



In-jet Fragmentation

in collaboration with Asmita Mukherjee and Werner Vogelsang

[Kaufmann, Mukherjee, Vogelsang: PRD92:054015 & PRD93:114021]



Outline

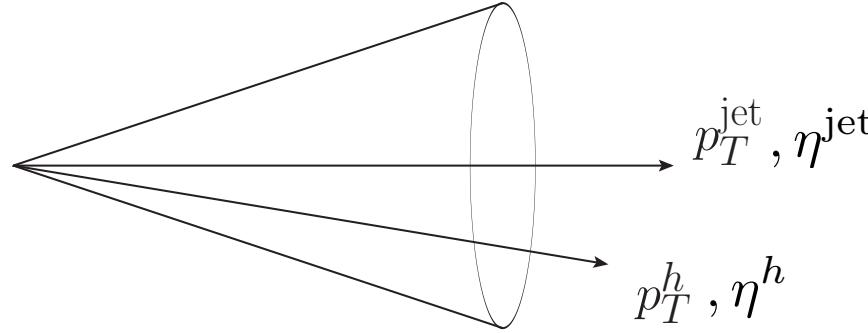
- Overview
- Calculation of the cross section
- Numerical Results
- Photon-in-jet
- Summary & Outlook



Process of interest

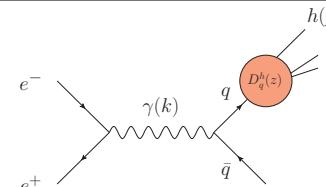
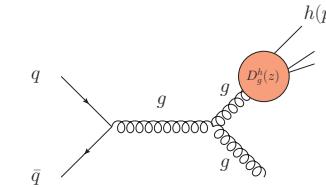
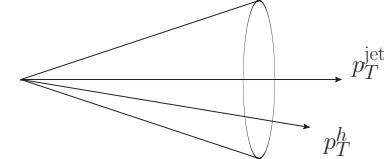
$$pp \rightarrow (\text{jet } h)X$$

where the hadron is observed INSIDE a fully reconstructed jet and is part of the jet



- Fragmentation functions
- Jet (sub-)structure
- Spin correlations (Collins effect)
- pA and AA scattering \rightsquigarrow “jet quenching”

Determination of FFs

Process	D_g^h ?	Direct scan?	
$e^+e^- \rightarrow hX$ $ep \rightarrow ehX$			 $z = \frac{2p_h \cdot k}{k^2}$
$pp \rightarrow hX$			 $\hat{p}_T = \frac{p_T^h}{z_c} \quad z_c^{\min} = \frac{2p_T^h}{\sqrt{S}} \cosh \eta$ $d\sigma \propto \int_{z_c^{\min}}^1 \frac{dz_c}{z_c^2} \ [f_a \otimes f_b \otimes d\hat{\sigma}_{ab}^c(\hat{p}_T, \dots)] \ D_c^h(z_c)$
$pp \rightarrow (\text{jet } h) X$			 $z = z_h \equiv \frac{p_T^h}{p_T^{\text{jet}}}$



A Warning...

Description of hadron-jet momentum correlation is not unique! Some possible definitions:



$$z_h = ?$$

$$\frac{p_T}{p_T^{\text{jet}}}$$

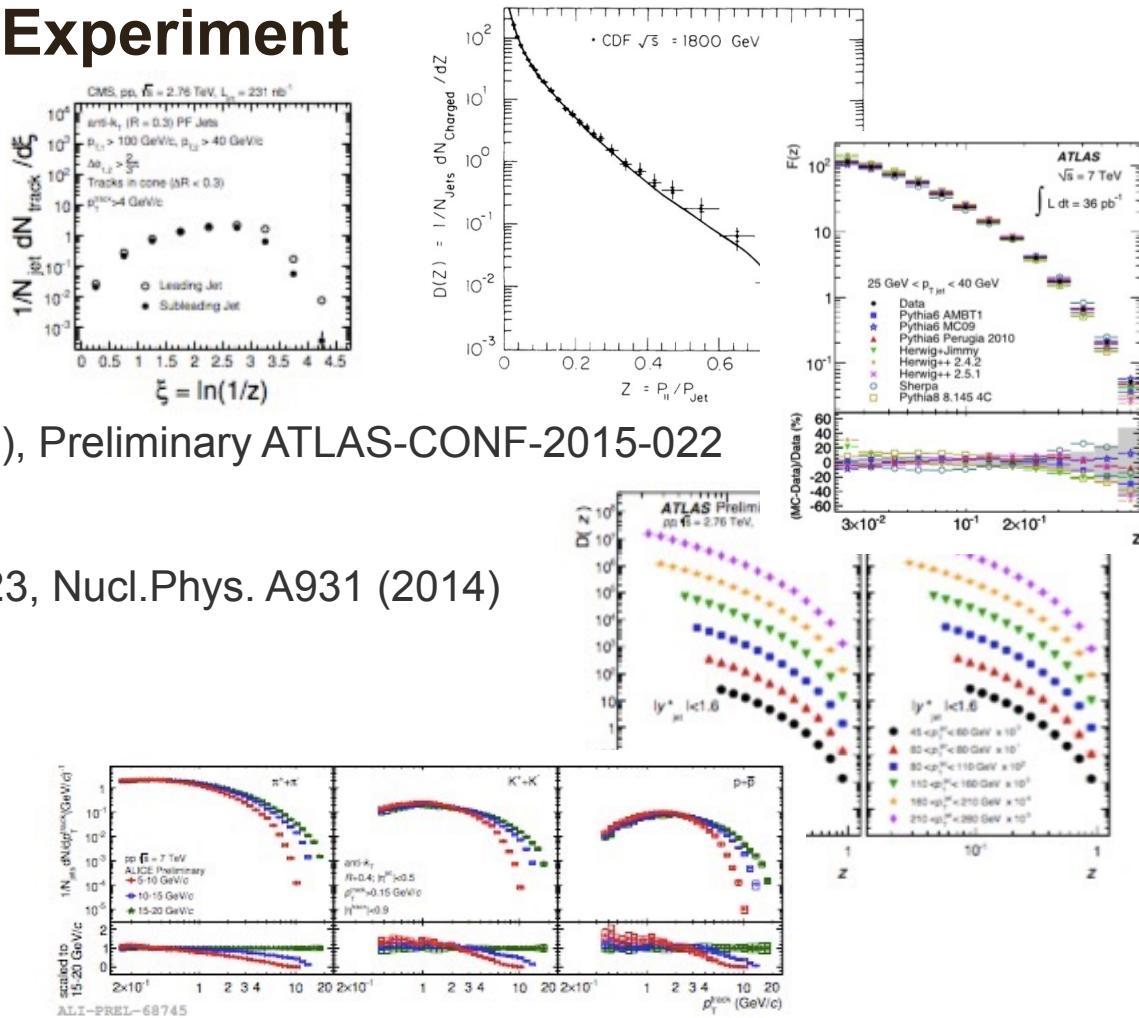
$$\frac{\vec{p}_T \cdot \vec{p}_T^{\text{jet}}}{|\vec{p}_T^{\text{jet}}|^2}$$

$$\frac{p_T}{p_T^{\text{jet}}} \cos \Delta R$$

$$\Delta R \equiv \sqrt{\Delta\eta^2 + \Delta\phi^2}$$

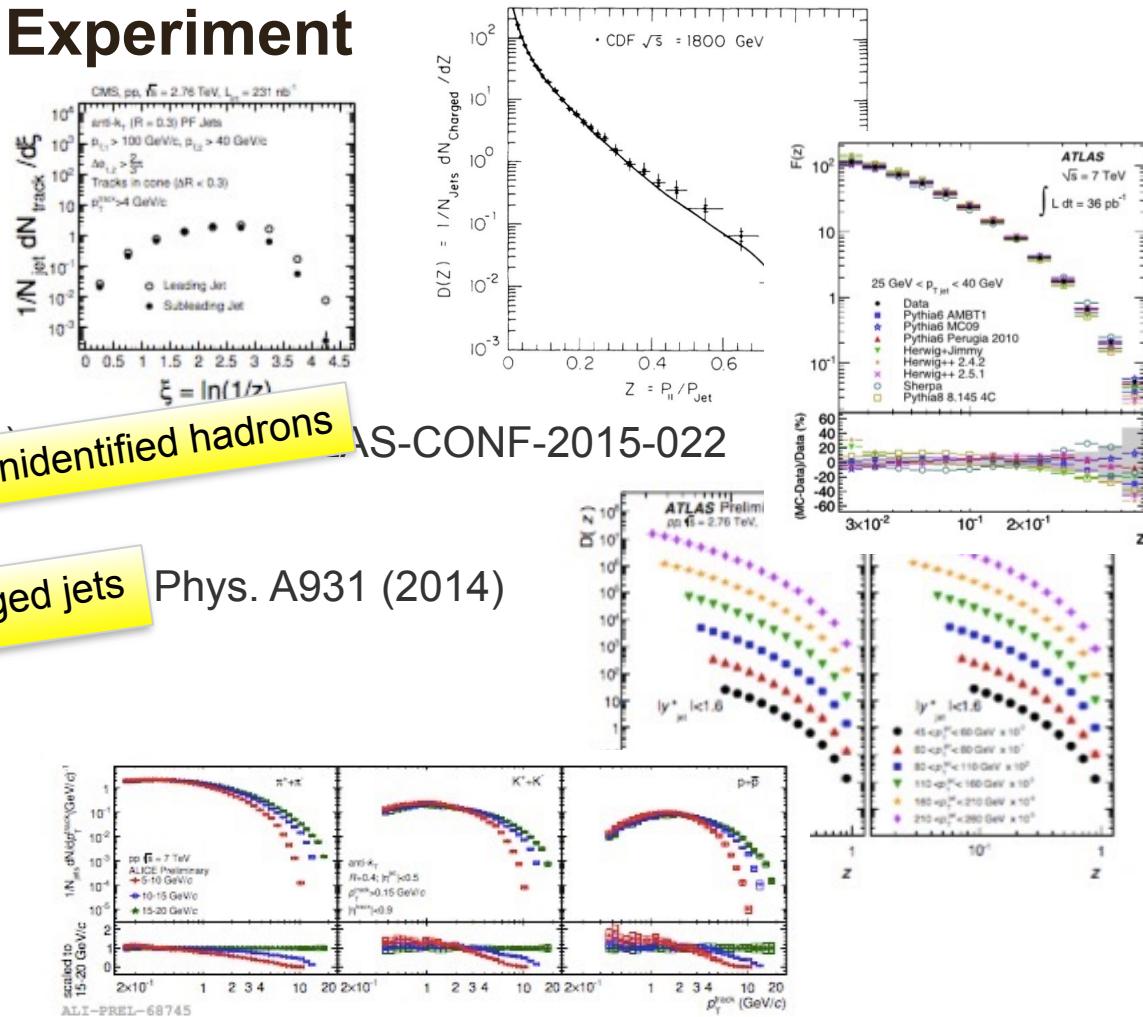
What's available - Experiment

- Tevatron
 - CDF: PRL65(1990)
- LHC
 - ATLAS: EPJ C71(2011), Preliminary ATLAS-CONF-2015-022
 - CMS: JHEP10(2012)
 - ALICE: arXiv:1408.5723, Nucl.Phys. A931 (2014)
- RHIC
 - STAR: in progress...



