Decay : $\tilde{g} \rightarrow \tilde{q}q$ with $\tilde{q} \rightarrow q\tilde{\chi}$ dominates $\rightarrow$ 0lepton + ≥ 4 jets + MET
- Discriminant variable : $M_{\text{Eff}} > 1200 \text{ GeV}$

Signal Region (SR) Definition: Lepton veto $p_T(e/\mu) > 20/10 \text{ GeV}$

Requirement	Channel	
	C 4j	
$E_T^{\text{miss}} [\text{GeV}] >$	160	Trigger-driven
$p_T(j_1) [\text{GeV}] >$	130	
$p_T(j_2) [\text{GeV}] >$	60	
$p_T(j_3) [\text{GeV}] >$	60	Pile-up driven
$p_T(j_4) [\text{GeV}] >$	60	
$p_T(j_5) [\text{GeV}] >$	–	
$p_T(j_6) [\text{GeV}] >$	–	QCD rejection -driven
$\Delta\phi(\text{jet}, E_T^{\text{miss}})_{\min} >$	0.4 ($i = \{1, 2, 3\}$)	
$E_T^{\text{miss}} / m_{\text{eff}}(Nj) >$	0.25 (4j)	
$m_{\text{eff}}(\text{incl.}) [\text{GeV}] >$	1200	Discriminating variable

$\rightarrow$ 3 main backgrounds: QCD, $Z \rightarrow \nu\nu$ +jets, $[t\bar{t} \rightarrow bWbW \rightarrow b\bar{b}jj \& W \rightarrow l\nu]$

Background estimation (3)

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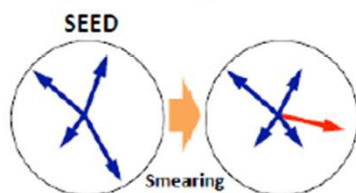
□ QCD/Multijets background (Data driven)

- Enter Signal Region because of fake MET or ν inside jet
- Can not trust MC + limited by MC stat
 - ✓ Compute jet response $R = p_T(\text{jet reco})/p_T(\text{jet true})$ and generate pseudo data to populate SR

1. Determine the jet response function R from dijet balance and 3-jets mercedes events (1jet aligned with MET)

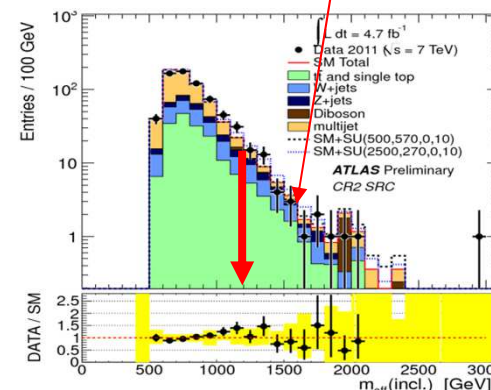
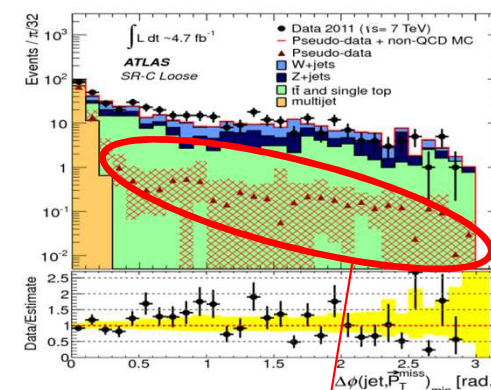
2. Take a control sample of multijets events with small MET.

3. Smear each jet by its response $R \rightarrow$ Pseudo-data



4. Normalize the shape obtained in a QCD enhanced region with low $\Delta\phi(\text{jet}, E_T^{\text{miss}}) < 0.4$

5. Propagate to signal region



Background estimation (4)

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□ $Z \rightarrow \nu\nu$ +jets (Data Driven)

- Enter SR because it is exactly signal like: Irreducible background !

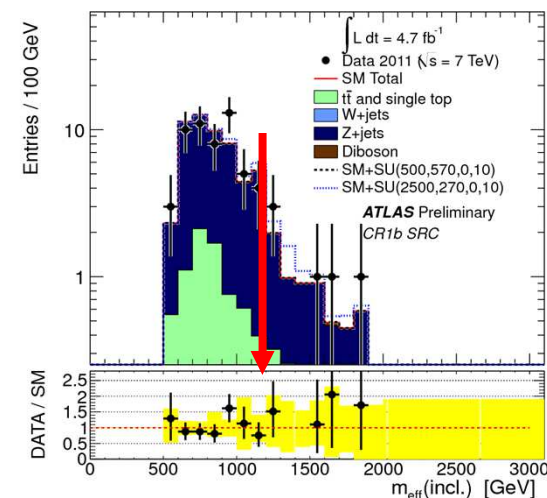
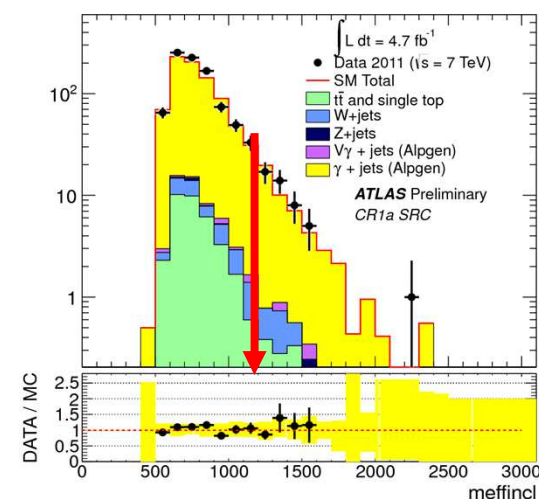
1. Use a close-by SM process: γ +jets

- ✓ Similar kinematic at $p_T \sim 400 \text{ GeV} \gg m_Z$
- Obtain a very pure sample
- ✓ Force the photon as MET
- ✓ Gain a factor ~ 3 in stat: $R = \sigma(Z+\text{jets})/\sigma(\gamma+\text{jets}) \sim 0.3$

2. Use a close-by SM process: $Z \rightarrow \ell\ell$ +jets

- ✓ More statistically limited (~ 10 times less than γ +jets)
- ✓ Will not consider it in the following

γ +4jets Control Region



Background estimation (5)

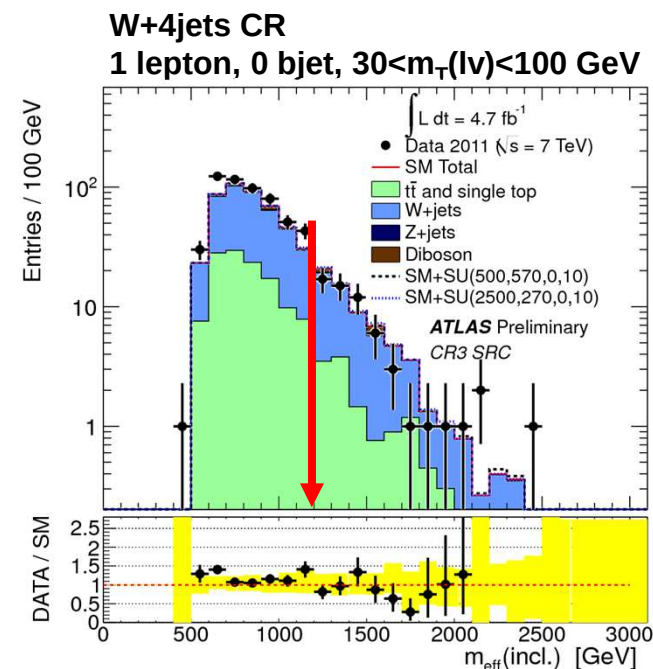
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□ $W \rightarrow l\nu$ + jets and $t\bar{t} \rightarrow b\bar{b}q\bar{q}$ (Semi Data-driven)

- Enter Signal Region because lepton is reconstructed as a jet, is τ , out of acceptance
- Have ν (real MET): can trust MC
 - ✓ Define enriched background “control” region (CR) by **reverting a cut** (Ex: ask 1lepton for 0lepton ch.)
 - ✓ Force the lepton as a jet (acceptable approximation)
- Look in the Control Region:
 - ✓ Monte Carlo should reproduce the data
 - ✓ High **Purity** ($N_{MC}^{oth} \sim \text{small}$), small Signal contamination
- Estimate N_{SR}^{bkg} **Transfer factor (c)** relying on MC shape:

$$N_{SR}^{Bkg} = \frac{N_{SR}^{MC}}{N_{CR}^{MC}} (N_{CR}^{data} - N_{CR}^{MC, others}) = N_{SR}^{MC} \left(\frac{N_{CR}^{data} - N_{CR}^{MC, others}}{N_{CR}^{MC}} \right)$$

Transfer Factor $c_{CR \rightarrow SR}$ **Scale factor** ($k \sim 1$)



➔ Systematics partially cancel in the ratio, but need small extrapolation ($c \sim 0.1-1$)

Background estimation (6)

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□ Summary: $SR=0\text{lepton} + \geq 4\text{jets} + \text{MET} + M_{\text{eff}}(\text{incl.}) > 1200 \text{ GeV}$

■ $t\bar{t}b\text{ar}+\text{jets}$

■ $W+\text{jets}$

■ $\gamma+\text{jets}$

■ QCD

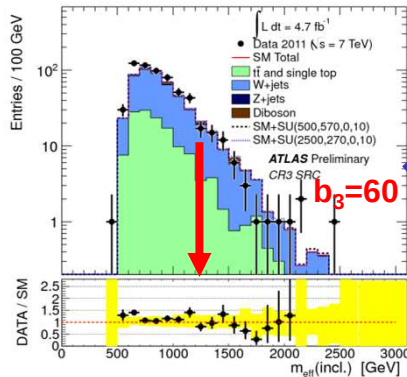
c = Transfer factor

p = purity

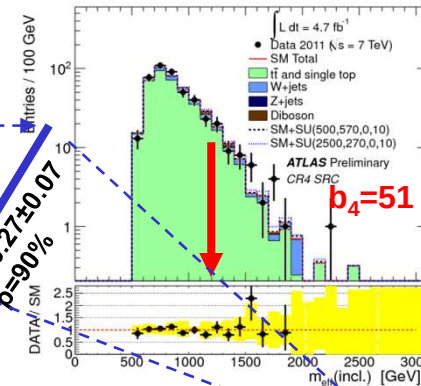
→ $CR \rightarrow SR$

--- $CRa \rightarrow CRb$

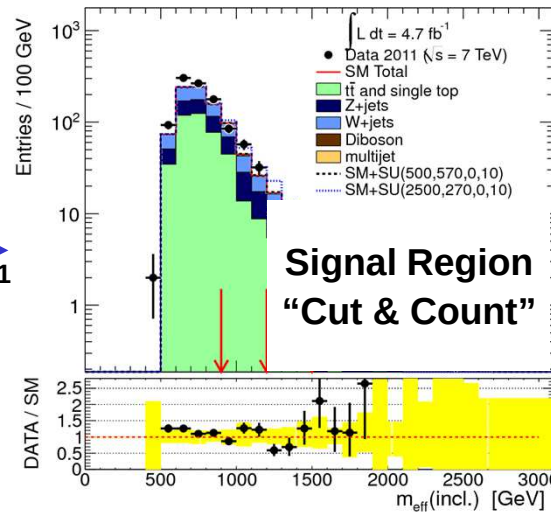
W+4jets CR3 (1 lepton, 0 bjet, mT)



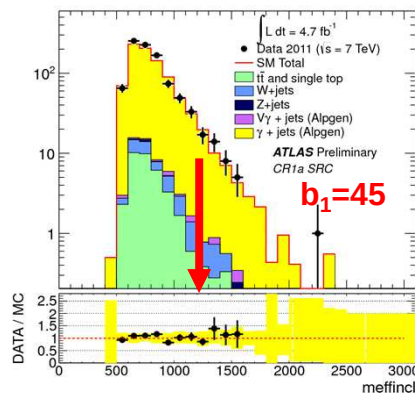
Top+4jets CR4 (1 lepton, 1 bjet, mT)



SR (0lepton+4jets+MET)



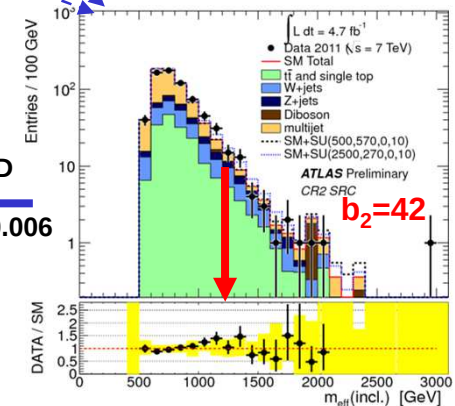
$\gamma+4\text{jets}$ CR1a



$\gamma \rightarrow Z$

$c=0.37 \pm 0.11$
 $p \sim 100\%$

QCD CR2 ($\Delta\phi(j, \text{MET}) < 0.4$)

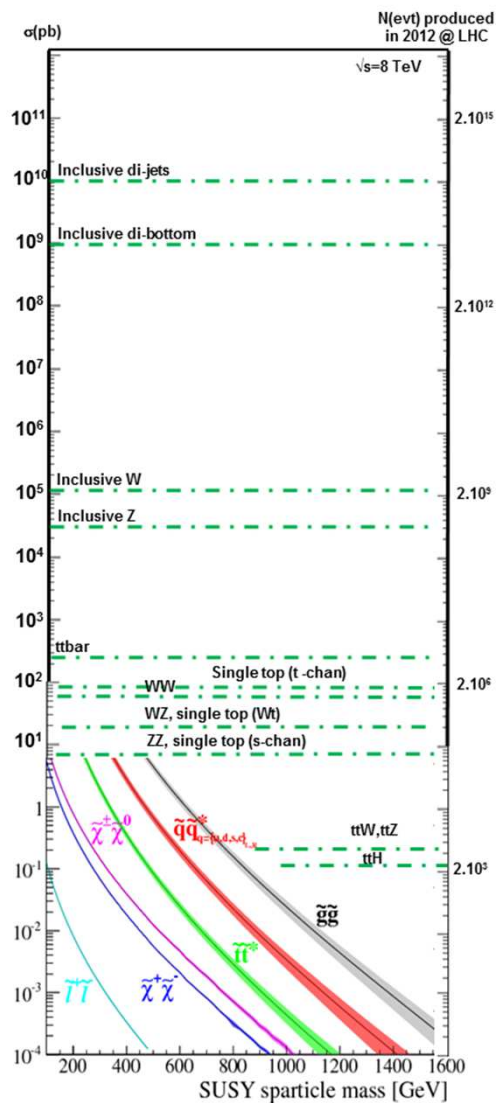


QCD $\leftarrow$ QCD

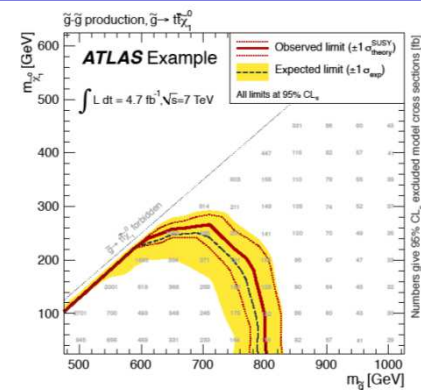
$c=0.0026 \pm 0.006$
 $p=15\%$

→ Errors contains exp. (Jet Energy scale, btagging) and theo. (PDF, scale) syst.

