

# Unpolarized and helicity distributions from global fits

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# Overview

## □ A PDF landscape

## □ Proton and neutron PDFs: the CJ15 fit

- Hadronic physics I : d/u ratio
- Hadronic Physics II : dbar / ubar ratio
- Nuclear Physics output: off-shell parton corrections
- High-energy: BSM searches

## □ Iterative Monte Carlo – the JAM approach

- “JAM15” helicity PDFs
- A strange puzzle
- “JAM-FF16” fragmentation functions

### *REFERENCES:*

\* *Accardi, PoS DIS2015 001 – “PDFs from protons to nuclei”*

\* *Accardi et al, PRDD93 (2016) 114017 – the CJ15 global fit*

\* *N.Sato et al, PRD93 (2016) 074005 – the JAM15 fit*

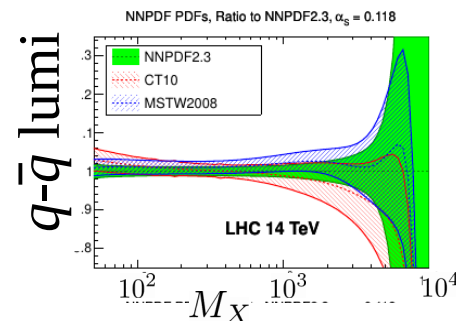
# A PDF landscape

# Why (n)(p)PDFs ?

Accardi – *Mod.Phys.Lett. A28 (2013) 35*  
 Forte and Watt – *Ann.Rev.Nucl.Part.Sci. 63 (2013) 291*

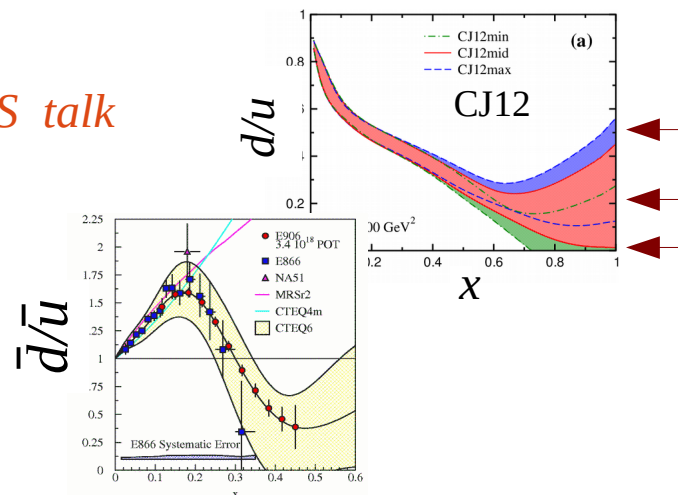
## High-energy (*large to small x*) → *J.Rojo's talk*

- Beyond the Standard Model searches
- NuTeV weak mixing angle
- Precision (Higgs) physics
- Small-x and gluonic “matter”



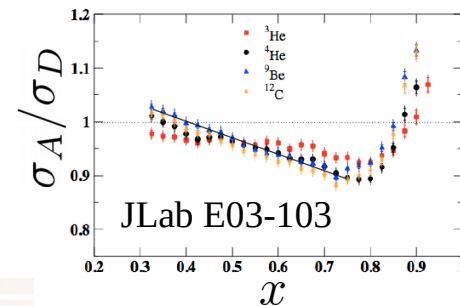
## Hadron structure (*large to medium x*) → *THIS talk*

- Effects of confinement on valence quarks
- $q - \bar{q}$  asymmetries; isospin asymmetry
- Strangeness, intrinsic charm
- Spin, parton orbital motion, color dynamics

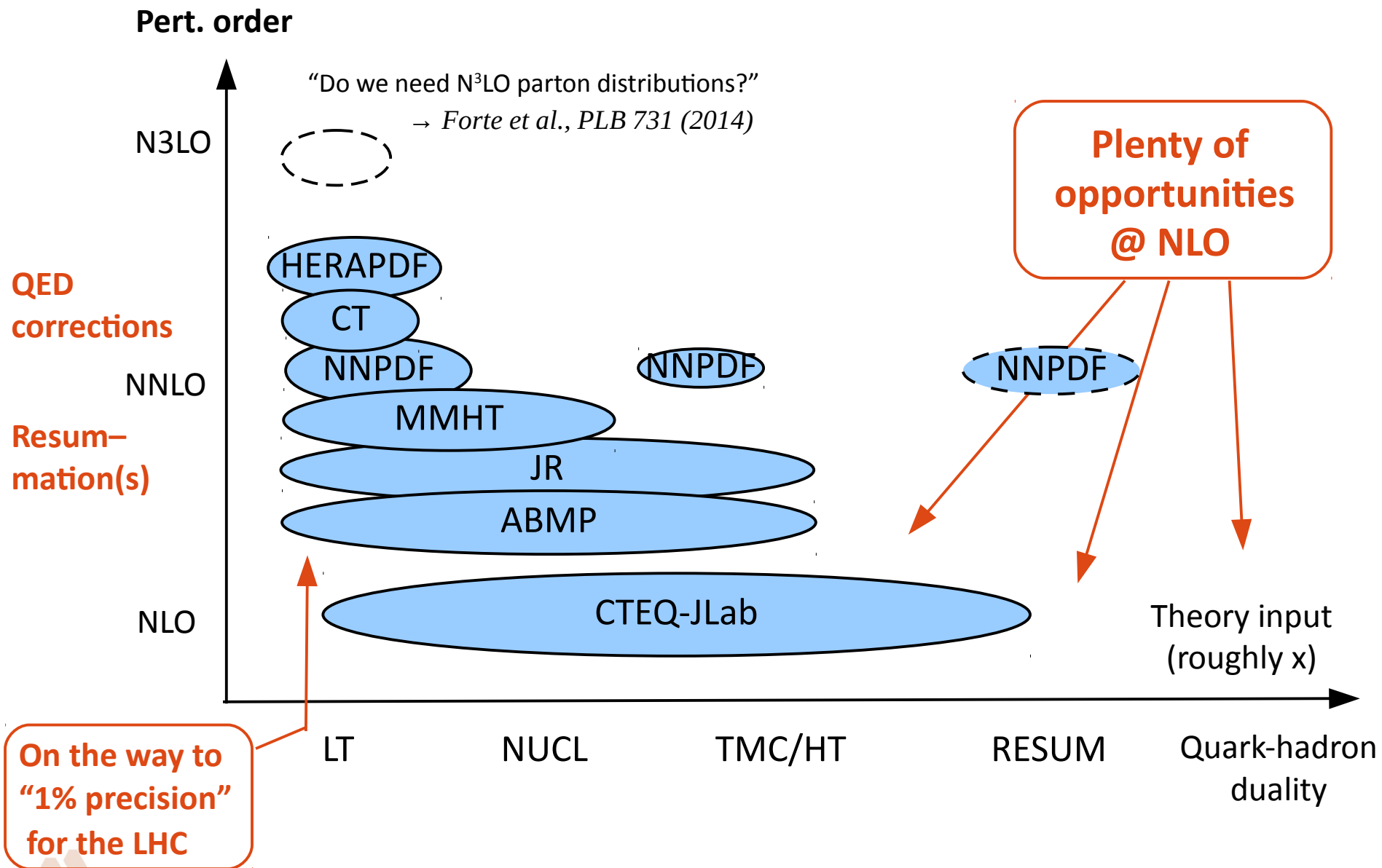


## Nuclear Physics → *Mostly in backup*

- Bound nucleons, EMC effect, SRC
- p+A and A+A collisions at RHIC / LHC
- Color propagation in nuclear matter

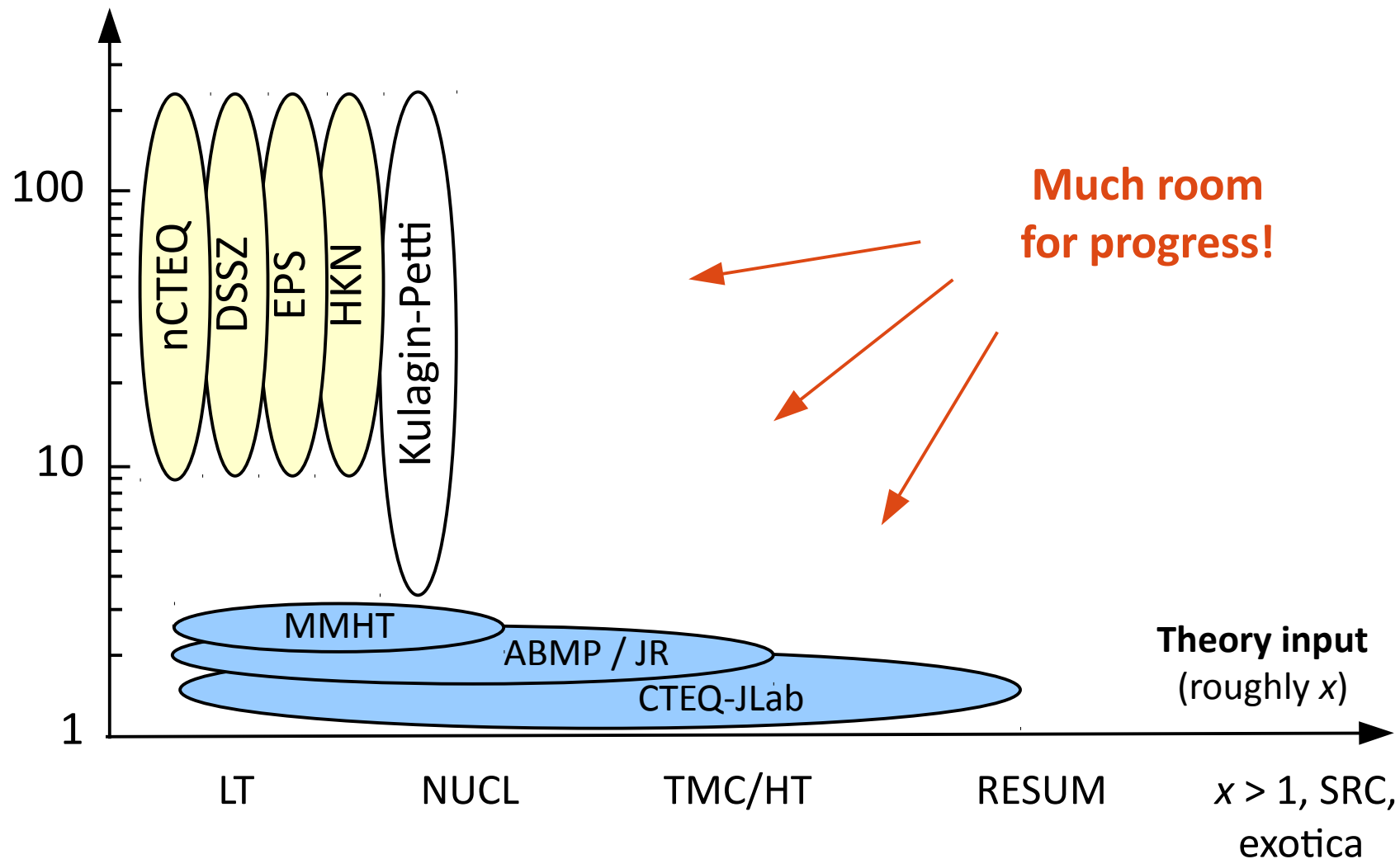


# A PDF landscape



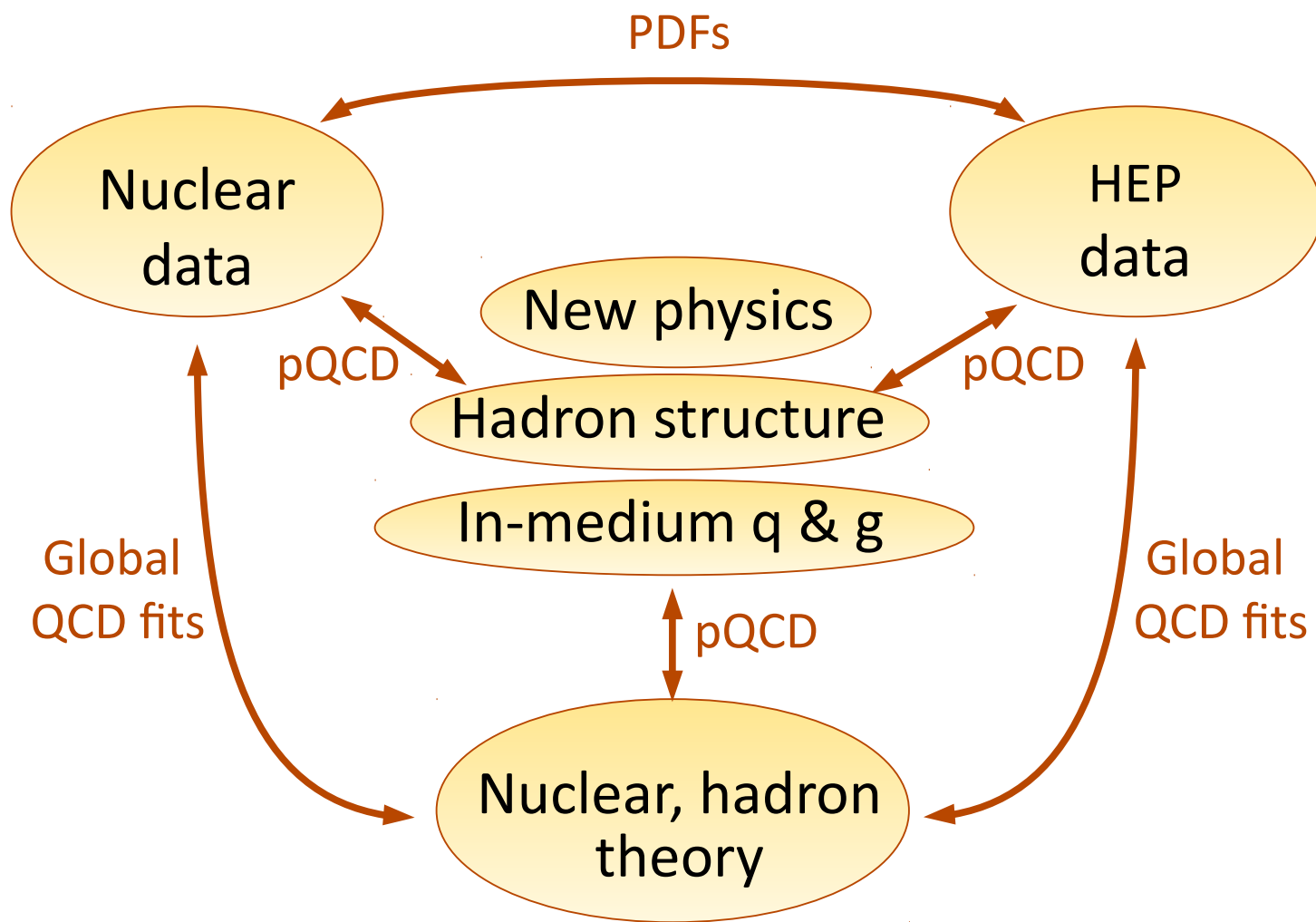
# A nPDF landscape

Atomic number

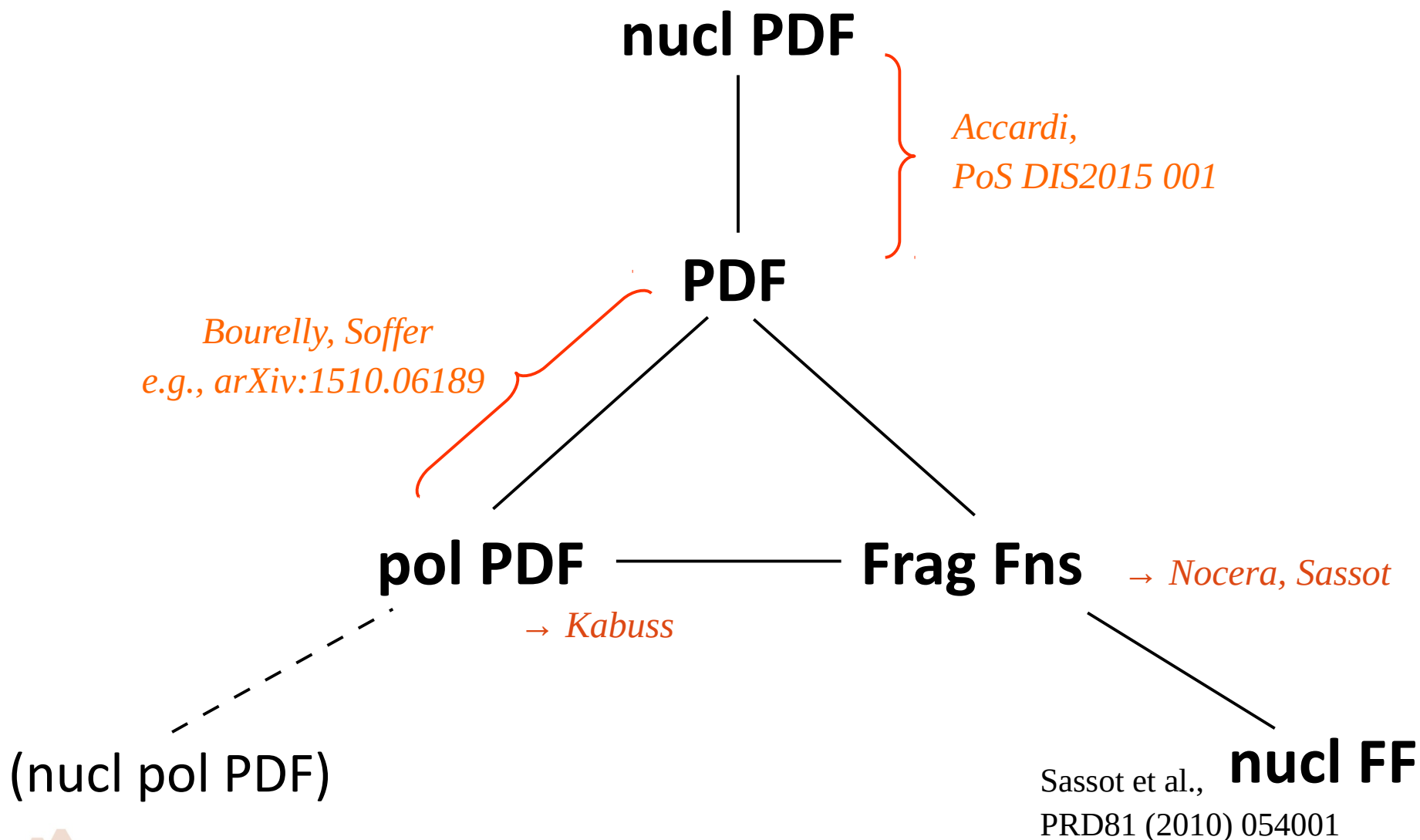


# Needs the betrothal of HEP and NUCL

- A global approach across subfields



# Other possible marriages



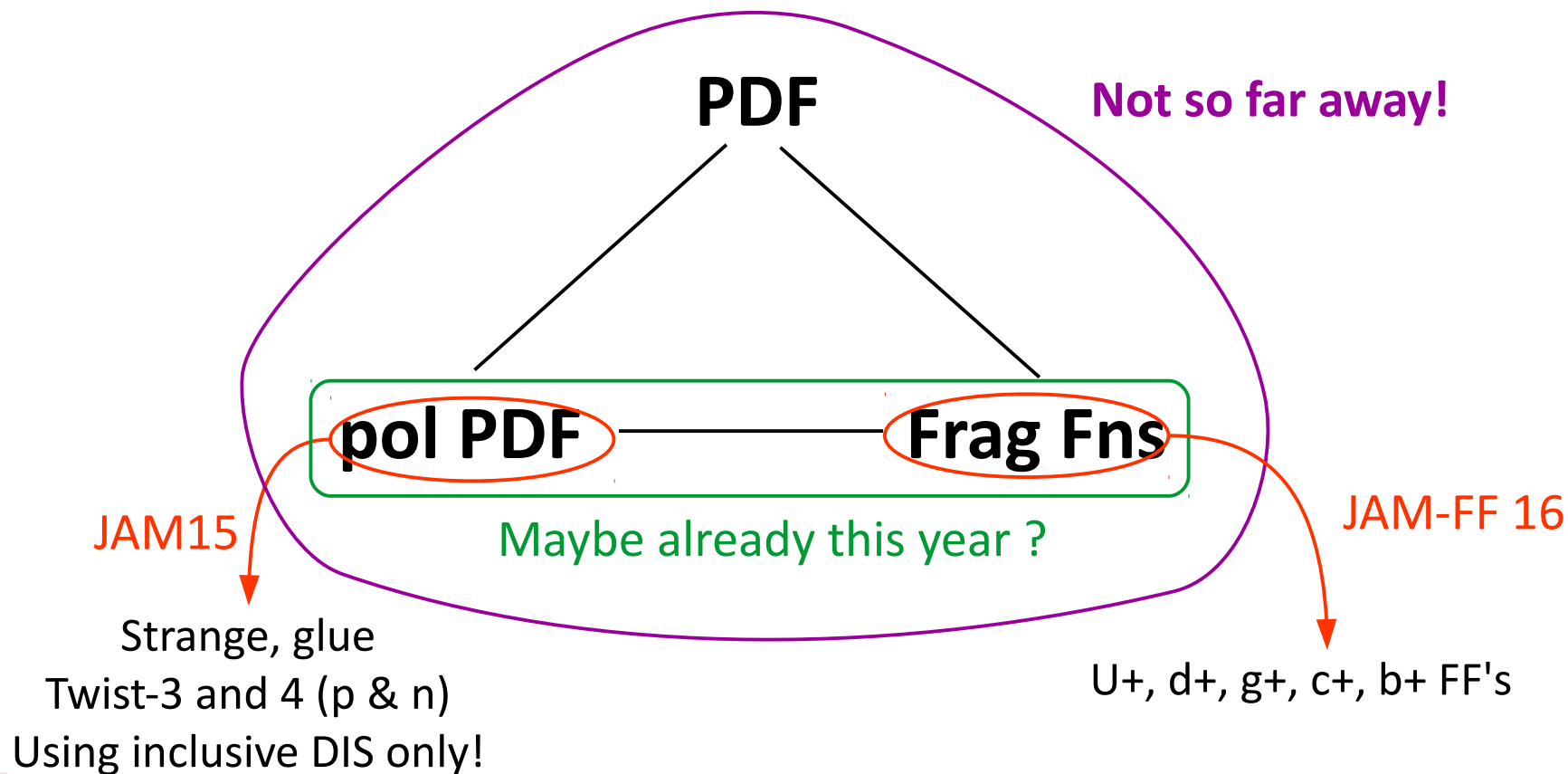


# New fitting methods

- More computing power, efficient implementations
    - New fitting, analysis methods
  - Traditional fits:
    - Detailed  $\chi^2$  scans, refined statistical analysis
  - Monte carlo fitting methods:
    - **NNPDF**: bootstrap + neural network fit
    - **JAM**: bootstrap + Iterative Monte Carlo (IMC) approach
      - *Sato, Ethier, Kuhn, Melnitchouk, Accardi (2015)*
- Large number of parameters, trustable uncertainty estimates
- Self organizing maps → *Liuti et al.*

# Iterative Monte Carlo approach

- ❑ Provides control over large number of parameters
- ❑ Maximizes extraction of physics information from data



# Proton and neutron PDFs - the CJ15 global fit -

*Accardi, Brady, Melnitchouk, Owens, Sato  
PRD93 (2016) 114017*

# The CTEQ-JLab global fits

## □ Collaborators:

- **Theory:** A.Accardi, W.Melnitchouk, J.Owens, N.Sato
- **Experiment:** E.Christy, C.Keppel, P.Monaghan

## □ All-x PDF global fits, focused on the “large” x region

- Maximize use of large-x data (esp. DIS)
- Include all relevant large-x / small- $Q^2$  theory corrections
- *Quantitatively evaluate theoretical systematic errors*
- *Use PDFs as tools for nuclear and particle physics*

## □ Latest public release: CJ15

- [www.jlab.org/cj](http://www.jlab.org/cj)
- Included in LHAPDF

# The CJ15 fit at a glance

	JLab & BONUS	HER MES	HERA I+II	Tevatron new W,Z	LHC	v+A di- $\mu$	Large-x treatment			
							Nucl.	HT TMC	Flex $d$	low-W DIS
<b>CJ15 *</b>	✓	✓	✓	✓	<i>in prog.</i>	✗	✓ ✓	✓	✓	✓
CT14			DIS 2016	✓ ✗	✓	✓				✓
MMHT14			✗✗✗	✓ ✗	✓	✓	✓			
NNPDF3.0					✓	✓		TMC only		
JR14	✓				✓	✓	✓	✓		
ABM15 **				✓ ✗	✓	✓				✓
HERAPDF2.0			✓	✗						

\* NLO only \*\* No jet data ✗ see 1503.05221 ✗✗ see 1508.06621 ✗✗ no reconstructed W