What's available - Experiment

- Tevatron
 - CDF dijet sample (90)
- LHC
 - ATLAS: EPJ C71(2011) 024002, ATLAS-CONF-2015-022
 - CMS: EPJC 71 (2011) 024002, CMS-CONF-2015-022
 - CMS: arXiv:1409.7571v1 [hep-ph] (2014)
- RHIC
 - STAR: in progress...





What's available - Theory

- e^+e^- collisions:
 - Jain, Procura, Waalewijn; JHEP05(2011) SCET
 - Procura, Waalewijn; PRD85(2012) SCET
 - ...
- pp collisions:
 - Arleo, Fontannaz, Guillet, Nguyen; JHEP04(2014) MC@NLO
 - Ritzmann, Waalewijn; PRD90(2014) SCET@NNLO
 - Kaufmann, Mukherjee, Vogelsang; PRD92(2015) NJA@NLO
 - Chien, Kang, Ringer, Vitev, Xing; JHEP05(2016) SCET@LO+NLL,R
 - Kang, Ringer, Vitev; arXiv:1606.07063 SCET@NLO+NLL,R
 - Dai, Kim, Leibovich; arXiv:1606.07411 SCET@NLO+NLL,R
 - ...



Outline

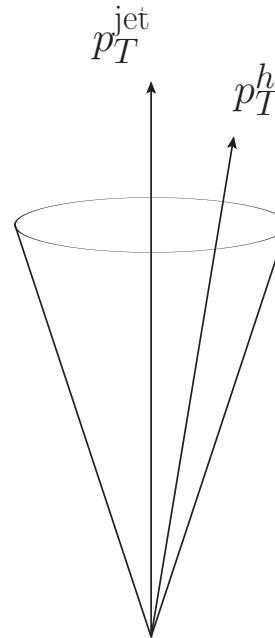
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- **Calculation of the cross section**
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Calculation of the cross section

[Kaufmann, Mukherjee, Vogelsang: PRD92:054015]

- NLO cross section for the process $pp \rightarrow (\text{jet } h) X$
- differential in jet's transverse momentum p_T^{jet} and rapidity η^{jet} , and the variable

$$z_h \equiv \frac{p_T^h}{p_T^{\text{jet}}}$$



Narrow Jet Approximation (NJA)

[Jäger, Stratmann, Vogelsang: PRD70:034010; Mukherjee, Vogelsang: PRD86:094009; Kaufmann, Mukherjee, Vogelsang: PRD91:034001]

- parameter which characterizes size of a jet

$$d_{ij} = \min \left((k_T^i)^{2p}, (k_T^j)^{2p} \right) \frac{(\phi_i - \phi_j)^2 + (\eta_i - \eta_j)^2}{R^2}$$

cone

$$(\phi_J - \phi_i)^2 + (\eta_J - \eta_i)^2 \leq R^2$$

- interested in rather narrow jets

$$\mathcal{R} \equiv R, \frac{1}{\beta} \quad \mathcal{R} \rightarrow 0$$

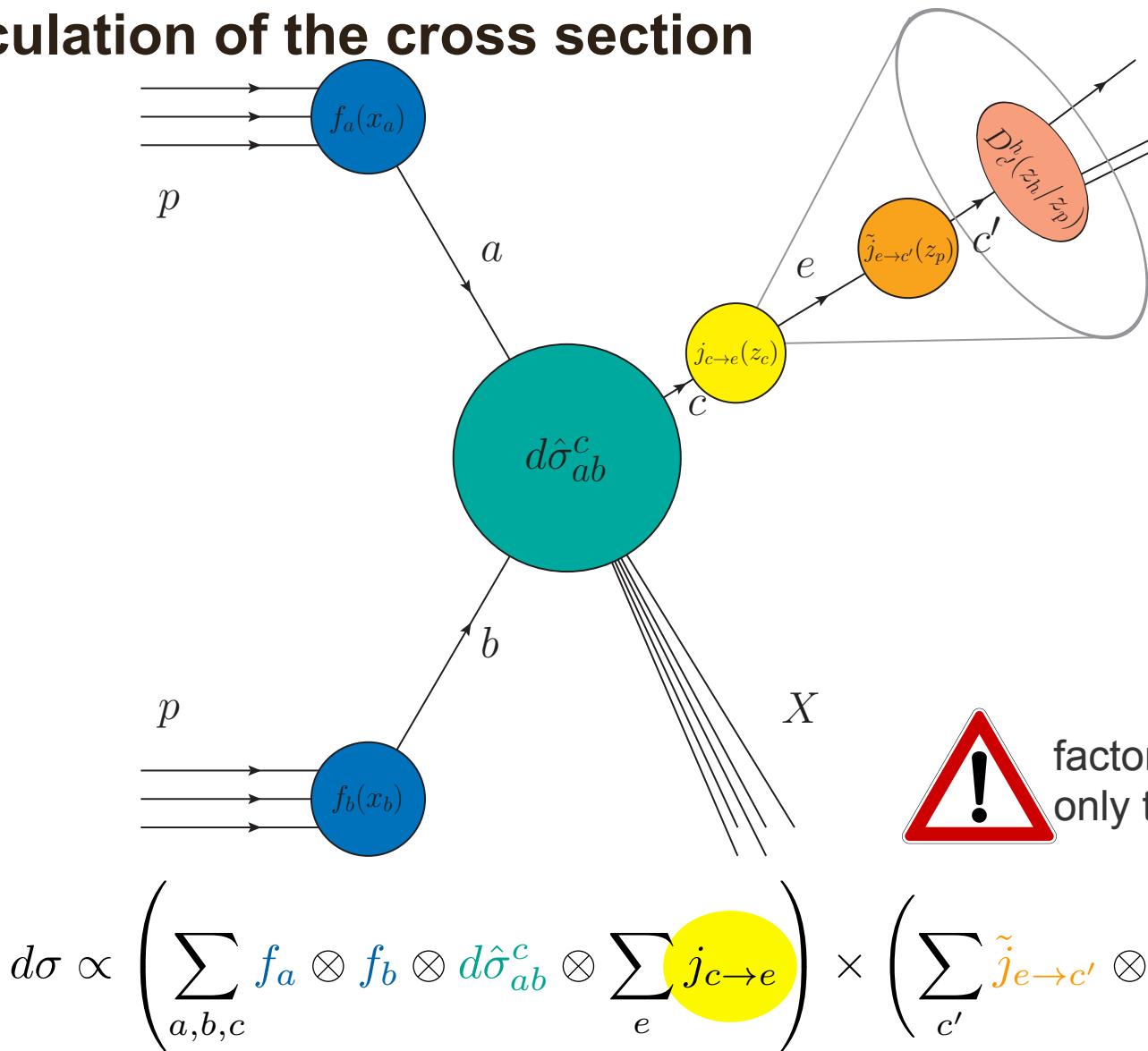
$$J_\beta(p) \equiv E - \beta \frac{m^2}{E}, \quad \beta > 1$$

J_{E_T}

- allows analytical calculation of partonic cross sections
- cross section $d\hat{\sigma} \propto \mathcal{A} \log \mathcal{R} + \mathcal{B} + \mathcal{O}(\mathcal{R}^2)$
- good approximation up to $\mathcal{R} \approx 0.7$



Calculation of the cross section



factorization “proved”
only to NLO!

Calculation of the cross section

$$\begin{aligned}
 \frac{d\sigma^{H_1 H_2 \rightarrow (\text{jet } h) X}}{dp_T^{\text{jet}} d\eta^{\text{jet}} dz_h} &= \frac{2p_T^{\text{jet}}}{S} \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a^{H_1}(x_a, \mu_F) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b^{H_2}(x_b, \mu_F) \\
 &\times \int_{z_c^{\min}}^1 \frac{dz_c}{z_c^2} \frac{d\hat{\sigma}_{ab}^c(\hat{s}, \hat{p}_T, \hat{\eta}, \mu_F, \mu'_F, \mu_R)}{vdv dw} \sum_e j_{c \rightarrow e} \left(z_c, \frac{\mathcal{R} p_T^{\text{jet}}}{\mu'_F}, \mu_R \right) \\
 &\times \sum_{c'} \int_{z_h}^1 \frac{dz_p}{z_p} \tilde{j}_{e \rightarrow c'} \left(z_p, \frac{\mathcal{R} p_T^{\text{jet}}}{\mu''_F}, \mu_R \right) D_{c'}^h \left(\frac{z_h}{z_p}, \mu''_F \right)
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Calculation of the cross section

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 \end{aligned}$$

single-inclusive
parton cross section

[Jäger,Stratmann,Vogelsang; PRD67(2003)]
[Aversa,Chiappetta,Greco,Guillet; Nucl.Phys.B327(1989)]

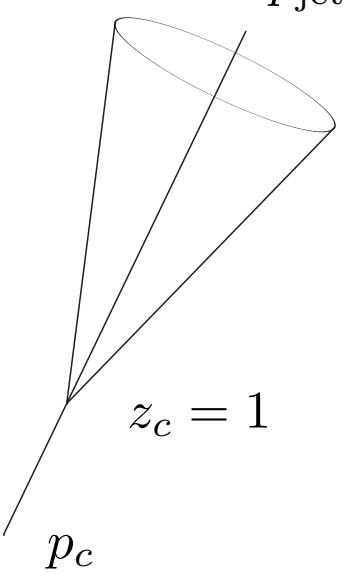
Calculation of the cross section

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 \end{aligned}$$

“jet-function”: formation of jet
with $p_{\text{jet}} = z_c p_c$

Calculation of the cross section

$$\frac{d\sigma^{H_1 H_2 \rightarrow (\text{jet } h) X}}{dp_T^{\text{jet}} d\eta^{\text{jet}} dz_h} = \frac{2p_T^{\text{jet}}}{S} \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a^{H_1}(x_a, \mu_F) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b^{H_2}(x_b, \mu_F)$$