Fit Results



OR



- 1) Need to suppress QCD / W,Z / top by $\sim 10^{10} / 10^5 / 10^2$
- 2) Estimate small remaining quantities
- 3) Fit results
- 4) Interpret the results if no excess

Fit Results (1)

ATLAS-CONF-2012-033

□ Building the likelihood

- Likelihood function : products of Poisson pdf* for SR and CR (as mutually exclusive) & syst.

$$L(n | \mu, b, \theta) = P_{SR} \times P_Z \times P_W \times P_{Top} \times P_{QCD} \times C_{syst}$$

n = Number of observed events in data

μ = SUSY signal strength to be tested

b = background

θ = systematics treated with Gaussian **Nuisance param.**

Free param.

- Inputs: Transfer factors** (**c**), data events in SR (**s**) and CR_j (**b_j**)

$$P_{SR} = P(n | \lambda_S(\mu, b, \theta)) = \mu \cdot c_{SR \rightarrow SR}(\theta) \cdot s + \sum_j c_{jR \rightarrow SR}(\theta) \cdot b_j$$

$$P_i = P(n | \lambda_i(\mu, b, \theta)) = \mu \cdot c_{SR \rightarrow iR}(\theta) \cdot s + \sum_j c_{jR \rightarrow iR}(\theta) \cdot b_j$$

$\lambda(\mu, b, \theta)$ = expected number of events

C_{CR,SR→SR}

Region	Main CR/Process			
	CR1a / Z/γ+jets	CR2 / QCD jets	CR4 / tτ+ Single Top	CR3 / W+jets
CR1a	1	0	0	0
CR2	0.1	1	0.39	0.2
CR4	0.0034	0	1	0.093
CR3	0.0078	0	0.32	1
SR	0.37	0.0026	0.27	0.2

→ Can correctly take the systematic correlation and cross-contamination into account by doing a simultaneous fit of all regions

*pdf=probability density function

**can be replaced by scale factor

Fit Results (2)

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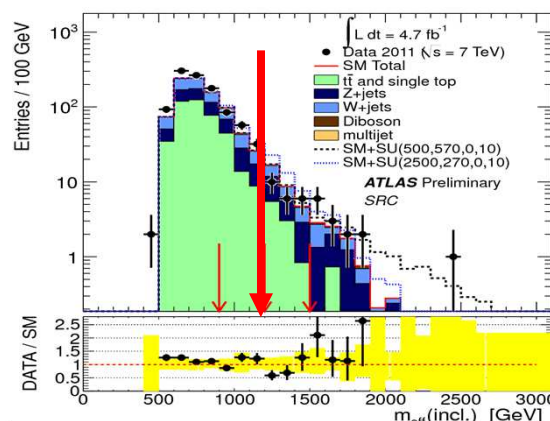
□ Background-only fit ($\mu=0$)

- Predict the background in the Signal Region (SR) by maximizing the likelihood
 - ✓ Give value to **free** and **nuisance** parameters [including statistical error]
 - ✓ Cross-checks of the extrapolation are done in “validation” region, close to SR
- Data in SR not in the fit + no signal contamination in CR (can be reproduced by theorists)

	Background in SR $c_{JR \rightarrow SR}(\theta) \cdot b_j$				Others	Total Background in SR
	Zvν+jets	QCD	W+jets	Top	Dibosons	SR
MC	16	0.01	11	10	1.7	39
Fit Output	17±6	0.02±0.03	8±3	12±5	1.7±0.9	39±9 [±5(stat)±7(syst)]

→ **25% error** (mainly from γ/Z acceptance, CR stat)

→ Observed **36** evts in Data. No Excess !
Now let's quantify it (discovery fit).



Fit Results (3)

ATLAS-CONF-2012-033

□ Discovery fit: background fit + N(data) in the signal region

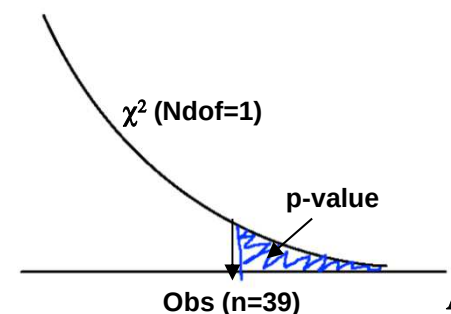
- **Goal:** Test the compatibility of data with background-only hypothesis in the signal region
- **Test statistic:** based on one-sided profile log likelihood ratio (a la Higgs)

$$\Lambda(\mu) = -2 \times [L(n | \mu, \hat{b}, \hat{\theta}) - L(n | \hat{\mu}, \hat{b}, \hat{\theta})] \sim \chi^2 \text{ dist with Ndof} = 1^* (\mu \geq 0)$$

Maximise L for a choice of μ (here $\mu=0$) *Maximise L for μ floating*

*In practice this approximation works well for sufficient stat ($n>5$). If not the case, use toys

- Compute a p-value :
 - ✓ Evaluate the probability for the observation to be signal-like
- In $\text{Olepton} + \geq 4\text{jets} + \text{MET} + M_{\text{eff}}(\text{incl.}) > 1200 \text{ GeV}$:



Predict 39 ± 9 and observe 36
p-value=0.6 (-0.2 σ). Compatible !

Fit Results (4)

ATLAS-CONF-2012-033

□ From discovery fit, can derive a model independent limit

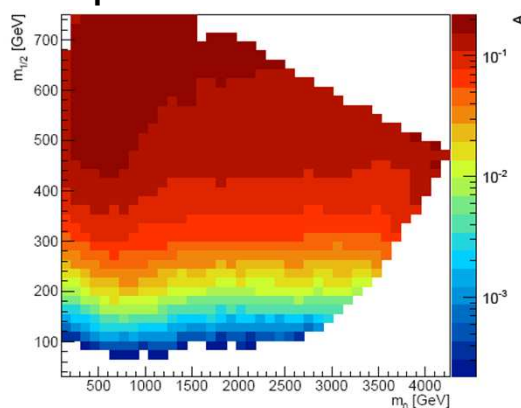
- Limit on visible cross-section of non-SM process: $\sigma_{\text{vis}} = \sigma \times A \times \epsilon$
- In $0\text{lepton} + \geq 4\text{jets} + \text{MET} + M_{\text{eff}}(\text{incl.}) > 1200 \text{ GeV}$:

Predict 39 ± 9 , observe 36 $\rightarrow$

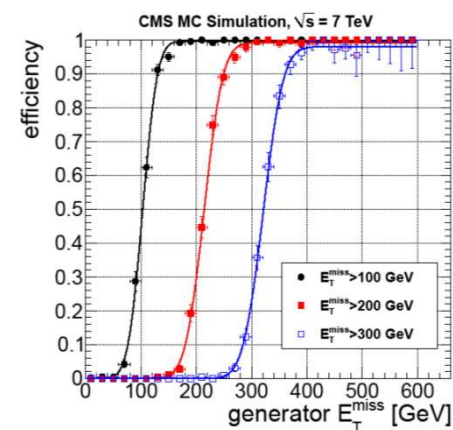
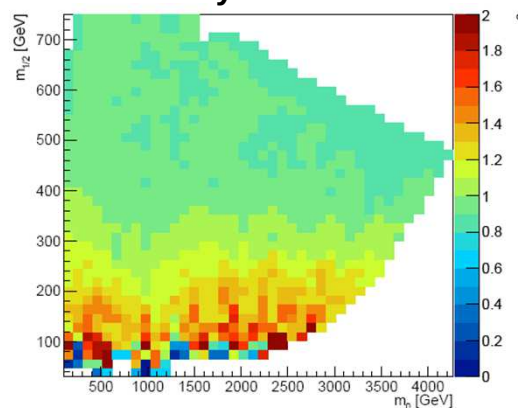
- $\rightarrow$ Exclude at 95%CL $N(\text{BSM}) \geq 18$ and $N/L = \sigma_{\text{vis}} > 3.7 \text{ fb}$
- $\rightarrow$ Expected to exclude $N(\text{BSM}) \geq 19$ and $N/L = \sigma_{\text{vis}} > 4.1 \text{ fb}$

- A and ϵ given for a well-defined SUSY model : Examples below

Acceptance of Truth cuts ~ 0.1 -10%



Efficiency wrt Truth ~ 1



$\rightarrow$ Result can be recasted in other models than the one considered

Fit Results (5)

ATLAS-CONF-2012-033

□ Exclusion fit : discovery fit with $\mu=1$ and a given SUSY Model

- Exclude (or not) a given SUSY model (with well defined cross-section).
- Test statistic:** based on one-sided profile log likelihood ratio (a la Higgs)

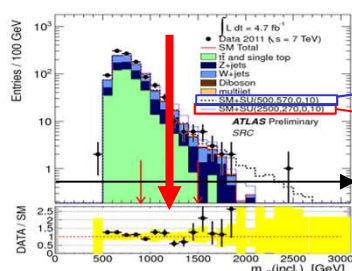
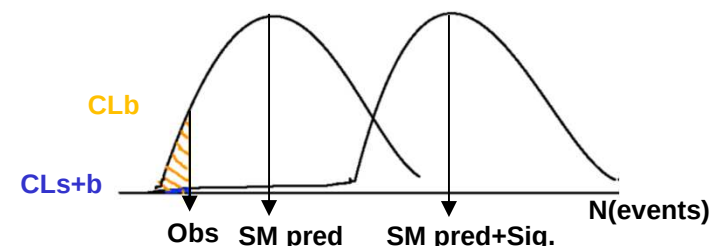
$$\Lambda(\mu) = -2 \times [L(n | \mu, \hat{\hat{b}}, \hat{\hat{\theta}}) - L(n | \hat{\mu}, \hat{b}, \hat{\theta})] \sim \chi^2 \text{ dist with Ndof} = 1^* (\mu \geq 0)$$

Maximise L for a choice of μ
(here $\mu=1$)

Maximise L for μ floating

*In practice this approximation works well for sufficient stat. ($n > 5$). If not the case, use toys

- Use CLs=CLs+b/CLb prescription (a la Higgs):
 - ✓ Protection from a down-fluctuation of the SM background
- In $0\text{lepton} + \geq 4\text{jets} + \text{MET} + M_{\text{eff}}(\text{incl.}) > 1200 \text{ GeV}$:
 - ✓ CLs **p-value** < 0.05 $\rightarrow$ exclude the model at 95 % CL



NS=5 $\rightarrow$ CLs p-value=0.4. Not excluded

NS=14 $\rightarrow$ CLs p-value=0.1. Not excluded

Fit Results (6)

□ Exclusion limits : a new standard ATLAS/CMS procedure (>June 2012)

- Ease the life of theorist by separating the signal theoretical and experimental systematics

Expected limit:



- **Central value:** all uncertainties included in the fit as nuisance parameters, except theoretical signal uncertainties (PDF,scales)

- **$\pm 1\sigma$ band** : $\pm 1\sigma$ results of the fit

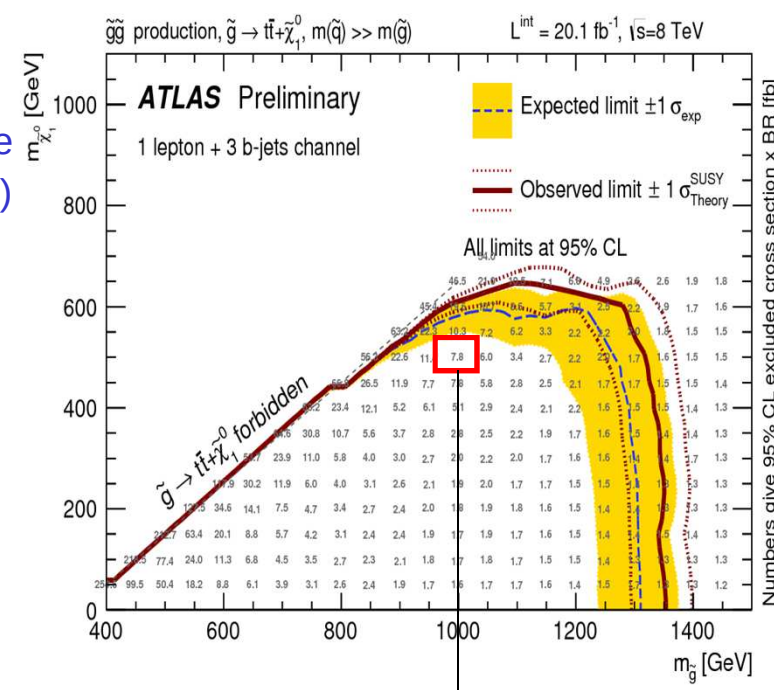
Observed limit:



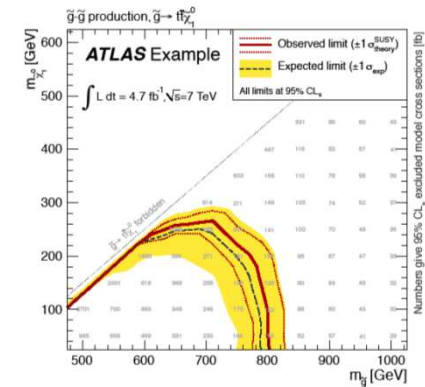
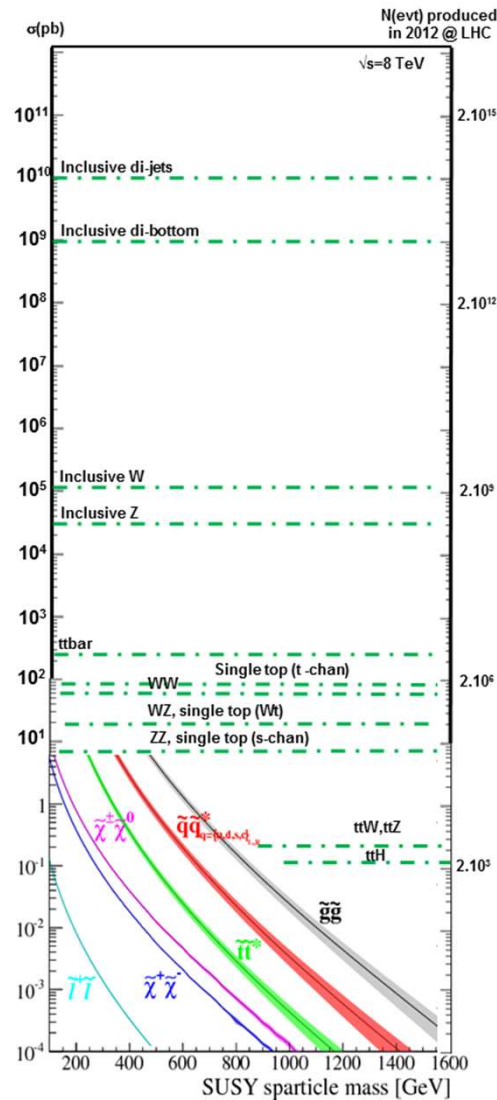
- **Central value:** Idem as for expected limit
- **$\pm 1\sigma$ band** : re-run and increase/decrease the signal cross section by the theoretical signal uncertainties (PDF, scales)