# New in CJ15

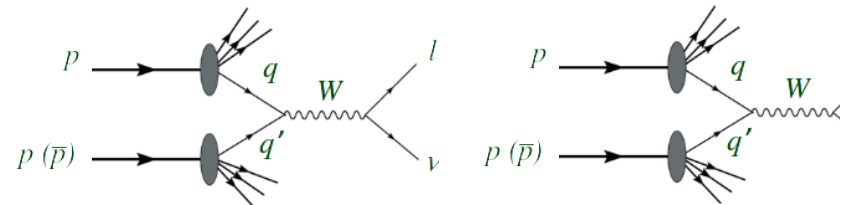
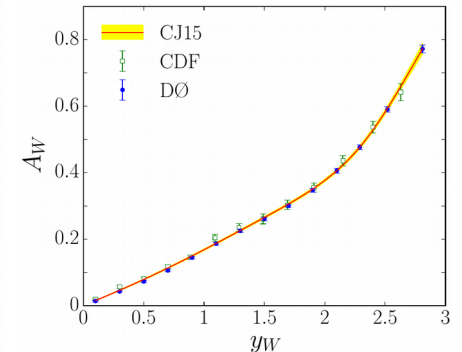
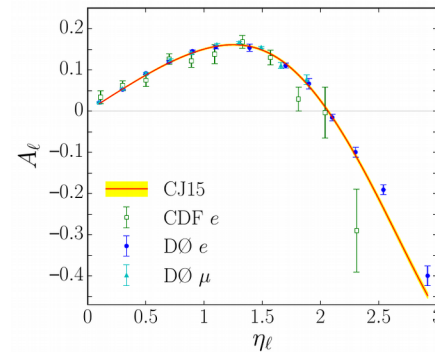
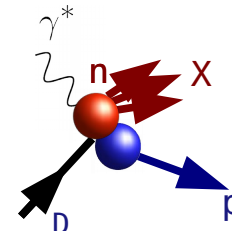
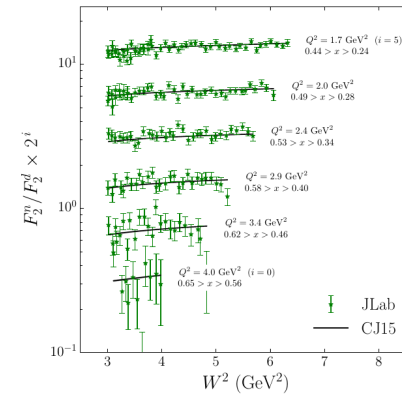
□ s-ACOT scheme for heavy flavors

□ New data:

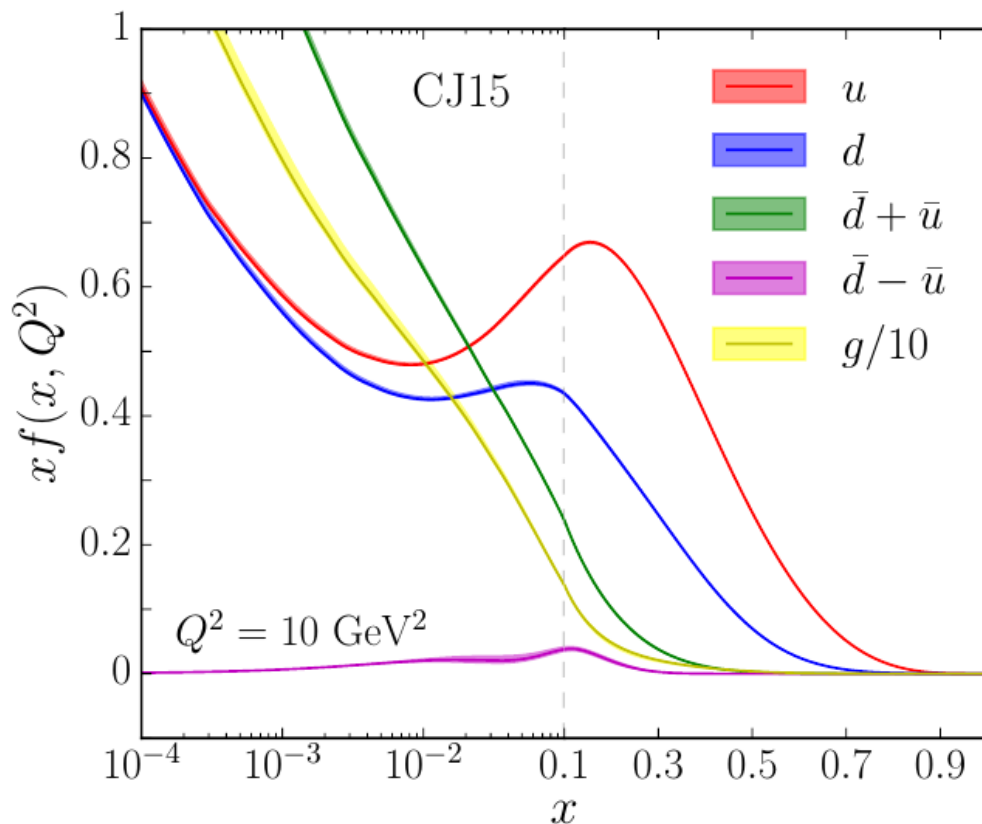
- BONUS spectator tagged DIS
- HERA I+II combination
- HERMES F2
- High-statistics W-boson charge asymmetries from DØ

□ New off-shell nucleon treatment in deuteron targets (DIS and DY)

- Parametrized vs. modeled → absorbs wave function uncertainty
- Comparison to extractions from deuteron and DIS on heavier targets



# CJ15 - PDFs



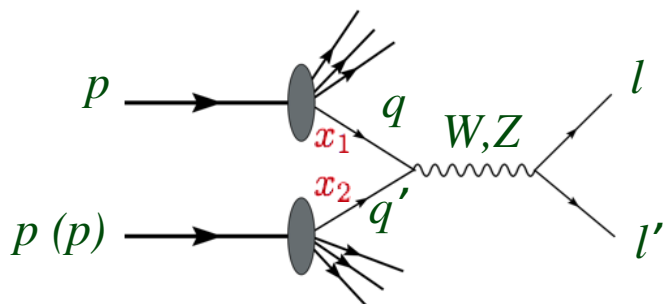
- Hessian error analysis
  - Correlated errors where available
- Error bands displayed for  $\Delta\chi^2 = 2.71$  (90% confidence level in a perfect, Gaussian world)

□ NLO fit gives  $\chi^2/\text{datum} = 1.04$

□ LO fit much worse – cannot accommodate  $Q^2$  dependence of data

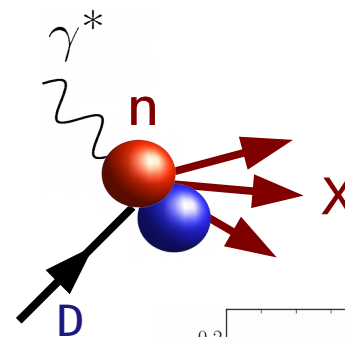
# NUCL / HEP symbiosis

- $W$  and  $Z \rightarrow$  constrain  $d$ -quark at largest  $x$  on proton targets

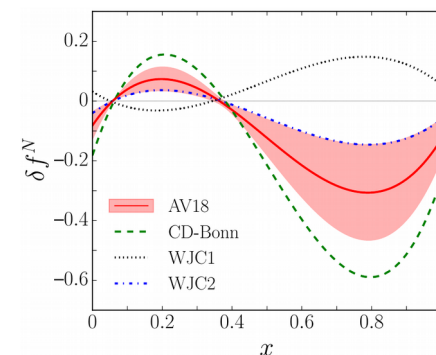


$$A_W(y) \approx \frac{d/u(x_2) - d/u(x_1)}{d/u(x_2) + d/u(x_1)}$$

- Compare to deuteron DIS
  - $\rightarrow$  constrain deuteron corrections
  - $\rightarrow$  **Off shell correction – first time in Deuteron!**



- Abundant DIS deuteron data
  - $\rightarrow$  precise  $u, d$  flavor separation





# NUCL / HEP symbiosis

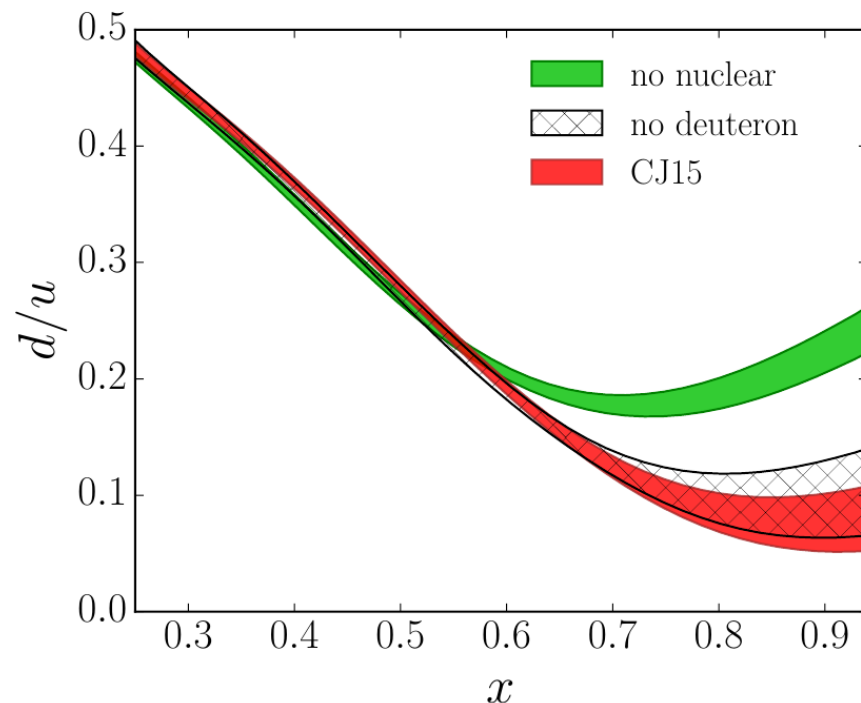
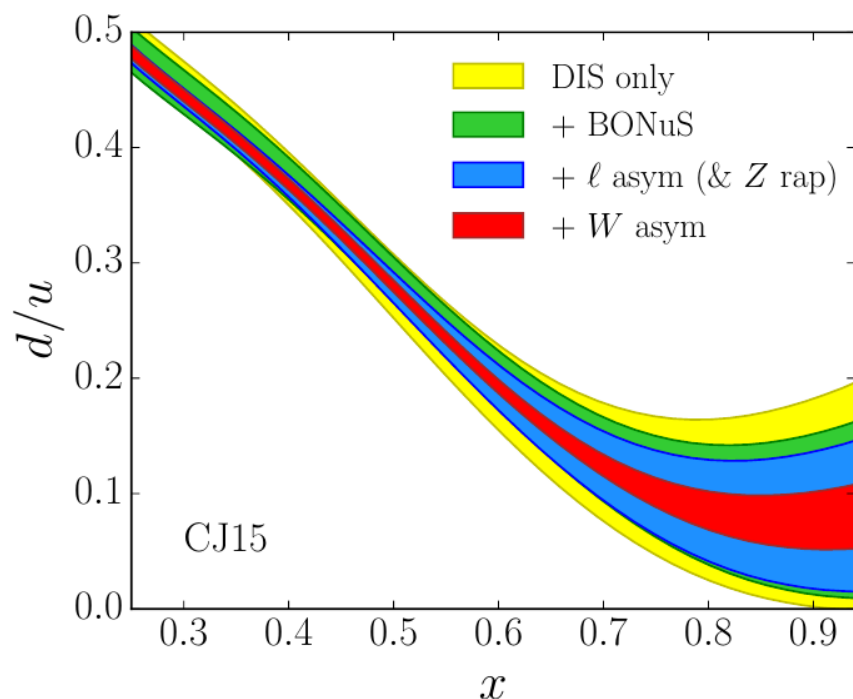
Observable	Experiment	# points	$\chi^2$			
			LO	NLO	NLO (OCS)	NLO (no nucl)
DIS $F_2$	BCDMS ( $p$ ) [81]	351	430	<b>438</b>	436	440
	BCDMS ( $d$ ) [81]	254	297	<b>292</b>	289	301
	SLAC ( $p$ ) [82]	564	488	<b>434</b>	435	441
	SLAC ( $d$ ) [82]	582	396	<b>376</b>	380	<b>507</b>
DIS $F_2$ tagged	Jefferson Lab ( $n/d$ ) [21]	191	218	<b>214</b>	213	219
$W$ /charge asymmetry	CDF ( $e$ ) [88]	11	11	<b>12</b>	12	13
	DØ ( $\mu$ ) [17]	10	37	<b>20</b>	19	29
	DØ ( $e$ ) [18]	13	20	<b>29</b>	29	14
	CDF ( $W$ ) [89]	13	16	<b>16</b>	16	14
	DØ ( $W$ ) [19]	14	39	<b>14</b>	15	<b>82</b>
$Z$ rapidity	CDF ( $Z$ ) [90]	28	100	<b>27</b>	27	26
	DØ ( $Z$ ) [91]	28	25	<b>16</b>	16	16
	⋮	⋮	⋮	⋮	⋮	⋮
$\chi^2$ /datum			1.33	<b>1.04</b>	1.04	1.09

# NUCL / HEP symbiosis

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Z rapidity	CDF ( $Z$ ) [90]	28	100	<b>27</b>	27	26

- ❑ Ignoring nuclear dynamics, SLAC( $d$ ) and DØ( $W$ ) pull  $d$  quark in opposite directions
  - **DØ (W) data determine nuclear corrections !!**
  - other asymmetries inconclusive by themselves
  - **BONUS data validate DØ(W) analysis**

# Hadronic physics output 1: d/u ratio



□ **d-quark determined by  $p+p \rightarrow W+X$**

□ **Nuclear corrections dominant at large  $x$**

- SLAC(d)'s statistical power used to fit the off-shell function...
- ... and to improve d/u flavor separation, esp. at  $x < 0.3$  (see backup)

# Hadronic physics output 1: d/u ratio

→  $d/u$  ratio at high  $x$  of interest for nonperturbative models of nucleon

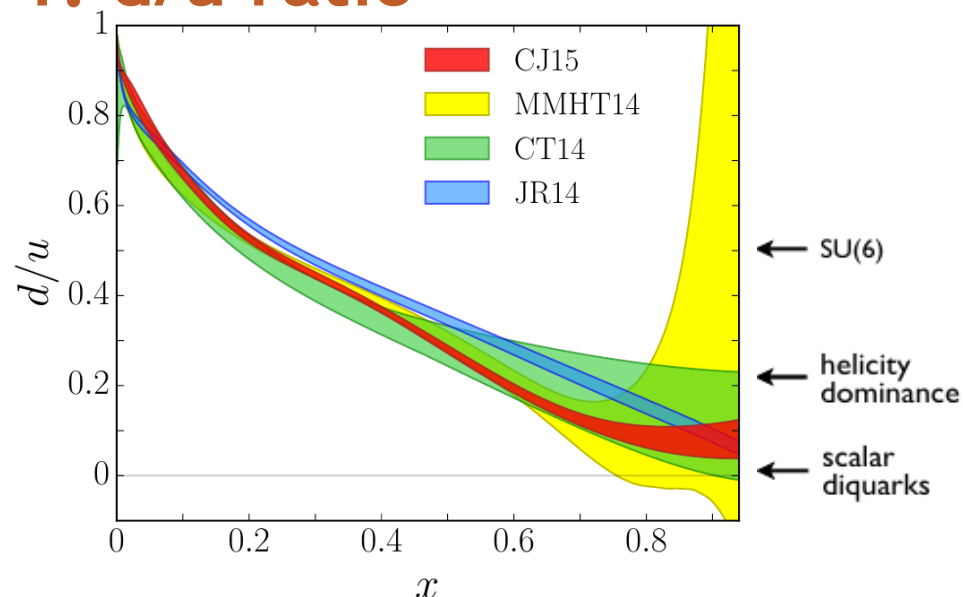
→ **CJ15:**

more flexible parametrization

$$d \rightarrow d + b x^c u$$

allows finite, nonzero  $x = 1$  limit

(standard PDF form gives 0 or  $\infty$  unless  $a_2^d = a_2^u$ )



**MMHT14:** fitted deuteron corrections  
standard  $d$  parametrization  
→ “UNDERCONSTRAINED”

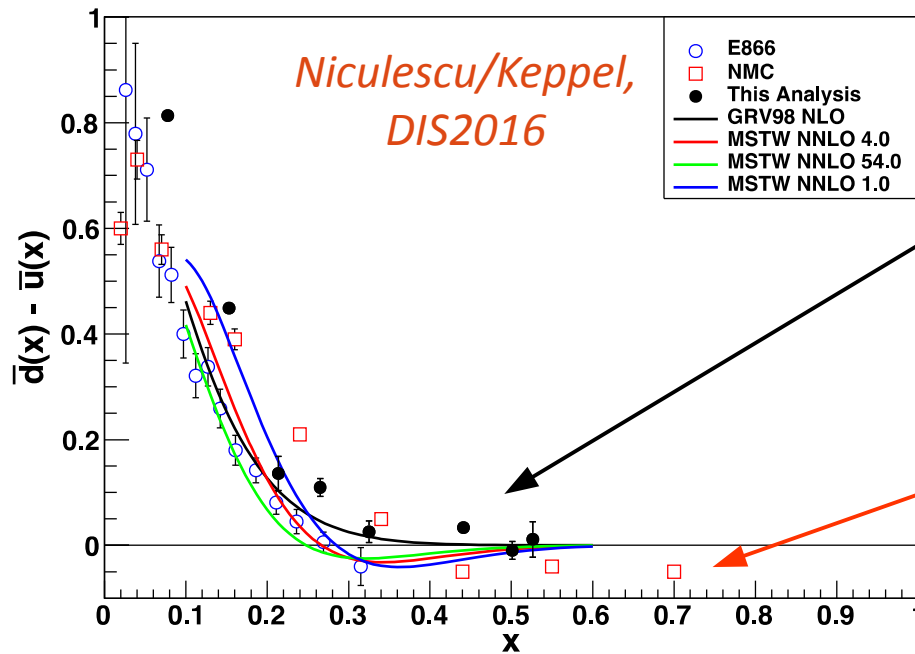
**JR14 (and ABM12):**

Similar deuteron corrections  
standard  $d$ ; no lepton/W asym.  
→ “OVERCONSTRAINED”

**CT14:**  $\beta_u = \beta_d \implies d/u$  finite  
No nuclear corrections

# Hadronic physics output 2: dbar/ubar

[AA, Keppel, Niculescu, DIS 2016]



*Niculescu/Keppel,  
DIS2016*

○ E866  
□ NMC  
● This Analysis  
— GRV98 NLO  
— MSTW NNLO 4.0  
— MSTW NNLO 54.0  
— MSTW NNLO 1.0

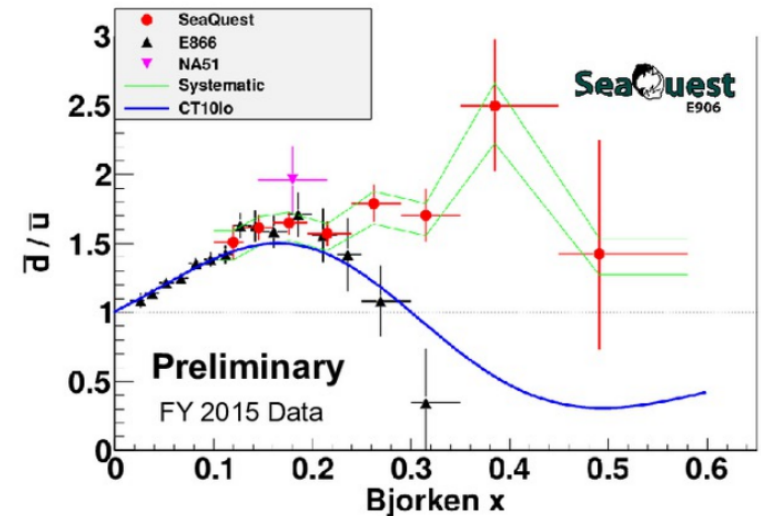
This analysis: use CJ15 to “remove” nuclear corrections

NMC = Peng et al. analysis of NMC data as in publication (use MSTW08 for  $u_v, d_v$ )

□ No evidence fo sign change at large x !

□ And SeaQuest agrees!

— presented for the first time a few days later



Bryan Kerns

*B.Kerns, DNP April 2016*

# Iterative Monte Carlo: the JAM approach

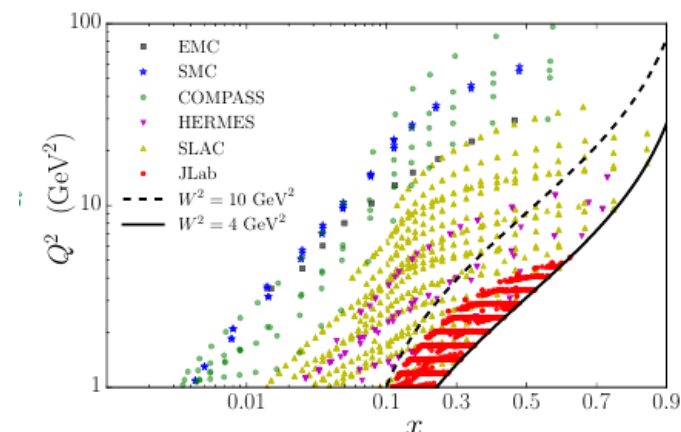
*Sato, Ethier, Melnitchouk, Kuhn, Accardi*

*PRD 2016 [arXiv:1601.07782] & work in progress*

# Fitting strategy

## Parametrization

- $xf(x) = Nx^a(1-x)^b(1+c\sqrt{x}+dx)$
- LT quark distributions  $\rightarrow \Delta u^+, \Delta d^+, \Delta s^+, \Delta g$
- T3 quark distributions  $\rightarrow D_u, D_d$
- T4 structure functions  $\rightarrow H_p, H_n$



**Chi-squared minimization**  $\rightarrow$  with correlated systematic uncertainties

$$\chi^2 = \sum_i \left( \frac{D_i - T_i \left( 1 - \sum_k r^k \beta_i^k / D_i \right)^{-1}}{\alpha_i} \right)^2 + \sum_k (r^k)^2$$

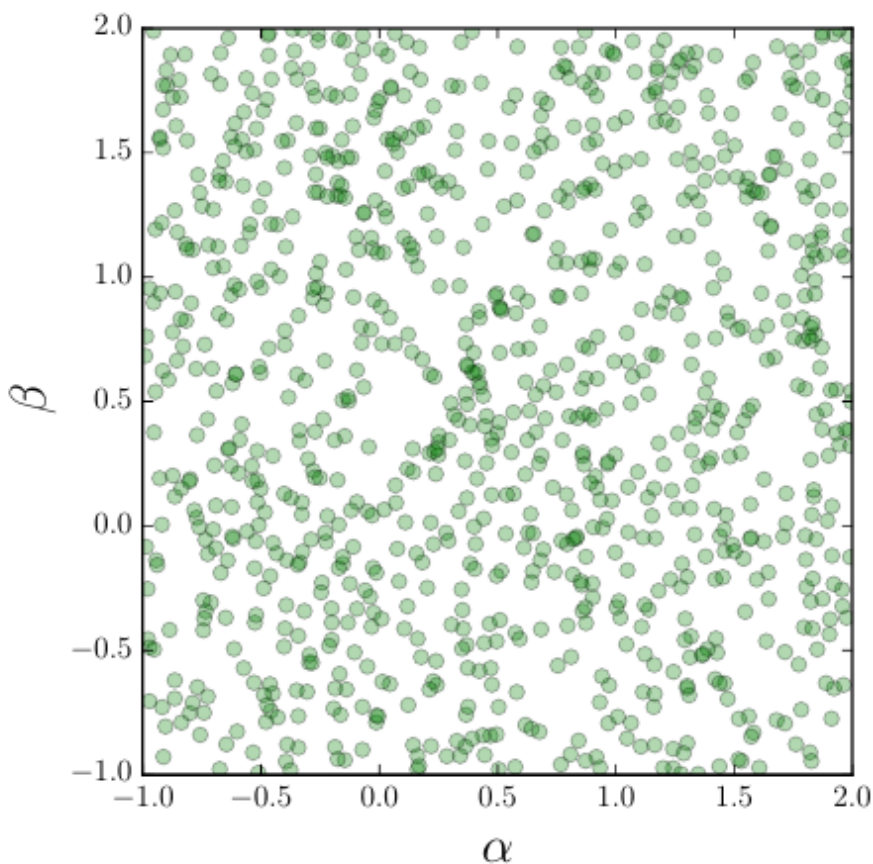
## Issues

- Stability in the moments (e.g.  $\Delta\Sigma^{(1)}$ )
- Is the solution given by a single fit unique?  $\rightarrow$  False minima
- Is over-fitting present in our fits?
- Which parameters should be fixed and at which value?
- Determination of uncertainty bands.