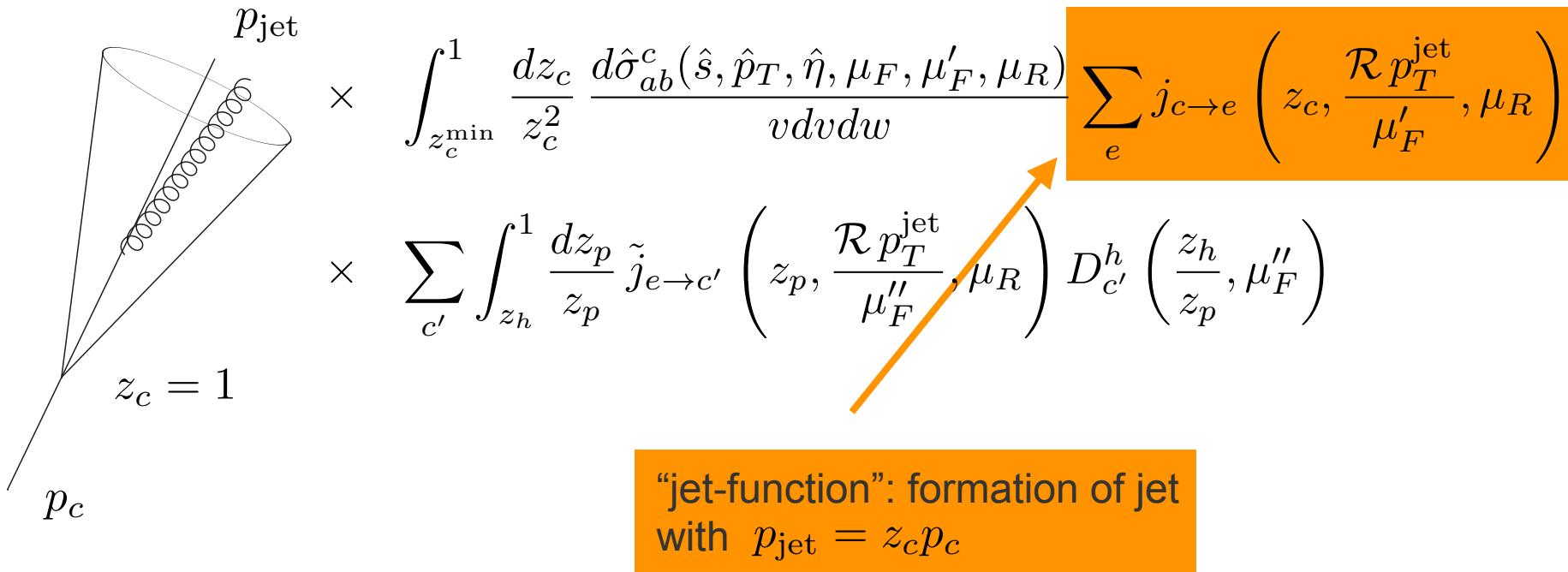
$$\times \int_{z_c^{\min}}^1 \frac{dz_c}{z_c^2} \frac{d\hat{\sigma}_{ab}^c(\hat{s}, \hat{p}_T, \hat{\eta}, \mu_F, \mu'_F, \mu_R)}{vdv dw} \sum_e j_{c \rightarrow e} \left(z_c, \frac{\mathcal{R} p_T^{\text{jet}}}{\mu'_F}, \mu_R \right)$$

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“jet-function”: formation of jet
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Calculation of the cross section

$$\frac{d\sigma^{H_1 H_2 \rightarrow (\text{jet } h) X}}{dp_T^{\text{jet}} d\eta^{\text{jet}} dz_h} = \frac{2p_T^{\text{jet}}}{S} \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a^{H_1}(x_a, \mu_F) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b^{H_2}(x_b, \mu_F)$$



p_{jet}

$z_c = 1$

p_c

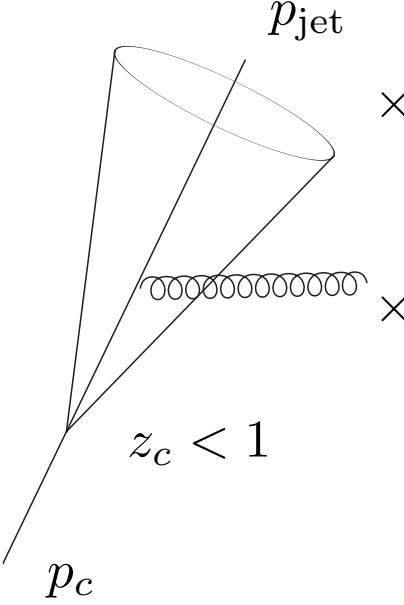
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“jet-function”: formation of jet with $p_{\text{jet}} = z_c p_c$

Calculation of the cross section

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“jet-function”: formation of jet
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Calculation of the cross section

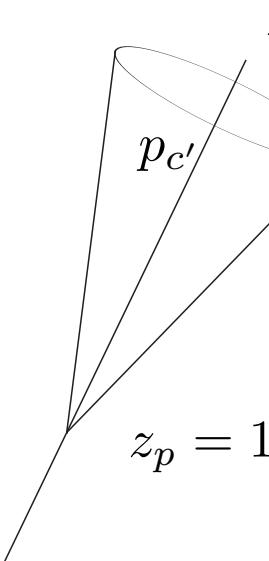
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 \end{aligned}$$



“jet-function”: partonic
fragmentation, $p_{c'} = z_p p_{\text{jet}}$

Calculation of the cross section

$$\frac{d\sigma^{H_1 H_2 \rightarrow (\text{jet } h) X}}{dp_T^{\text{jet}} d\eta^{\text{jet}} dz_h} = \frac{2p_T^{\text{jet}}}{S} \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a^{H_1}(x_a, \mu_F) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b^{H_2}(x_b, \mu_F)$$



$$\times \int_{z_c^{\min}}^1 \frac{dz_c}{z_c^2} \frac{d\hat{\sigma}_{ab}^c(\hat{s}, \hat{p}_T, \hat{\eta}, \mu_F, \mu'_F, \mu_R)}{vdv dw} \sum_e j_{c \rightarrow e} \left(z_c, \frac{\mathcal{R} p_T^{\text{jet}}}{\mu'_F}, \mu_R \right)$$

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↑

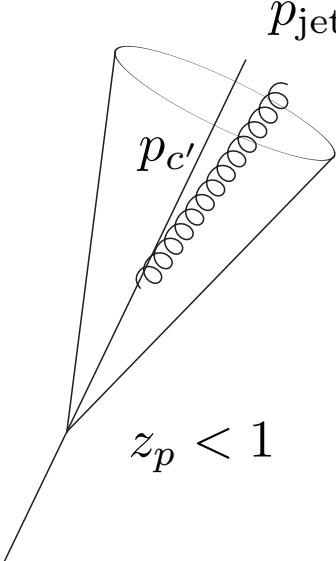
“jet-function”: partonic fragmentation, $p_{c'} = z_p p_{\text{jet}}$

Calculation of the cross section

$$\frac{d\sigma^{H_1 H_2 \rightarrow (\text{jet } h) X}}{dp_T^{\text{jet}} d\eta^{\text{jet}} dz_h} = \frac{2p_T^{\text{jet}}}{S} \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a^{H_1}(x_a, \mu_F) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b^{H_2}(x_b, \mu_F)$$

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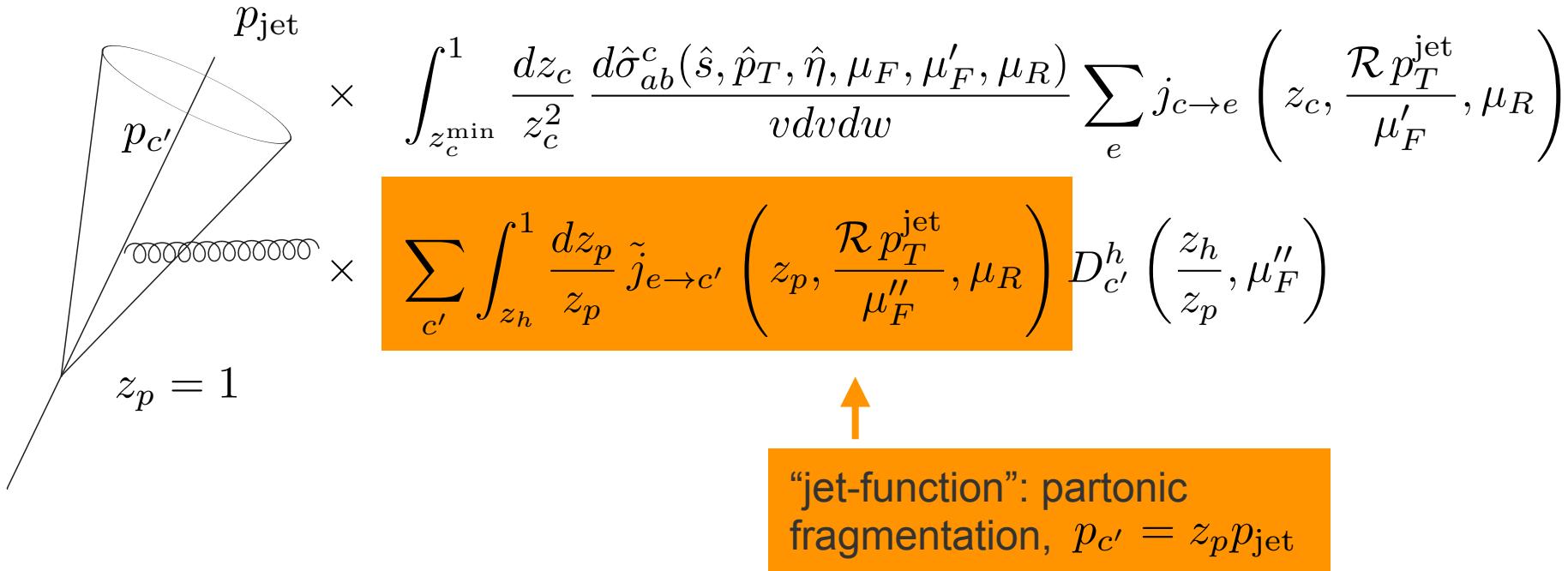


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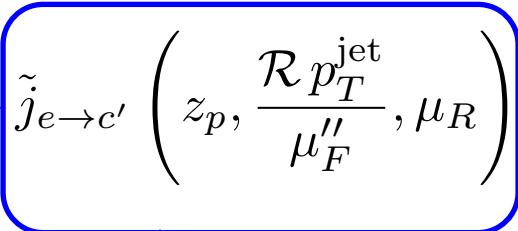