Excluded Model Production Cross section (σ_{excl})

e.g $M(\tilde{g})=1 \text{ TeV}$ $\sigma(\tilde{g}\tilde{g})\sim 30 \text{ fb}$ If $\sigma_{\text{excl}} < \sigma(\tilde{g}\tilde{g})$



Interpretation

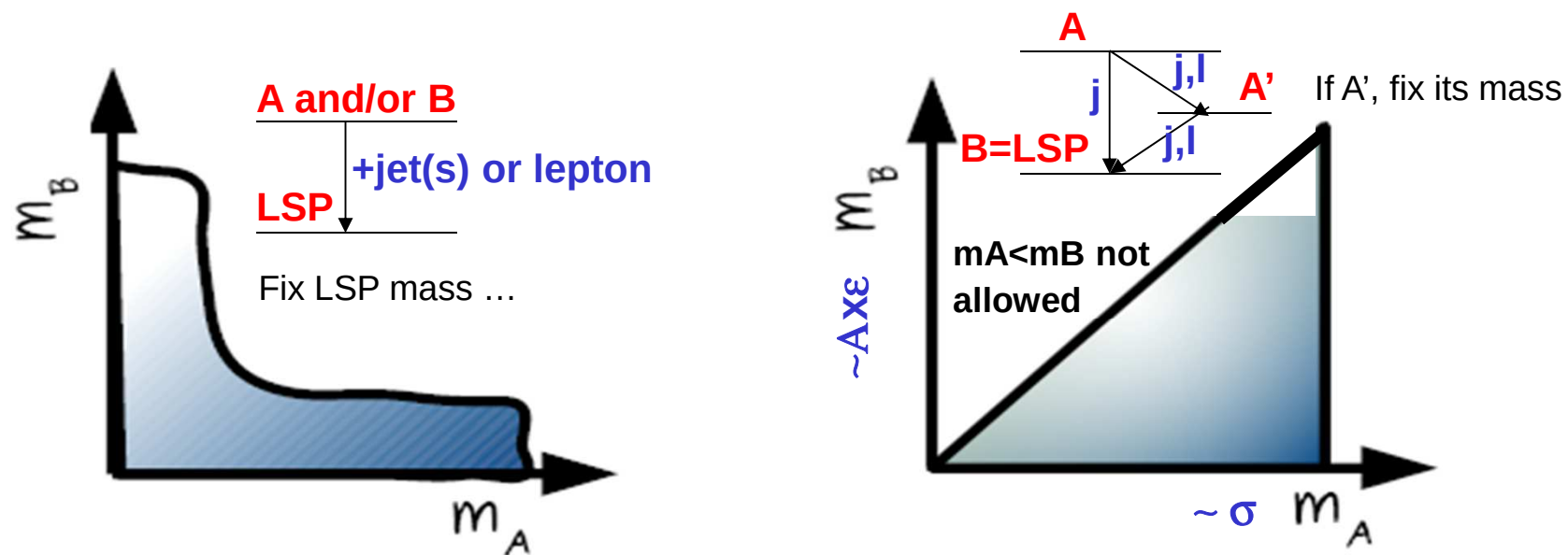


- 1) Estimate small remaining quantities
- 2) Need to suppress QCD / W,Z / top by $\sim 10^{10} / 10^5 / 10^2$
- 3) Fit results
- 4) Interpret the results if no excess

Interpretation (1)

❑ Derive a limit in a simplified decay chain Model (SMS)

- Well suited for **natural SUSY** and **direct production** (not a SUSY model !):
 - ✓ 29 sparticles → 2 or 3, **decoupled** all other particles, force a specific decay mode (BR=100% !)
- Assumptions on the chirality and nature of particle “arbitrary”



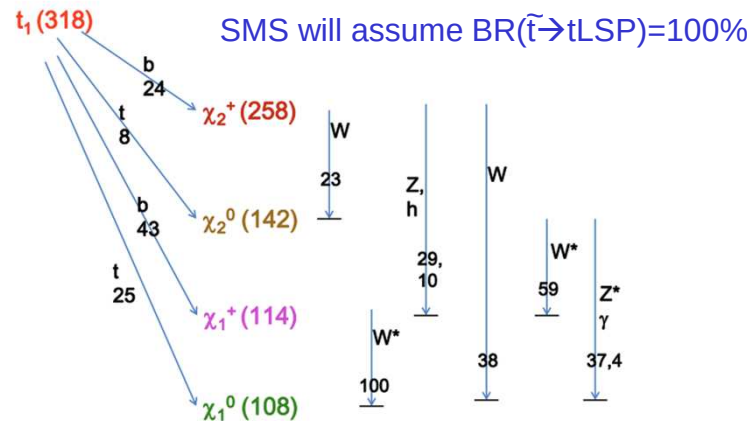
→ Very helpful also to design analyses. Possible to recast in other models

Interpretation (2)

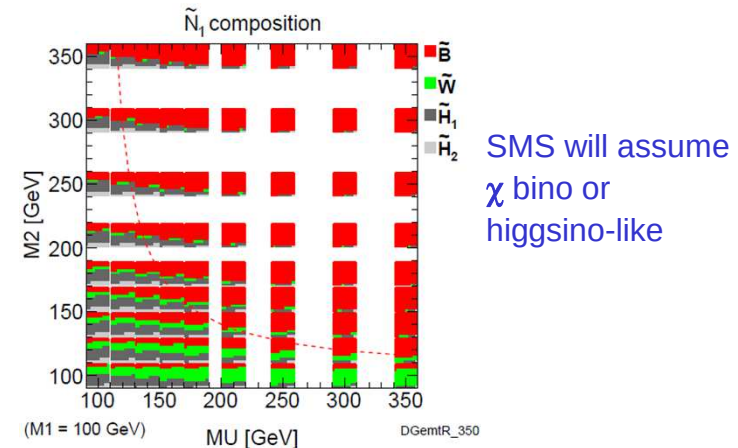
❑ **Derive a limit in a simplified MSSM**

- Reduce number of SUSY parameters from 105 (MSSM) to 19, i.e. “manageable”:
 - ✓ Well justified assumptions → no arbitrary fixing of the BR decay
 - ✓ “Standard” exp. constraints
- Recover the SUSY complexity → can track missing features of SMS in “simple” cases

Direct Stop production

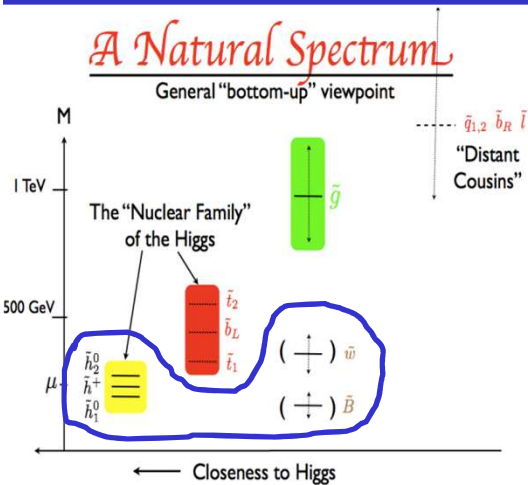


Direct Gaugino production

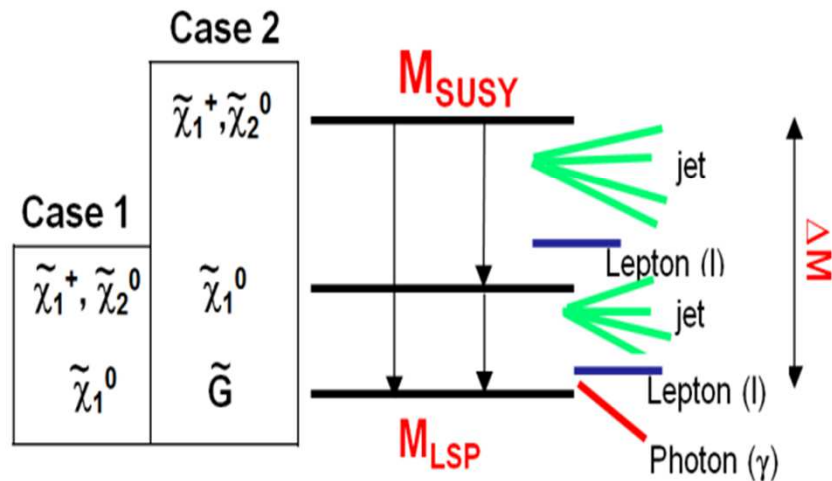


→ Complementary approach to Simplified Model

Lecture Part Ib

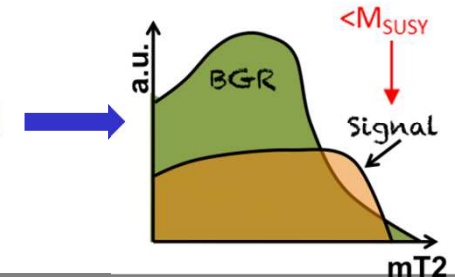


A: MSSM
B: Natural $>10\%$
C1: **RPC**
C2: LSP nature
C3: Spectra opening



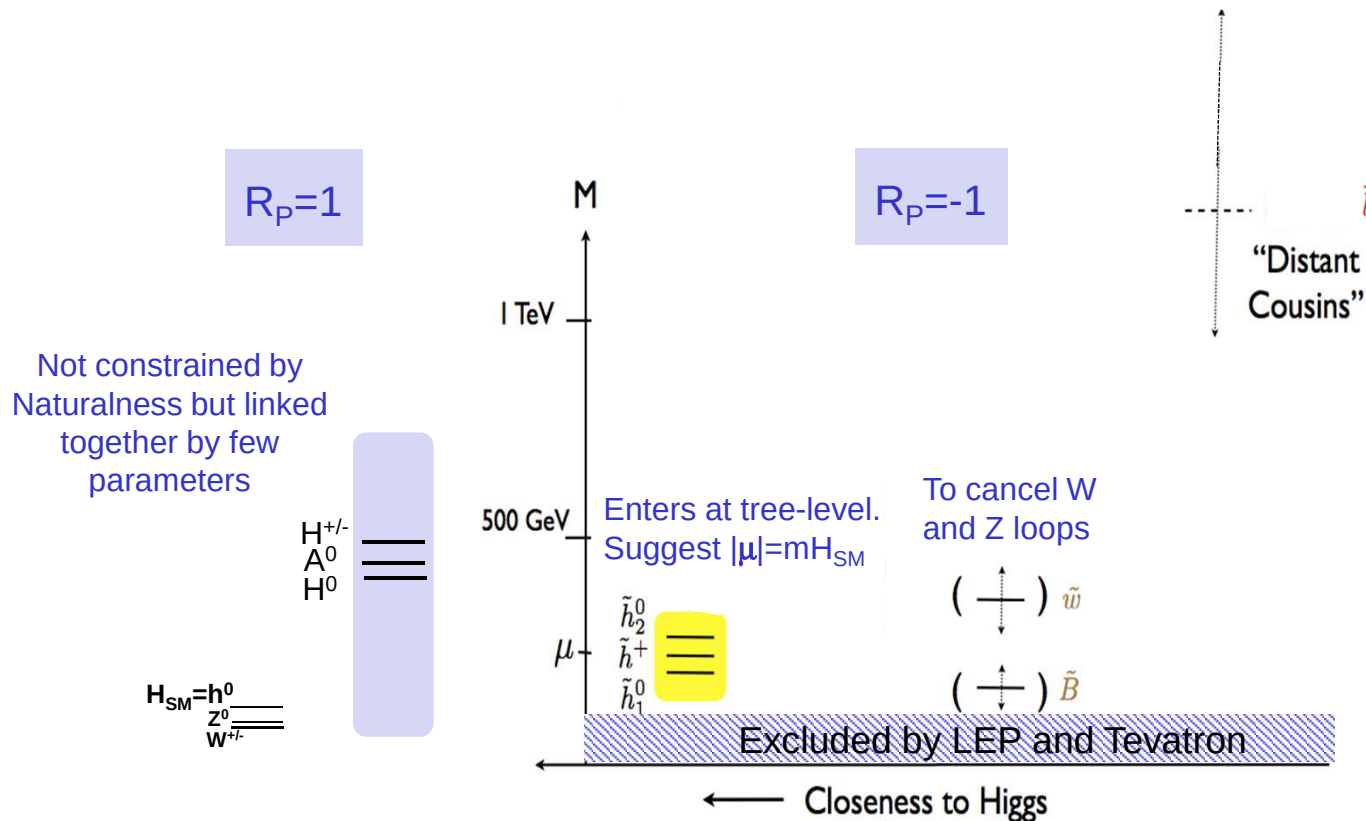
0/1/2 jets
+ E_T^{miss}
+ 1/2/3/4 lepton
+ 0/1/2 photon

- Absolute value of vectorial sum: $E_T^{\text{miss}} \sim \Delta M$
- Number of jets (can be jet-veto !)
- Endpoint at $m_{T2} \sim (\mathbf{M}_{\text{SUSY}}^2 - \mathbf{M}_{\text{LSP}}^2) / \mathbf{M}_{\text{SUSY}}$ [[hep-ph/9906349,0304226](https://arxiv.org/abs/hep-ph/9906349)]



Parameters of the EW sector (1)

Before looking for SUSY particles, let's look at the EW sector at large



Parameters of the EW sector (2)

$R_P=1$

1- Standard Model

$$\phi_{\text{Higgs}} = \left(\frac{1}{\sqrt{2}} (v + H^0 + iH^0) \right)$$

$$\begin{aligned} W^0, B, \text{Im}(H^0) &\longrightarrow \gamma, Z \quad (\theta_W \sim 30^\circ) \\ W^+, W^-, H^+, H^- &\xrightarrow{\text{EWSB}} W^+, W^- \\ \text{Re}H^0 &\longrightarrow H \end{aligned}$$

1 param.
(m_H)

$$V(\phi^\dagger \phi) = -\frac{m_H^2}{2} \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2$$

$$v = \sqrt{m_H^2 / (2\lambda)} = (\sqrt{2} G_F)^{-1/2} \simeq 246 \text{ GeV}$$

2- SUSY

$$H_d = \begin{pmatrix} (v_d + \phi_d^0 + i\chi_d^0)/\sqrt{2} \\ \phi_d^- \end{pmatrix} \quad H_u = \begin{pmatrix} \phi_u^+ \\ (v_u + \phi_u^0 + i\chi_u^0)/\sqrt{2} \end{pmatrix}$$

Names	Spin	P_R	Gauge Eigenstates	Mass Eigenstates
Gauge Bosons	1	+1	W^+, W^-, W^0, B	W^+, W^-, Z^0, γ
Higgs bosons	0	+1	$H_u^0, H_d^0, H_u^+, H_d^-$	$h^0, H^0, A^0, H^\pm$

$$\theta_W + 2 \text{ angles } \alpha, \beta \quad \cos^2(\beta - \alpha) = \frac{m_{h^0}^2 (m_Z^2 - m_{h^0}^2)}{m_A^2 (m_{H^0}^2 - m_{h^0}^2)}, \quad \tan\beta = v_u/v_d$$