**Solution**  $\rightarrow$  MC approach; iterate until convergence  $\rightarrow$  **data driven uncertainties**

# Iterative Monte Carlo (IMC) analysis

Toy example  $\rightarrow$  fitting 2 model parameters  $\alpha, \beta$



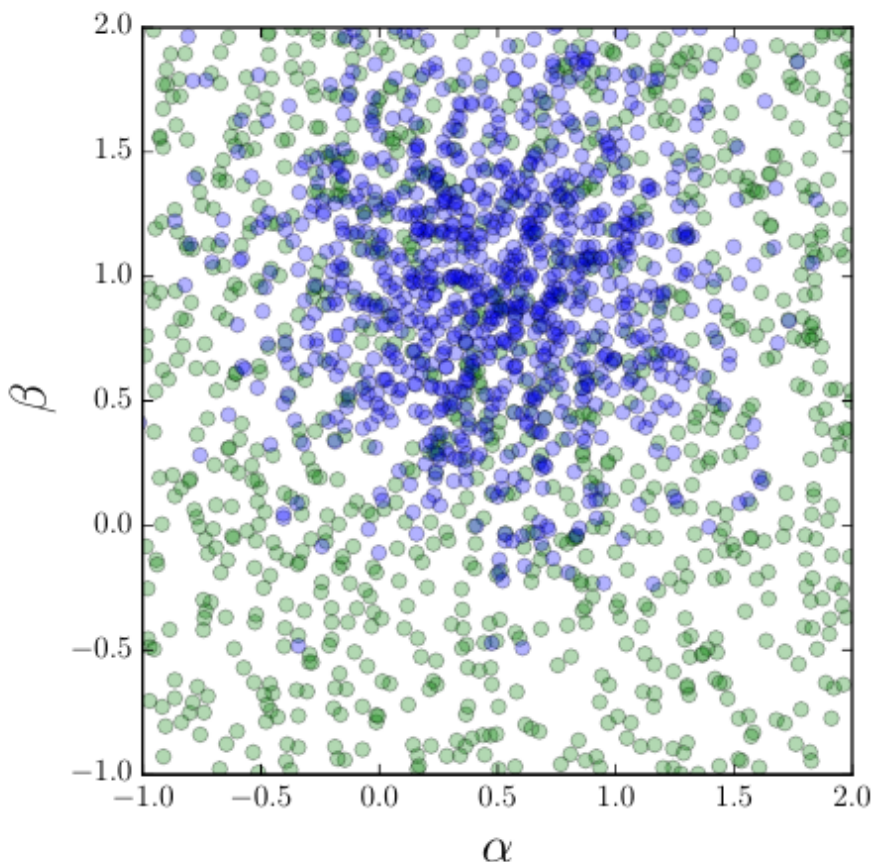
## I. Flat sampling

Initial priors  $\{(\alpha, \beta)\}$



# Iterative Monte Carlo (IMC) analysis

Toy example  $\rightarrow$  fitting 2 model parameters  $\alpha, \beta$



## I. Flat sampling

Initial priors  $\{(\alpha, \beta)\}$

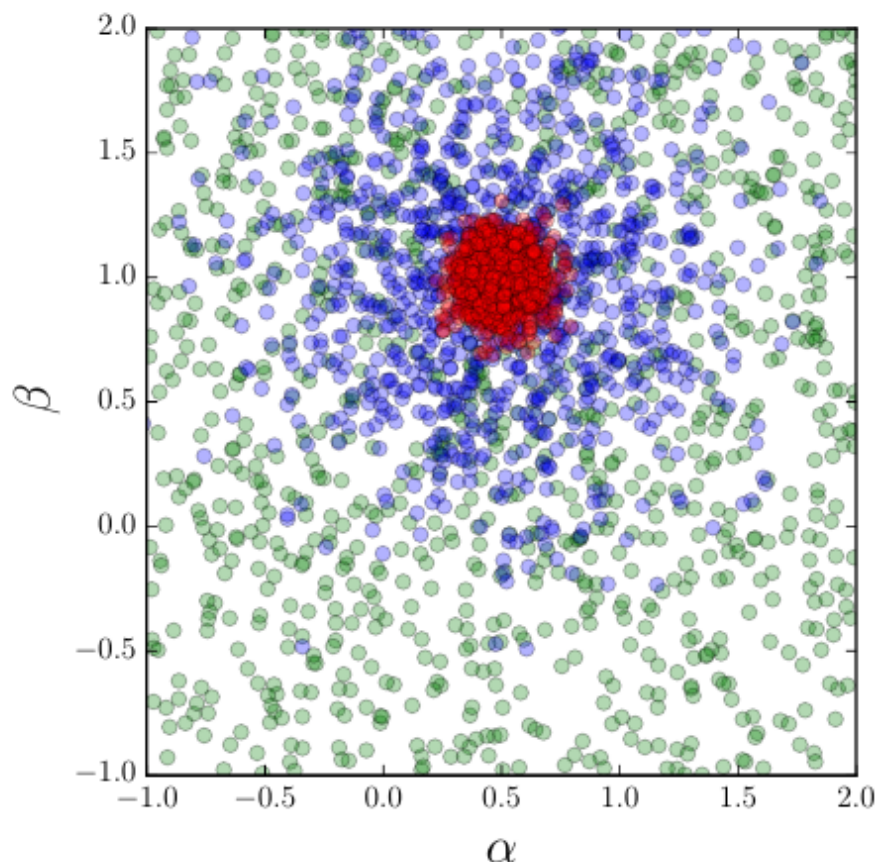
## II. First iteration

priors  $\{(\alpha, \beta)\}$

posteriors  $\{(\alpha, \beta)\}$

# Iterative Monte Carlo (IMC) analysis

Toy example  $\rightarrow$  fitting 2 model parameters  $\alpha, \beta$



## I. Flat sampling

Initial priors  $\{(\alpha, \beta)\}$

## II. First iteration

priors  $\{(\alpha, \beta)\}$

posteriors  $\{(\alpha, \beta)\}$

## III. Second iteration

priors  $\{(\alpha, \beta)\}$

posteriors  $\{(\alpha, \beta)\}$

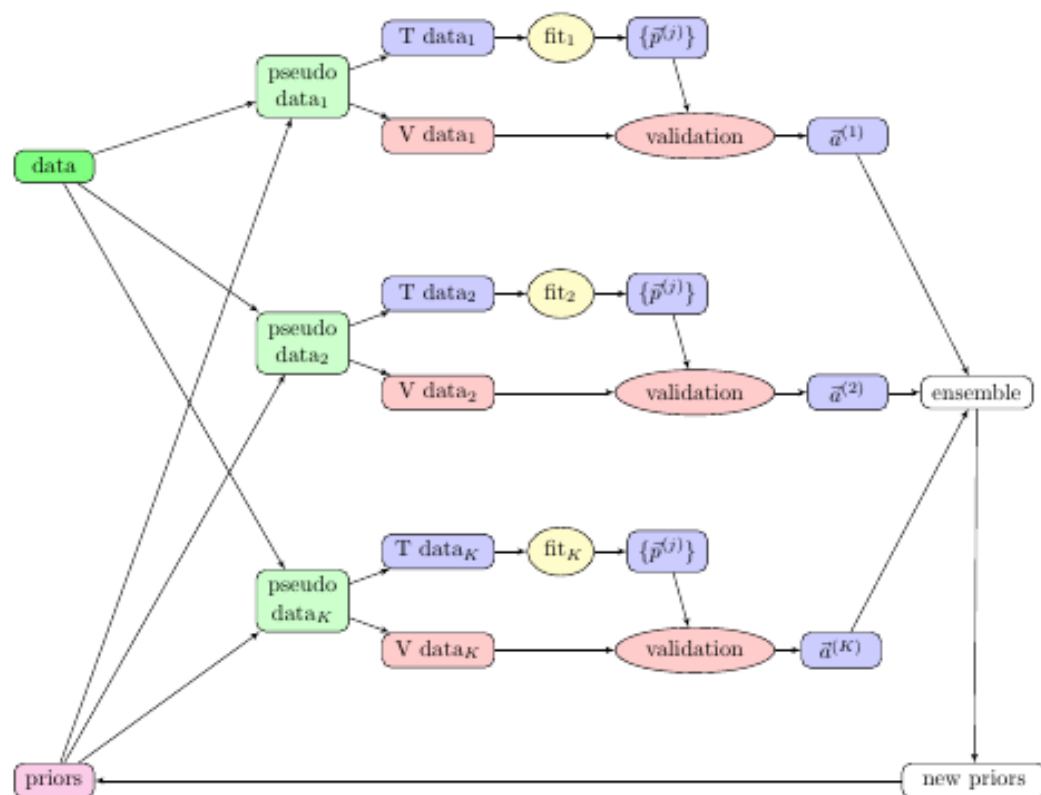
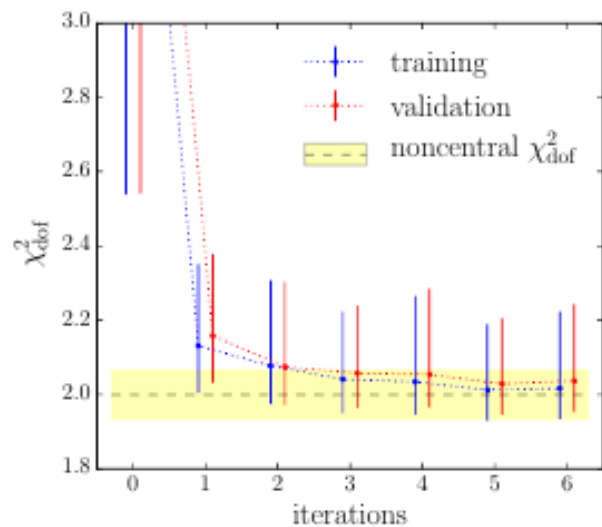
... until convergence

# Iterative Monte Carlo (IMC) analysis

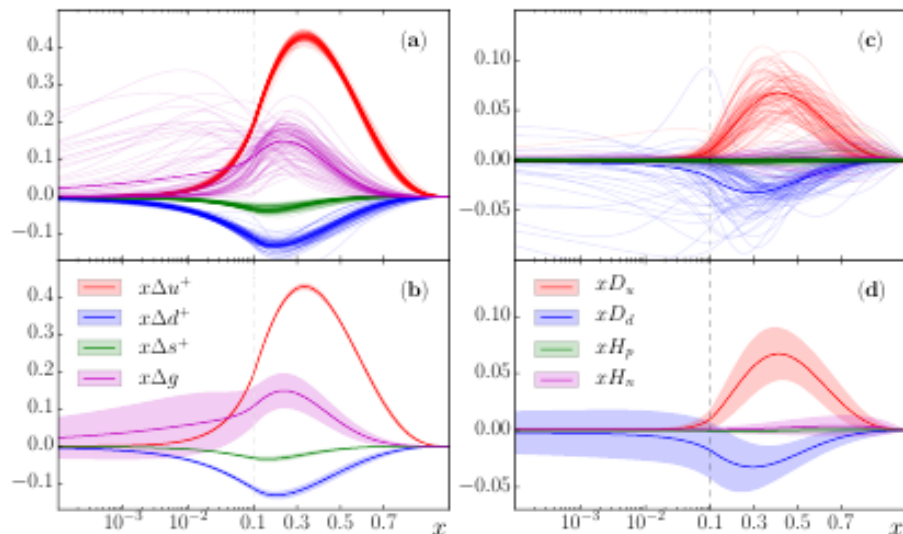
## Each iteration

- Generate pseudo data sets via data resampling
- Random data partition  $\rightarrow$  Training & Validation
- Fit the training set
- Validation

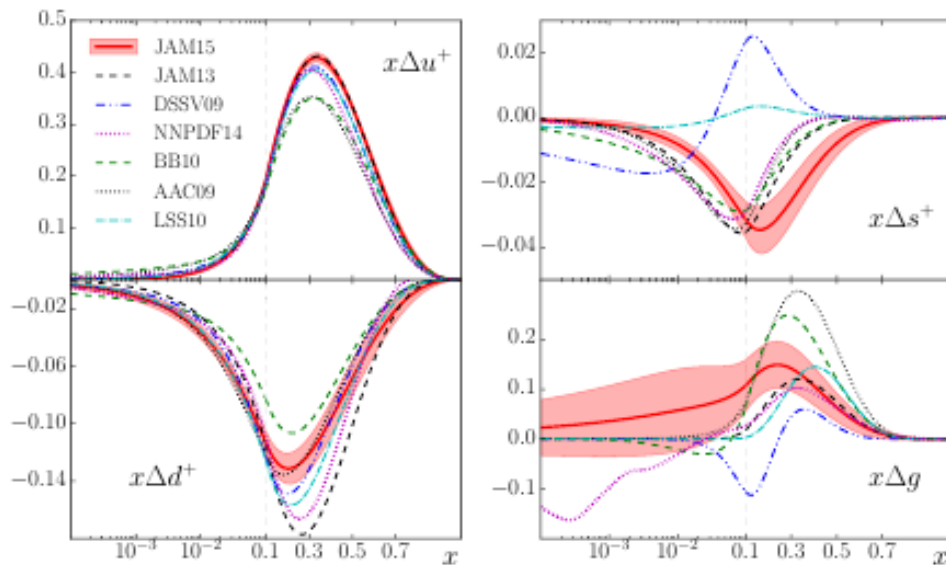
# of fits: 10000



# The JAM15 polarized PDFs

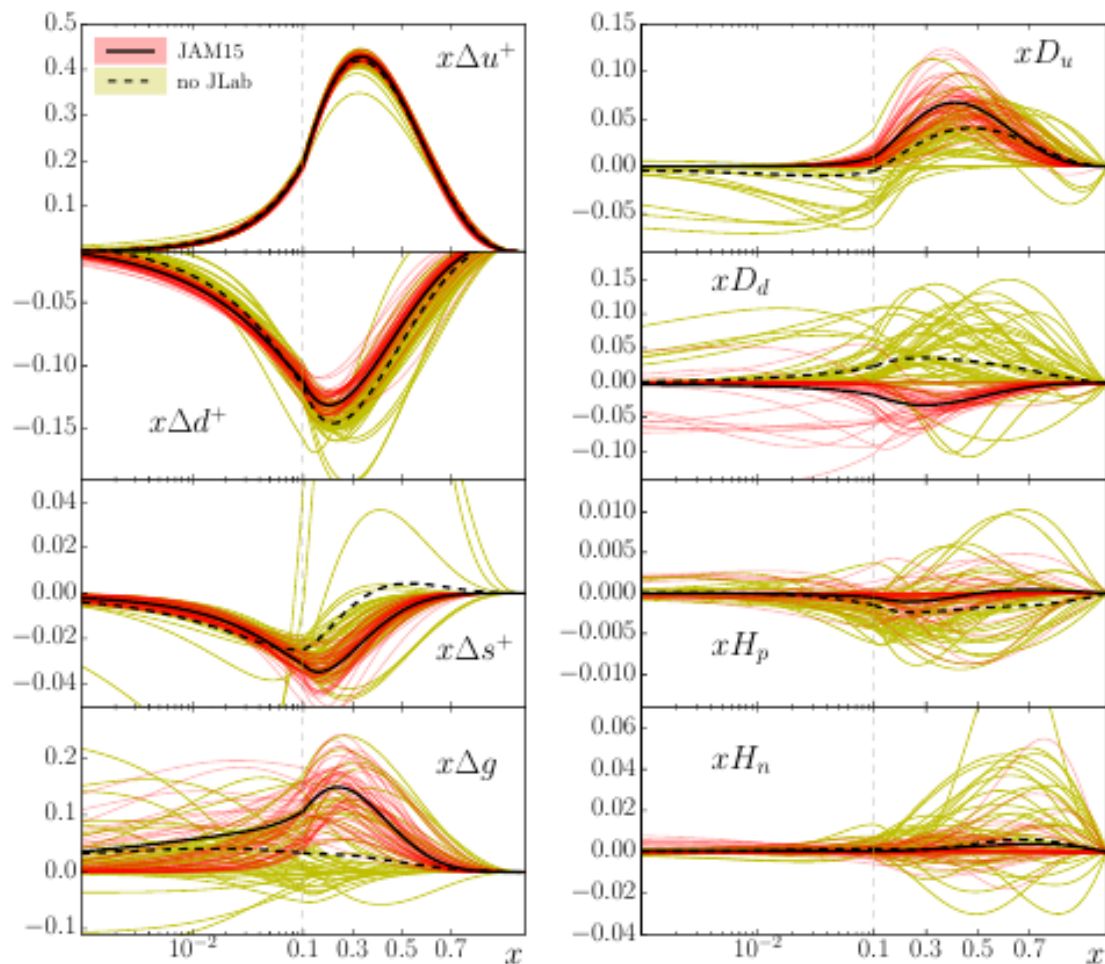


moment	truncated	full
$\Delta u^+$	$0.82 \pm 0.01$	$0.83 \pm 0.01$
$\Delta d^+$	$-0.42 \pm 0.01$	$-0.44 \pm 0.01$
$\Delta s^+$	$-0.10 \pm 0.01$	$-0.10 \pm 0.01$
$\Delta \Sigma$	$0.31 \pm 0.03$	$0.28 \pm 0.04$
$\Delta G$	$0.5 \pm 0.4$	$1 \pm 15$
$d_2^p$	$0.005 \pm 0.002$	$0.005 \pm 0.002$
$d_2^n$	$-0.001 \pm 0.001$	$-0.001 \pm 0.001$
$h_p$	$-0.000 \pm 0.001$	$0.000 \pm 0.001$
$h_n$	$0.001 \pm 0.002$	$0.001 \pm 0.003$



- Significant constraints on  $\Delta s^+$  and  $\Delta g$
- Non zero T3 quark distributions
- T4 contribution to  $g_1$  consistent with zero
- **Negative  $\Delta s^+$**
- JAM15  $\Delta g$  compatible with recent DSSV fits.

# Impact of JLab data



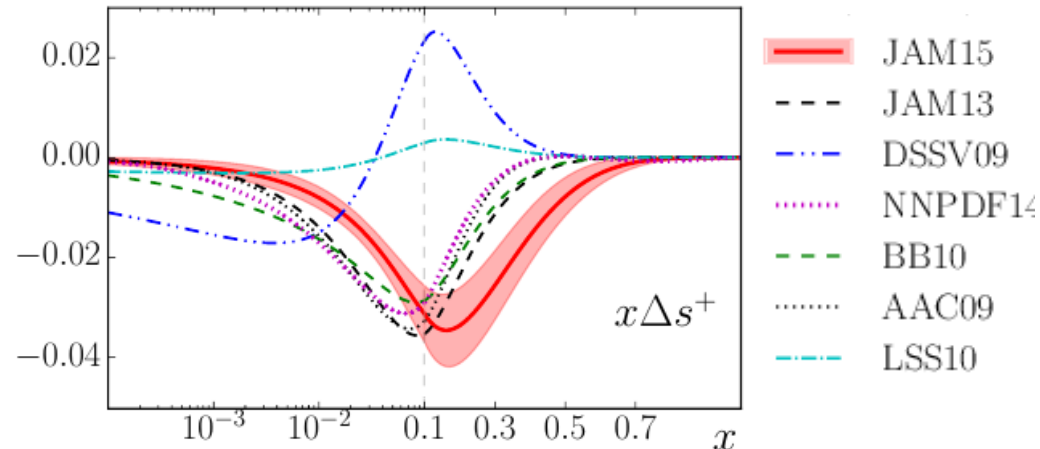
- JLab data  $\rightarrow 0.1 < x < 0.7$
- Constraints on small  $x$  from large  $x \rightarrow$  weak baryon decay constraints
- Large uncertainties in  $\Delta s^+$ ,  $\Delta g$  removed by JLab data
- Non vanishing T3 quark distributions
- T4 distributions consistent with zero



# A strange puzzle

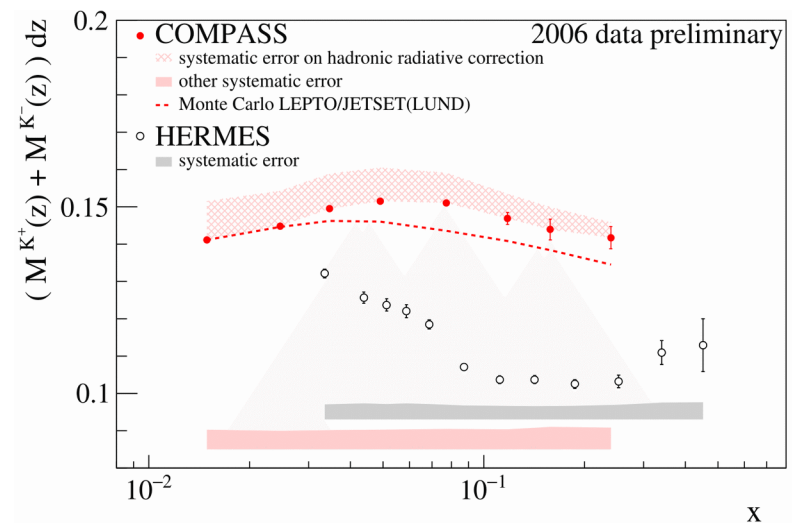
## □ $\Delta s$ : positive or negative?

- Depends on kaon FF used in SIDIS calculations!



## □ $s$ : large or small?

- Difference in size reconciled by removing hadron mass corrections  
→ see backup
- Extraction of  $s(x)$  strongly affected on kaon FF systematic uncertainty



# The JAM FF 2016 fit

## □ Kaon FF too uncertain, correlated to strange PDF in SIDIS

- Cannot take kaon FF off the shelf
- Need in-house extraction

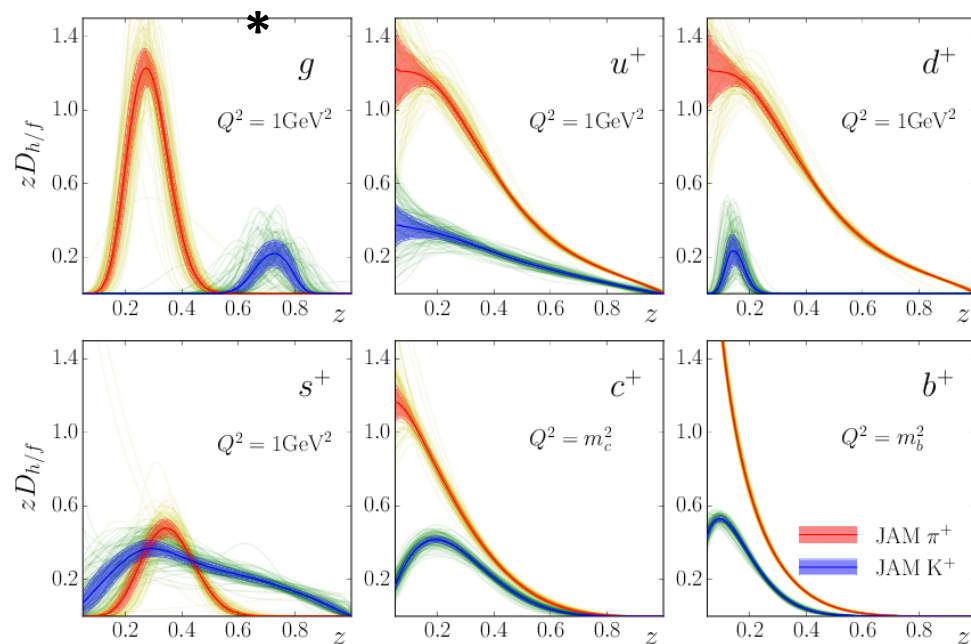
## □ Iterative MC approach

- Only SIA data used :  
npts=245,  $\chi^2 = 305.2$

## □ Strange-to-kaon FF:

- Between HKNS and DSS
- Expect combined DIS/SIDIS analysis to give negative  $\Delta s$   
(But wait and see!)