Leading order

$$\begin{aligned}
 \frac{d\sigma^{H_1 H_2 \rightarrow (\text{jet } h) X}}{dp_T^{\text{jet}} d\eta^{\text{jet}} dz_h} &= \frac{2p_T^{\text{jet}}}{S} \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a^{H_1}(x_a, \mu_F) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b^{H_2}(x_b, \mu_F) \\
 &\times \int_{z_c^{\min}}^1 \frac{dz_c}{z_c^2} \frac{d\hat{\sigma}_{ab}^c(\hat{s}, \hat{p}_T, \hat{\eta}, \mu_F, \mu'_F, \mu_R)}{vdv dw} \sum_e j_{c \rightarrow e} \left(z_c, \frac{\mathcal{R} p_T^{\text{jet}}}{\mu'_F}, \mu_R \right) \\
 &\times \sum_{c'} \int_{z_h}^1 \frac{dz_p}{z_p} \tilde{j}_{e \rightarrow c'} \left(z_p, \frac{\mathcal{R} p_T^{\text{jet}}}{\mu''_F}, \mu_R \right) D_{c'}^h \left(\frac{z_h}{z_p}, \mu''_F \right)
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Leading order

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 \end{aligned}$$



$\delta_{ec'} \delta(1 - z_p)$ $\delta_{ce} \delta(1 - z_c)$

Leading order

$$\begin{aligned}
 \left. \frac{d\sigma^{H_1 H_2 \rightarrow (\text{jet } h) X}}{dp_T^{\text{jet}} d\eta^{\text{jet}} dz_h} \right|_{\text{LO}} &= \frac{2p_T^{\text{jet}}}{S} \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a^{H_1}(x_a, \mu_F) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b^{H_2}(x_b, \mu_F) \\
 &\quad \times \frac{d\hat{\sigma}_{ab}^{c, \text{Born}}(\hat{s}, \hat{p}_T, \hat{\eta}, \mu_F, \mu'_F, \mu_R)}{vdv dw} \times D_c^h(z_h, \mu''_F) \\
 &\propto D_c^h(z_h, \mu''_F).
 \end{aligned}$$

pQCD vs SCET

- Consistent at NLO
- Instead of the factorized structure with $j_{c \rightarrow e}(z_c)$ and $\tilde{j}_{e \rightarrow c'}(z_p)$ consider the combinations

$$\mathcal{K}_{c \rightarrow c'}(z, z_p; \lambda, \kappa) \equiv \sum_e j_{c \rightarrow e}(z, \lambda) \tilde{j}_{e \rightarrow c'}(z_p, \kappa)$$

- Same functions found in SCET $\mathcal{J}_{cc'} = \mathcal{K}_{c \rightarrow c'}$ [Dai,Kim,Leibovich; arXiv:1606.07411]
[Kang,Ringer,Vitev; arXiv:1606.06732]
- semi-inclusive fragmenting jet functions (FJFs)

$$\mathcal{G}_q^h(z, z_h, \omega_J, \mu_G) = \int_{z_h}^1 \frac{dz'_h}{z'_h} \left[\mathcal{J}_{qq}(z, z'_h, \omega_J, \mu_G) D_q^h \left(\frac{z_h}{z'_h}, \mu_G \right) + \mathcal{J}_{qg}(z, z'_h, \omega_J, \mu_G) D_g^h \left(\frac{z_h}{z'_h}, \mu_G \right) \right]$$

$$\mathcal{G}_g^h(z, z_h, \omega_J, \mu_G) = \int_{z_h}^1 \frac{dz'_h}{z'_h} \left[\mathcal{J}_{gg}(z, z'_h, \omega_J, \mu_G) D_g^h \left(\frac{z_h}{z'_h}, \mu_G \right) + \sum_{i=q, \bar{q}} \mathcal{J}_{gi}(z, z'_h, \omega_J, \mu_G) D_i^h \left(\frac{z_h}{z'_h}, \mu_G \right) \right]$$

pQCD vs SCET

[Kang,Ringer,Vitev; arXiv:1606.06732]

- DGLAP like evolution equations in (z, μ) -space. z_h is unaffected

$$\frac{d}{d \log \mu^2} \begin{pmatrix} \mathcal{G}_S^h(z, z_h, \omega_J, \mu) \\ \mathcal{G}_g^h(z, z_h, \omega_J, \mu) \end{pmatrix} = \frac{\alpha_s(\mu)}{2\pi} \begin{pmatrix} P_{qq}(z) & 2N_f P_{gq}(z) \\ P_{qg}(z) & P_{gg}(z) \end{pmatrix} \otimes \begin{pmatrix} \mathcal{G}_S^h(z, z_h, \omega_J, \mu) \\ \mathcal{G}_g^h(z, z_h, \omega_J, \mu) \end{pmatrix}$$

with the singlet given by

$$\mathcal{G}_S^h(z, z_h, \omega_J, \mu) = \sum_{i=q, \bar{q}} \mathcal{G}_i^h(z, z_h, \omega_J, \mu) = 2N_f \mathcal{G}_q^h(z, z_h, \omega_J, \mu)$$

- $\Rightarrow \log \mathcal{R}$ resummation
- Beyond LL: techniques analogue to evolution of FFs with NLO kernels, i.e. expansion around LO solution (PEGASUS)

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- $\Rightarrow \log \mathcal{R}$ resummation
 - Beyond LL: techniques due to evolution of FFs with NLO kernels, i.e. expansion around **'PEGASUS'**
- Details in talk by Daniele Anderle*



Outline

- Overview
- Calculation of the cross section
- **Numerical Results**
- Photon-in-jet
- Summary & Outlook



Numerical Results

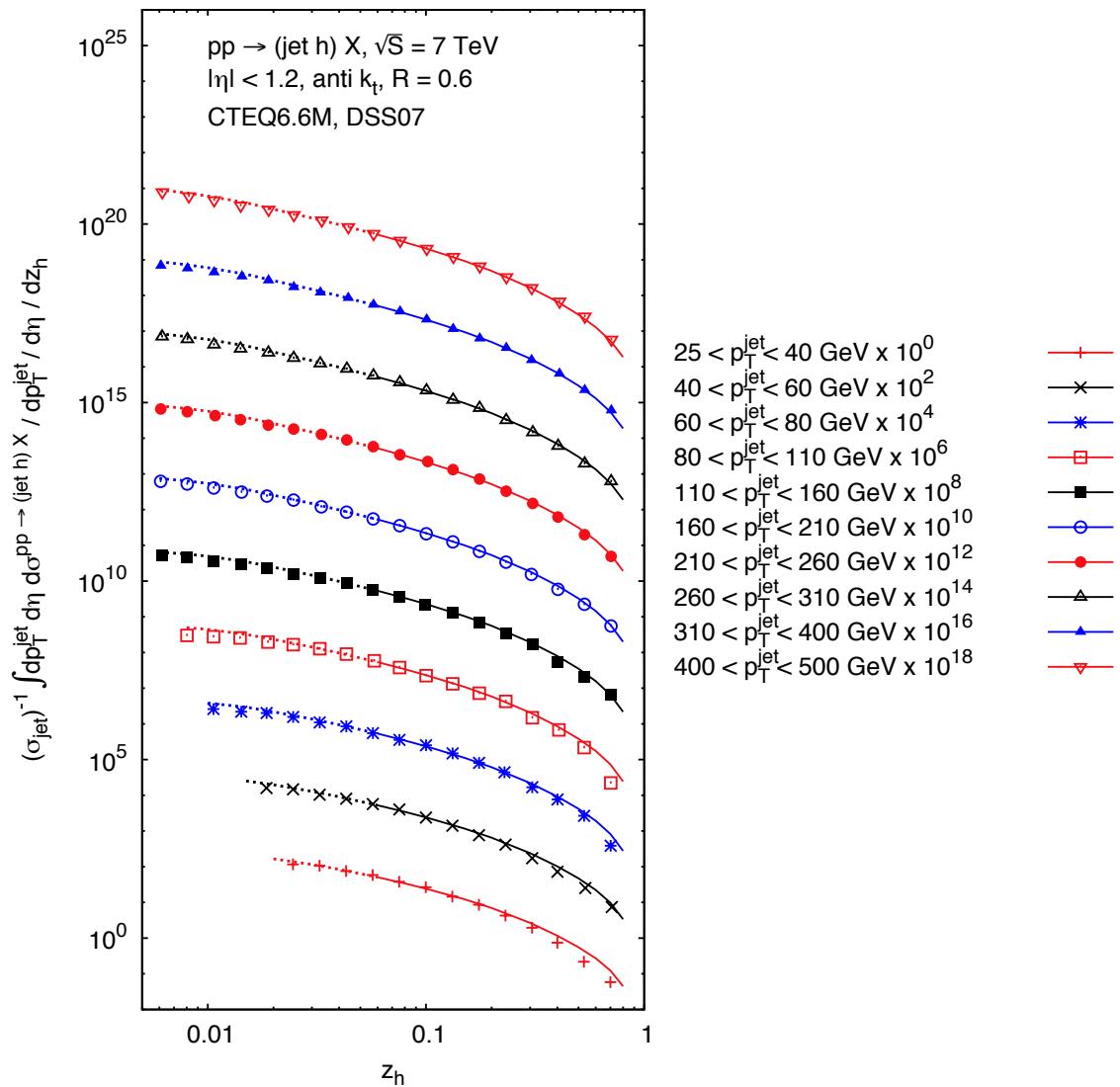
- Usually, cross section differential in z_h and normalized to the single-inclusive jet spectrum, i.e.

$$F(z_h) \equiv \left(\int dp_T^{\text{jet}} d\eta^{\text{jet}} \frac{d\sigma^{\text{jet } X}}{dp_T^{\text{jet}} d\eta^{\text{jet}}} \right)^{-1} \times \int dp_T^{\text{jet}} d\eta^{\text{jet}} \frac{d\sigma^{(\text{jet } h)X}}{dp_T^{\text{jet}} d\eta^{\text{jet}} dz_h}$$

- Many theoretical uncertainties cancel out, e.g. choice of PDFs or initial state factorization scale