$m_H + 1 \text{ mass}$ m_A Masses of other Higgses are related to m_A at tree-level

$R_P=-1$

2- SUSY: Each **gauge** field has a partner with S-1/2 in the vector multiplet

$$\begin{aligned} c_W &= \cos\theta_W \\ s_W &= \sin\theta_W \\ c_\beta &= \cos\beta \\ s_\beta &= \sin\beta \end{aligned} \quad \begin{pmatrix} M_1 & 0 & -c_\beta s_W m_Z & s_\beta s_W m_Z \\ 0 & M_2 & c_\beta c_W m_Z & -s_\beta c_W m_Z \\ -c_\beta s_W m_Z & c_\beta c_W m_Z & 0 & -\mu \\ s_\beta s_W m_Z & -s_\beta c_W m_Z & -\mu & 0 \end{pmatrix}$$

Mixing

Bino, Wino, Higgsino $\rightarrow$ Neutralinos

Names	Spin	P_R	Gauge Eigenstates	Mass Eigenstates
neutralinos	1/2	-1	$\tilde{B}^0, \tilde{W}^0, \tilde{H}_u^0, \tilde{H}_d^0$	$\tilde{N}_1, \tilde{N}_2, \tilde{N}_3, \tilde{N}_4$
charginos	1/2	-1	$\tilde{W}^\pm, \tilde{H}_u^\pm, \tilde{H}_d^\pm$	$\tilde{C}_1^\pm, \tilde{C}_2^\pm$

Charginos

$$\begin{pmatrix} M_2 & \sqrt{2} s_\beta m_W \\ \sqrt{2} c_\beta m_W & \mu \end{pmatrix}$$

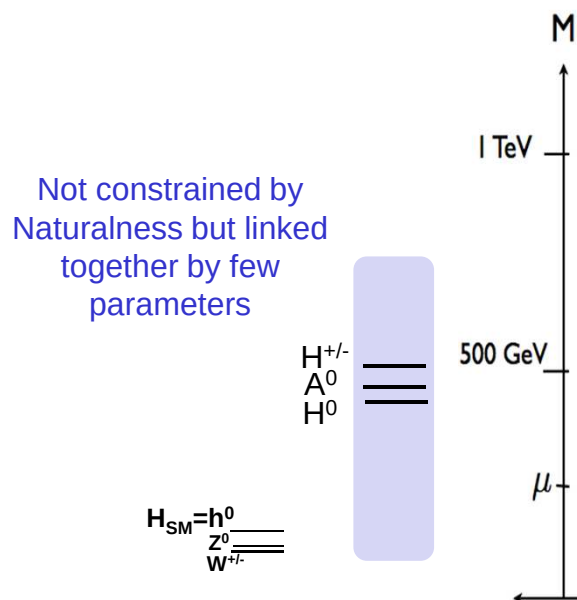
Masses of Gauge Eigenstates

Bino Wino Higgsino

4 param. $M_1, M_2, \mu, \tan\beta$

SUSY Higgses

$R_p=1$



1- Standard Model:

- $m_H=125$ GeV.
- Assume h^0 Higgs discovered (*most probable*)

2- SUSY

$$H_d = \begin{pmatrix} v \\ +i\chi_d^0/\sqrt{2} \end{pmatrix} \quad H_u = \begin{pmatrix} \phi_u^+ \\ (v_u + \phi_u^0 + i\chi_u^0)/\sqrt{2} \end{pmatrix}$$

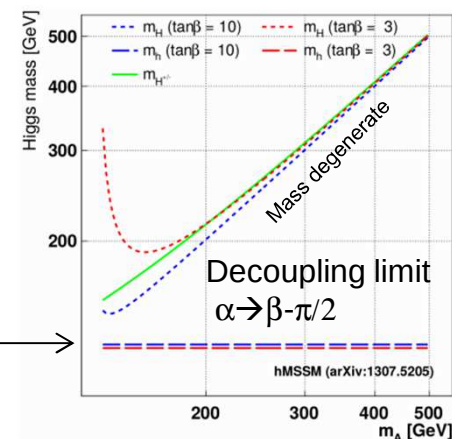
EWSB

Names	Spin	P_R	Gauge Eigenstates	Mass Eigenstates
Gauge Bosons	1	+1	W^+, W^-, W^0, B	W^+, W^-, Z^0, γ
Higgs bosons	0	+1	$H_u^0, H_d^0, H_u^+, H_d^-$	$h^0, H^0, A^0, H^\pm$

$$m_A^*, \tan\beta=v_u/v_d$$

hMSSM, 1307.5205

$$m_{h^0}=125 \text{ GeV}$$

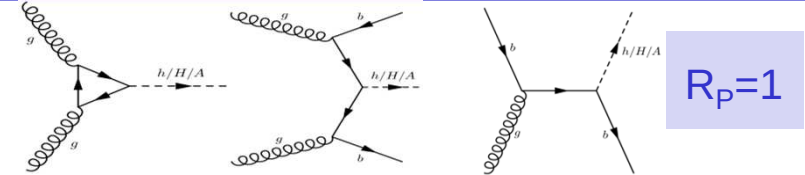


SUSY Higgses (1)

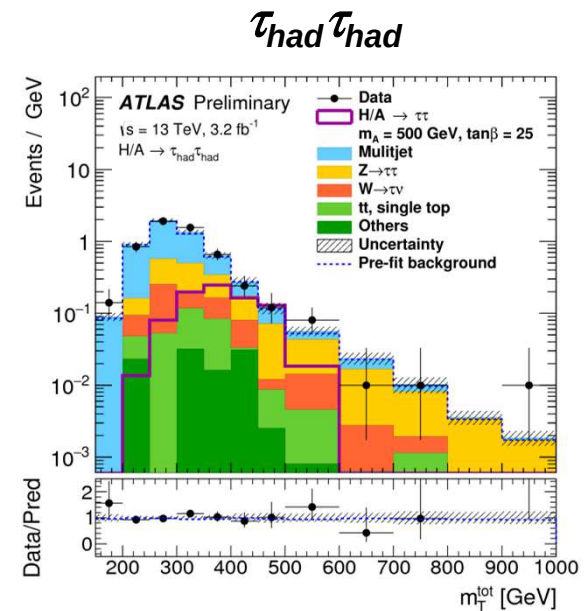
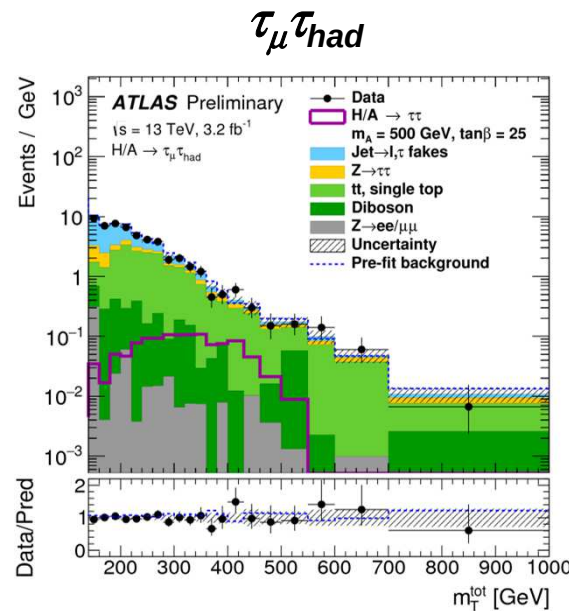
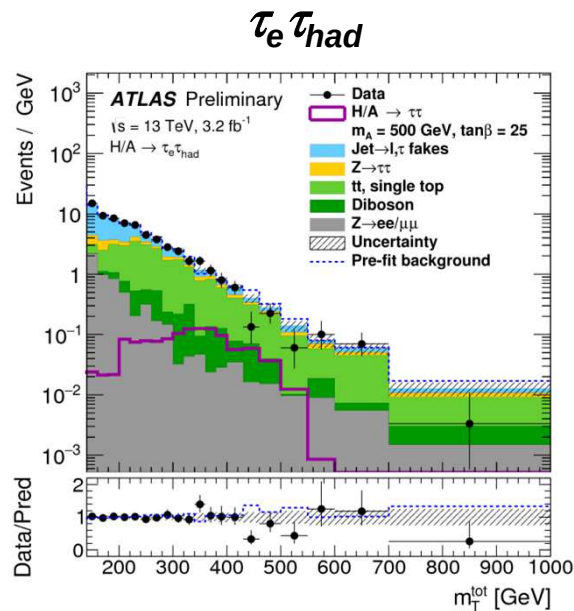
ATLAS-CONF-2015-061

□ Neutral Higgs searches : $\phi=A/H$

- Production: $gg \rightarrow \phi$ or $gg \rightarrow b\phi, bb\phi$
- Coupling (*prefer down-type fermions*) : $\phi \rightarrow \tau\tau$ (10%), bb (90%)
 - ✓ Choose $\tau\tau$ since bb more challenging experimentally
- Fit discriminant variable m_T ($\tau\tau$ system) in 3 channels (could be further split in b-tag and bjet in principle)



$$m_T^{\text{tot}} = \sqrt{m_T^2(E_T^{\text{miss}}, \tau_1) + m_T^2(E_T^{\text{miss}}, \tau_2) + m_T^2(\tau_1, \tau_2)},$$



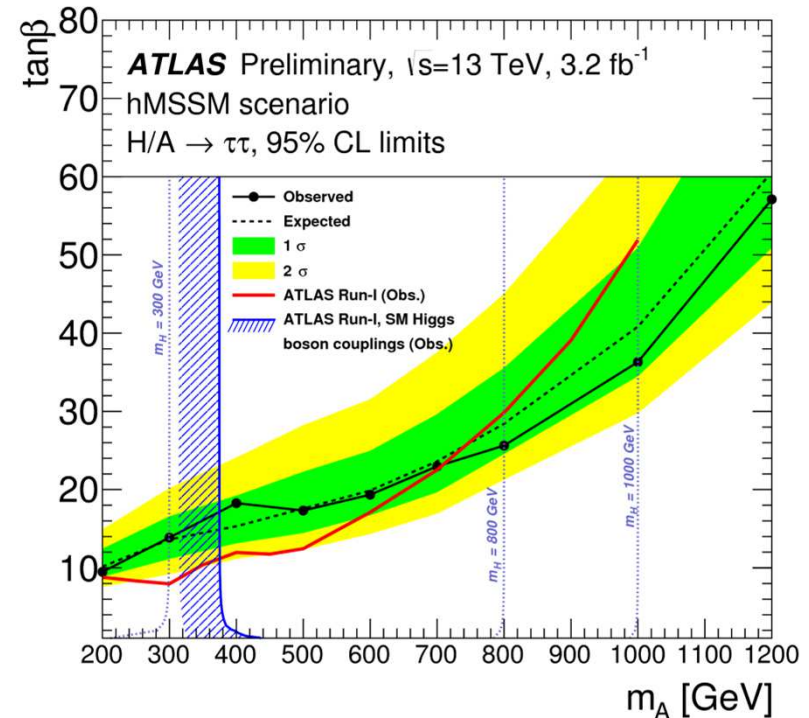
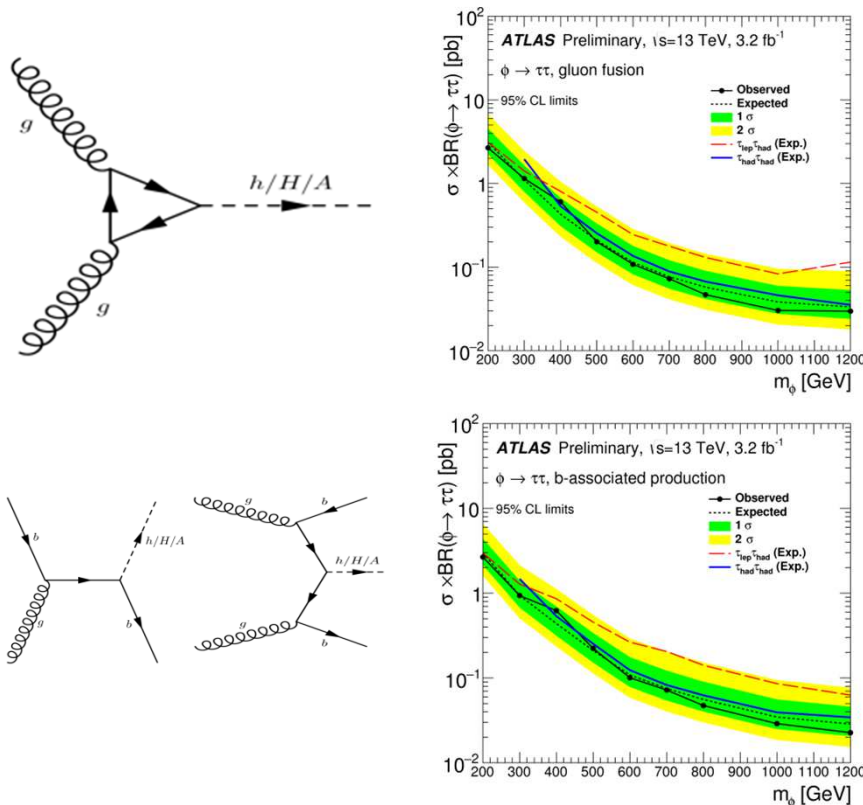
SUSY Higgses (2)

ATLAS-CONF-2015-061

□ Combine all channels $\phi=A/H \rightarrow \tau_{\text{had}}\tau_{\text{had}}, \tau_{\text{had}}\tau_{\text{lep}}$

$R_P=1$

▪ $\tau_{\text{had}}\tau_{\text{had}}$ more sensitive at high ϕ mass (>400 GeV)



Clearly favor high mass neutral SUSY Higgses in hMSSM

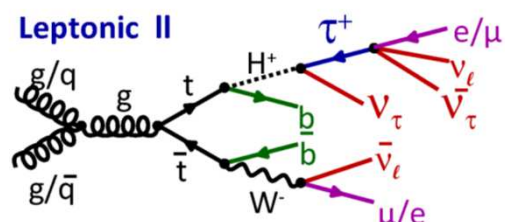
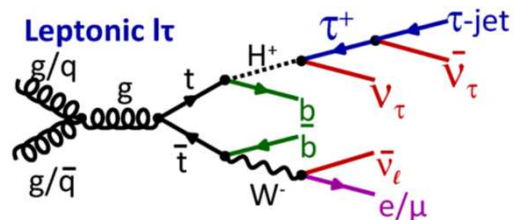
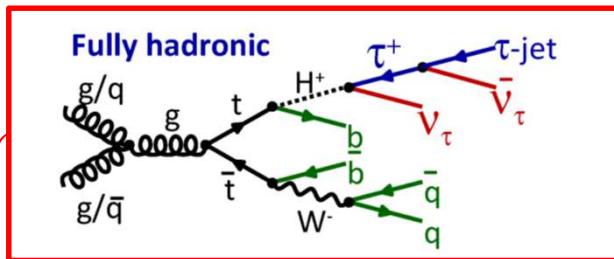
SUSY Higgses (3)