## (nearly final) JAM-FF16



\* not the official name, yet

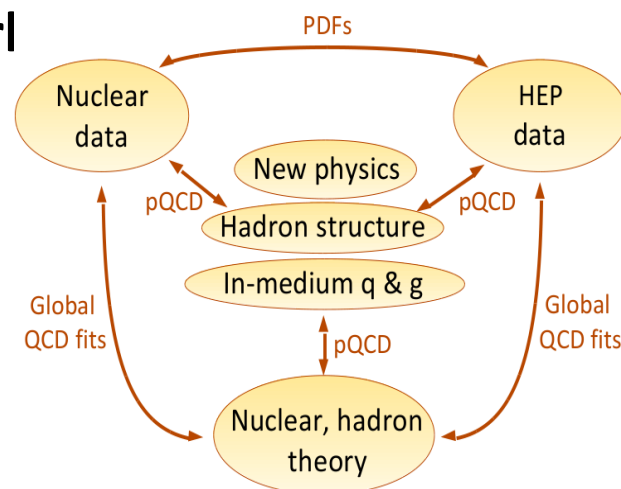
# Conclusions

## □ Entering a new precision era in large-x PDFs

- Most groups are finally on board
  - Much to be learned from each other
- New data (now and in the future), new fitting approaches
- Conquering nuclear corrections – time for threshold resummation ?
- Conquering the world: towards unified PDF+pPDF+FF fits
  - To the benefit of humankind (e.g.,  $\Delta s$ )

## □ High-energy and nuclear physics need to work

- Progress in hadron / nuclear structure
- Precision PDFs for BSM searches
- Make the most of LHC and JLab 12
- Prepare for the EIC





# Appendix: details on CJ15 fits

# CJ15 - data set

$$W^2 > 3.5 \text{ GeV}^2 \implies x \lesssim 0.85$$

$$\text{BONUS } F_2^n / F_2^d \implies x \lesssim 0.65$$

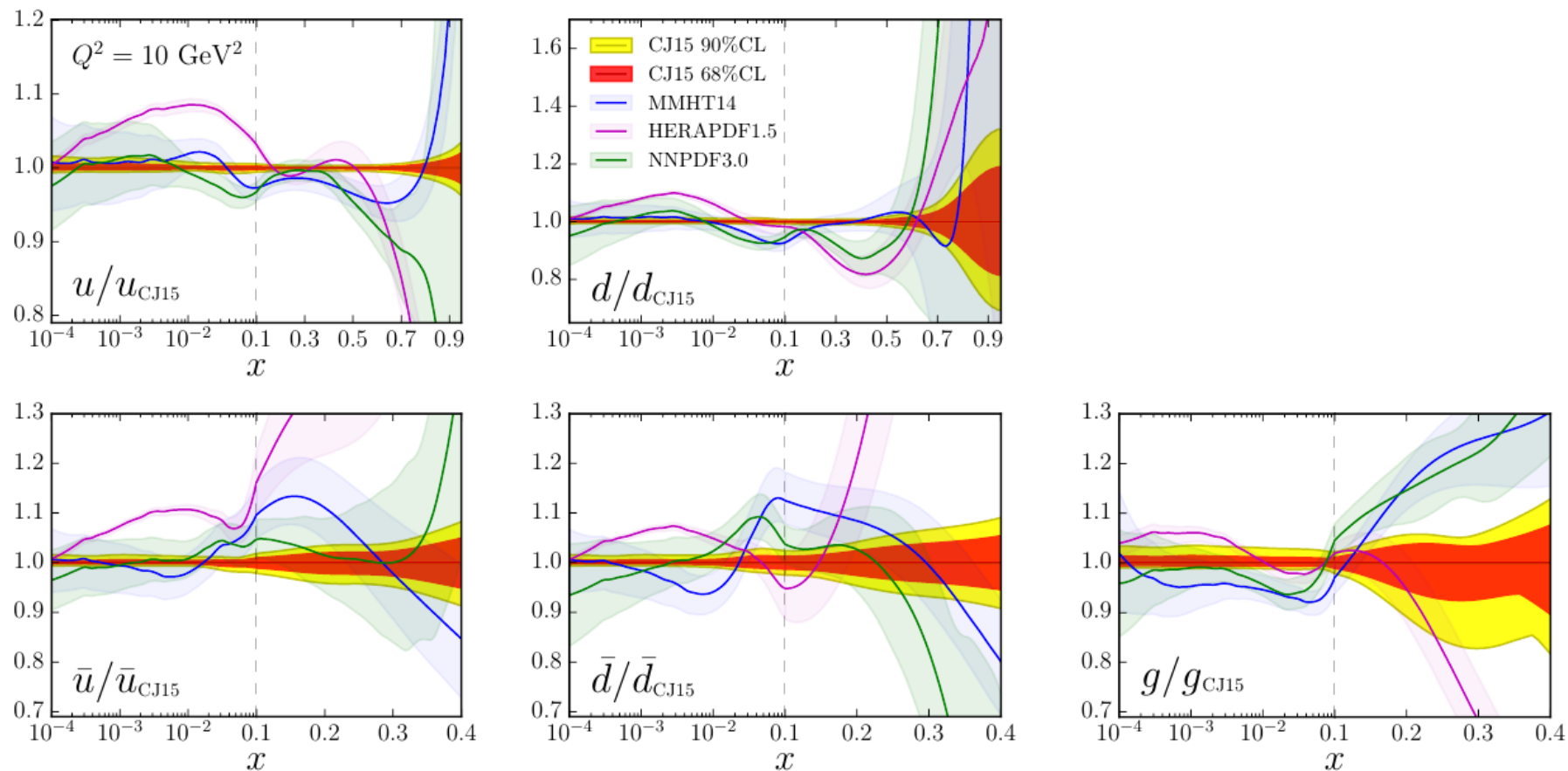
HERA I+II

$$\text{D0 } A_\ell : x \lesssim 0.5$$

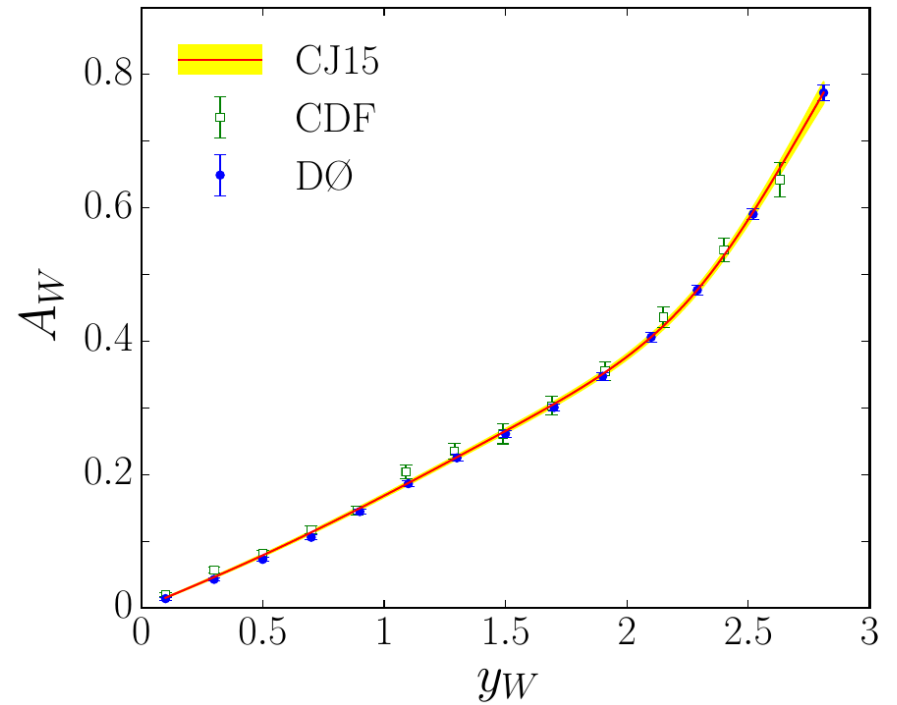
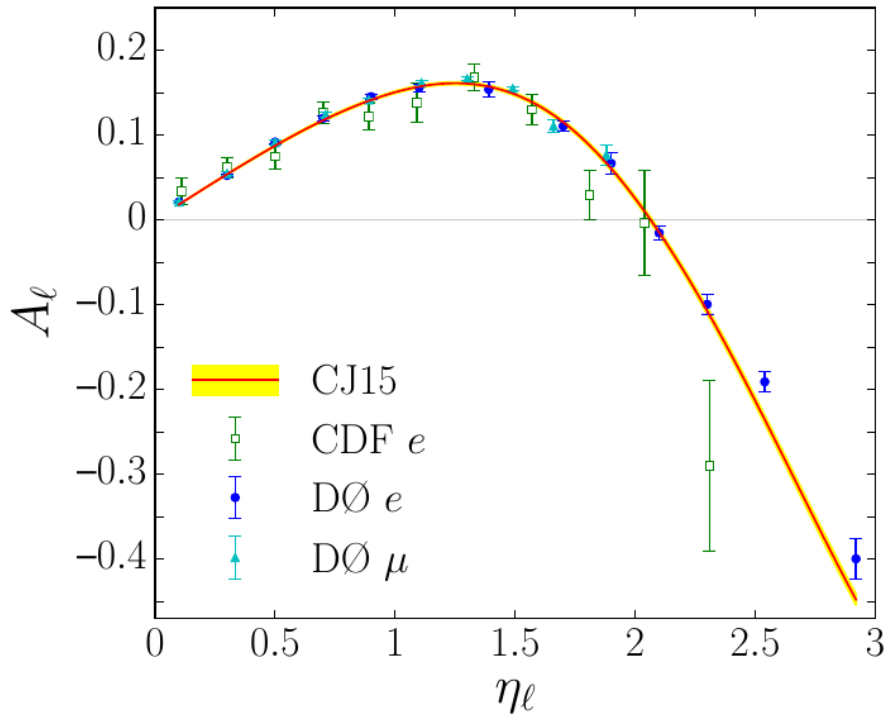
$$\text{D0 } A_W : x \lesssim 0.85$$

Observable	Experiment	# points	$\chi^2$				
			LO	NLO	NLO (OCS)	NLO (no nucl)	NLO (no nucl/D0)
DIS $F_2$	BCDMS ( $p$ ) [81]	351	430	<b>438</b>	436	440	427
	BCDMS ( $d$ ) [81]	254	297	<b>292</b>	289	301	301
	SLAC ( $p$ ) [82]	564	488	<b>434</b>	435	441	440
	SLAC ( $d$ ) [82]	582	396	<b>376</b>	380	507	466
	NMC ( $p$ ) [83]	275	431	<b>405</b>	404	405	403
	NMC ( $d/p$ ) [84]	189	179	<b>172</b>	173	174	173
	HERMES ( $p$ ) [86]	37	56	<b>42</b>	43	44	44
	HERMES ( $d$ ) [86]	37	51	<b>37</b>	38	36	37
	Jefferson Lab ( $p$ ) [87]	136	166	<b>166</b>	167	177	166
	Jefferson Lab ( $d$ ) [87]	136	131	<b>123</b>	124	126	130
DIS $F_2$ tagged	Jefferson Lab ( $n/d$ ) [21]	191	218	<b>214</b>	213	219	219
DIS $\sigma$	HERA (NC $e^-p$ ) [85]	159	325	<b>241</b>	240	247	244
	HERA (NC $e^+p$ 1) [85]	402	966	<b>580</b>	579	588	585
	HERA (NC $e^+p$ 2) [85]	75	184	<b>94</b>	94	94	93
	HERA (NC $e^+p$ 3) [85]	259	307	<b>249</b>	249	248	248
	HERA (NC $e^+p$ 4) [85]	209	348	<b>228</b>	228	228	228
	HERA (CC $e^-p$ ) [85]	42	44	<b>48</b>	48	45	49
	HERA (CC $e^+p$ ) [85]	39	56	<b>50</b>	50	51	51
	Drell-Yan	ES66 ( $pp$ ) [29]	121	148	<b>139</b>	139	145
	ES66 ( $pd$ ) [29]	129	207	<b>145</b>	143	158	157
W/charge asymmetry	CDF ( $e$ ) [88]	11	11	<b>12</b>	12	13	14
	DØ ( $\mu$ ) [17]	10	37	<b>20</b>	19	29	28
	DØ ( $e$ ) [18]	13	20	<b>29</b>	29	14	14
	CDF ( $W$ ) [89]	13	16	<b>16</b>	16	14	14
Z rapidity	DØ ( $W$ ) [19]	14	39	<b>14</b>	15	82	—
	CDF ( $Z$ ) [90]	28	100	<b>27</b>	27	26	26
jet	DØ ( $Z$ ) [91]	28	25	<b>16</b>	16	16	16
	CDF (run 2) [92]	72	33	<b>15</b>	15	23	25
$\gamma$ +jet	DØ (run 2) [93]	110	23	<b>21</b>	21	14	14
	DØ 1 [94]	16	17	<b>7</b>	7	7	7
	DØ 2 [94]	16	34	<b>16</b>	16	17	17
	DØ 3 [94]	12	34	<b>25</b>	25	24	25
	DØ 4 [94]	12	76	<b>13</b>	13	13	13
total		4542	5894	<b>4700</b>	4702	4964	4817
total + norm			6022	<b>4708</b>	4710	4972	4826
$\chi^2/\text{datum}$			1.33	<b>1.04</b>	1.04	1.09	1.07

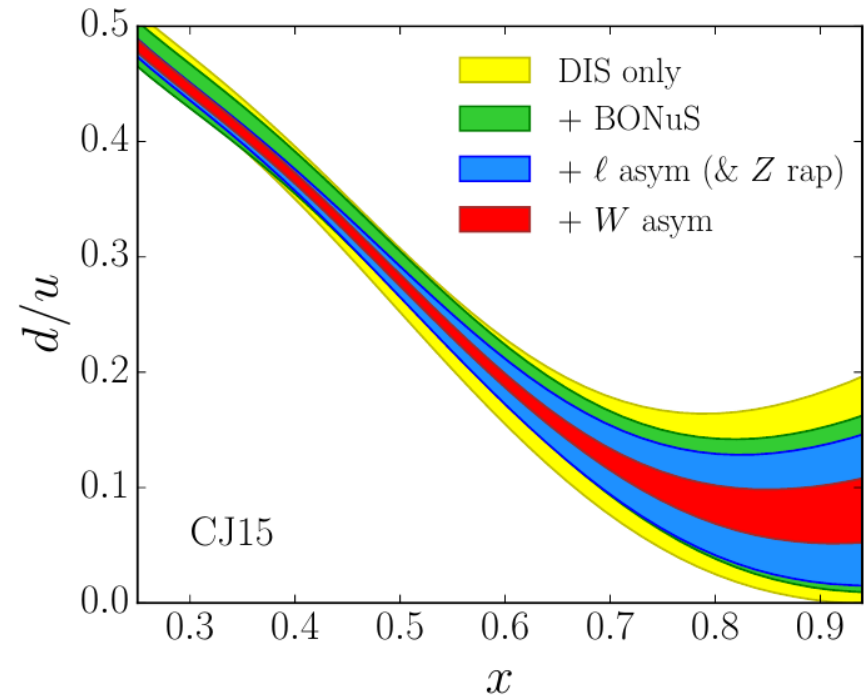
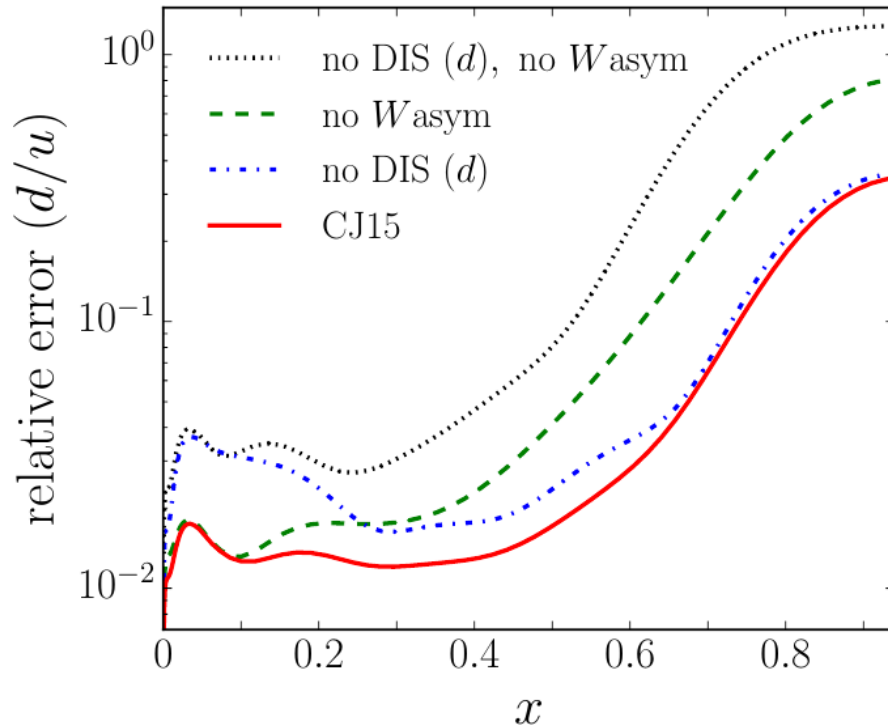
# CJ15 vs. others



# W-lepton and W asymmetry at Tevatron



# What fits what?



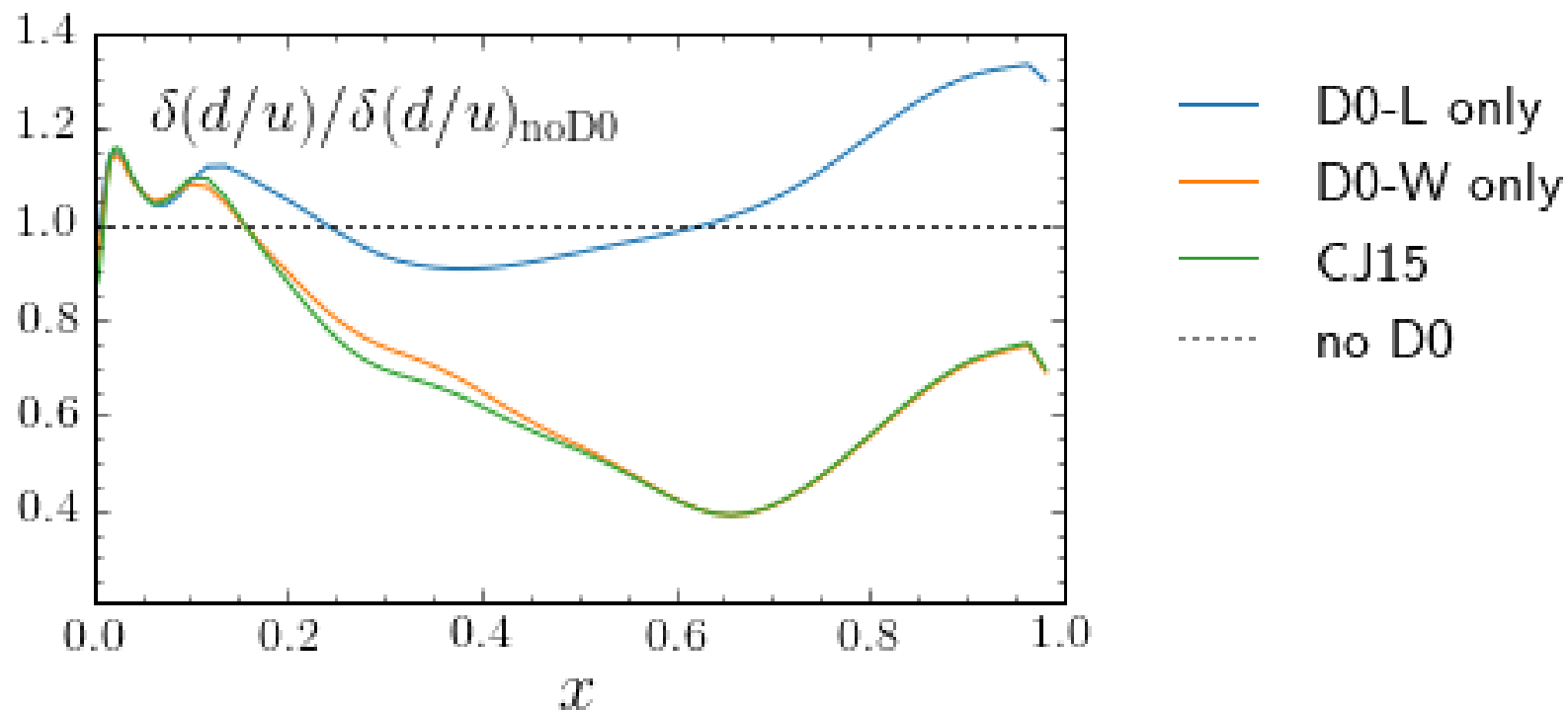
## Large $x > 0.3$ :

- D0's  $W$ -asymmetry determines the  $d$ -quark
- SLAC( $d$ )'s statistical power used to fit the off-shell function

## Moderate $x < 0.3$ :

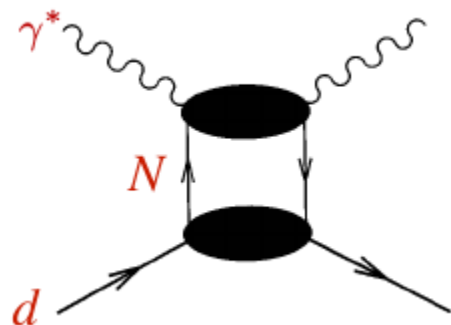
- SLAC( $d$ ) enables precise  $d/u$  flavor separation

# W-lepton and W asymmetry at Tevatron



# Nuclear corrections

- At large  $x$ , DIS dominated by incoherent scattering from individual nucleons



$$q^d(x, Q^2) = \int \frac{dz}{z} dp^2 f_{N/d}(z, p^2) \tilde{q}^N(x/z, p^2, Q^2)$$

nucleon momentum distribution in  $d$  ("smearing function")

PDF in bound (off-shell) nucleon

$$\rightarrow z = \frac{p \cdot q}{p_d \cdot q} \approx 1 + \frac{p_0 + \gamma p_z}{M} \left[ p_0 = M + \varepsilon, \quad \varepsilon = \varepsilon_d - \frac{\vec{p}^2}{2M} \right]$$

momentum fraction of  $d$  carried by  $N$

$$\rightarrow \text{at finite } Q^2, \text{ smearing function depends on } \gamma = \sqrt{1 + 4M^2 x^2 / Q^2}$$

- Offshell expansion; parametrize first order coefficient,  $x_1$  fixed with valence sum rule

$$\tilde{q}^N(x, p^2) = q^N(x) \left[ 1 + \frac{(p^2 - M^2)}{M^2} \delta q^N(x) \right]$$

$$\delta q^N = C_N(x - x_0)(x - x_1)(1 + x - x_0) \quad \int_0^1 dx \delta q^N(x) (q^N(x) - \bar{q}^N(x)) = 0$$

# Nucleon off-shellness constrained by D0 data (!)

- The “wrong” nuclear corrections creates tension between DIS(D) and W asym
  - The fits then choses the “right” one

- **Deuteron to nucleon “EMC” ratio**  $D/(p+n)$

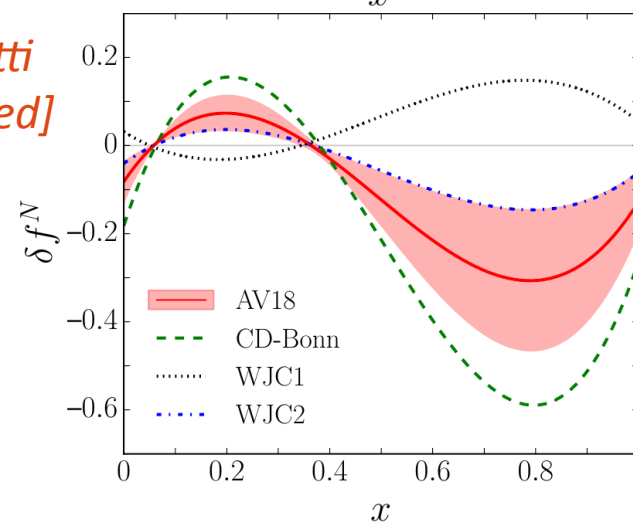
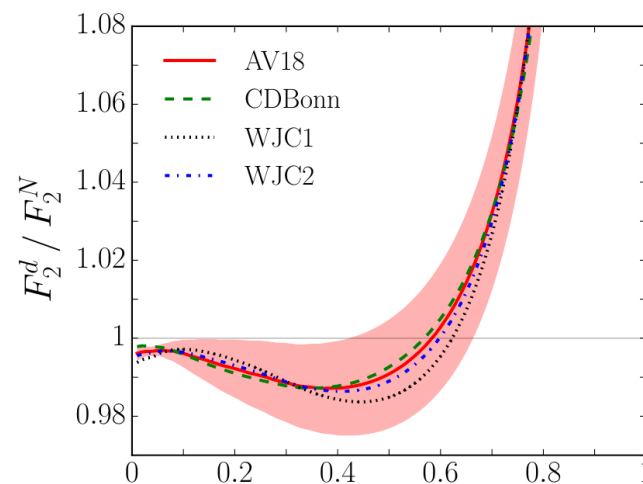
- Stable w.r.t. choice of nucleon w.fn.
- (WJC1 disfavored  $\chi^2$ -wise)
- No evidence for antishadowing

- **Off shell correction – first time in Deuteron!**

- Good statistical precision!
- Magnitude compensates for wave function's missing / excessive strength
- Physical result or fitting away other physics?