Numerical Results

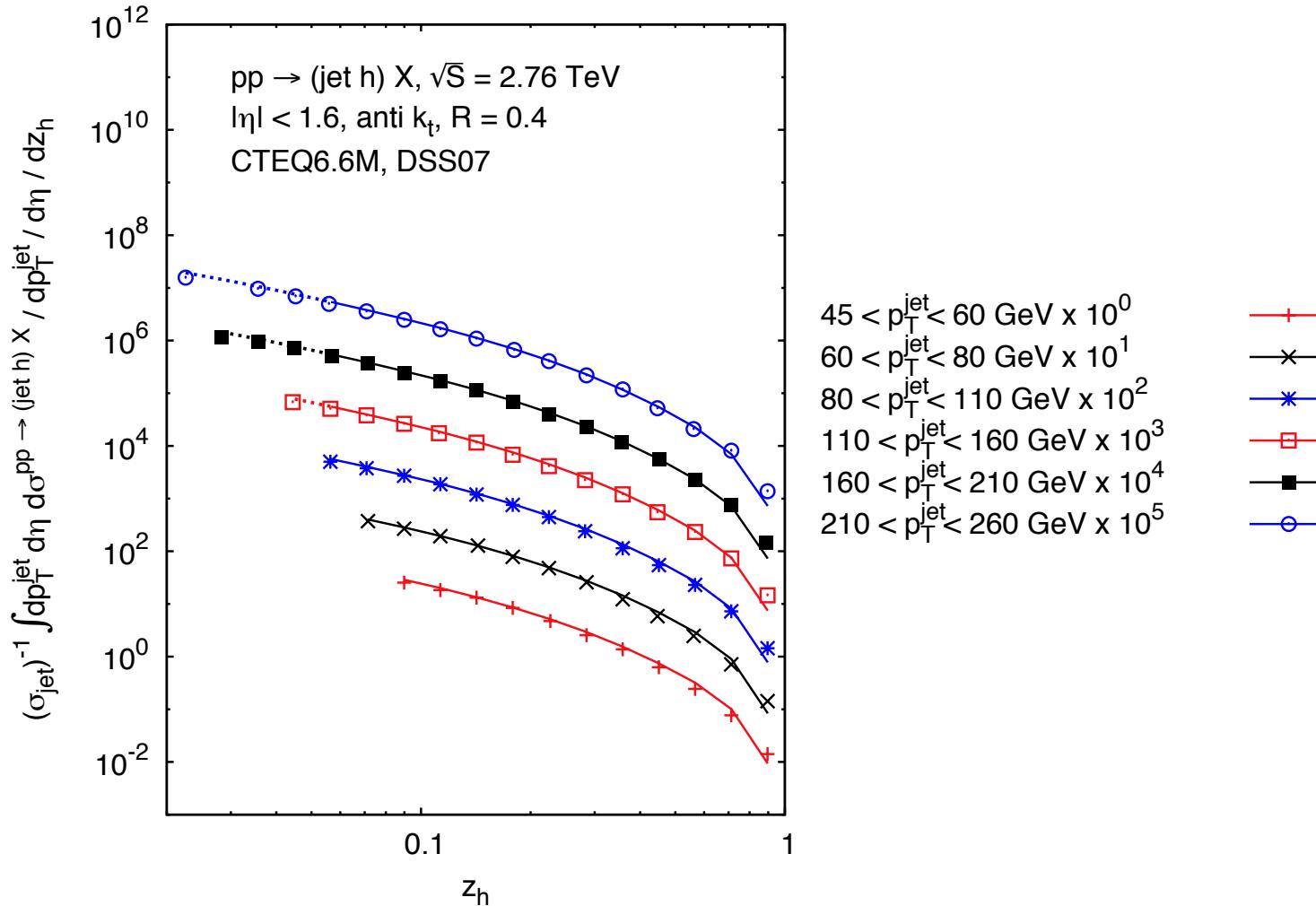
[Aad et al: Eur.Phys.J. C71,1795 (2011)]





Numerical Results

[ATLAS-CONF-2015-022]





Outline

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Photon-in-Jet

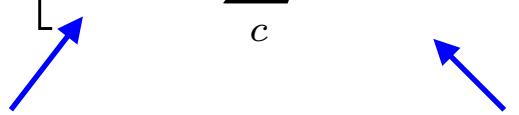
[Kaufmann, Mukherjee, Vogelsang: PRD93:114021]

- Poorly known photon fragmentation functions $D_c^\gamma(z, \mu_F)$
- process $pp \rightarrow (\gamma \text{ jet})X$
- “democratic jet algorithm”: photon is part of the jet!
- additional direct part, however not @LO for $z_\gamma < 1$

$$d\sigma = \sum_{a,b} f_a \otimes f_b \otimes \left[d\hat{\sigma}_{ab}^\gamma + \sum_c d\hat{\sigma}_{ab}^c \otimes D_c^\gamma \right]$$

direct

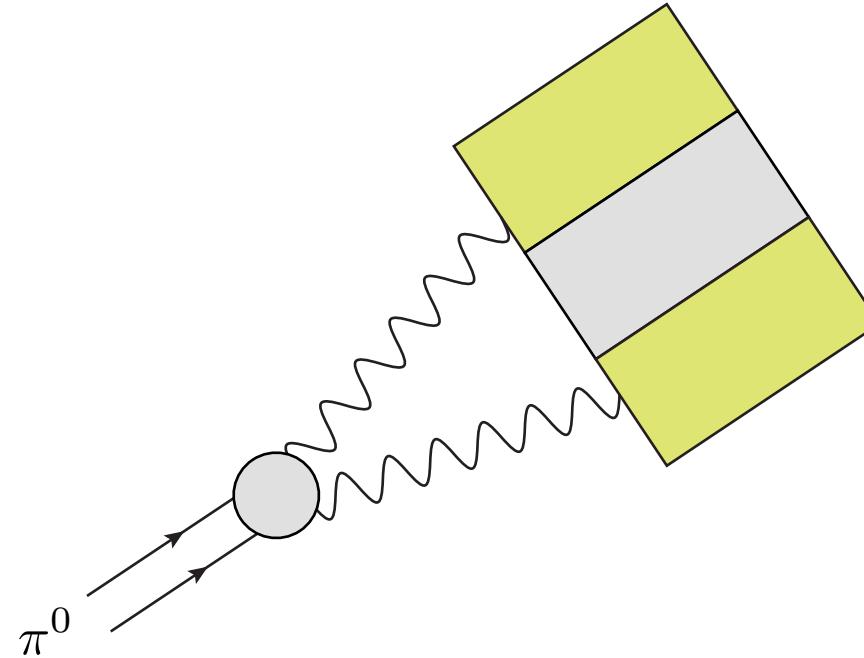
fragmentation





Background $\pi^0 \rightarrow \gamma\gamma$ ($\approx 98.8\%$)

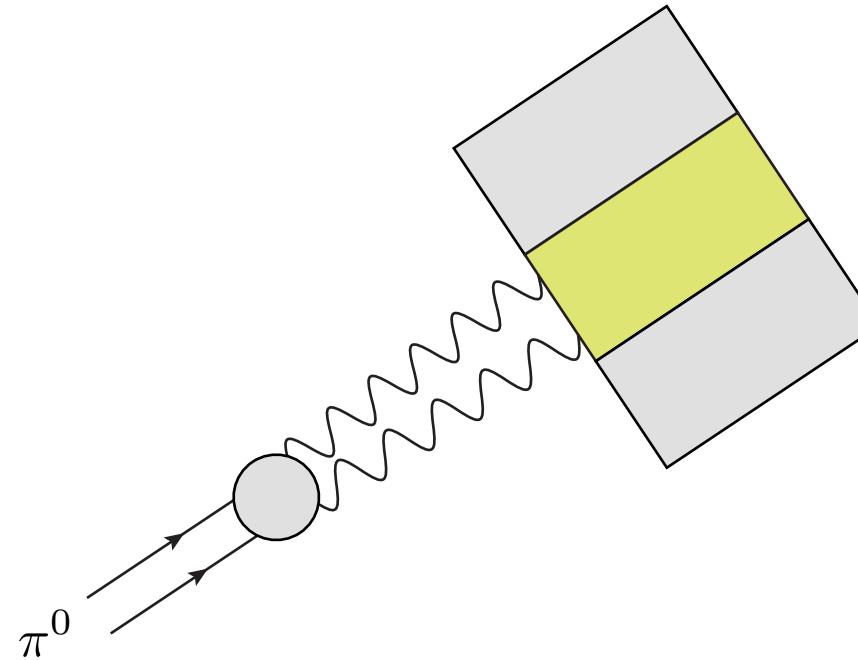
- As long as both photons are detected, pion may be reconstructed.





Background $\pi^0 \rightarrow \gamma\gamma$ ($\approx 98.8\%$)

- 2 collinear photons mimic 1 e.m. signal

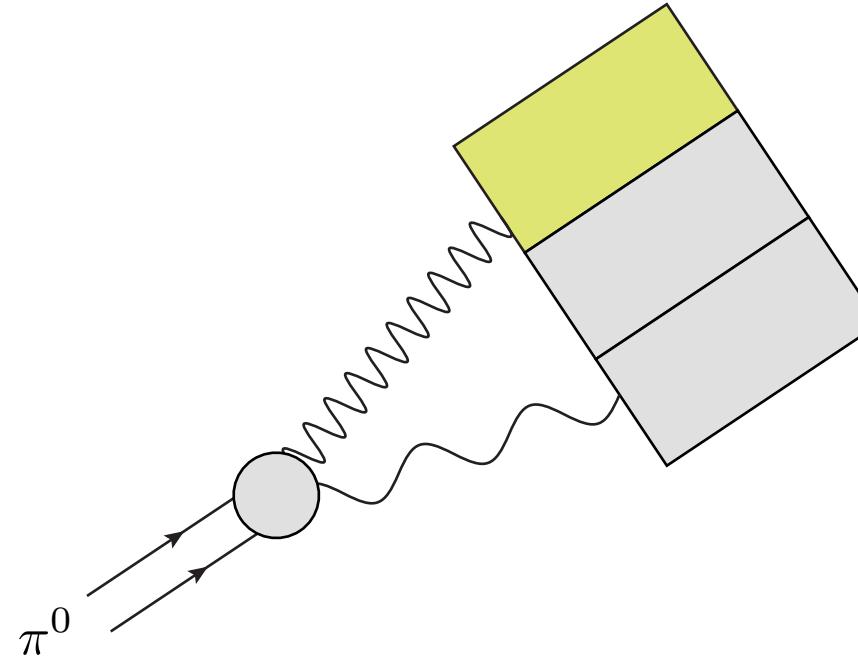


- with $\delta\eta = \delta\phi \approx 0.01$ relevant for $p_T^\pi > 20 \text{ GeV}$



Background $\pi^0 \rightarrow \gamma\gamma$ ($\approx 98.8\%$)

- Decay asymmetric in energy, one photon below energy threshold



- Energy threshold $\varepsilon < 10 \text{ MeV}$ required \rightsquigarrow main experimental challenge