□ Charged Higgs : $H^{+/-}$

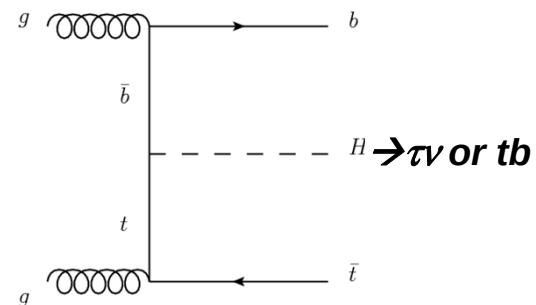
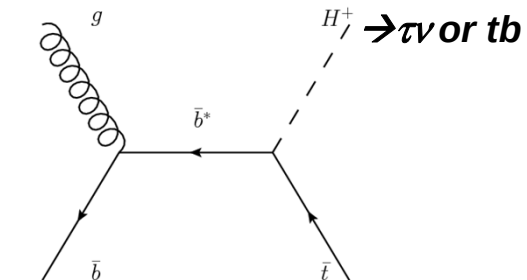
$R_p=1$

- An elementary **charged scalar** particle : clearly indicate new physics beyond SM !
- Coupling (*prefer down-type fermions*) : BR ($H^{+/-} \rightarrow b\bar{t}$) if opened, BR ($H^{+/-} \rightarrow \tau\nu$) for $\tan\beta > 1$
- Production: $m_{H^{+/-}} < m_t$: $t\bar{t}$ events

$m_{H^{+/-}} > m_t$: Associated prod. with b, t



→ **Most powerful mode (see 1204.2760)**



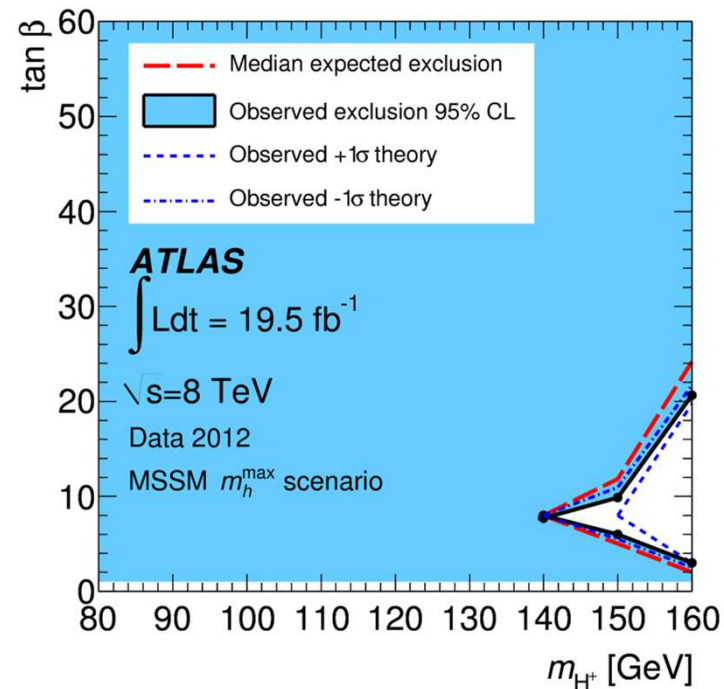
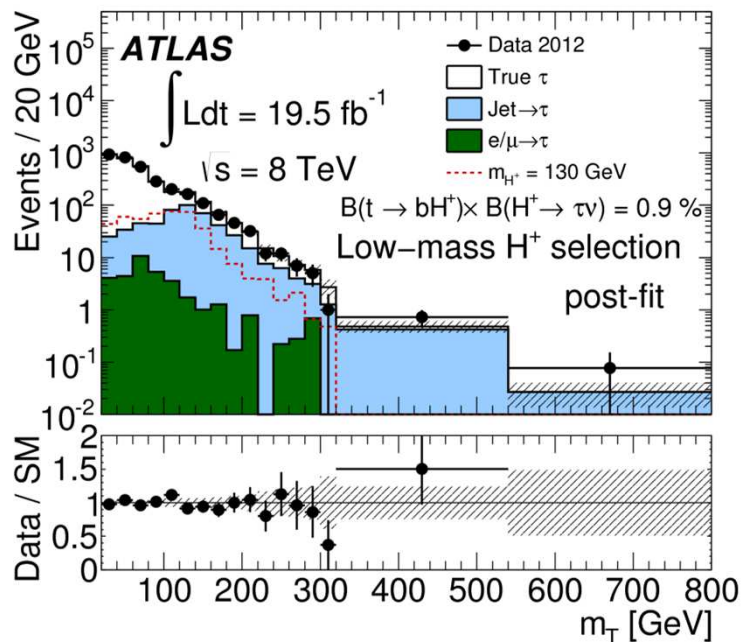
SUSY Higgses (4)

ATLAS: 1412.6663
CMS: 1508.07774

❑ Charged Higgs ($H^{\pm} \rightarrow \tau\nu$) and $m_H < m_t$

$R_p=1$

- =1 hadronic tau, ≥ 1 b-jet, ≥ 4 jets [$W \rightarrow jj$], $MET > 65$ GeV : veto on electron/muon
- Background dominated by fake taus [Data-driven fake factor method]
- Fit discriminant variable $m_T(\tau\nu \text{ system})$: signal will appear as a broad excess



Close the low mass region with Run1 ...

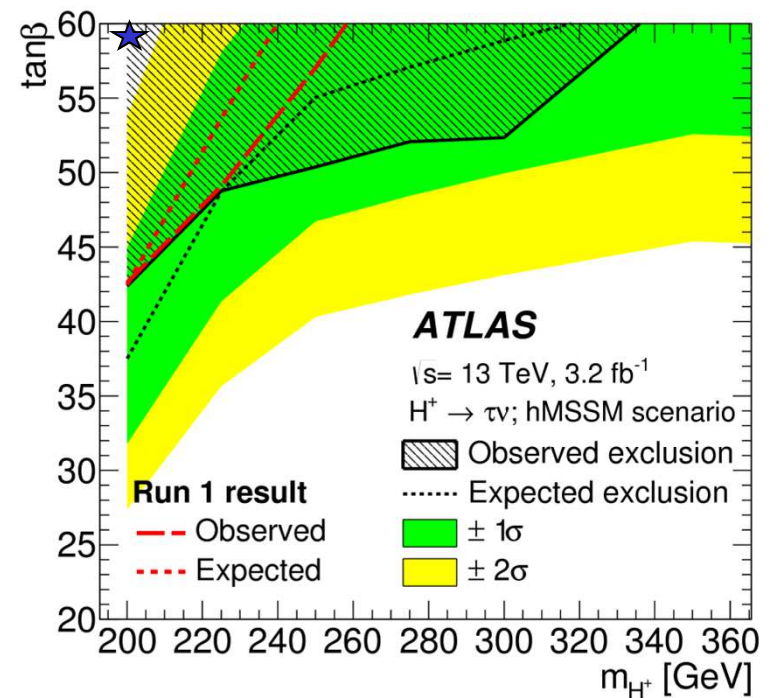
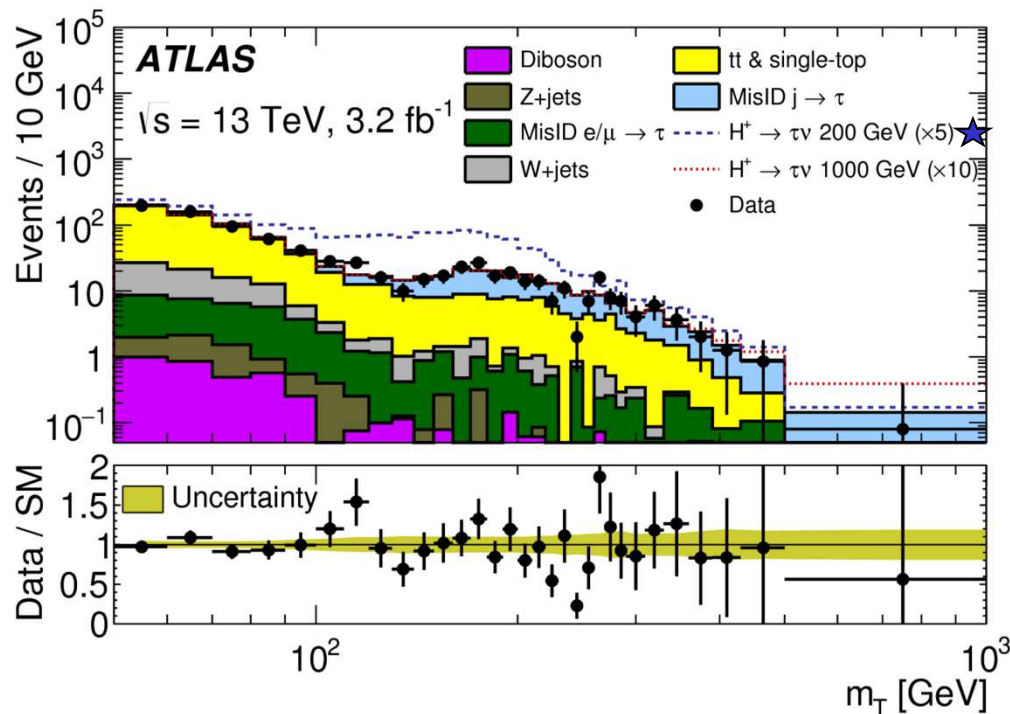
SUSY Higgses (5)

ATLAS: 1603.09203

❑ Charged Higgs ($H^{\pm} \rightarrow \tau\nu$) and $m_H > m_t$

$R_p=1$

- =1 hadronic tau, ≥ 1 b-jet, ≥ 3 jets [$W \rightarrow jj$], $MET > 150$ GeV : veto on electron/muon
- Background dominated by fake taus [Data-driven fake factor method]
- Fit discriminant variable m_T ($\tau\nu$ system) : signal will appear as a broad excess



... and start to be sensitive to higher mass with Run2

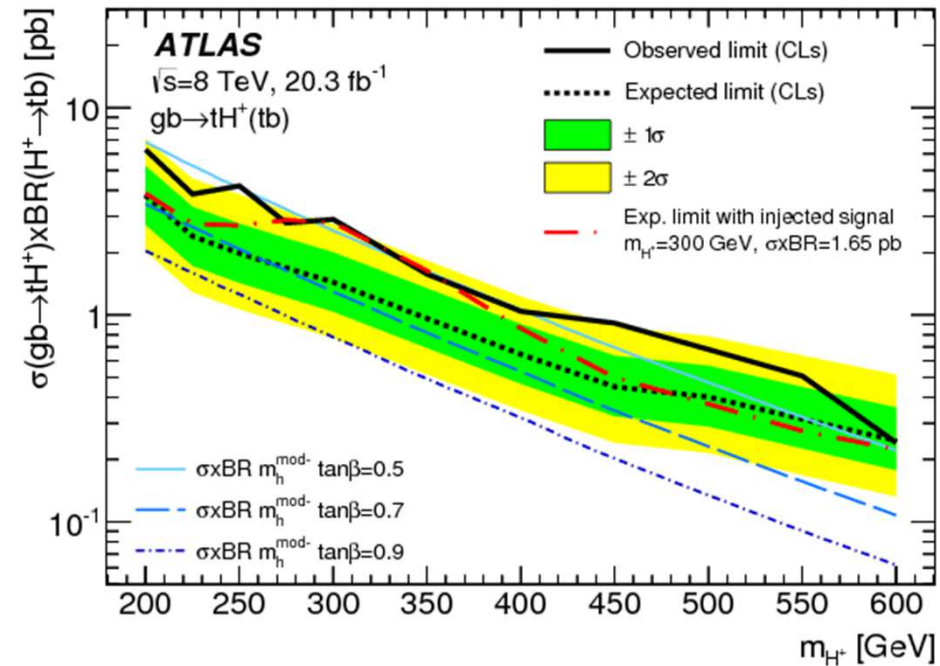
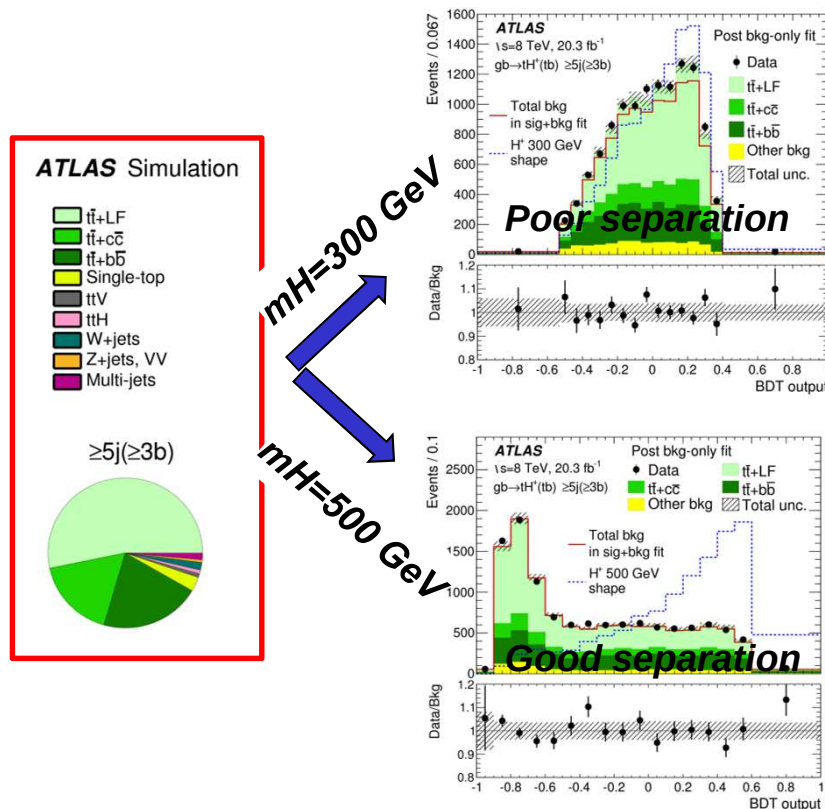
SUSY Higgses (6)

ATLAS: 1512.03704
CMS: 1508.07774

$R_p=1$

❑ Charged Higgs ($H^{\pm} \rightarrow tb$) and $m_H > m_t$

- Very poor signal over background around 300 GeV (even with BDT) ...
- Systematic dominated (by unknowns on $t\bar{t}b$ modelling)



SUSY Higgses (7)

□ Summary

$R_p=1$

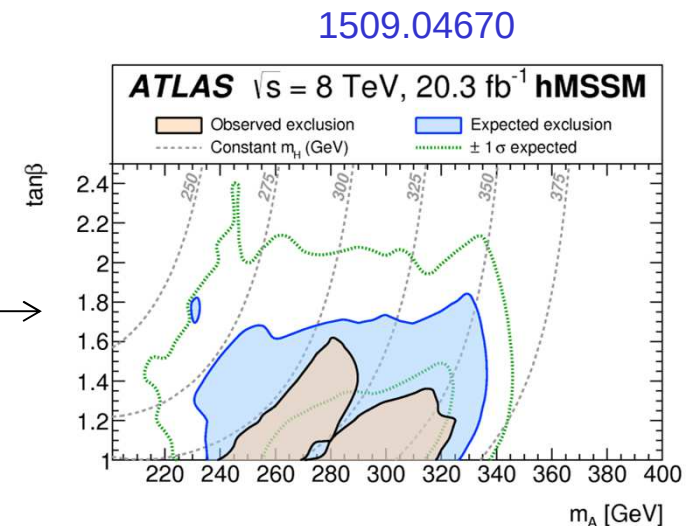
- Reasonable to assume that h^0 is the boson discovered (even if a bit heavy for MSSM)
- SUSY Higgses excluded at low masses and typically $m(A^0/H^0, H^{\pm}) > 200\text{-}300\text{ GeV}$
- Other related searches :