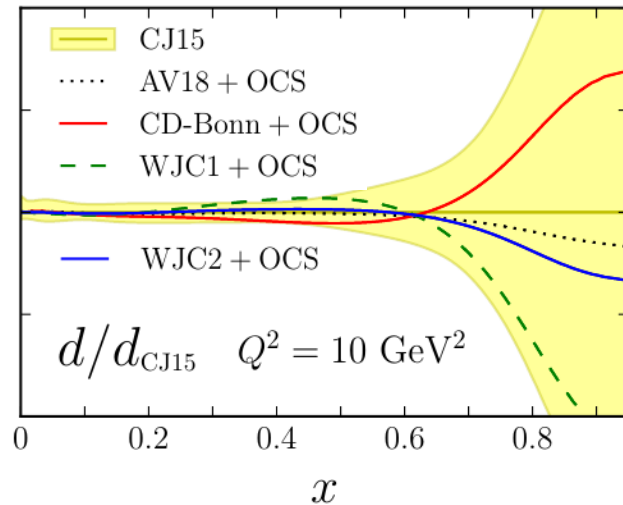
→ also: R. Petti  
[WG1 – Wed]

$$\delta f^N = C(x - x_0)(x - x_1)(1 + x_0 - x)$$





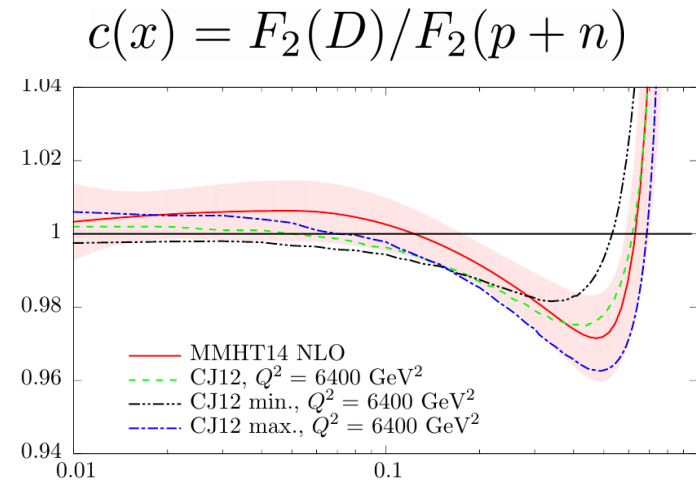
# Cross checks



- Fit with with a 1-parameter model of the off-shell effects
  - Obtain compatible  $d$  quarks

OCS = Off-shell Covariant Spectator model

- MMHT14 parametrize the whole nuclear effect
  - Obtain similar result
  - (but cannot explore the nuclear dynamics)



MMHT14, EPJ C75 (2015) 204

## Hadronic physics output 2: dbar/ubar

- Peng et al. suggest NMC F2 data indicate negative dbar/ubar

J.C. Peng *et al.*, PLB 736 (2014), 411-414

$$\bar{d}(x) - \bar{u}(x) = \frac{1}{2}[u_v(x) - d_v(x)] - \frac{3}{2x}[F_2^p(x) - F_2^n(x)]$$

- But extract  $F_2^n$  without accounting for nuclear corrections
- Note also:
  - Ambiguities in choice of valence or total up and down above
  - NLO effects not negligible when dbar/ubar  $\sim 0$

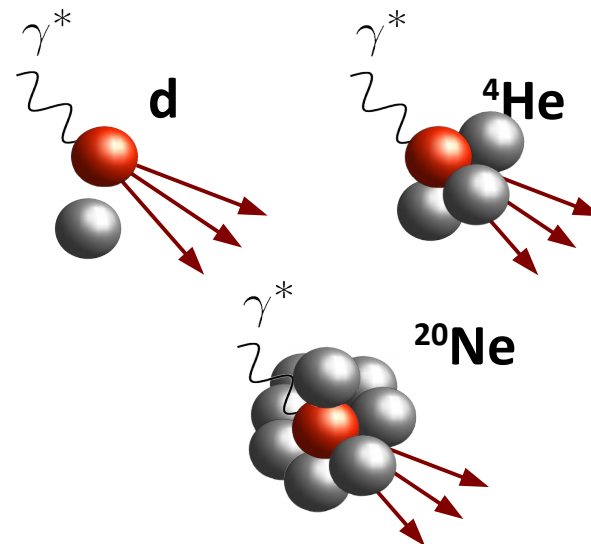
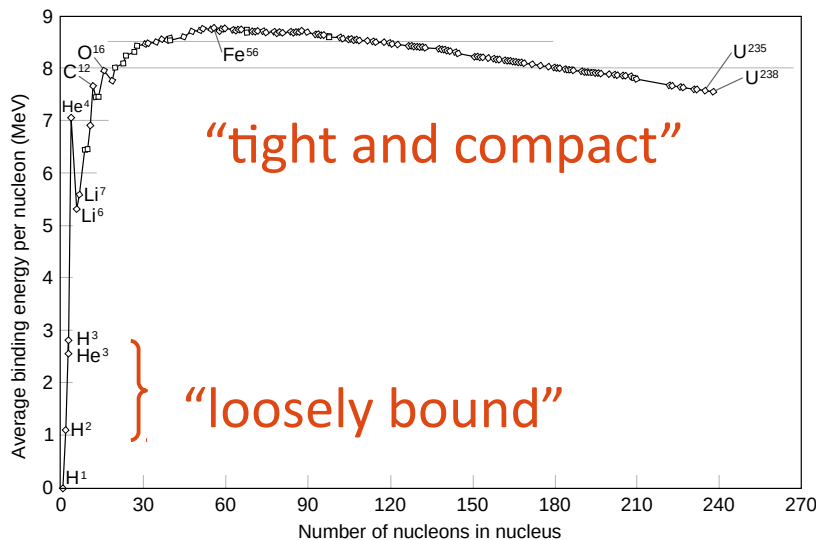
- Using CJ15 fit, can “remove” nuclear effects → *Niculescu & Keppel, DIS2016*

$$F_2^n \equiv F_2^d(\text{measured}) \times \frac{F_2^n}{F_2^d}(\text{CJ15})$$

# Nuclear physics output

❑ **QUESTION:** Does the nuclear environment affect the off-shell behavior of a nucleon?

- For example, partial deconfinement [Close, Jaffe, Roberts (1985)]

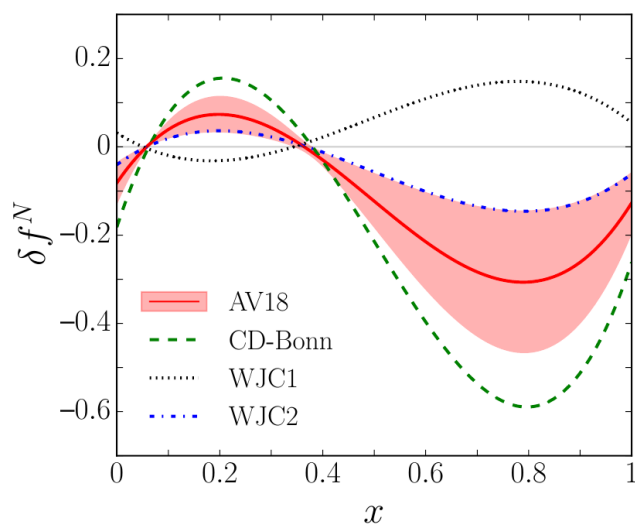


$$\delta q = \delta q(x; A) ??$$

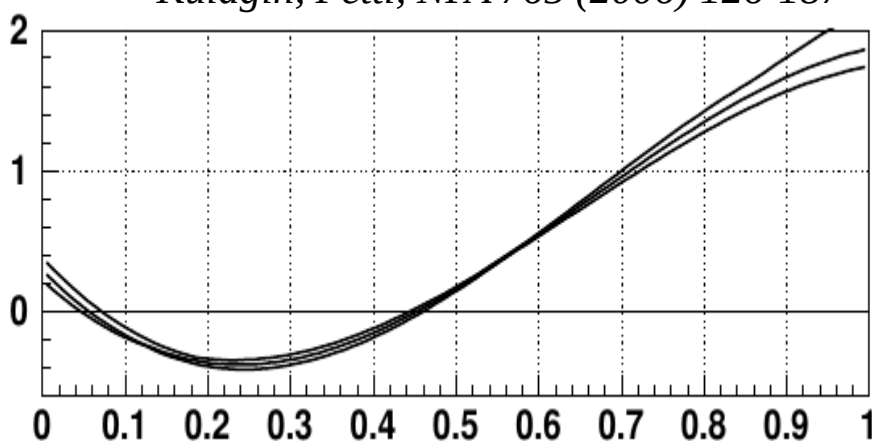
# Nuclear physics output

## □ Compare to Kulagin-Petti fit to e+A collisions

- Same functional form (but different normalization)



*Kulagin, Petti, NPA 765 (2006) 126-187*

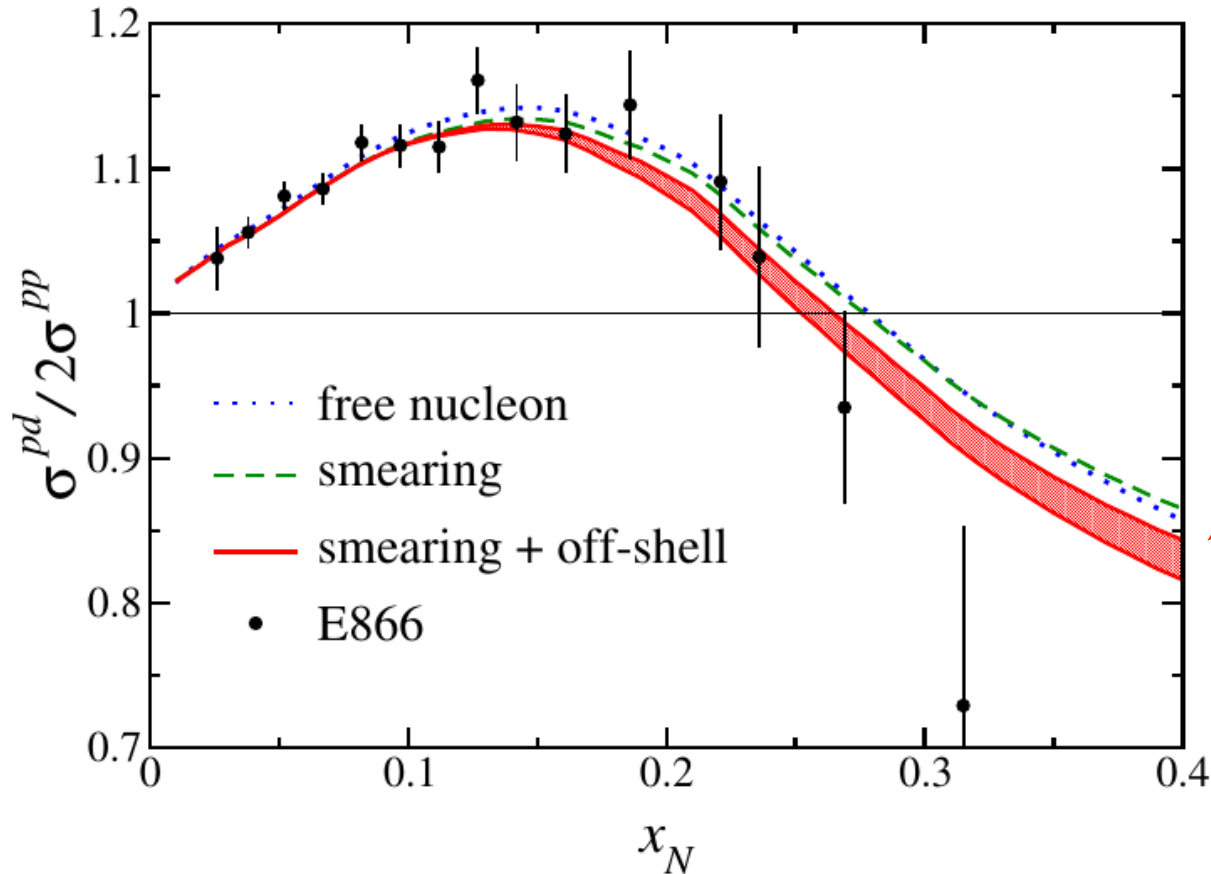


## □ Different shape and size

- no nuclear universality ??  $\delta f^N$
- too hard nuclear spectral function at large momentum ??
- wrong parametrization? → **To be investigated**

# Nuclear corrections...

*Ehlers, AA, Brady, Melnitchouk, PRD90 (2014)*

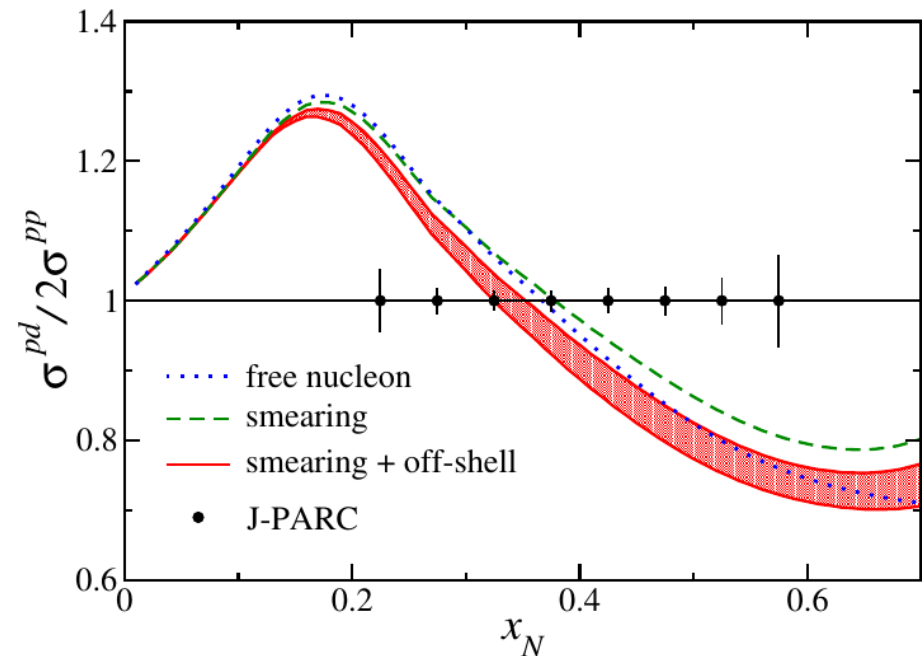
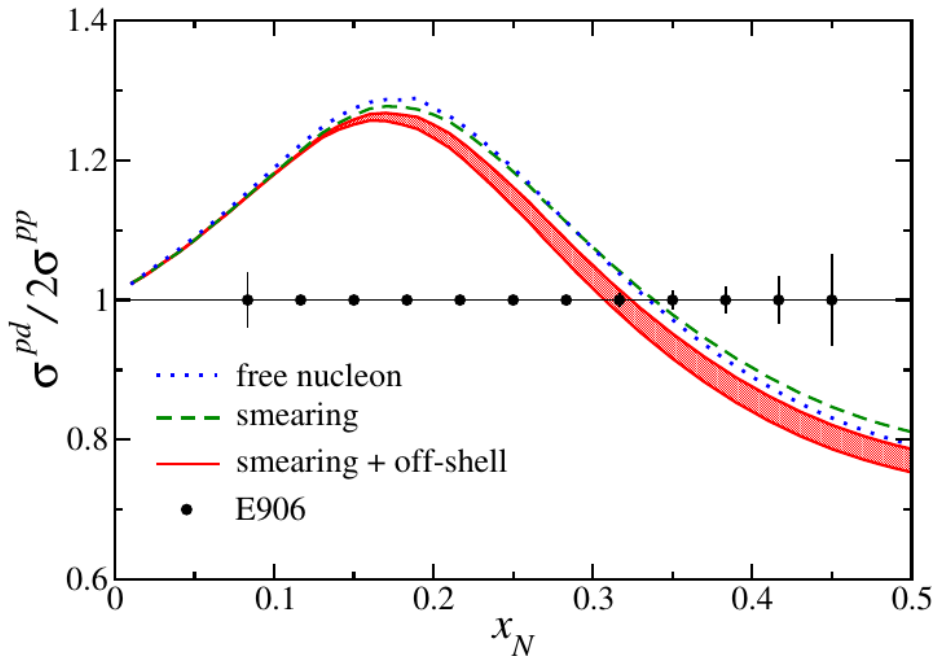


**Red band:**  
combined wave fn.  
& off-shell model  
uncertainty

Off-shell corrections help make  $d\bar{u}u$  stay positive

# Future DY reaches into large- $x$

*Ehlers, AA, Brady, Melnitchouk, PRD90 (2014)*

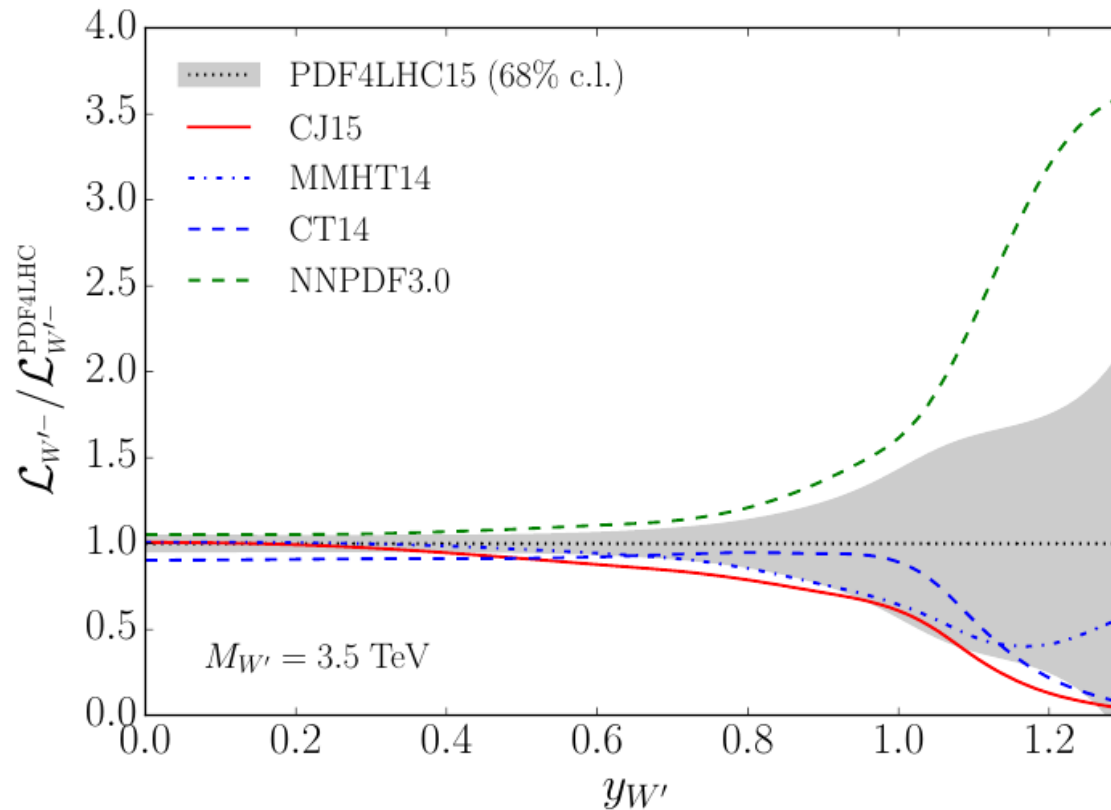


□ **E906/Sea Quest:** off-shell effects even more important

□ **J-PARC:** can cross-check nuclear smearing vs. DIS

# BSM physics output

→ see also: *R.Placakyte [WG1 – Tue]*



$$\mathcal{L}_{W'^-} = \frac{2\pi G_F}{3\sqrt{2}} x_1 x_2 \left[ \cos^2 \theta_C (\bar{u}(x_2) d(x_1) + \bar{c}(x_2) s(x_1)) + \sin^2 \theta_C (\bar{u}(x_2) s(x_1) + \bar{c}(x_2) d(x_1)) \right] + (x_1 \leftrightarrow x_2)$$

# Strangeness and strangeness asymmetry

$$s^\pm(x) = s(x) \pm \bar{s}(x) \quad [s^\pm] = \int_0^1 dx x s^\pm(x)$$

□ In pre-LHC fits, mostly constrained by  $\nu$ +A data

- CCFR inclusive DIS
- NuTeV muon pair production
- NOMAD and CHORUS

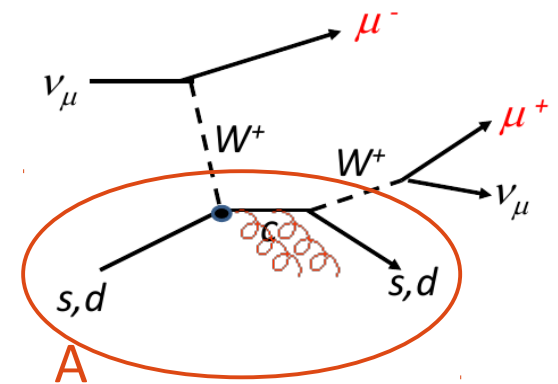
□ Nuclear corrections again...

– Initial state nuclear wave-function mods

- Partly under control using nPDFs
- But: double counting!! → either use in nPDF or in PDF fits !

– Final state propagation of the charm quark / D meson

- Not under theoretical / phenomenological control  
(*cf.* heavy quark “puzzle” in A+A at RHIC, LHC)

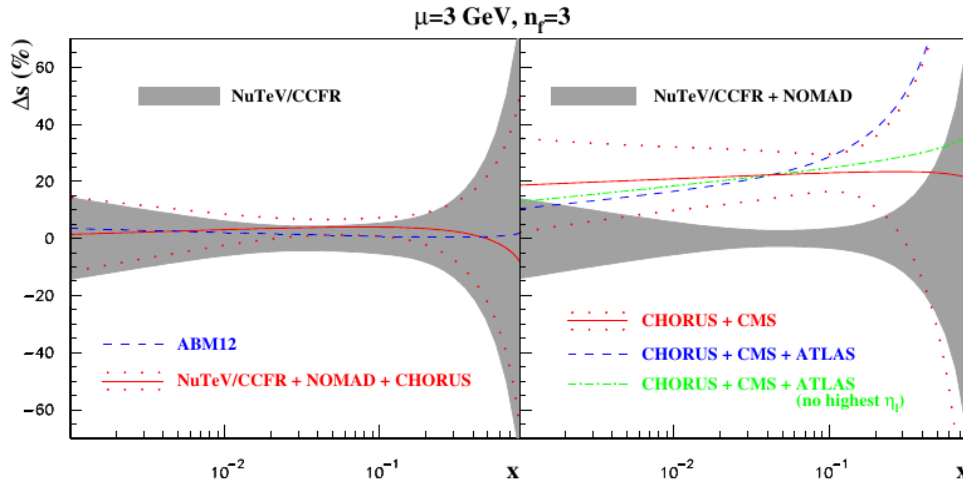




# Strange tensions

□  $\nu+A \rightarrow \text{dimuons}$  vs.  $p+p \rightarrow W+c$  at LHC

Alekhin et al., arXiv:1404.6469



$$g s_p \rightarrow W c$$

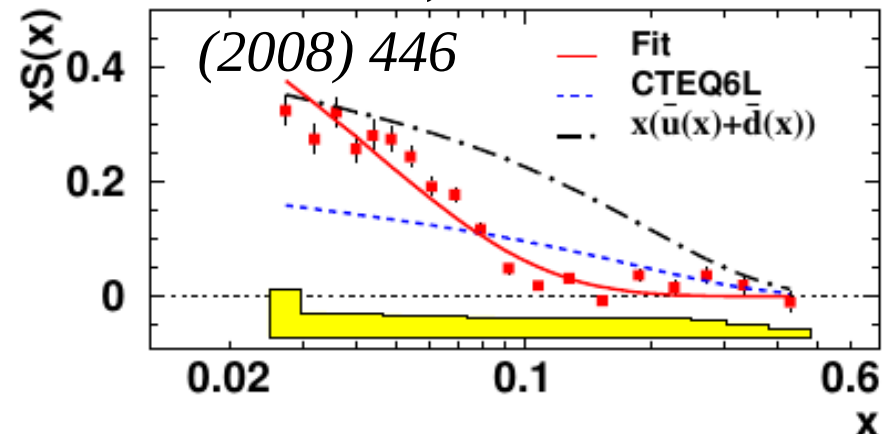
$$\nu s_A \rightarrow \mu^- \mu^+ \nu_\mu s$$

FSI ?