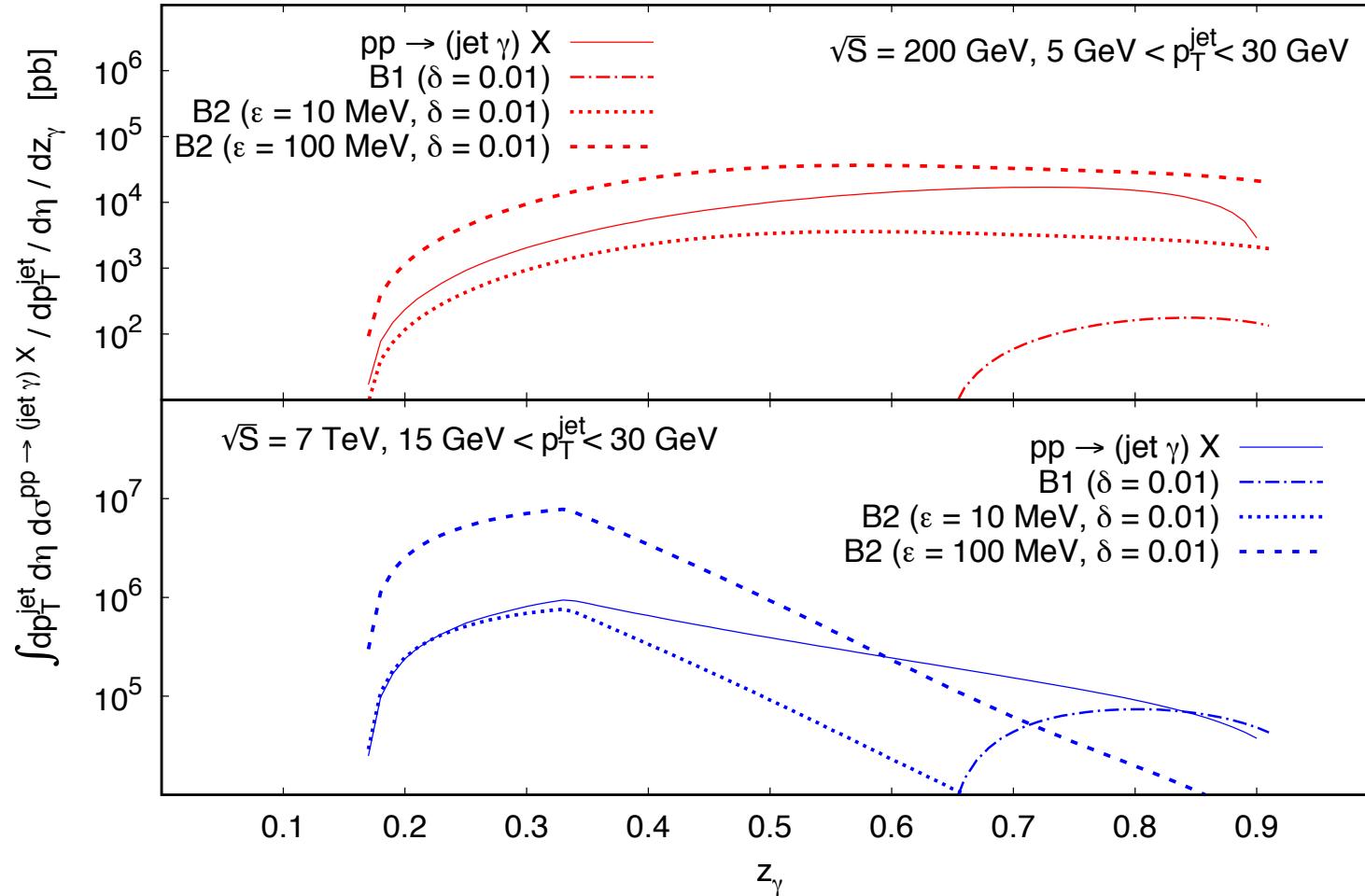
Background - Note

[Kaufmann, Mukherjee, Vogelsang: PRD93:114021]

- Most photon observables suffer from π^0 background
- Background calculation not limited to photon-in-jet
- Corresponding pion cross section is needed, e.g. background for single-inclusive photon productions requires single-inclusive pion cross section, etc.

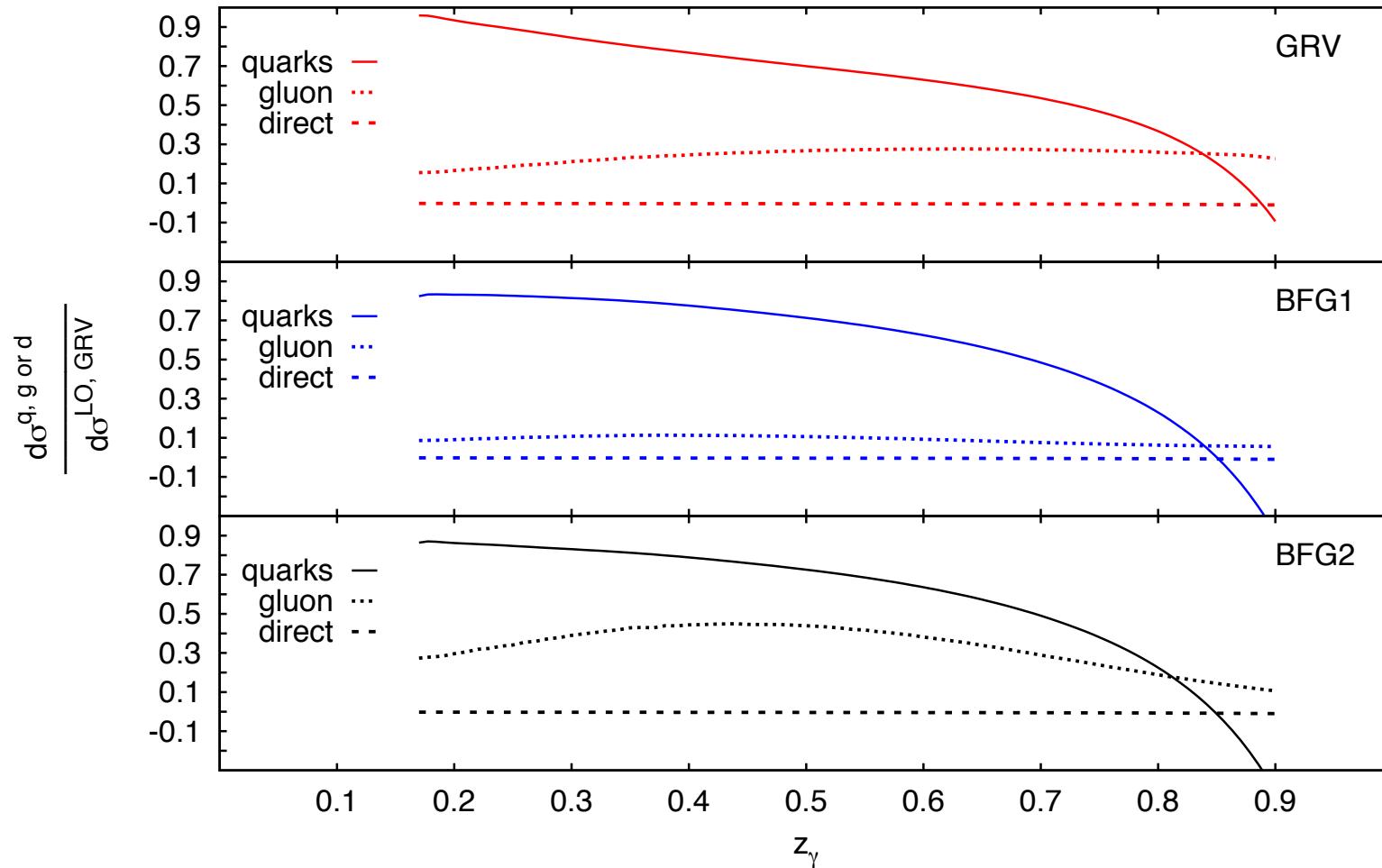
Numerical Results

$|\eta| < 1, p_T > 5 \text{ GeV, anti } k_t, R = 0.6, \text{CT10, GRV / DSS14}$



Numerical Results

$\sqrt{S} = 200 \text{ GeV}$, $5 \text{ GeV} < p_T^{\text{jet}} < 30 \text{ GeV}$, $|y| < 1$, $p_T > 5 \text{ GeV}$, anti k_t , $R = 0.6$, CT10





Summary

- $pp \rightarrow (\text{jet } h/\gamma) X$ are promising for $D_c^{h/\gamma}(z)$
- large amount of data
- various theoretical calculations in different frameworks
- Different definitions for z_h , charged jets, exclusive/inclusive,...



Outlook

- Detailed comparison Experiment vs. Theory — Need for higher precision?
- new (global) FF fits including the jet+hadron data
- Jet+Photon? An experimental challenge...



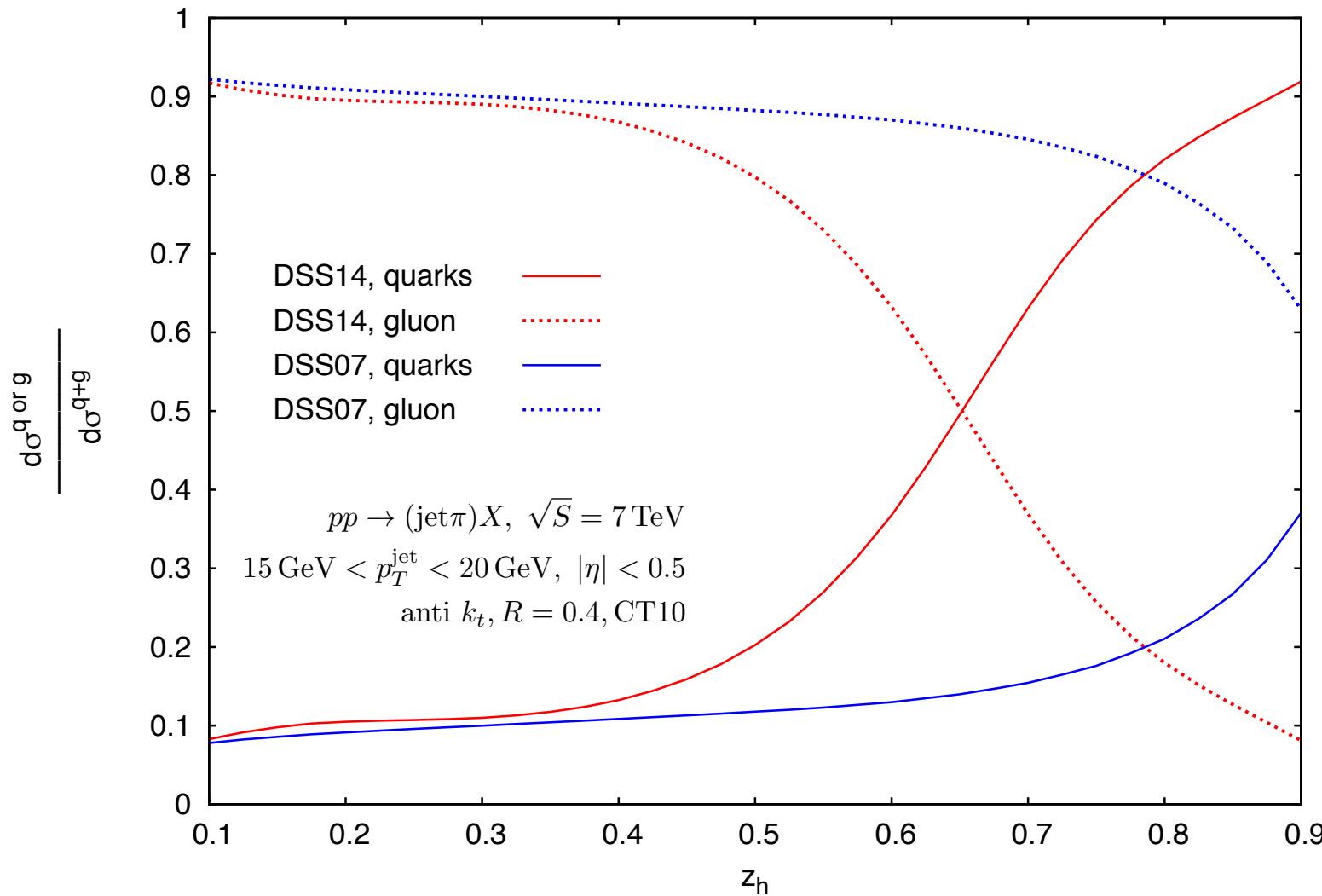
Thank you.