✓ VBF $H^\pm$

✓ Add another neutral scalar **a** like in NMSSM:

- $h \rightarrow 2a \rightarrow 2b2\mu$ [CMS-PAS-HIG-14-041]
- $h \rightarrow 2a \rightarrow 4\mu$ [1506.00424]
- $h \rightarrow 2a \rightarrow 2\mu 2\tau$ [CMS-PAS-HIG-15-011, 1505.01609]
- $h \rightarrow 2a \rightarrow 4\tau$ [CMS-PAS-HIG-14-022]
- $h \rightarrow 2a \rightarrow 4\gamma$ [1509.05051]

✓ diHiggs production, e.g. $H \rightarrow hh \rightarrow bb\tau\tau$



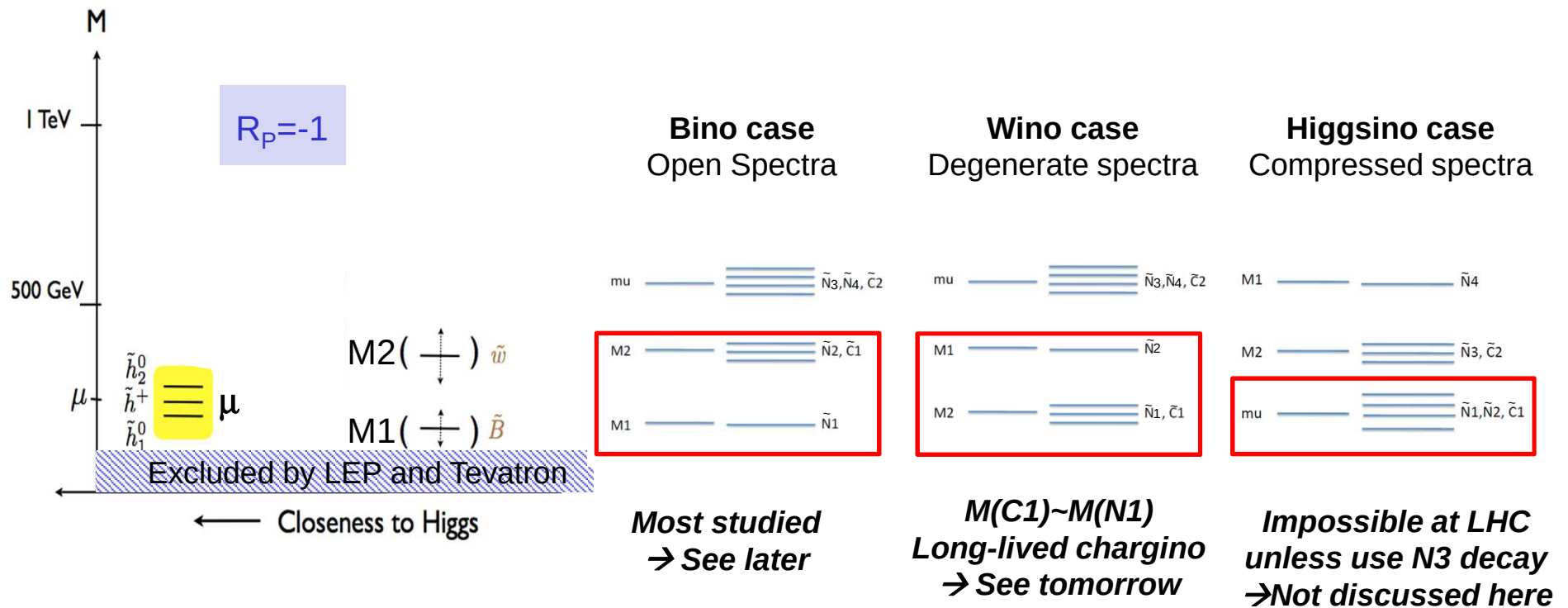
In the following, I assume $A^0/H^0/H^{\pm}$ decoupled

✓ It is conservative: If they happen to be lighter EWKinos cross-section will increase.

EWKinos

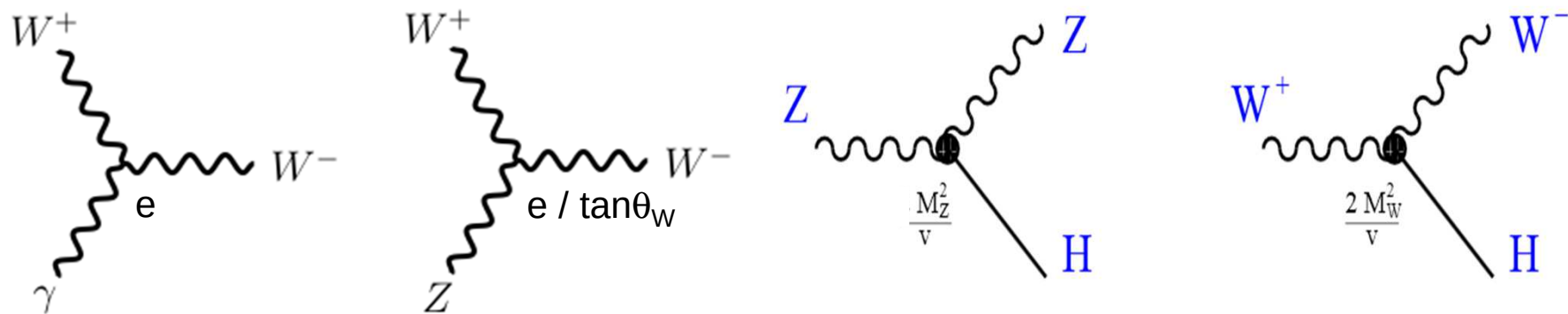
Will mainly focus on lightest states ($\tilde{C}1$, $\tilde{N}1$, $\tilde{N}2$)

A: MSSM
B: Natural >10%
B1: **RPC**
B2: **LSP nature=N1**
B3: Unspecified Spectra



EWKinos (1)

□ Reminder : Triple Gauge coupling in the SM EWK sector



- $\lambda_{WW\gamma} \sim 0.6 \lambda_{WWZ}$
- $\lambda_{ZZH} \sim 0.5 \lambda_{WWH}$
- Neutral gauge coupling forbidden at tree level
- $H\gamma\gamma$ forbidden at tree level

Simple rule: Keep one SM particle and change the 2 others with
 $W \leftrightarrow \tilde{C}1$, $\gamma (H) \leftrightarrow \tilde{N}1$, $Z (H) \leftrightarrow \tilde{N}2$ can get a feel for SUSY couplings

EWKinos (2)

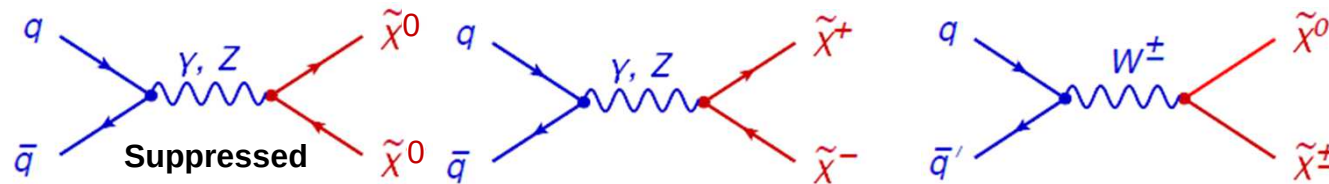
Bino case
Open Spectra

μ ——— $\tilde{N}_3, \tilde{N}_4, \tilde{C}_2$

M2 ——— $\tilde{N}_2, \tilde{C}_1$
M1 ——— $\tilde{N}_1$

□ How to produce them ?

- Gauginos mediated by γ, Z, W (marginally via $h^0/H^0/A^0/H^{\pm-}$)



Note: Coupling depends on Ewkino flavor: (W-Bino, Z/ γ -Wino) ~ 0 , (W/Z/ γ -Higgsino) $\sim$ small, (W-Wino, Z/ γ -Bino) $\sim$ large

□ How they decay (and assumptions) ?

- Most natural is via Gauge bosons and Higgses $\rightarrow$ 1- Only decay via W/Z
- On shell or off-shell depends $\Delta M = M(N_2, C_1) - M(N_1) \rightarrow$ 2- On-shell W/Z $\Delta M > M(Z, W)$
- LHC = hadronic collider $\rightarrow$ 3- W/Z leptonic decays ($l = e, \mu$)

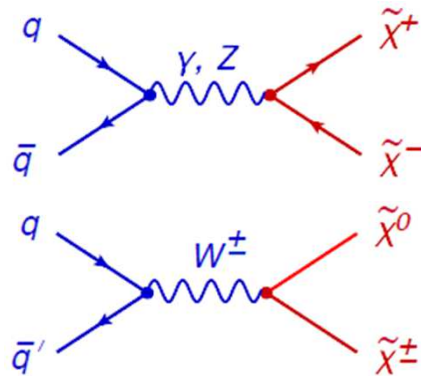
	N1	N2	C1
N1	N1N1		
N2	ZN1N1	ZZN1N1	
C1	WN1N1	WZN1N1	WWN1N1



	N1	N2	C1
N1	ISR-jet		
N2	2l+MET	4l+MET	
C1	1l+MET	3l+MET	2l+MET

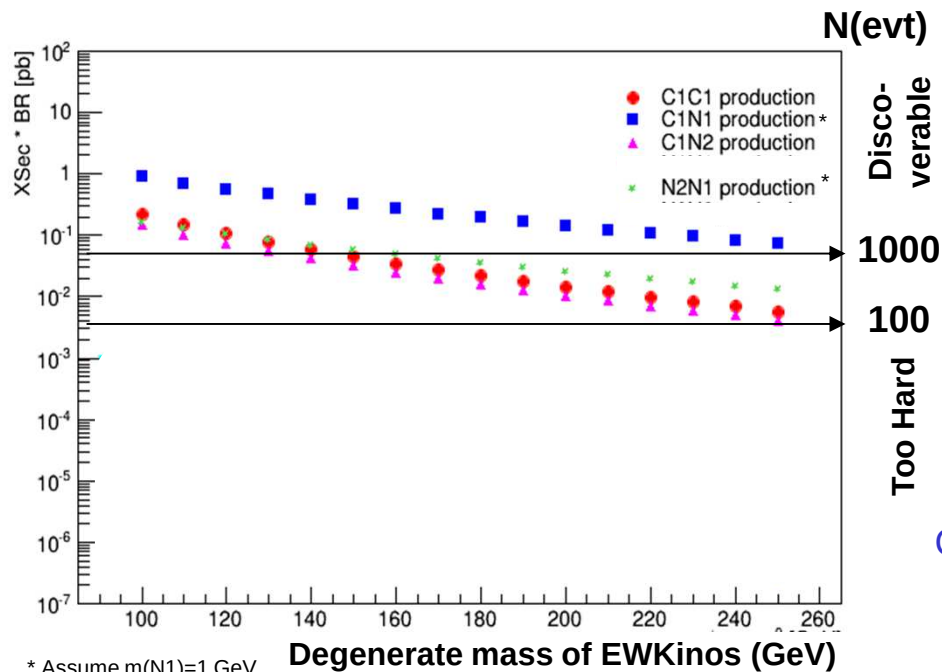
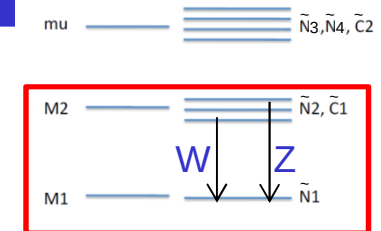
EWKinos (3)

Bino case
Open Spectra



No neutralino mass limit @LHC!

	N1	N2	C1
N1	ISR jet		
N2	2l+MET	4l+MET	
C1	1l+MET	3l+MET	2l+MET



	Main EW Background		
	Type	$\sigma \times BR$ (pb)	$\sigma/\sigma_{SUSY}@100$ GeV
N2N1*	Z $\rightarrow ll$	2200	2200/0.17 ~ 13000
C1N1*	W $\rightarrow l\nu$	25000	25000/1 ~ 25000
C1N2	WZ $\rightarrow ll\nu ll$	0.3	0.3/0.12 ~ 2.5
C1C1	WW $\rightarrow ll\nu ll$	2.6	2.6/0.25 ~ 10

Only C1C1 and C1N2 can be probed at LHC-8 (20 fb⁻¹)

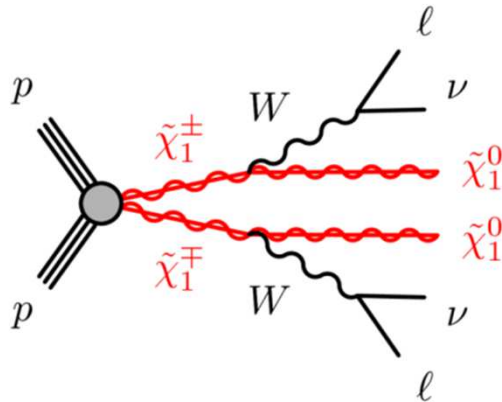
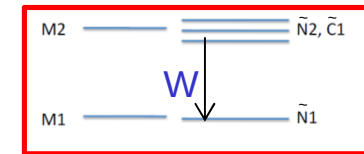
EWKinos (4)