□ Kaons in  $e+p$  at HERMES

- But.. fragmentation functions uncertainty

HERMES, PLB 666

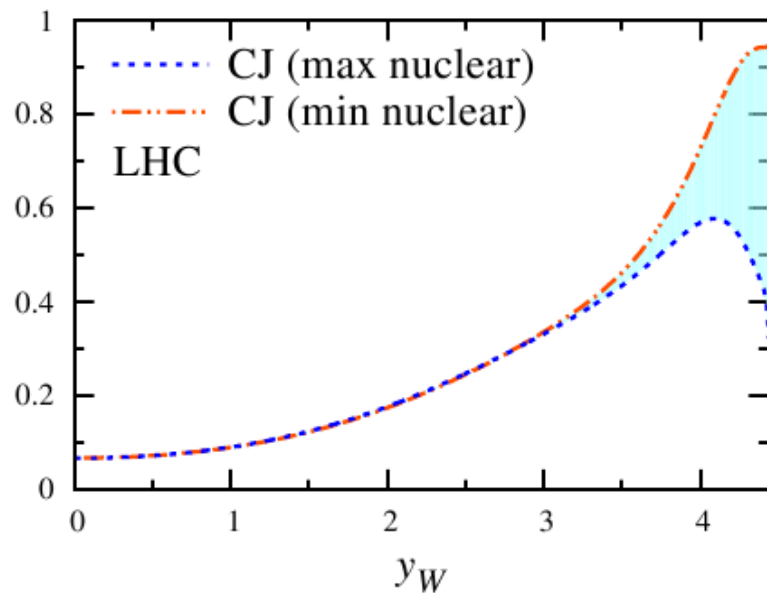


# W charge asymmetry at LHC

Brady, Accardi, Melnitchouk, Owens, *JHEP* 1206 (2012) 019

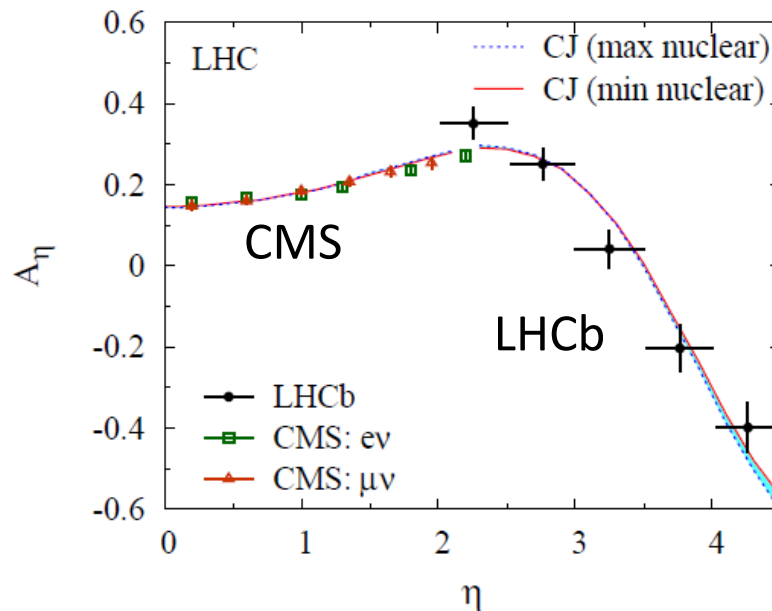
## Directly reconstructed W:

- highest sensitivity to large  $x$



## From decay lepton $W \rightarrow l + \nu$ :

- smearing in  $x$

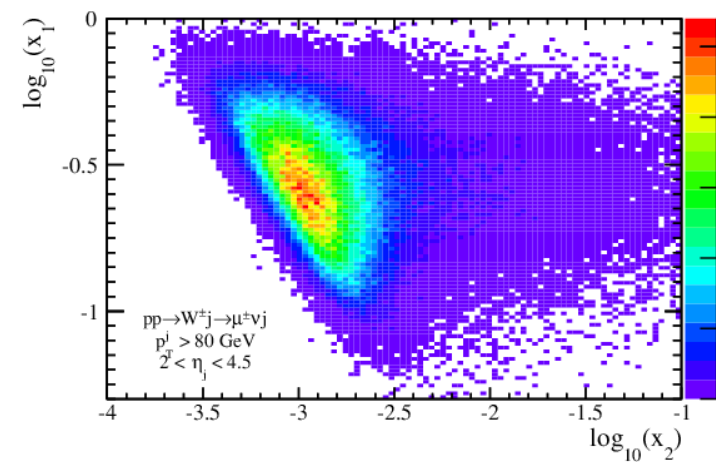
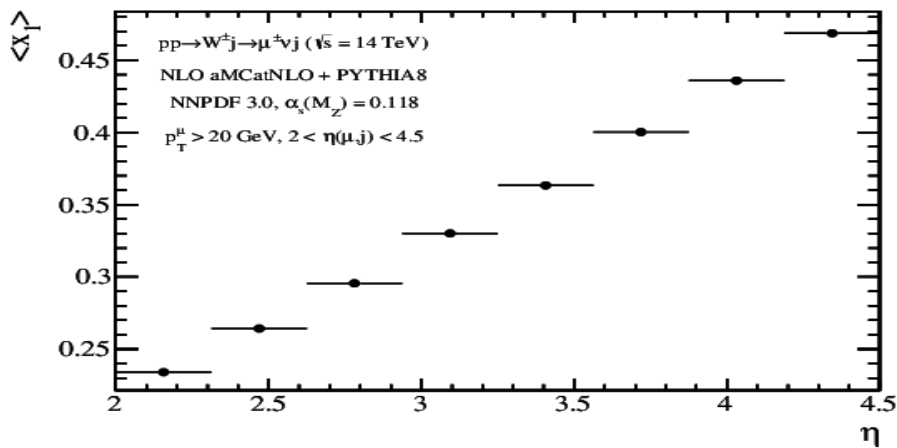
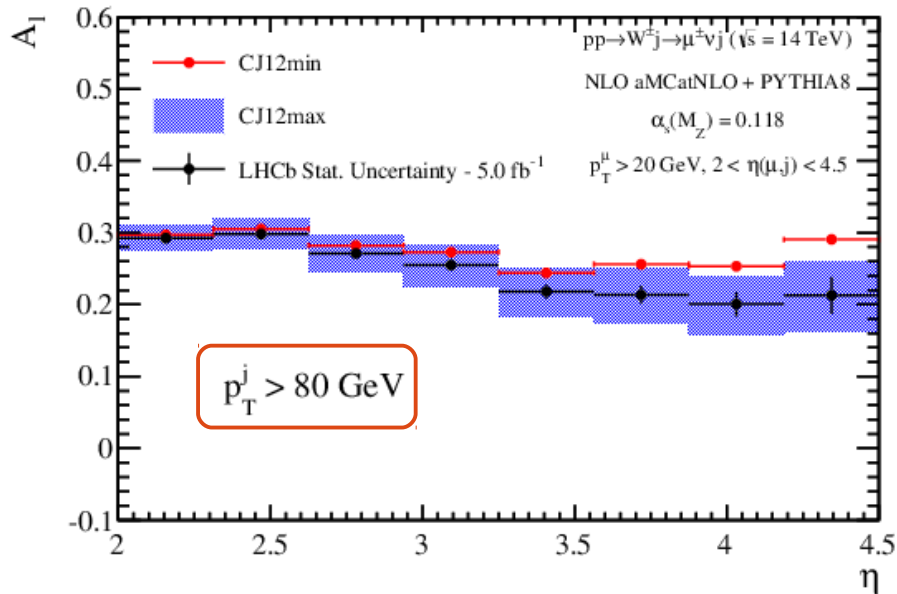


## ❑ Would be nice to reconstruct W at

- LHCb – But I am told “too many holes”...
- RHIC – how high in rapidity?
- AFTER@LHC ??

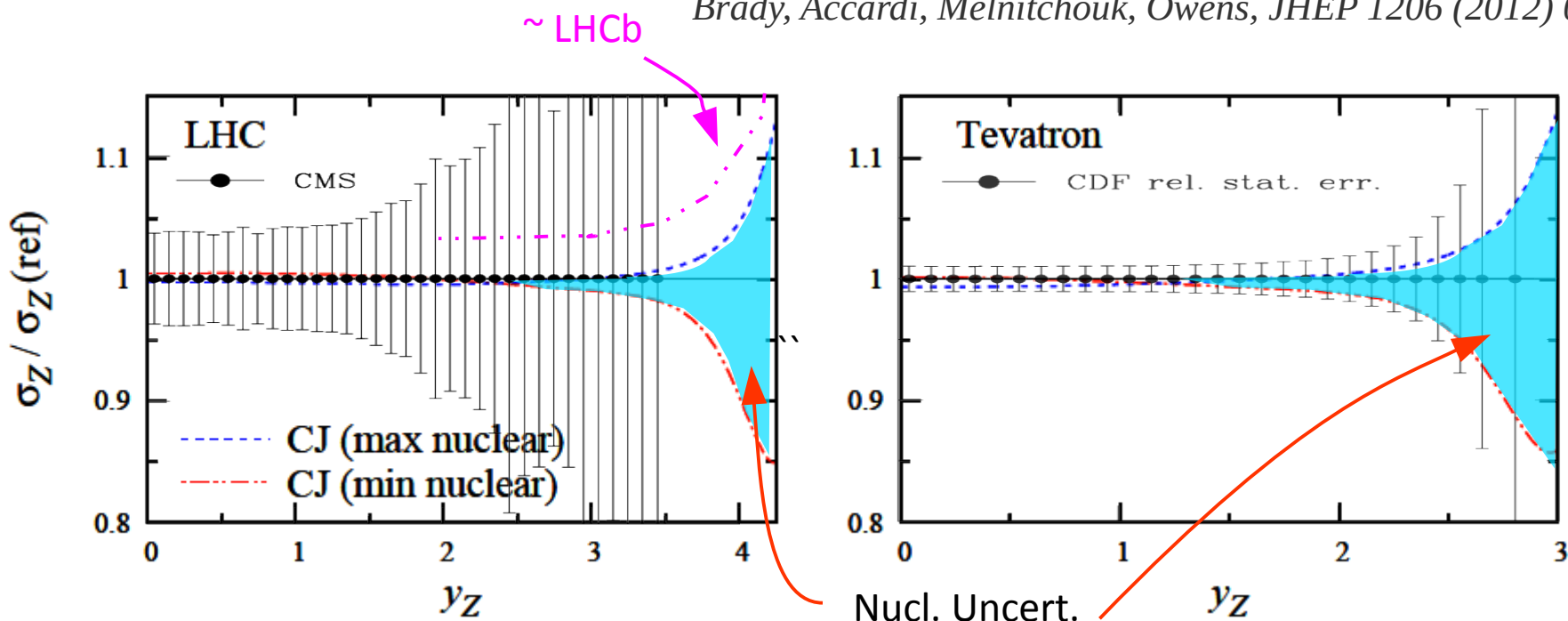
# W+c at LHCb

Farry and Gauld, PRD 93 (2016) 014008



# Z rapidity distribution

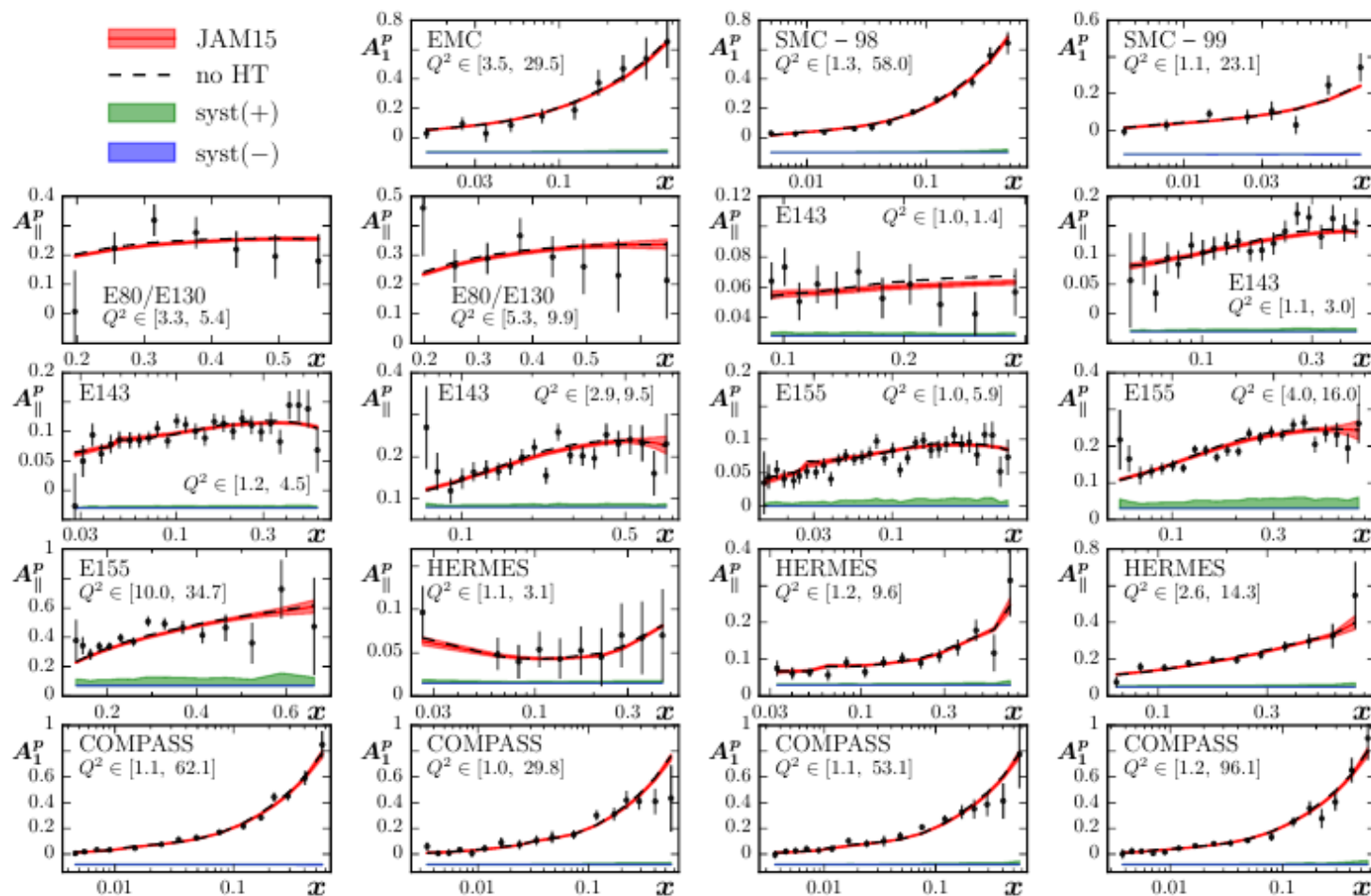
Brady, Accardi, Melnitchouk, Owens, *JHEP* 1206 (2012) 019



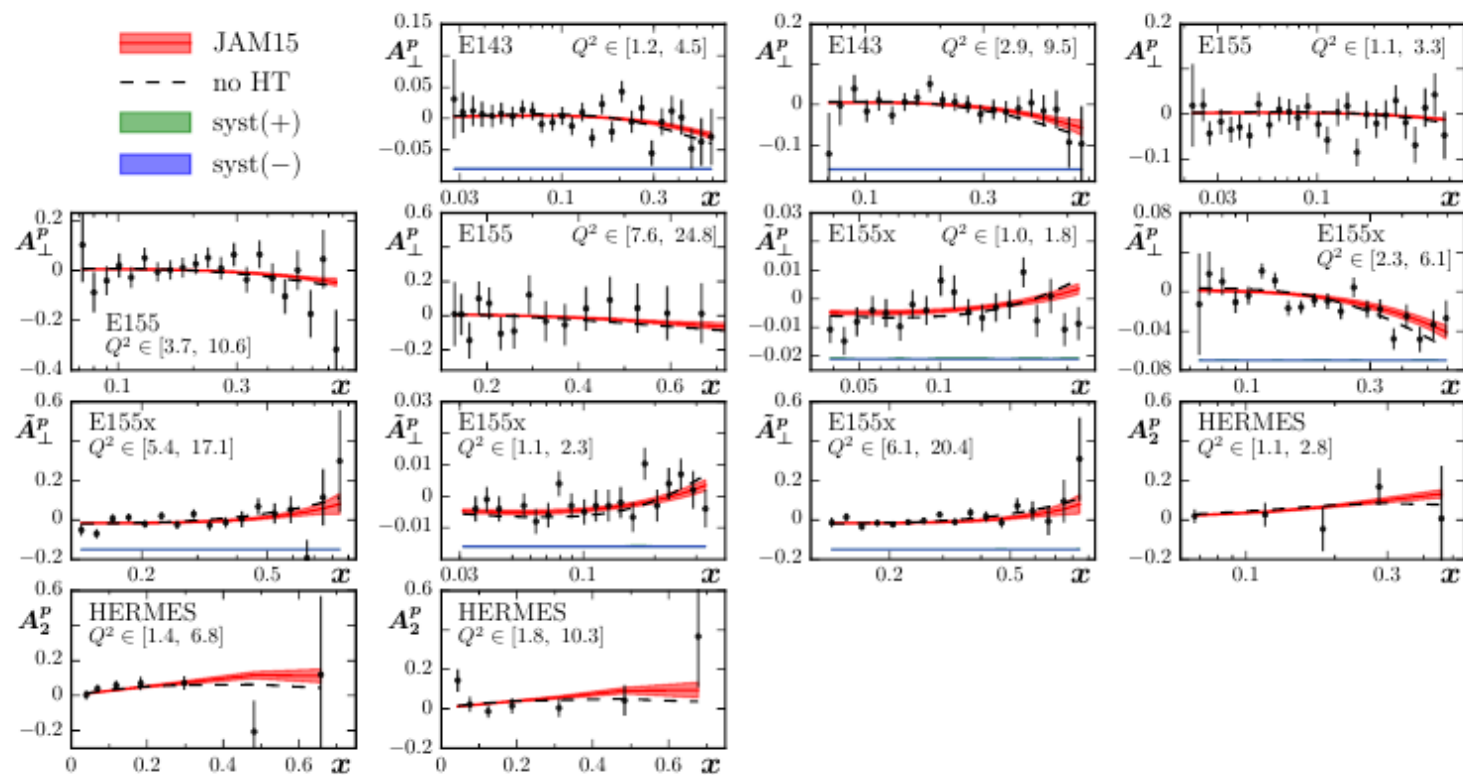
- ❑ Direct Z reconstruction is unambiguous in principle, but:
  - Needs better than 5-10% precision at large rapidity
  - Experimentally achievable?
    - At LHCb? RHIC? AFTER@LHC?
    - Was full data set used at Tevatron?