Backup Slides



Numerical Results



Calculation of the cross section

$$\begin{aligned}
 \frac{d\sigma^{H_1 H_2 \rightarrow (\text{jet } h) X}}{dp_T^{\text{jet}} d\eta^{\text{jet}} dz_h} &= \frac{2p_T^{\text{jet}}}{S} \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a^{H_1}(x_a, \mu_F) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b^{H_2}(x_b, \mu_F) \\
 &\times \int_{z_c^{\min}}^1 \frac{dz_c}{z_c^2} \frac{d\hat{\sigma}_{ab}^c(\hat{s}, \hat{p}_T, \hat{\eta}, \mu_F, \mu'_F, \mu_R)}{vdv dw} \sum_e j_{c \rightarrow e} \left(z_c, \frac{\mathcal{R} p_T^{\text{jet}}}{\mu'_F}, \mu_R \right) \\
 &\times \sum_{c'} \int_{z_h}^1 \frac{dz_p}{z_p} \tilde{j}_{e \rightarrow c'} \left(z_p, \frac{\mathcal{R} p_T^{\text{jet}}}{\mu''_F}, \mu_R \right) D_{c'}^h \left(\frac{z_h}{z_p}, \mu''_F \right)
 \end{aligned}$$

Calculation of the cross section

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analytical

Scales and Cancellations

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 \end{aligned}$$



Scales and Cancellations

- μ'_F is an artifact of the organization
- $d\hat{\sigma}_{ab}^c(\mu'_F)$ due to collinear factorization
- cancellation with $j(\mu'_F)$ because we require a jet in the final state
- powerful check, analytical AND numerical!
- μ''_F because of final state hadron
- needed to subtract collinear singularities when two partons jointly form the jet together
- $\log \mu''_F$ are standard scale logs that compensate evolution of FFs

Jet Functions

$$j_{a \rightarrow b}(z, \lambda) \equiv \delta_{ab} \delta(1-z) - \frac{\alpha_s}{2\pi} \left[\{P_{ba}(z) \log(\lambda^2(1-z)^2)\}_+ - P_{ba}^{(\epsilon)}(z) + \delta_{ab} \delta(1-z) I_a^{\text{algo}} \right]$$

- splitting functions for in $D = 4 - 2\epsilon$ dimensions:

$$\tilde{P}_{ij}(z) = P_{ij}(z) + \epsilon P_{ij}^{(\epsilon)}$$

- $\{\dots\}_+$ regularizes in the diagonal case, e.g. for P_{qq}

$$2C_F(1+z)^2 \left(\frac{\log(1-z)}{1-z} \right)_+ + P_{qq}(z) \log \lambda^2$$

- algorithm dependent quantities I_j^{algo} (pure numbers)

$$\bullet \quad \lambda = \frac{\mathcal{R} p_T^{\text{jet}}}{\mu'_F}$$

Jet Functions

$$\tilde{j}_{a \rightarrow b}(z_p, \kappa) \equiv \delta_{ab} \delta(1 - z_p) + \frac{\alpha_s}{2\pi} \left[\{ P_{ba}(z_p) \log(\kappa^2 (1 - z_p)^2) \}_+ - P_{ba}^{(\epsilon)}(z_p) + \delta_{ab} \delta(1 - z_p) I_a^{\text{algo}} + \mathcal{I}_{ba}^{\text{algo}}(z_p) \right]$$

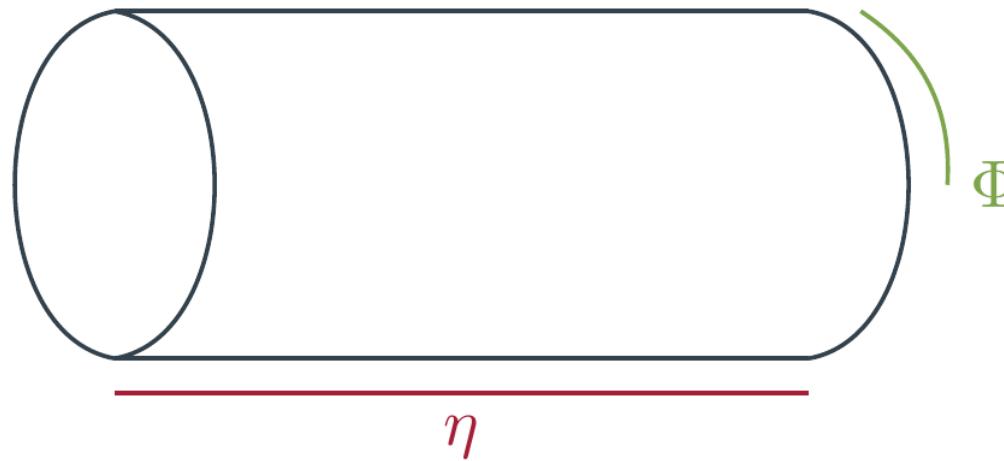
- algorithm dependent functions

$$\mathcal{I}_{c'c}^{\text{algo}}(z) = \begin{cases} 2P_{c'c}(z) \log\left(\frac{z}{1-z}\right) \Theta(1/2 - z) & \text{cone algorithm ,} \\ 2P_{c'c}(z) \log z & (\text{anti-})k_t \text{ algorithm ,} \\ P_{c'c}(z) \left[\log(z) + \log\left(\frac{z}{1-z}\right) \Theta(1/2 - z) \right] & J_{E_T} \text{ algorithm .} \end{cases}$$