ATLAS: 1403.5294

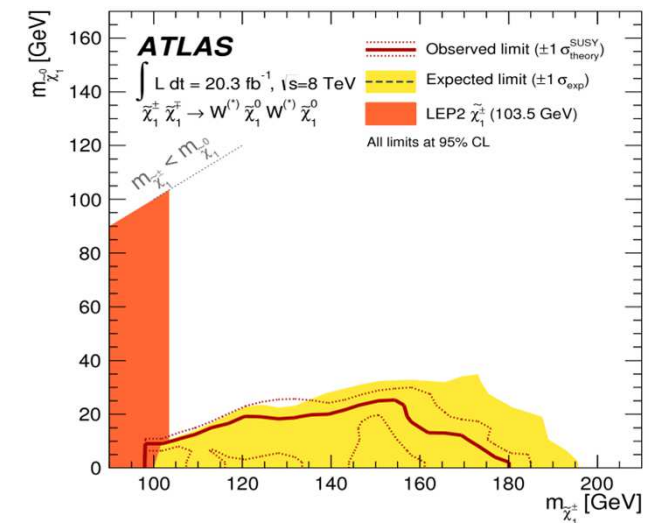
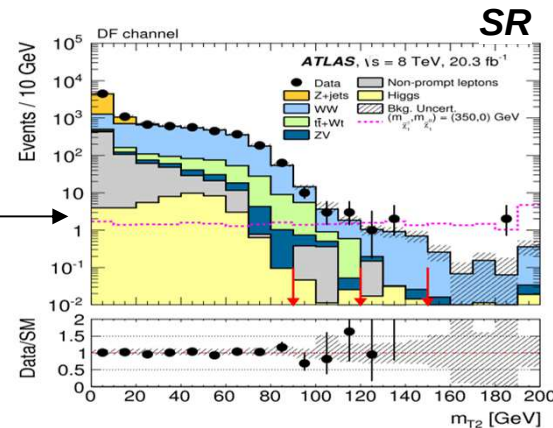
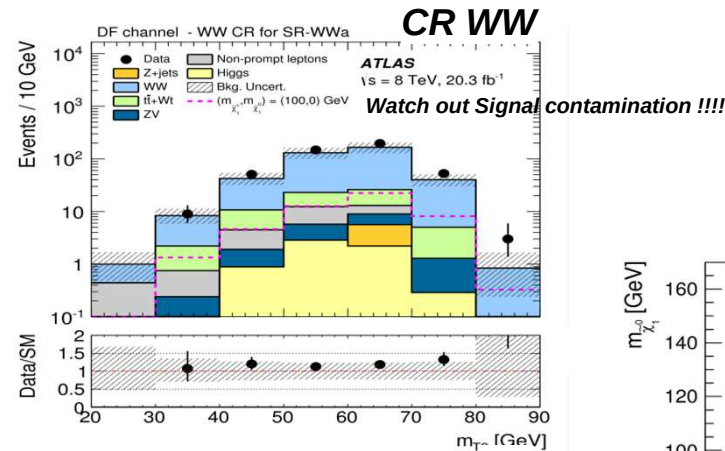
Bino case
Open Spectra

❑ Search for Direct Chargino (C1C1)

mu ———— $\tilde{N}_3, \tilde{N}_4, \tilde{C}_2$



- Expected signal is flat !
- WW steeply falling after 90 GeV

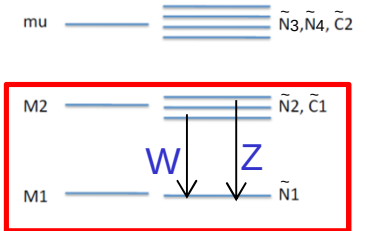


Can barely exclude
SUSY models !

EWKinos (5)

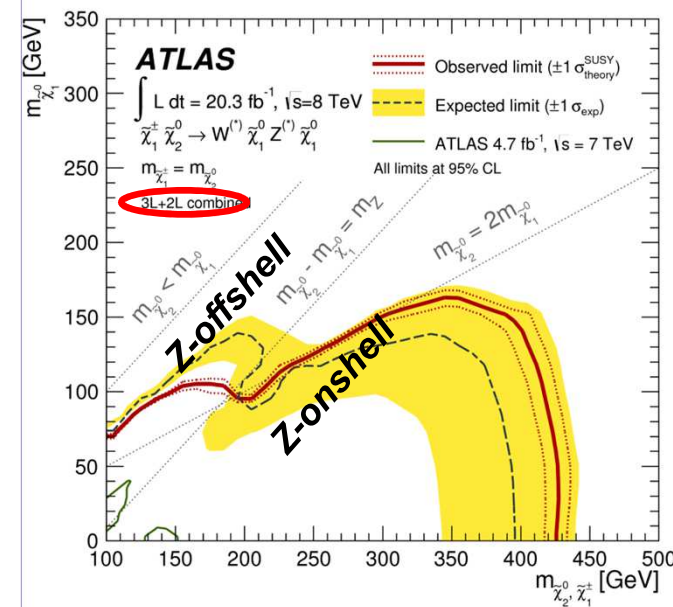
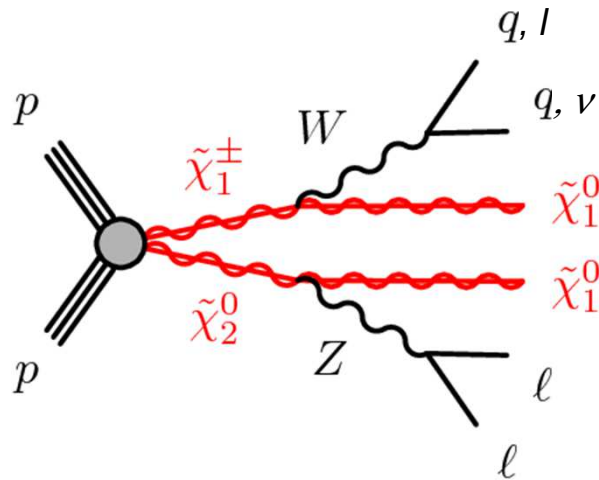
ATLAS: 1402.7029+1403.5294
CMS:1405.7570

Bino case
Open Spectra



□ Search for associate production (C1N2)

- Similar experimental challenge, but dominant background is now WZ
- Consider 3-lepton signal regions: w/wo Z veto, w/wo W veto, w/wo taus
- Can combine with 2l+2j

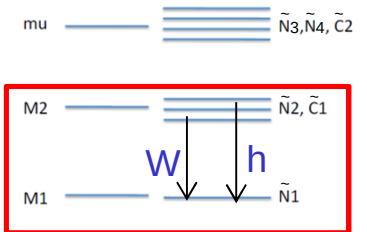


→ Exclude $m(\chi_1^\pm)=m(\chi_2^0)<420$ GeV for $m(\text{LSP})<80$ GeV

EWKinos (6)

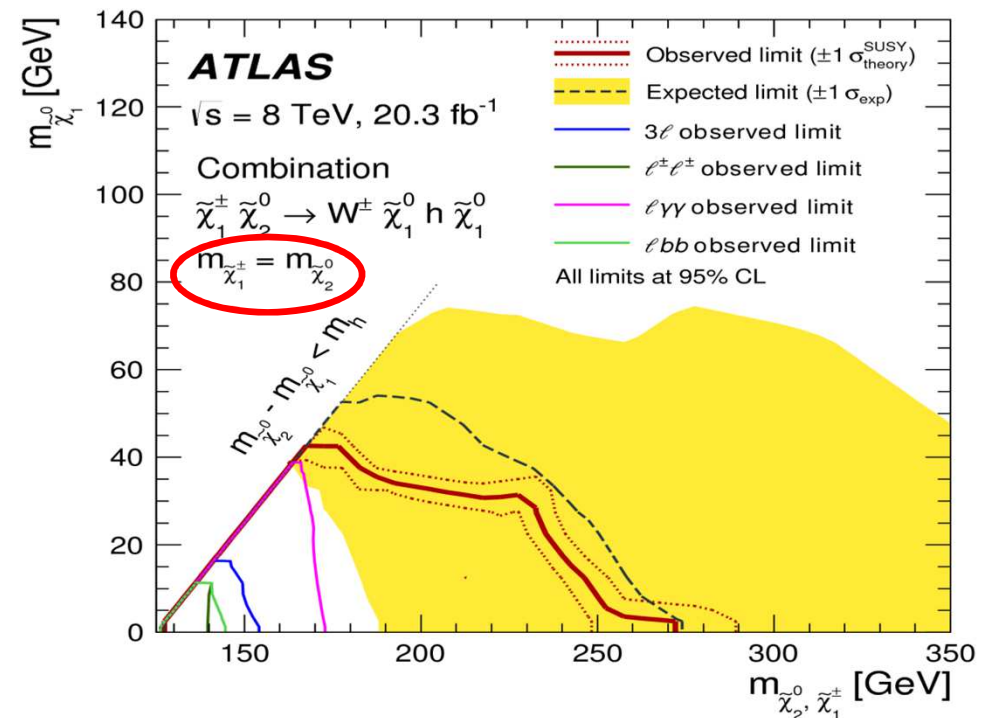
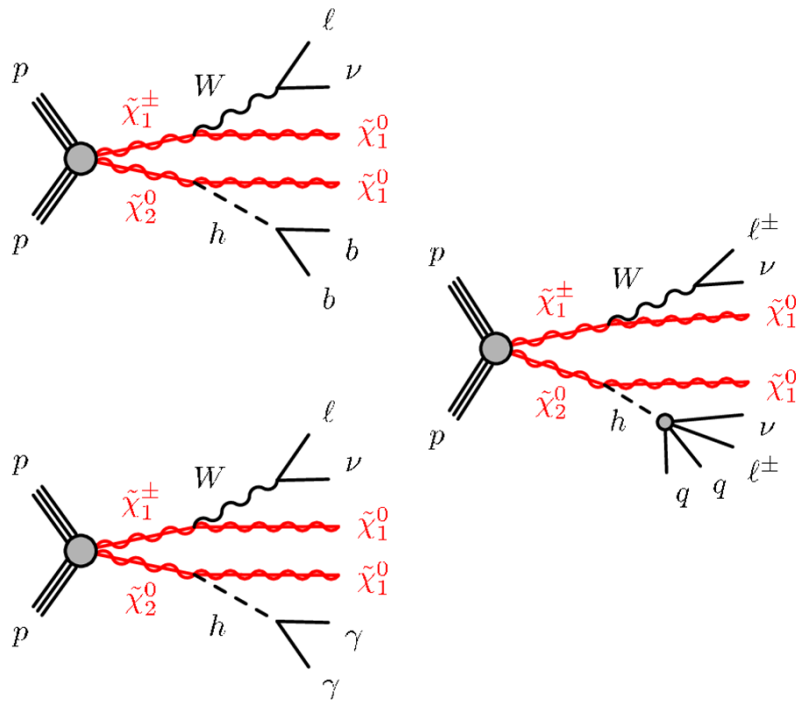
ATLAS: 1501.07110
CMS: 1409.3168

Bino case
Open Spectra



EWKinos production with Higgs in the cascade

- 3 signal regions considered, with very different final states
- $1l + \gamma\gamma$ is stronger



EWKinos (7)

A: MSSM

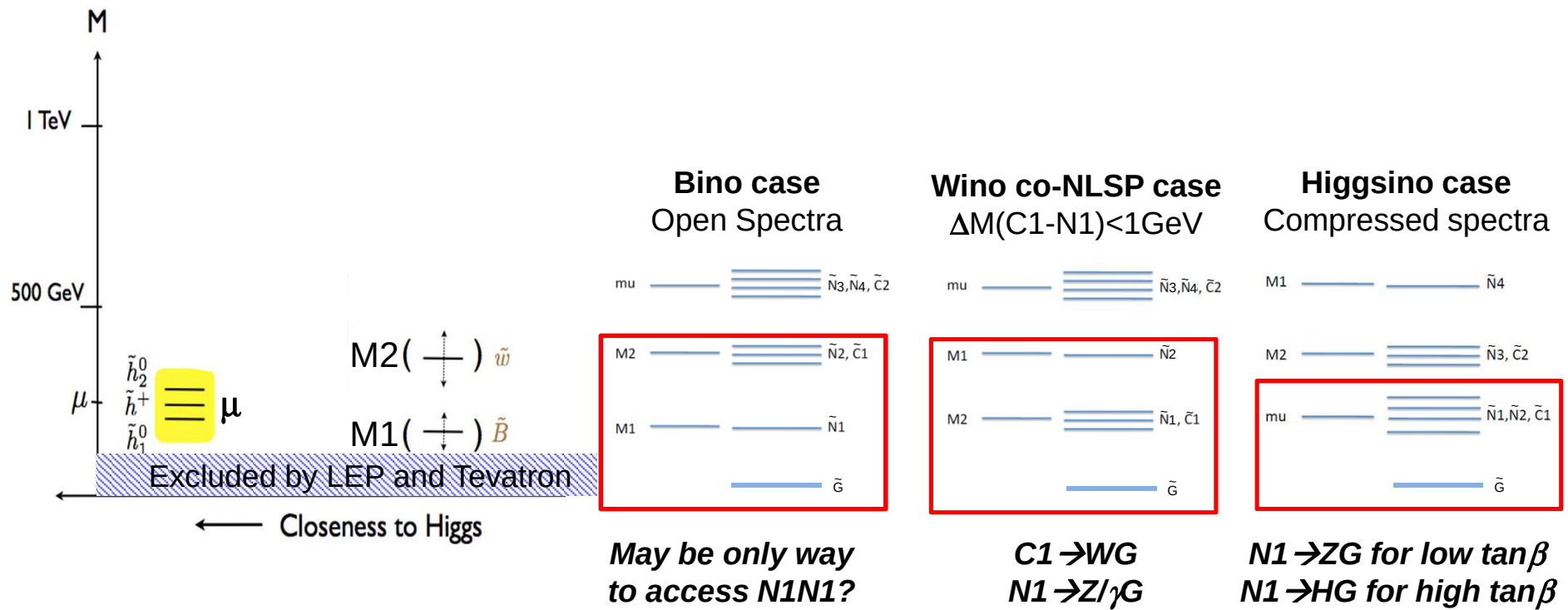
B: Natural >10%

C1: RPC

C2: LSP nature=G

C3: Unspecified Spectra

Main change



EWKinos (9)

ATLAS: 1507.05493

Wino co-NLSP case
 $\Delta M(C1-N1) < 1 \text{ GeV}$

Can access the Wino co-NLSP case

- C1N2 has very low cross-section
- Decay $C1 \rightarrow W^* N1$ is longer than $\sim 0.1 \text{ mm}$... and $C1 \rightarrow W \tilde{G}$ dominates !

	N1	C1
C1	$W\gamma GG$	$WWGG$
Final states	$1l+1\gamma+\text{MET}$	$2l+\text{MET}$
Bkg type	$W\gamma \rightarrow l\nu\gamma$	$WW \rightarrow l\nu l\nu$
$\sigma \times \text{BR (pb)}$	12	2.6
$\sigma/\sigma_{\text{SUSY}}@100 \text{ GeV}$	12/1~12	2.6/0.25~10

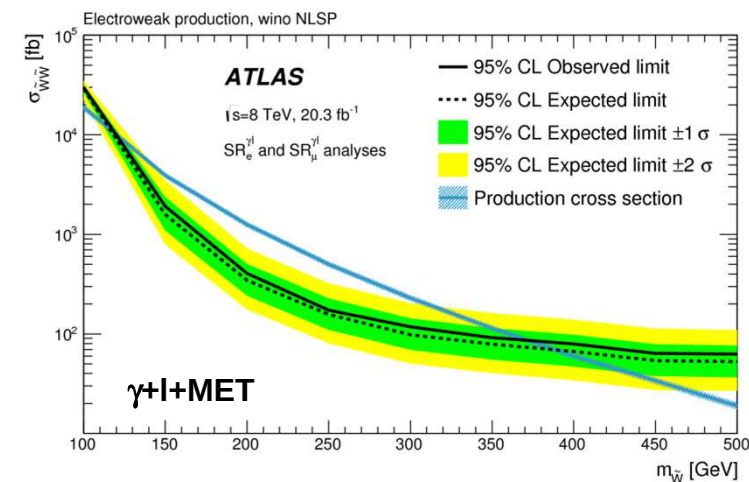
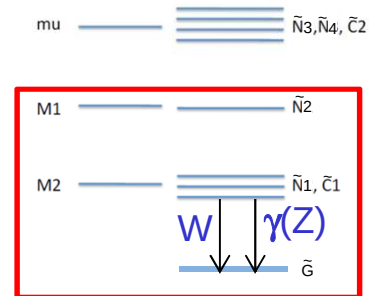
For $\text{LSP}=\chi_1^0$ was

25000

10

Now possible to cover it!