# Appendix: JAM15 polarized PDF

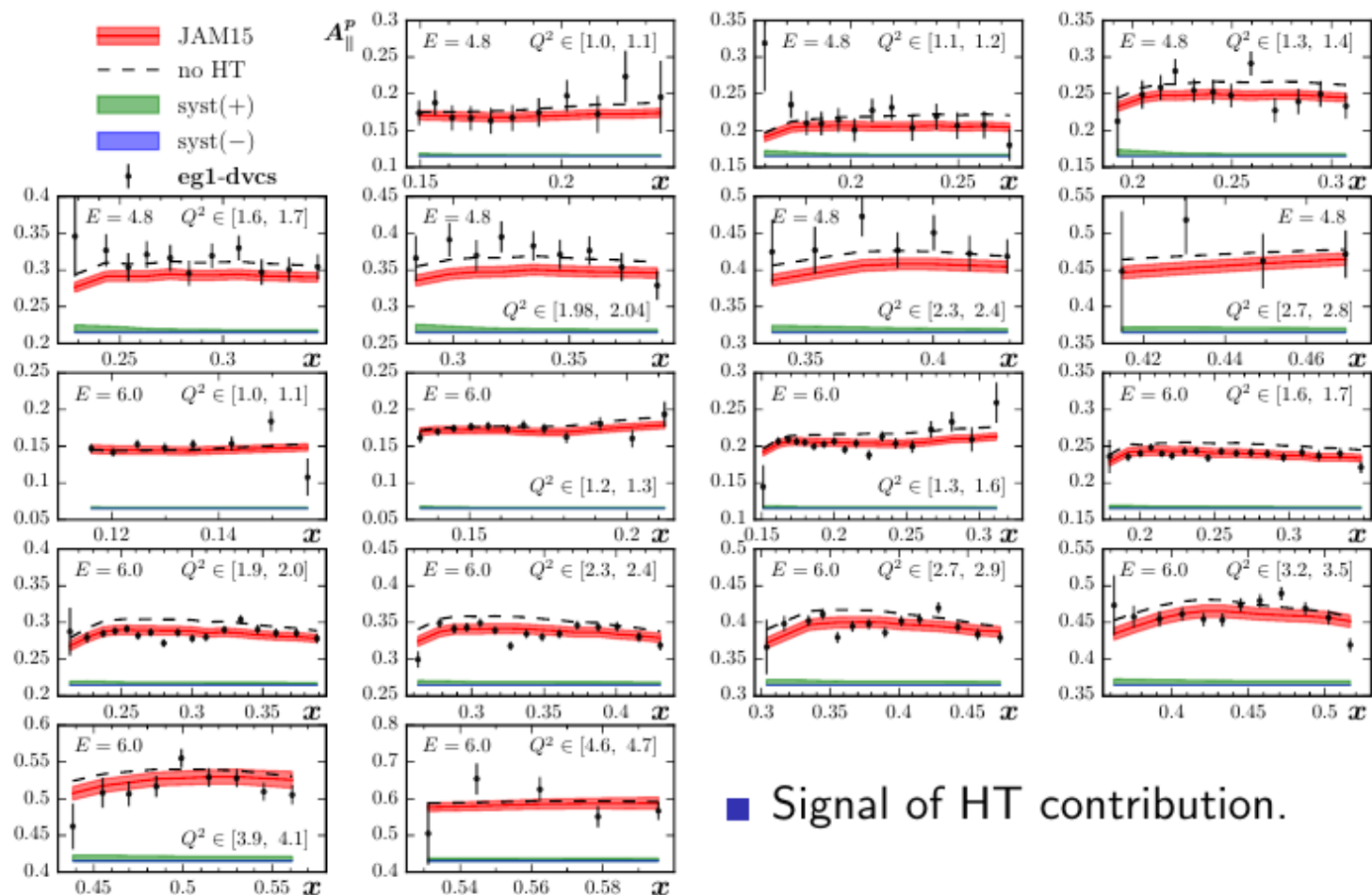
# Data vs. theory: proton



# Data vs. theory: proton

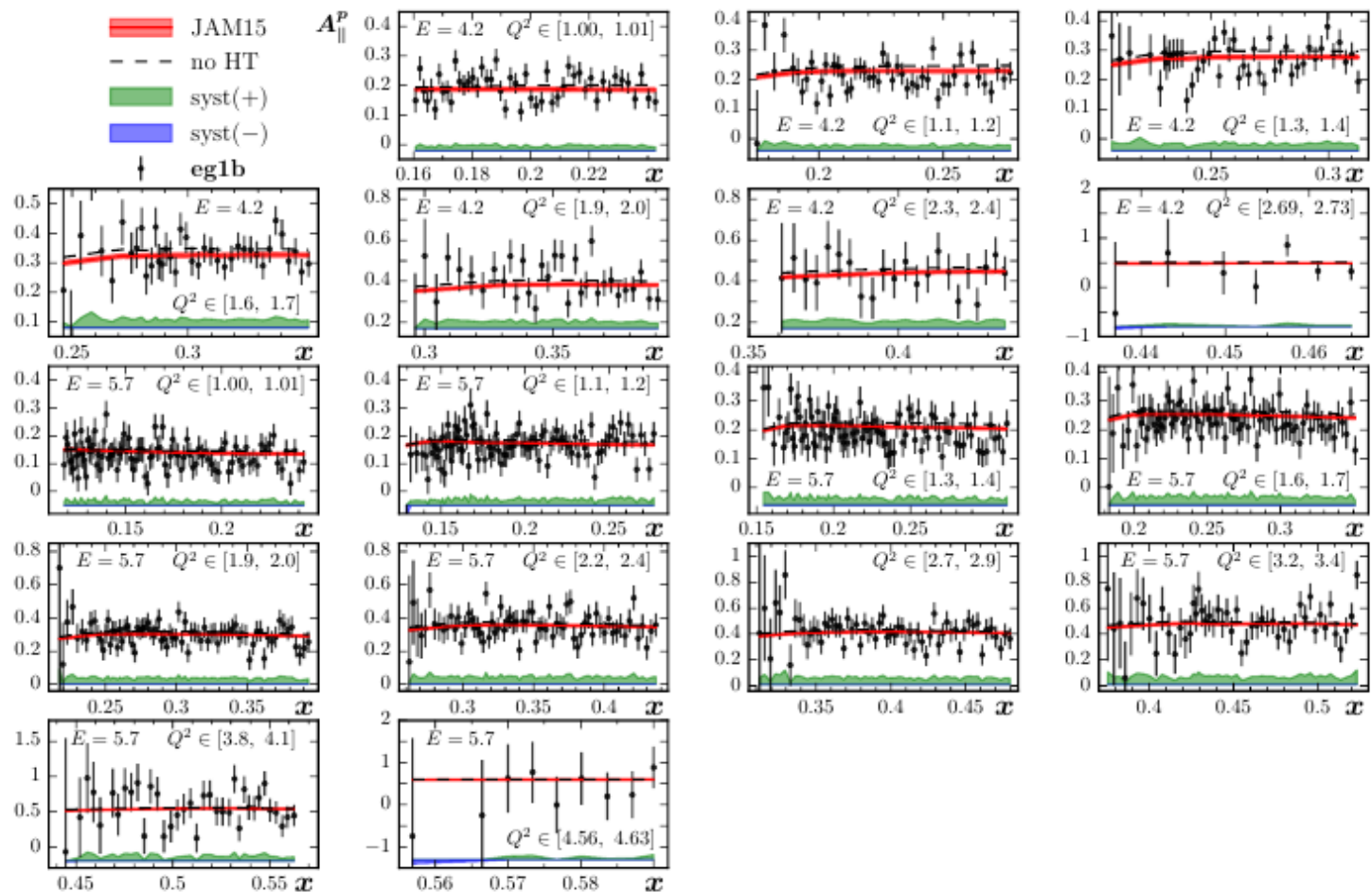


# Data vs. theory: proton EG1b-DVCS

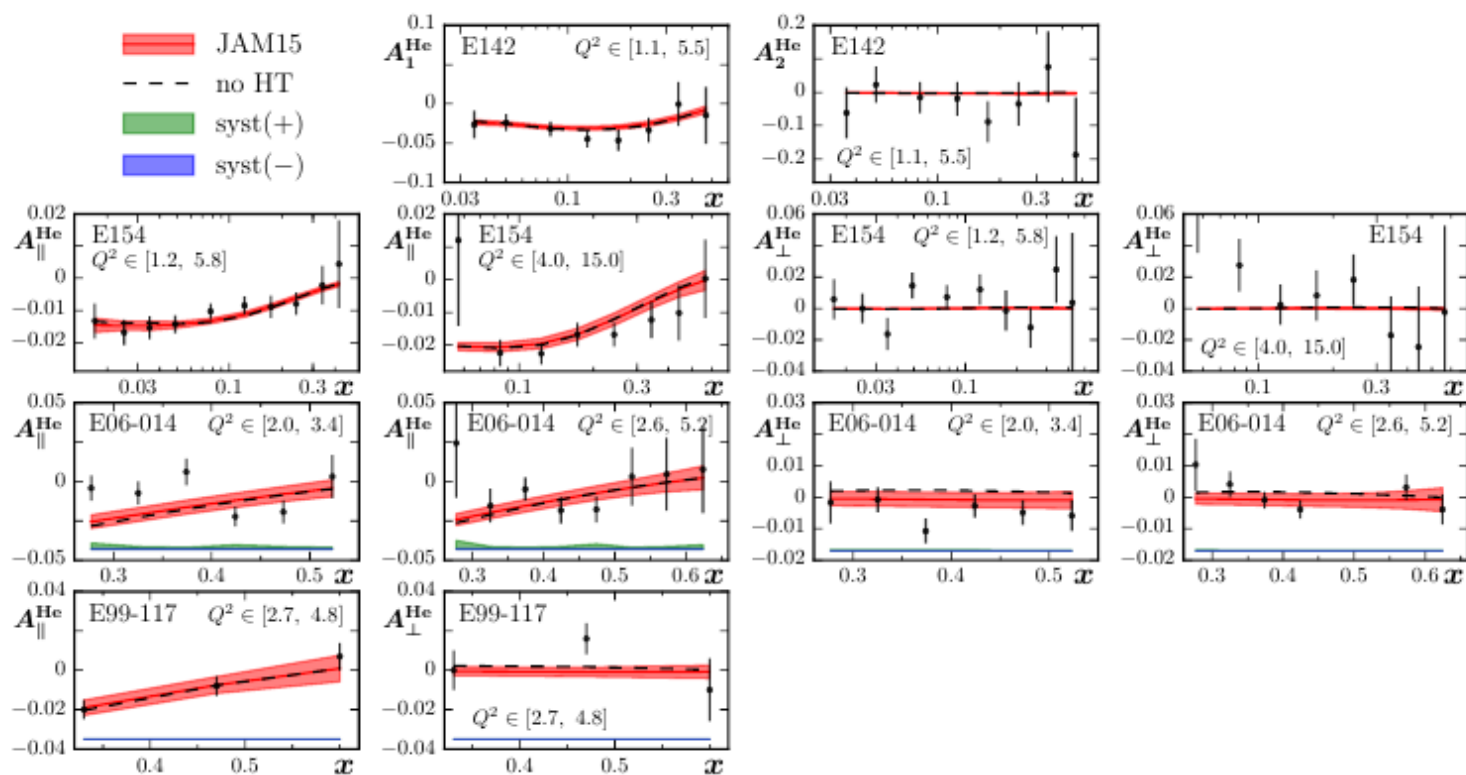




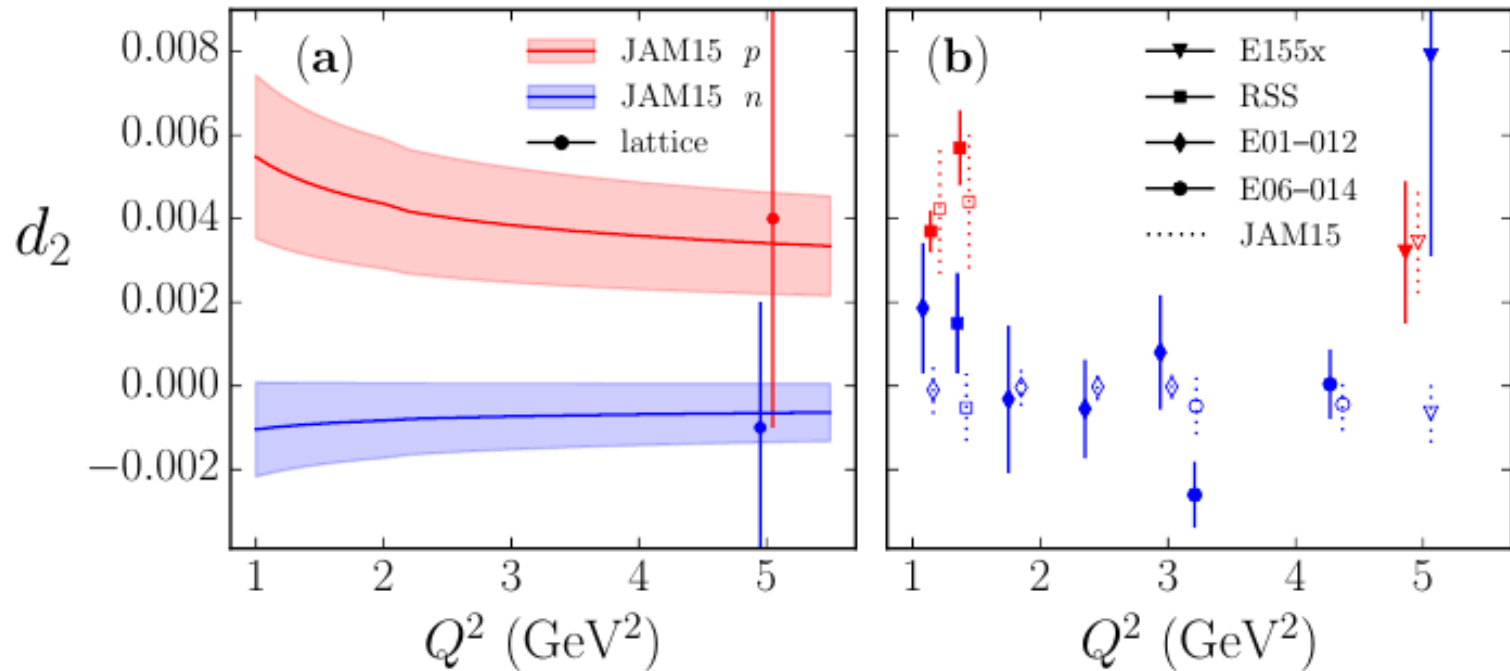
# Data vs. theory: proton EG1b



# Data vs. theory: $^3\text{He}$



# $d_2$ matrix element

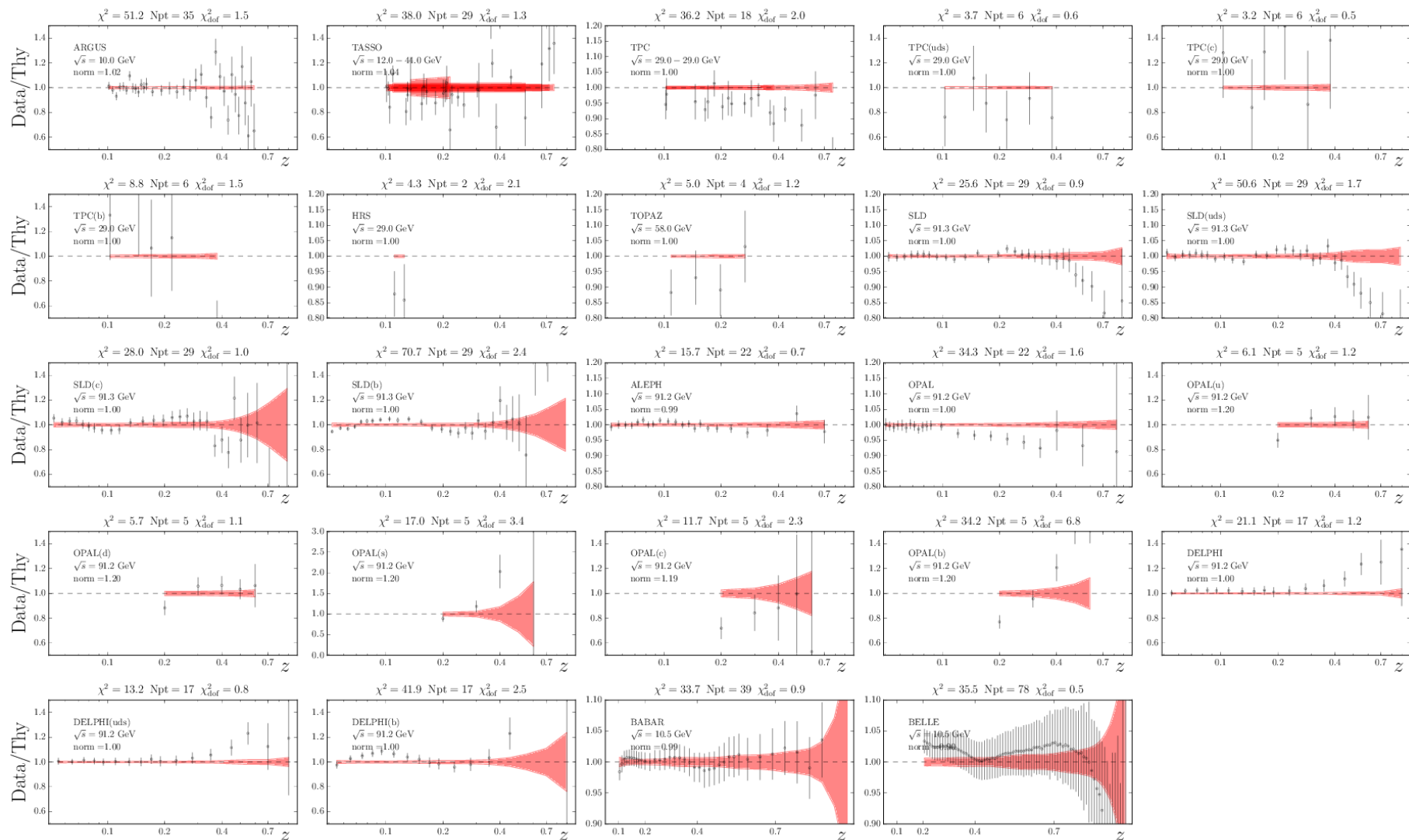


- $d_2(Q^2) \equiv \int_0^1 dx x^2 [2g_1^{T3}(x, Q^2) + 3g_2^{T3}(x, Q^2)]$
- $d_2$  is related to “color polarizability” or the “transverse color force” acting on quarks.
- Existing measurements of  $d_2$  are in the resonance region (contains TMC T4 and beyond.)
- Agreement with data indicates quark-hadron duality

# Appendix: preliminary JAM FFs

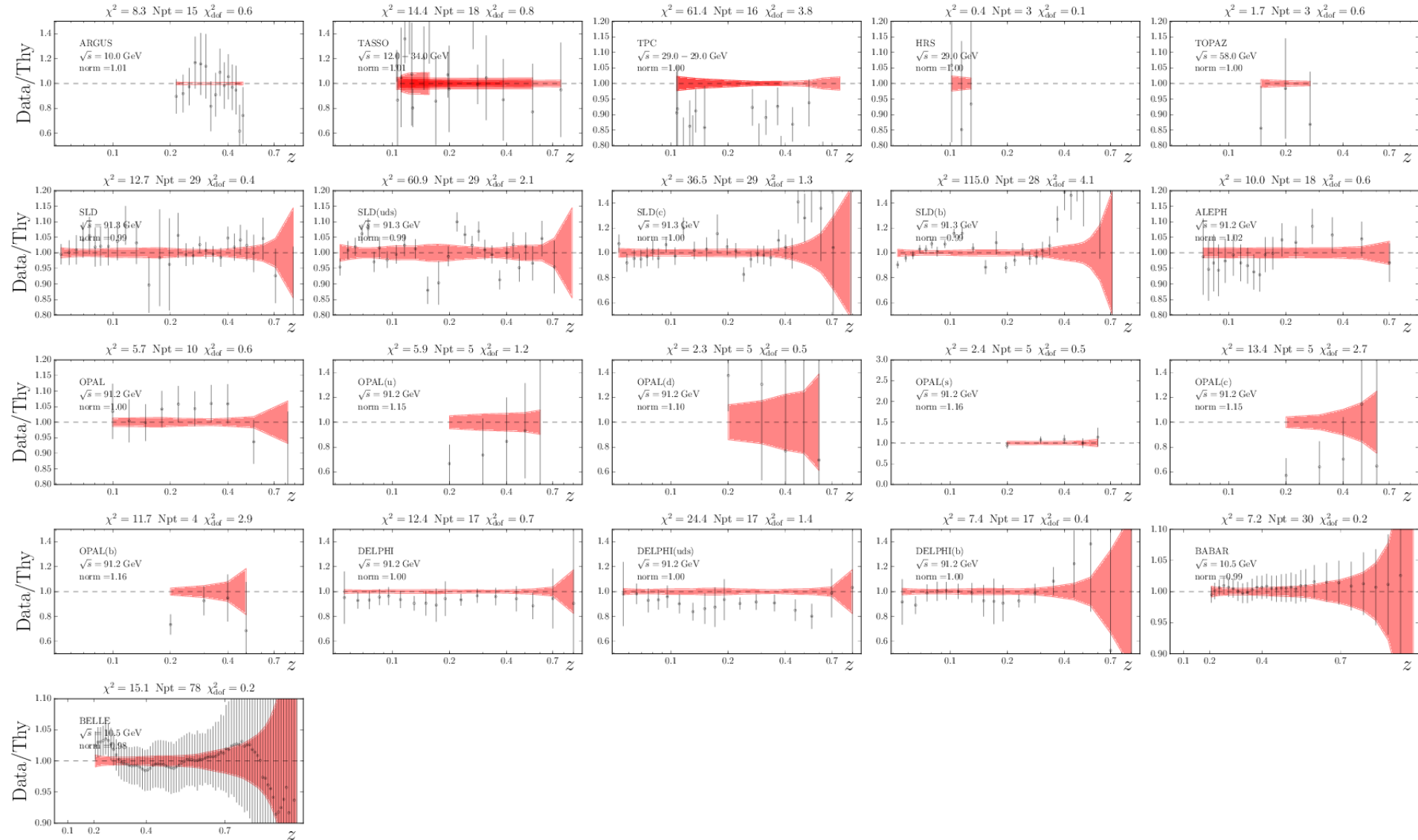
# Pion fit vs. data

Sato, Ethier, Melnitchouk, Accardi – in progress



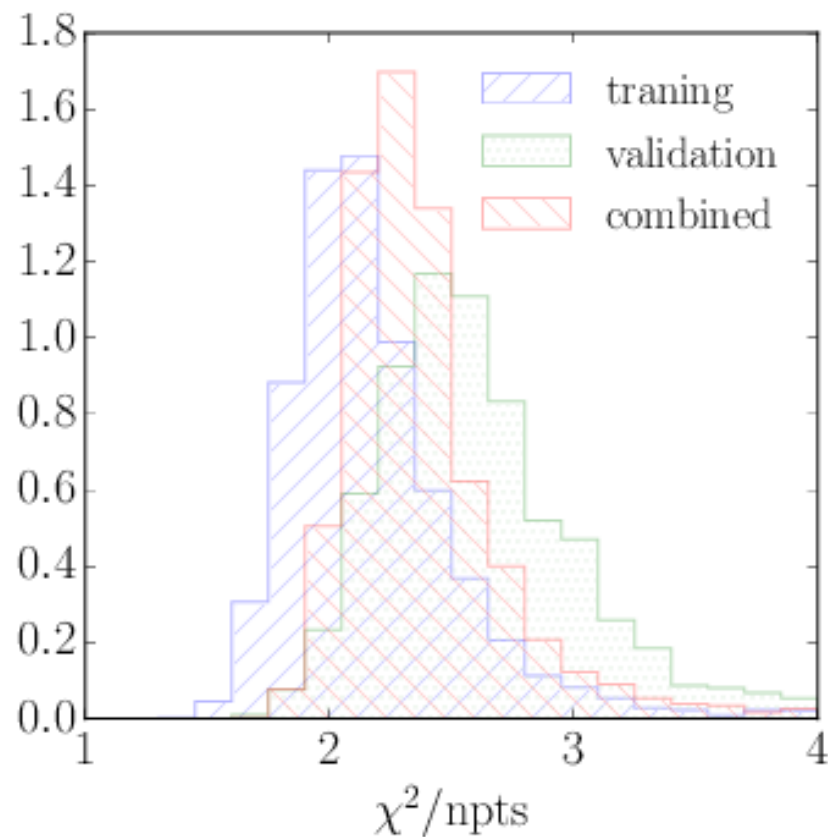
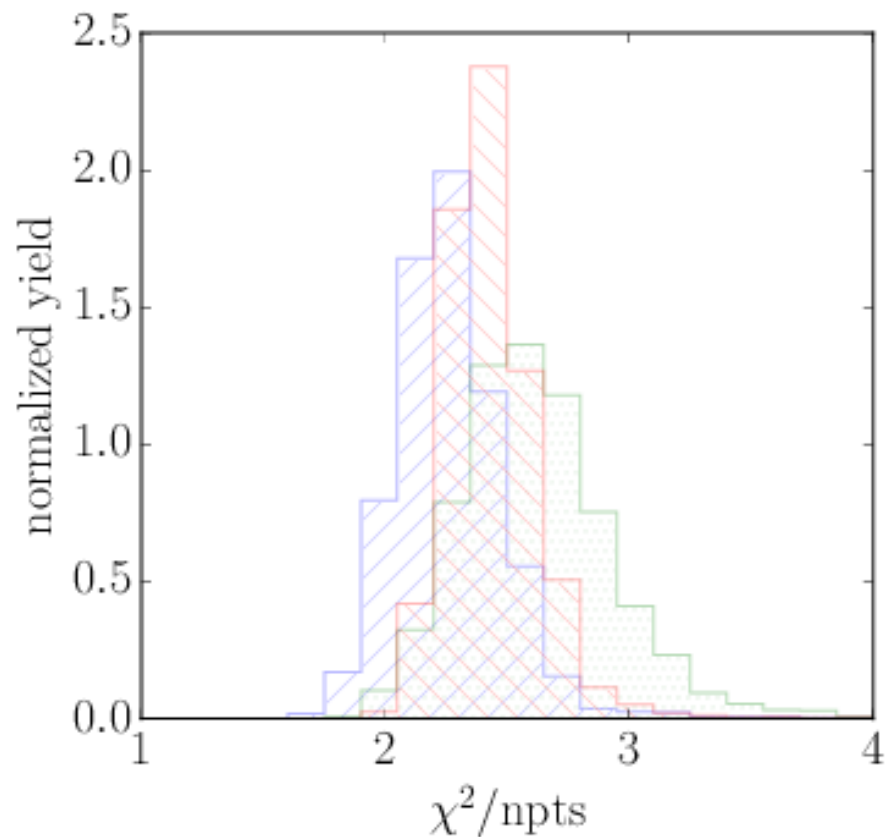
# Kaon fit vs. data

Sato, Ethier, Melnitchouk, Accardi – in progress



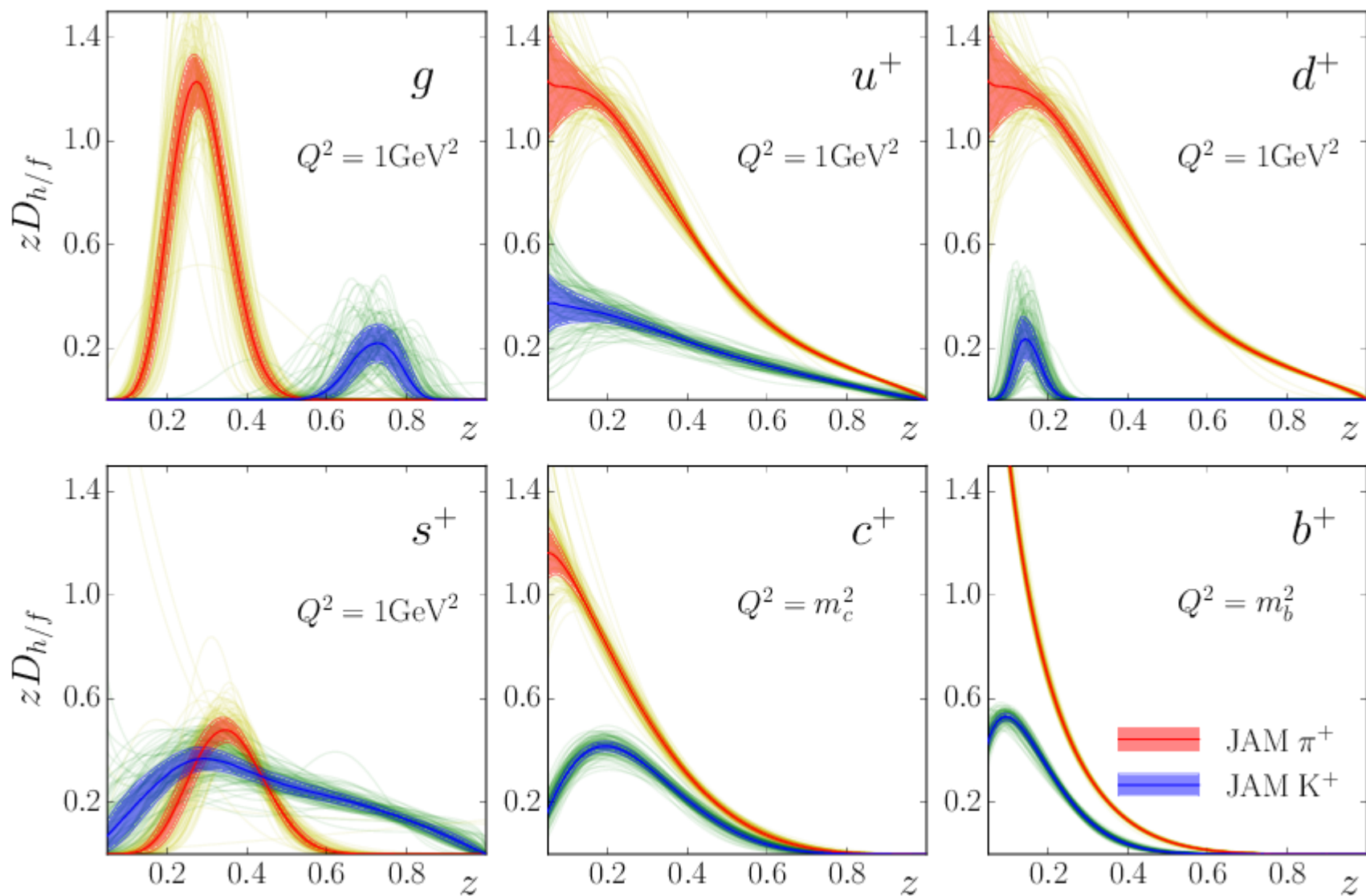
# Chi<sup>2</sup> distribution (pions)