- $\kappa = \frac{\mathcal{R} p_T^{\text{jet}}}{\mu_F''}$



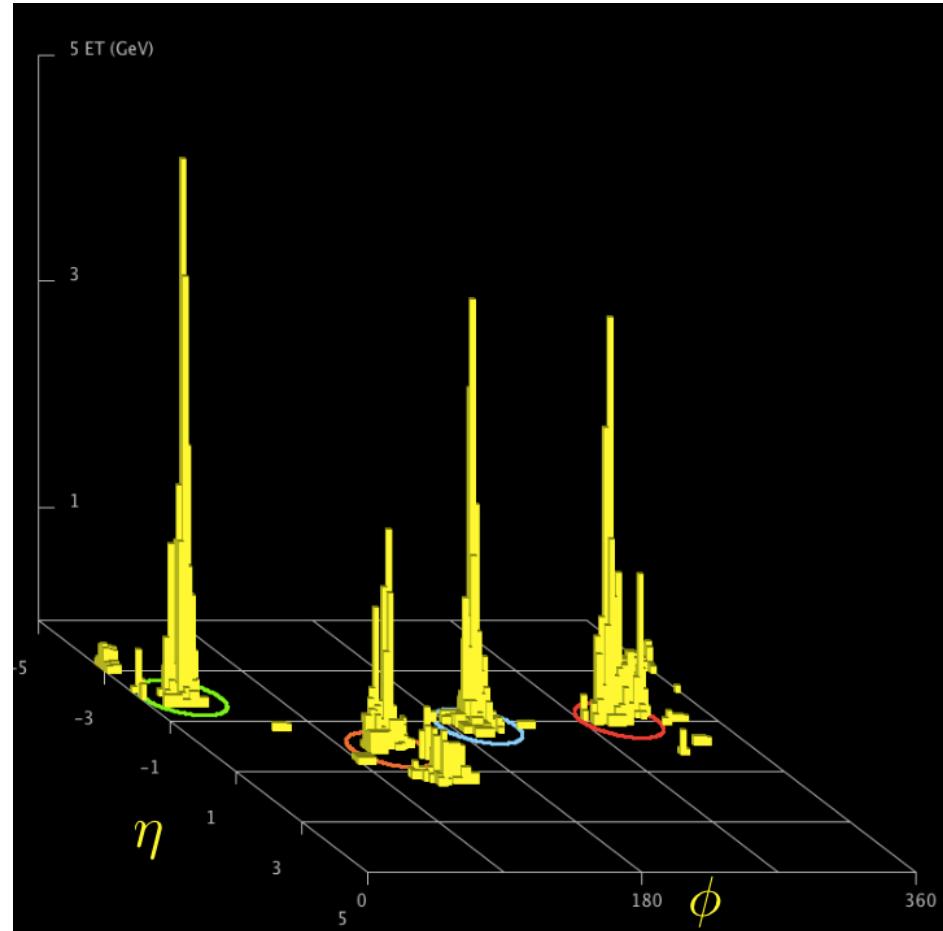
Rapidity and Azimuth



$$\eta = -\log \tan \frac{\theta}{2}, \quad [0, \pi] \rightarrow \mathbb{R}$$

Jet Algorithms

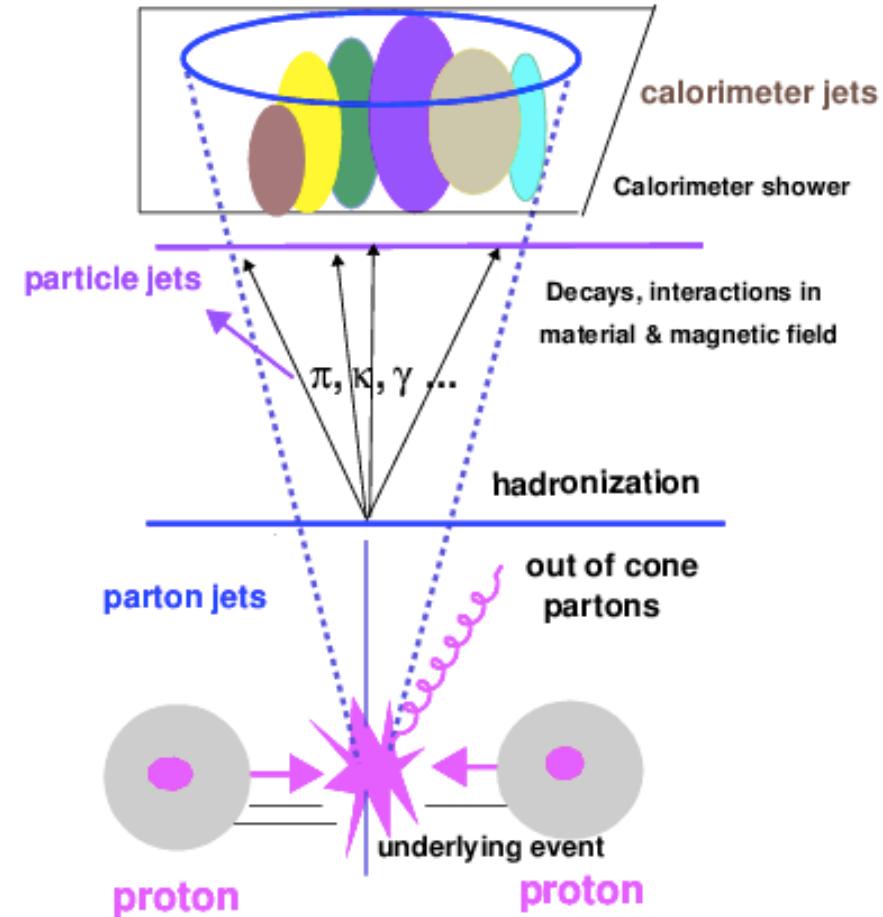
- no uniquely defined objects
- definition needed
- fulfill experimental and theoretical needs
- must be valid for all “levels”
- some examples:
 - cone
 - (anti-) k_T
 - J_{E_T}
 - ...



[Soper, CTEQ Summerschool 2013]

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[CMS-PAS-QCD-08-005]



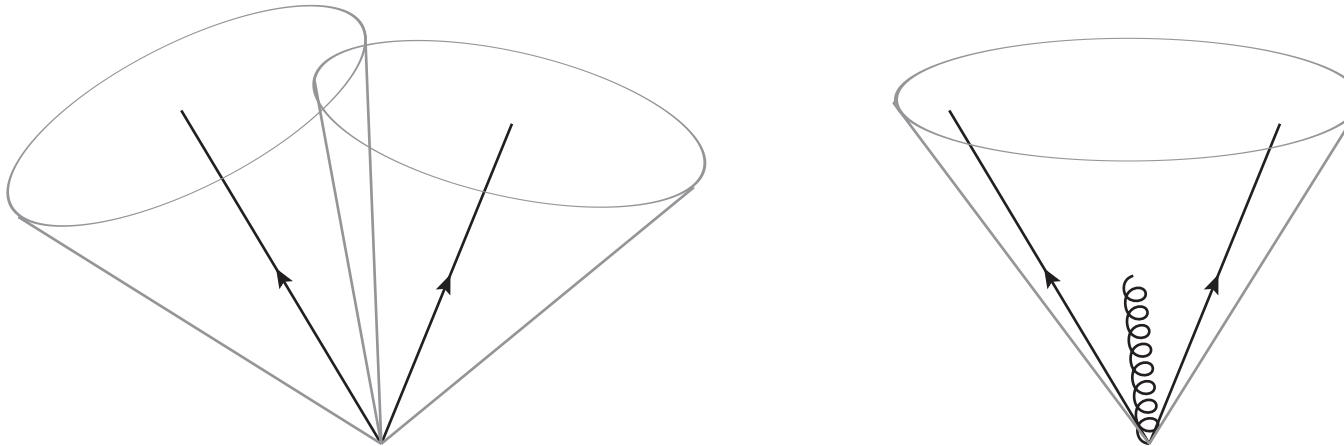
Cone algorithms

- all particles inside a geometrical cone

$$i \in \text{Jet} \Leftrightarrow (\phi_J - \phi_i)^2 + (\eta_J - \eta_i)^2 \leq R^2$$

Problem

- iterative search (seeds) for “stable” cones \Rightarrow not IRC safe





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- improvements (midpoint,...) only shift to higher orders
- IR safe: exact seedless algorithms, solve mathematical problem
- naive implementation: unusable $\mathcal{O}(N2^N)$
- improved implementation: SISCone $\mathcal{O}(N^2 \log N)$ [Salam, Soyez: JHEP0705:086]

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“size”

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kT-Algorithms

- distance between two objects

$$d_{ij} \equiv \min \left((k_T^i)^{2p}, (k_T^j)^{2p} \right) \frac{(\phi_i - \phi_j)^2 + (\eta_i - \eta_j)^2}{R^2}$$

- distance between object and beam

$$d_{iB} \equiv (k_T^i)^{2p}$$

- parameter p specifies algorithm:
 - $p = 1$ the kT algorithm [Ellis, Soper: PRD.48.3160],[Catani, Dokshitzer, Seymour, Webber: NuclPhys.B406]
 - $p = 0$ Cambridge-Aachen [Dokshitzer, Leder, Moretti, Webber: JHEP9708:001]
 - $p = -1$ anti-kT [Cacciari, Salam, Soyez: JHEP0804:063]

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kT-Algorithms

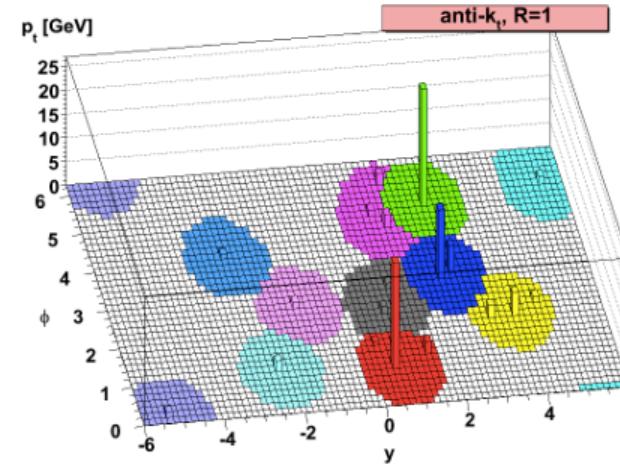
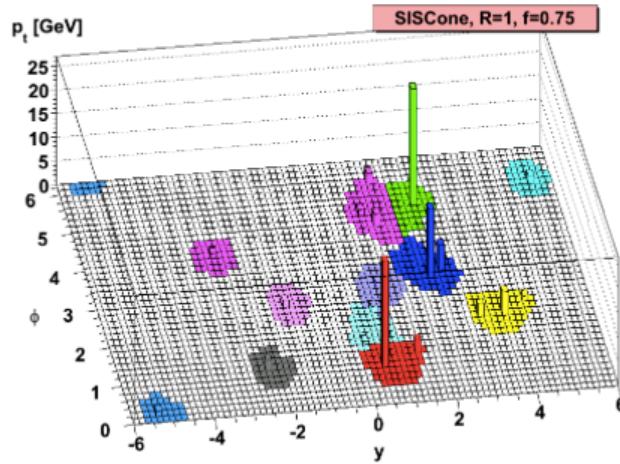
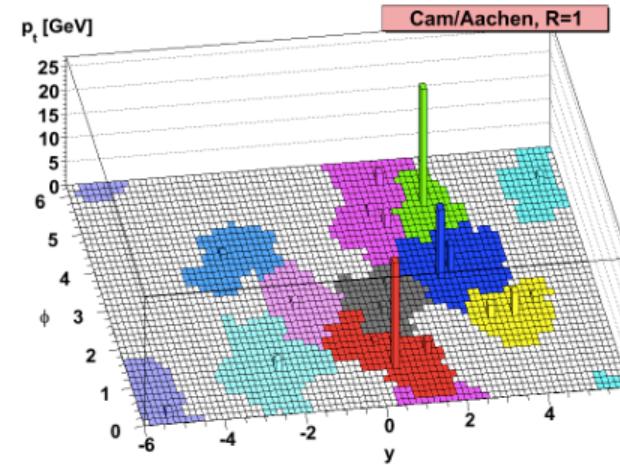
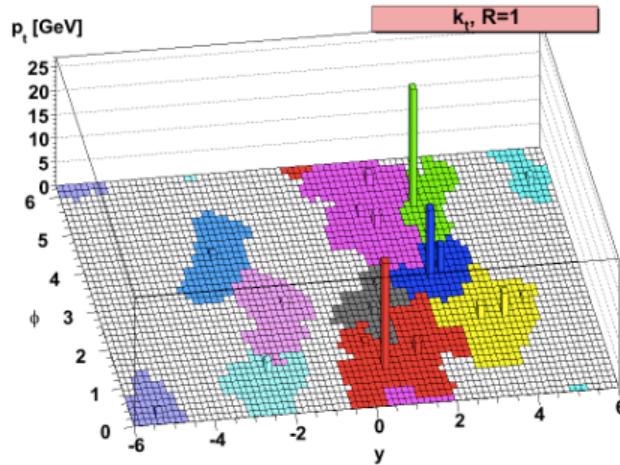
Iterative procedure

- initially: list of all particles
- while list not empty
 - for all objects in list: compute distances
 - if minimum distance is beam: object is a jet, remove from list
 - else, merge the two objects



kT- vs cone algorithms

[Cacciari, Salam, Soyez: JHEP0804:063]





Maximized Jet Function

[Georgi: arXiv:1408.1161; Kaufmann, Mukherjee, Vogelsang: PRD91:034001]

- jet as a sum of particles, increases jet energy
- not gain too much mass = be collimated

$$m^2 \equiv p_{\text{jet}}^2 = 2p_j \cdot p_k = 2E_j E_k (1 - \cos \theta_{jk})$$

- define jet function

$$J_\beta(p) \equiv E - \beta \frac{m^2}{E} , \quad \beta > 1$$

- increases with increasing energy
- decreases with increasing mass squared over energy
- highest J jet maximizes the jet function
- iterative procedure

Maximized Jet Function

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“inverse size”

- increases with increasing energy
- decreases with increasing mass squared over energy
- highest J jet maximizes the jet function
- iterative procedure