See $\text{LSP}=N1$ (p.54)



EWKinos (10)

CMS: 1409.3168

Higgsino case
Compressed spectra

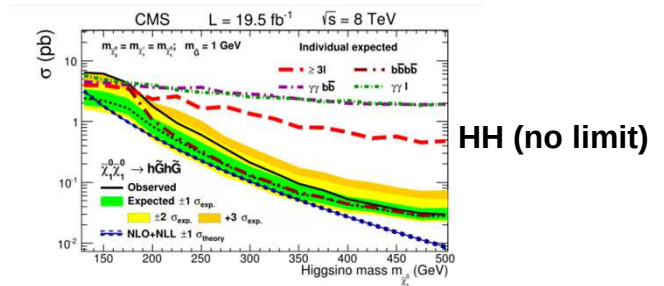
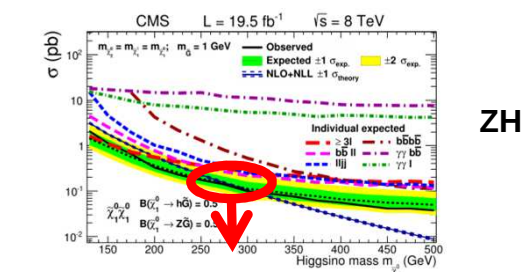
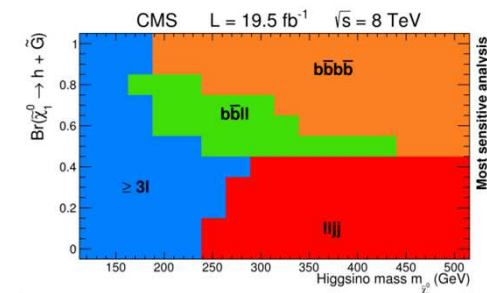
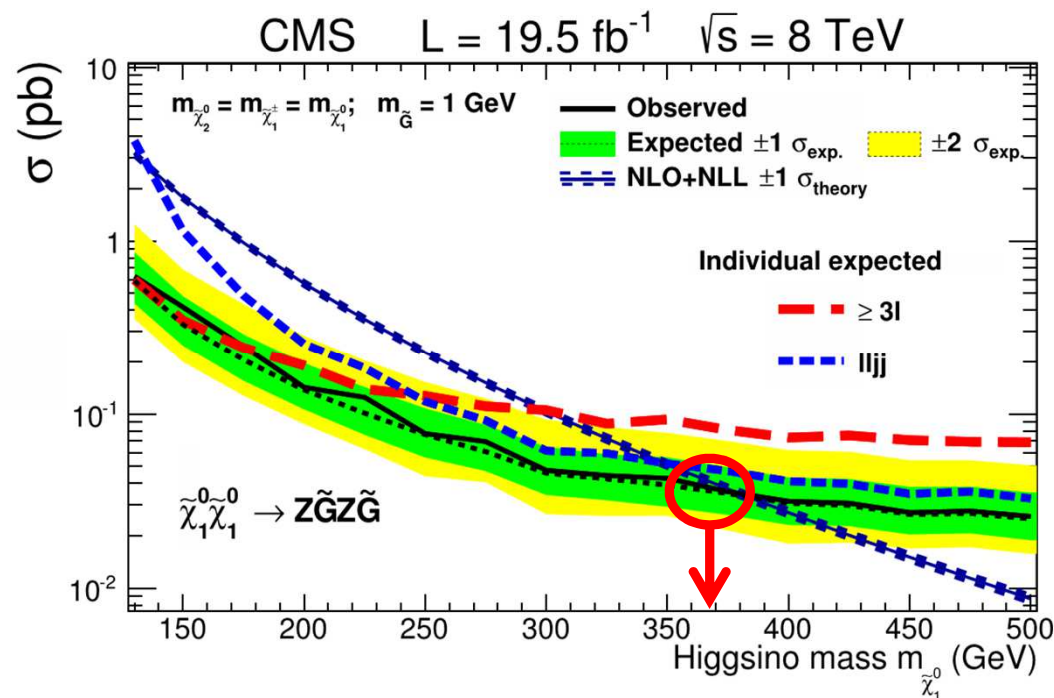
M1 $\longrightarrow \tilde{N}4$

M2 $\longrightarrow \tilde{N}3, \tilde{C}2$

mu $\longrightarrow \tilde{N}1, \tilde{N}2, \tilde{C}1$
 $Z/H \longrightarrow \tilde{G}$

Have sensitivity to the Higgsino case !

- N1, C1, N2 degenerate. Assume N1 is still the NLSP
- Can combine several channel to gain sensitivity



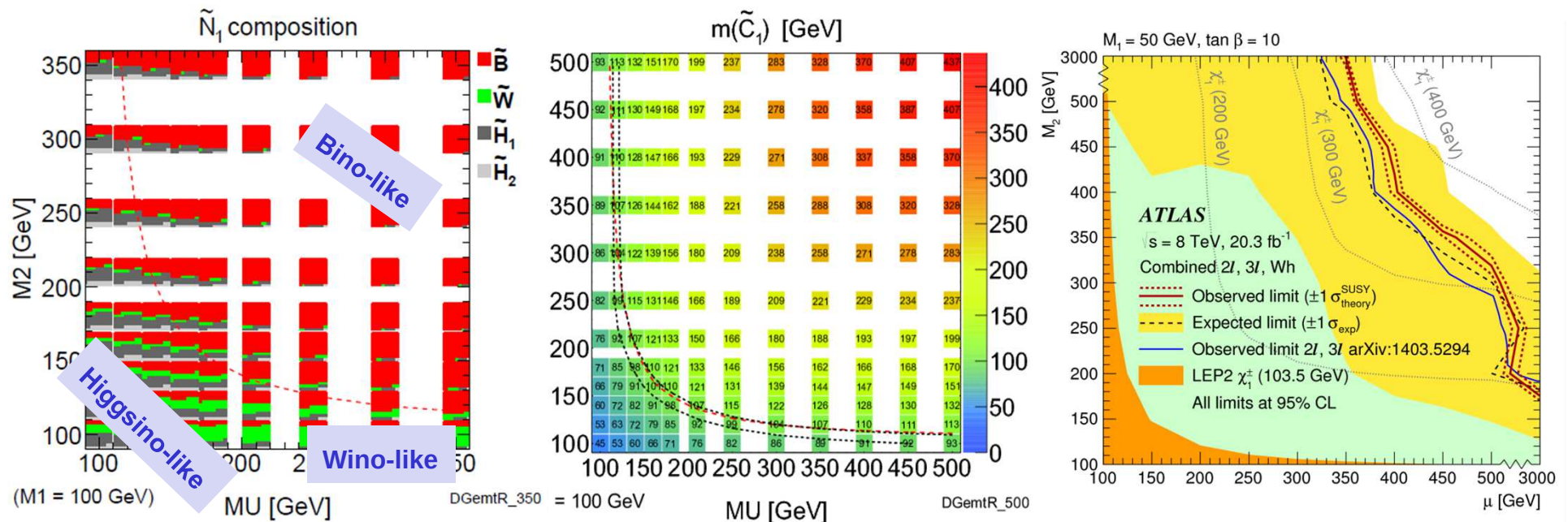
A strong constraint for naturalness

EWKinos (11)

ATLAS: 1509.07152
CMS-PAS-SUS-15-010

□ pMSSM = systematic scan of 19 parameters MSSM

- We have treated 3 extreme cases: Bino, Wino, Higgsino with BR=100%
 - ✓ Sometimes even SUSY impossible models ! (*N1 Higgs decay in Bino models*)
- Can consider SUSY-motivated case and all the intermediate cases with pMSSM



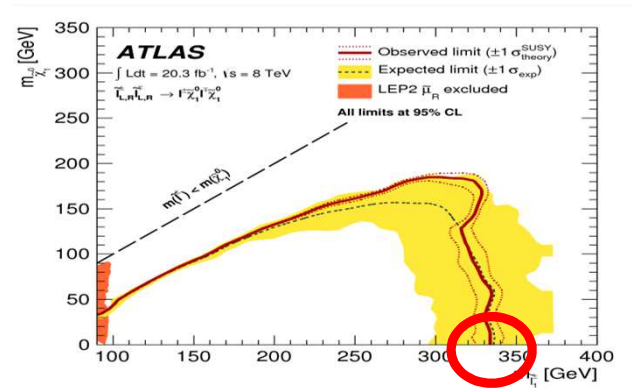
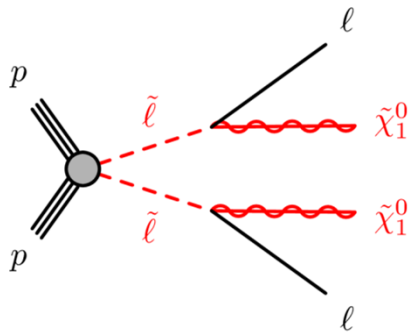
But not always easy to have a clear picture because of 4D !

Sleptons

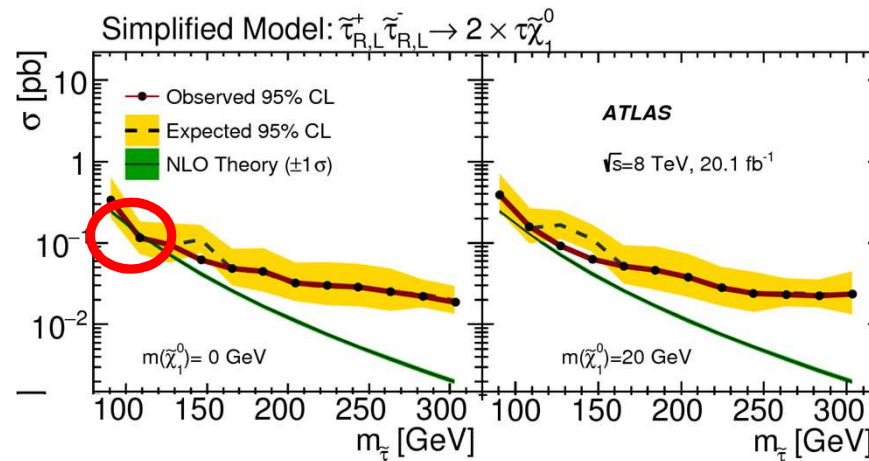
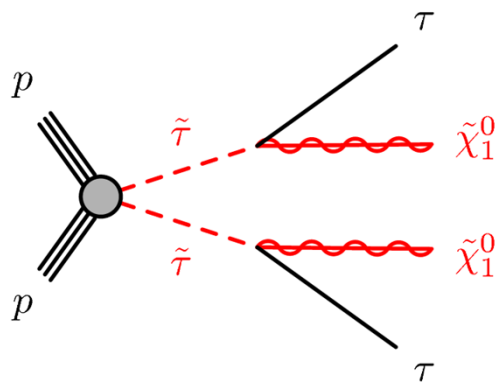
ATLAS: 1403.5294
CMS: 1405.7570
ATLAS-SUSY-2013-028

□ Charged sleptons govern by 5 parameters $(m_{eR}=m_{eL}, m_{\mu R}=m_{\mu L}, m_{\tau R}, m_{\tau L}, \theta_\tau)^*$

- Have same final states as $C1C1 \rightarrow WWN1N1$ so can be exclude beyond LEP

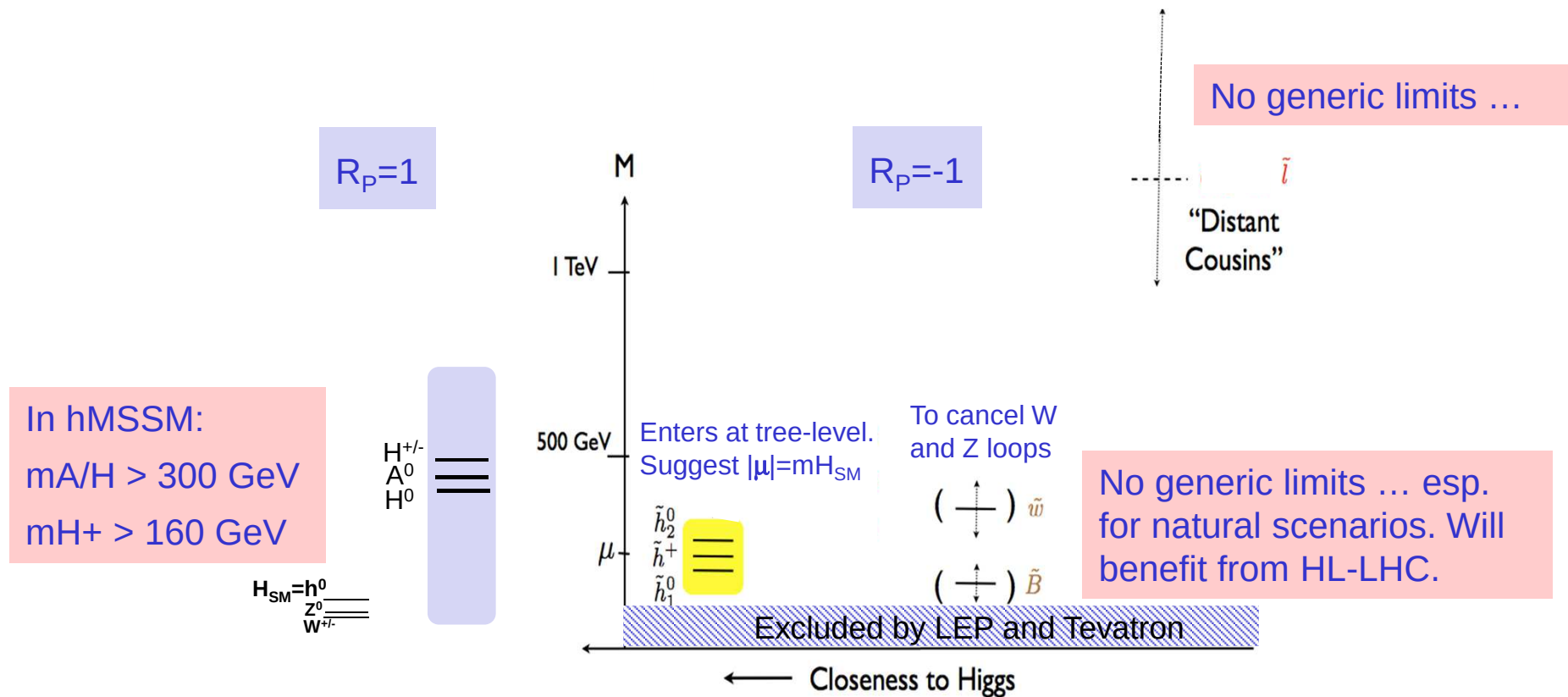


- $m(\tau) \sim 1.8$ GeV heavy, low mass stau conceivable but extremely challenging experimentally !



*Sneutrinos in RPC not directly accessible at LHC

Summary of Part Ib



SPARE