Sato, Ethier, Melnitchouk, Accardi – in progress



# Pions vs. Kaons

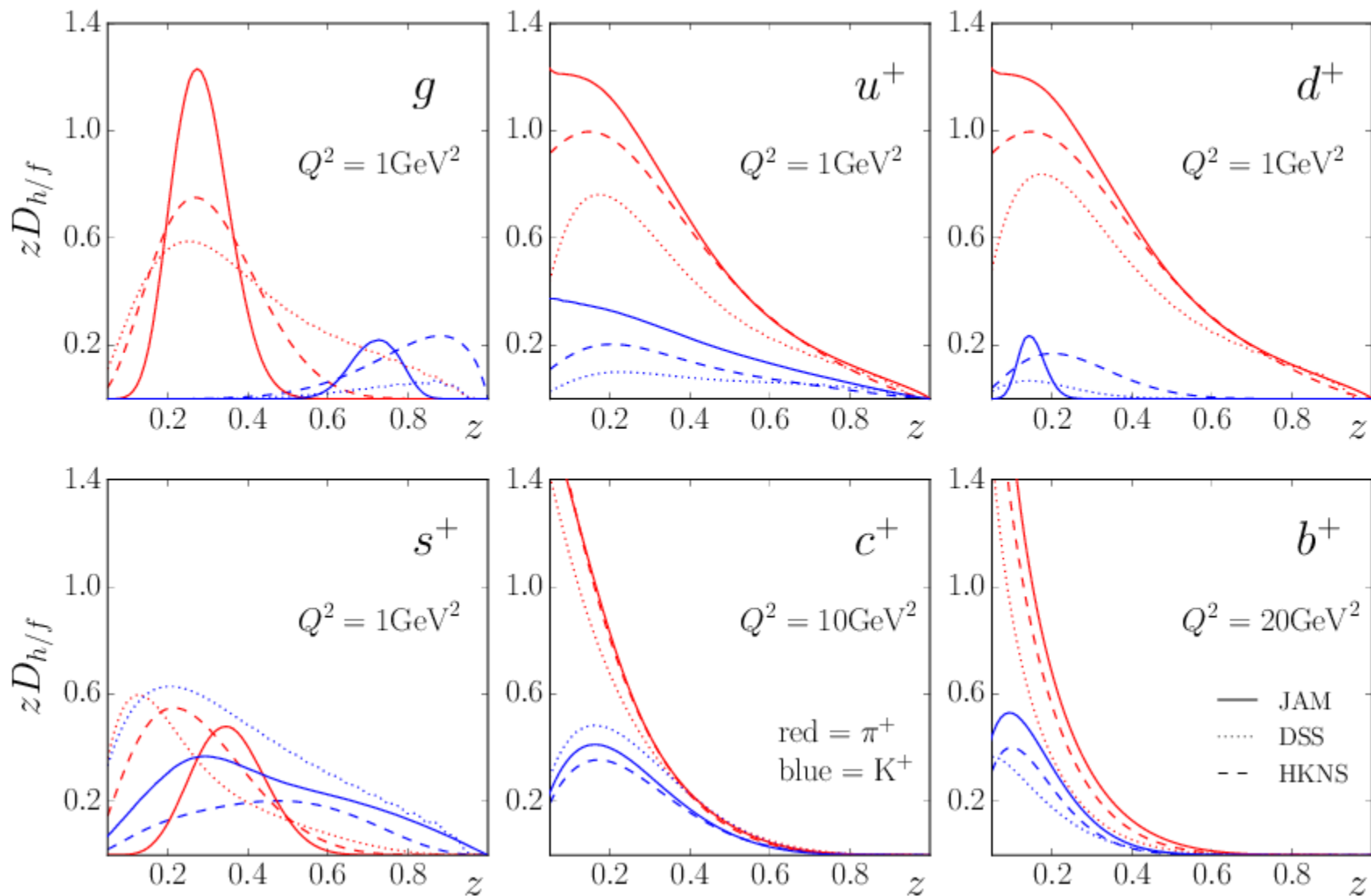
Sato, Ethier, Melnitchouk, Accardi – in progress





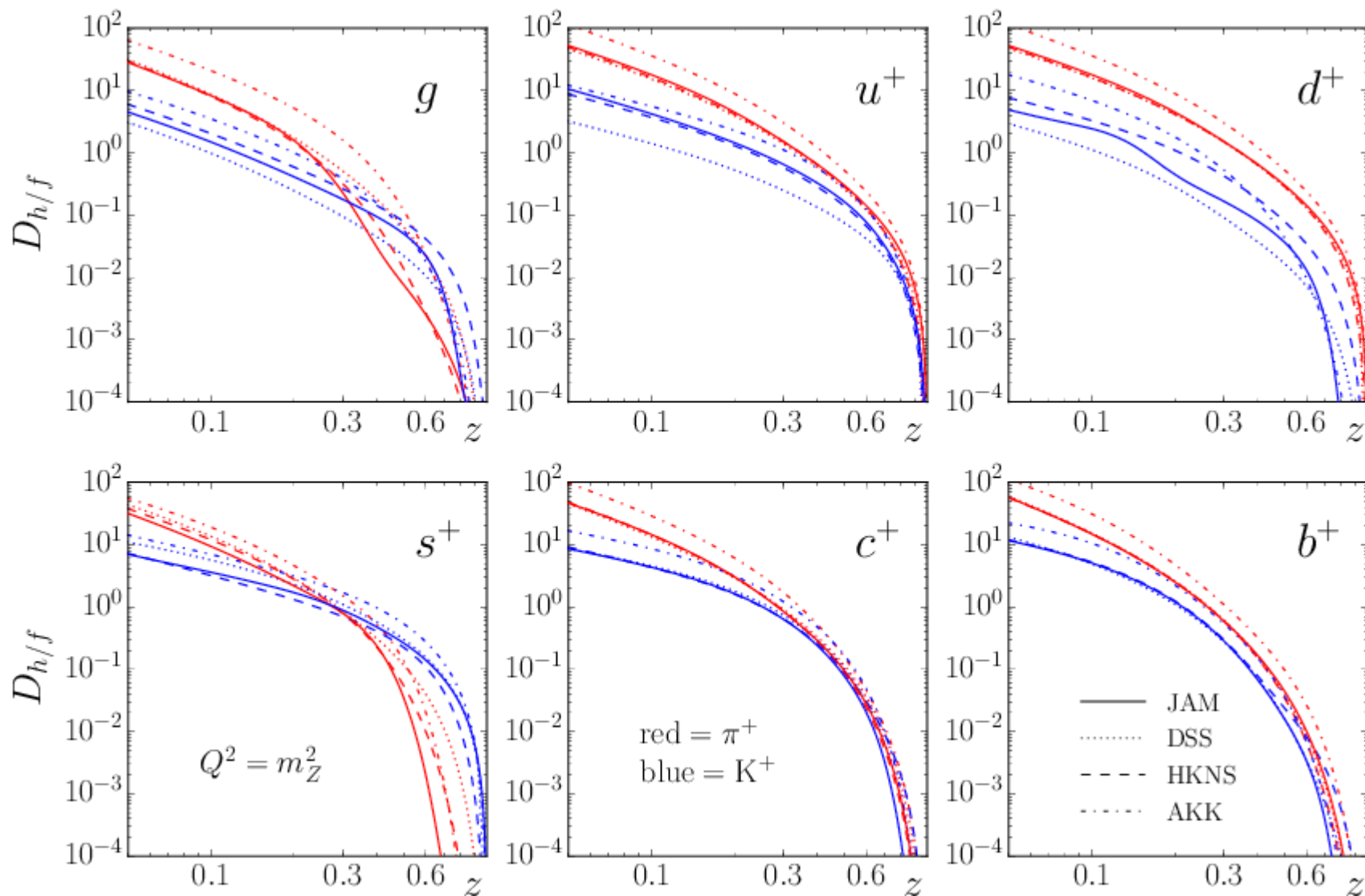
# JAM-FF vs. others

Sato, Ethier, Melnitchouk, Accardi – in progress



# JAM-FF vs. others

Sato, Ethier, Melnitchouk, Accardi – in progress

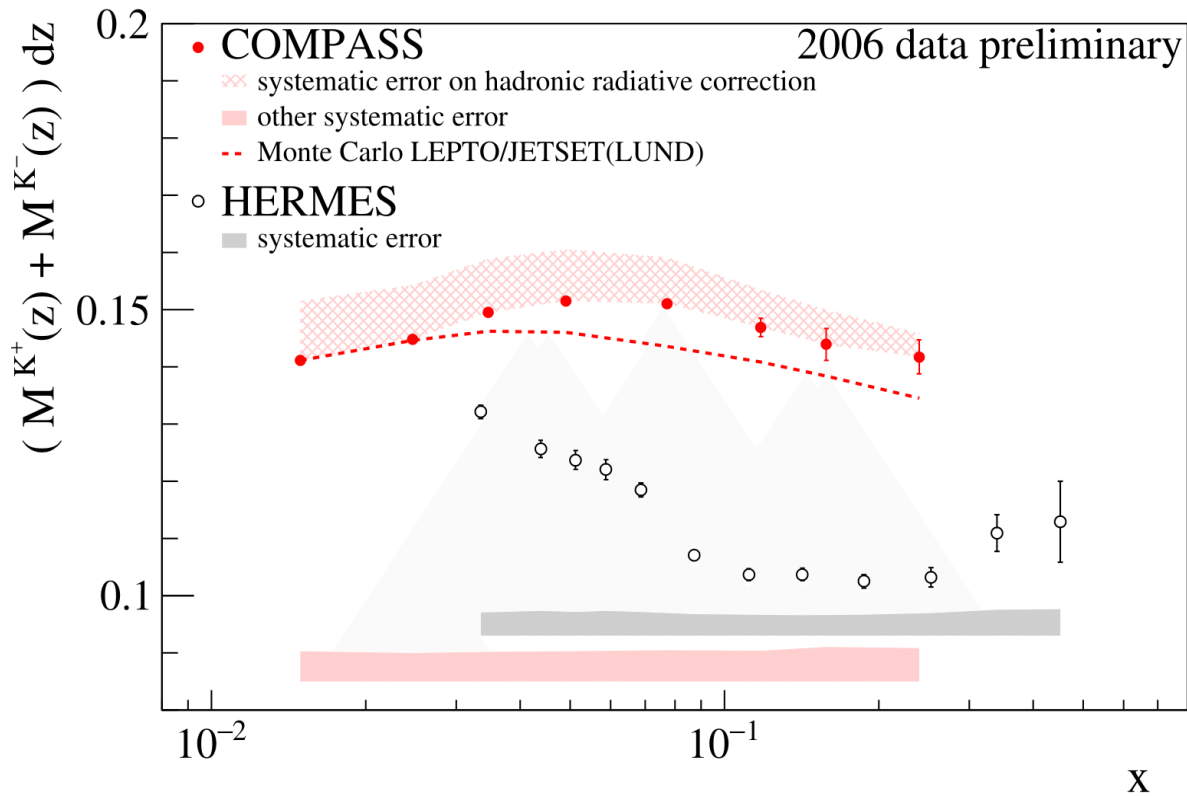


# Appendix: Hadron Mass Corrections, HERMES vs. COMPASS

*J.Guerrero, J.Ethier, AA, S.Casper, W.Melnitchouk, JHEP 1509 (2015) 169*

*J. Guerrero, A.Accardi, work in progress*

# In brief...



$$M_{\text{exp}}^K(x_B, Q^2) = \frac{\int_{0.2}^{0.8 (0.85)} \frac{dN^K}{dx_B dQ^2 dz_h}}{\frac{dN^{DIS}}{dx_B dQ^2}}$$

Strange quark “tagging”

# Hadron Mass corrections in brief

With a massive target nucleon, and massive hadron in final state:

- Respecting partonic 4-momentum conservation also requires a “massive” fragmenting parton with  $\tilde{k}'^2 \geq m_h^2/\zeta_h$ .
- $x_B$  and  $z_h$  mixed in the new scaling variables  $\xi_h$  and  $\zeta_h$

Parton model: 
$$M^{K(0)}(x_B, Q^2) = \frac{\sum_q e_q^2 q(x_B, Q^2) \int dz_h D_q^h(z_h, Q^2)}{\sum_q e_q^2 q(x_B, Q^2)}$$
  
(ignoring mass effects)

With HMCs: 
$$M^K(x_B, Q^2) = \frac{\int dz_h \sum_q e_q^2 q(\xi_h, Q^2) D_q^h(\zeta_h, Q^2)}{\sum_q e_q^2 q(\xi, Q^2)}$$

$$\xi_h = \frac{2x_B}{1 + \sqrt{1 + 4x_B^2 M^2/Q^2}} \left(1 + \frac{m_h^2}{\zeta_h Q^2}\right)$$

$$\zeta_h = \frac{z_h}{2} \frac{\xi}{x_B} \left(1 + \sqrt{1 - \frac{4x_B^2 M^2 m_{h\perp}^2}{z_h^2 Q^4}}\right)$$

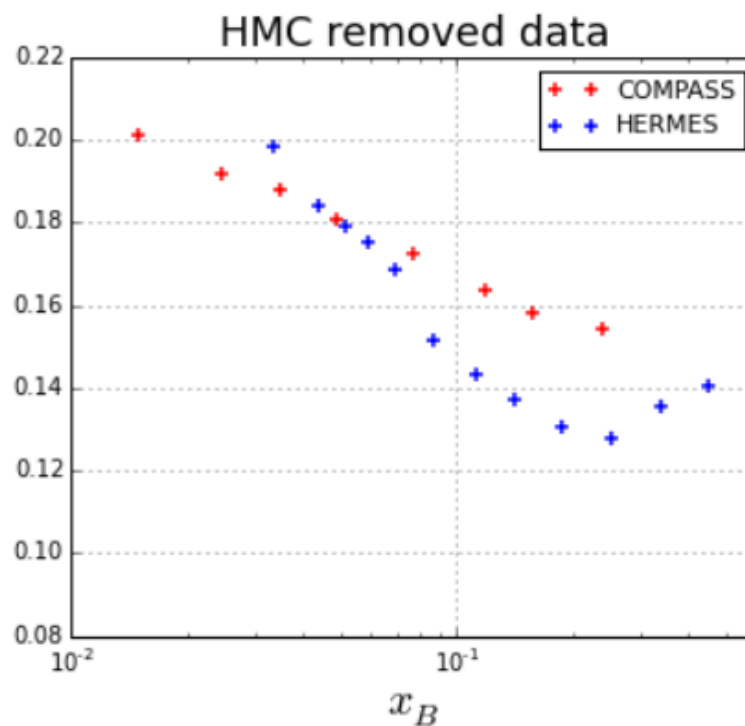
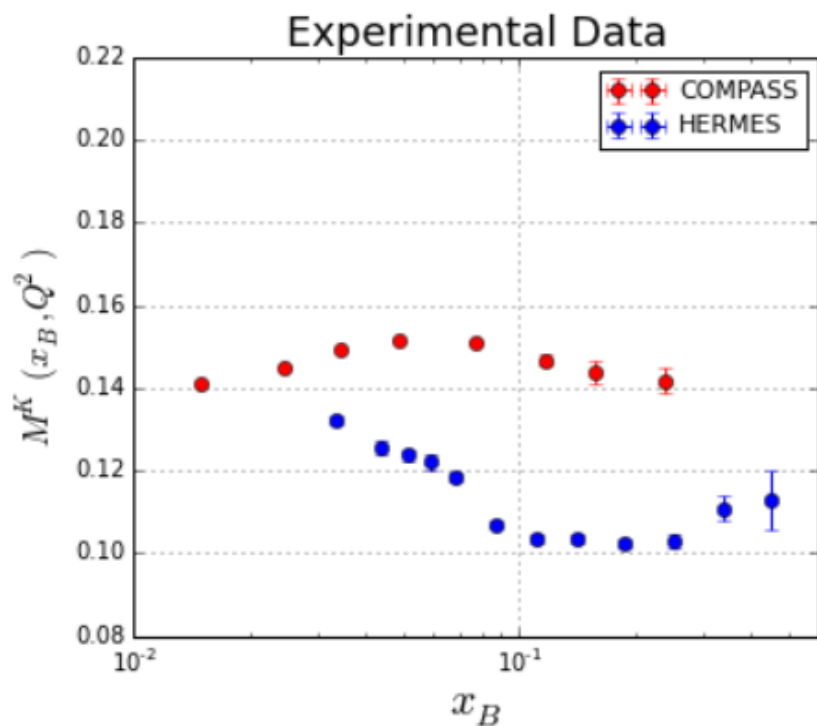
$$Q(x) \equiv u(x) + \bar{u}(x) + d(x) + \bar{d}(x)$$

$$S(x) \equiv s(x) + \bar{s}(x)$$

$$\mathcal{D}_Q^K(z) \equiv 4D_u^K(z) + D_d^K(z)$$

$$\mathcal{D}_s^K(z) \equiv 2D_s^K(z)$$

# COMPASS vs. HERMES



“Remove” HMCs

$$M_{\text{exp}}^{K(0)} \equiv M_{\text{exp}}^K \frac{M^{K(0)}}{M^K} \Bigg|_{\text{theory}}$$

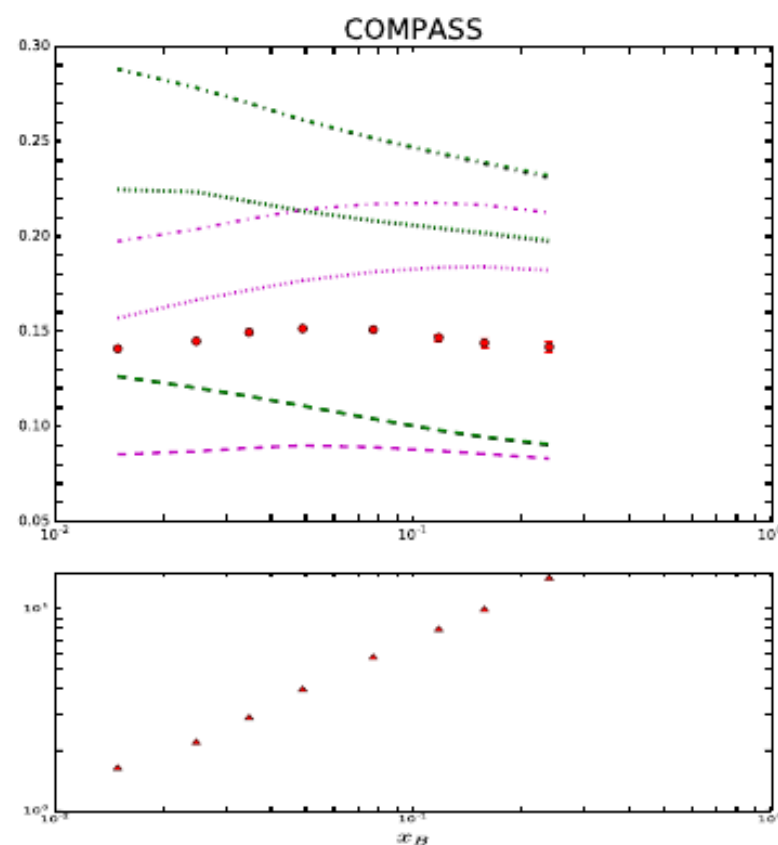
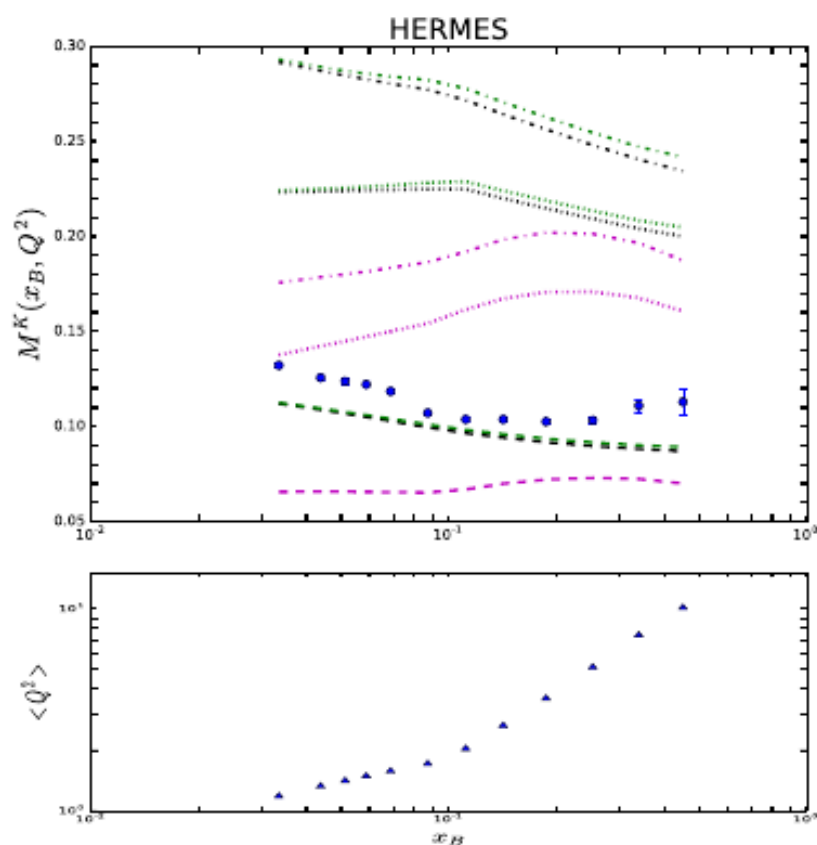
With massless hadrons

With HMC

→ data become directly comparable to parton model

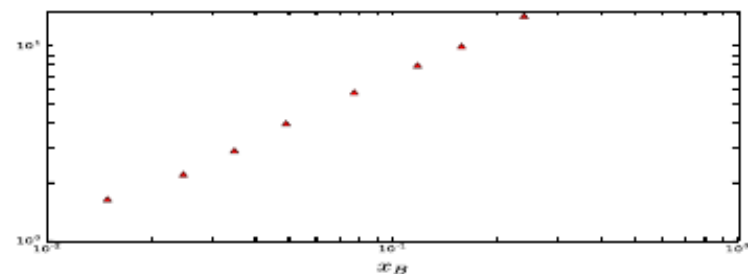
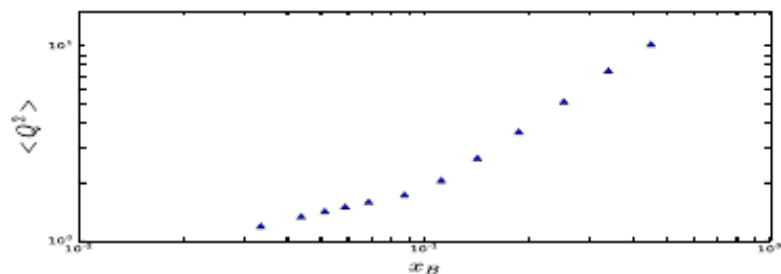
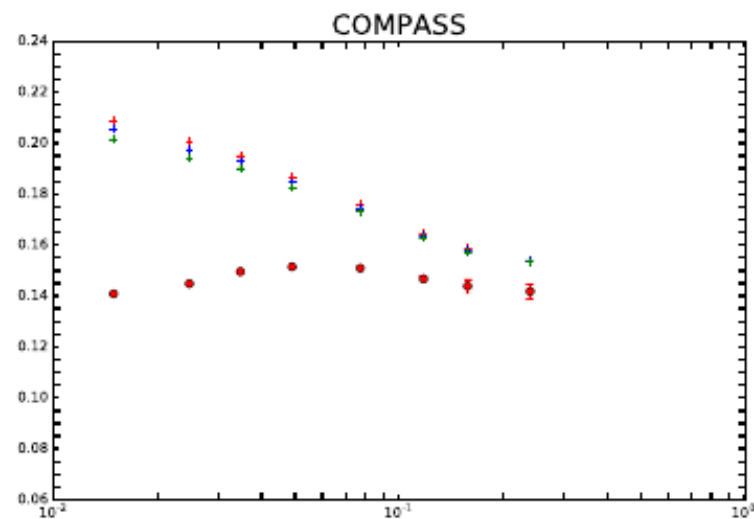
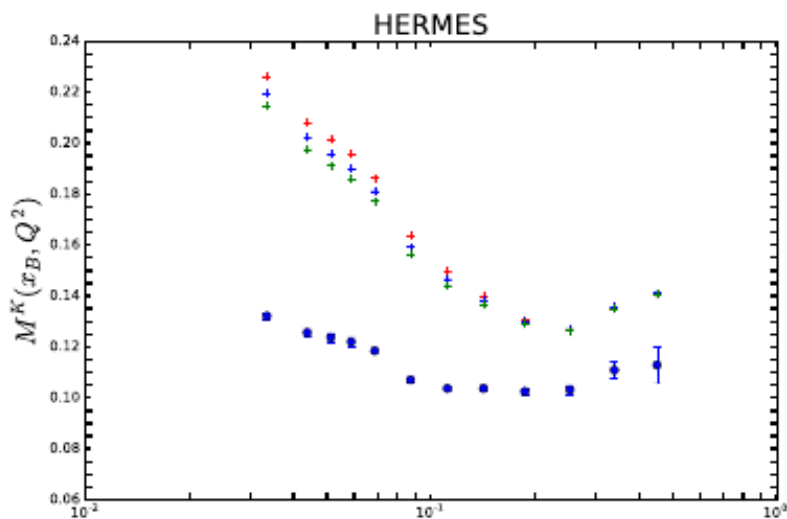
# Fragmentation function systematics

- PDF: CTEQ14
- FFs: DSS (dashed lines), HKNS(dashed-dotted lines), KKP(dotted lines)
- Data (dots) vs. Theory (Bjorken limit, Albino prescription, Our prescription)



# Fragmentation function systematics

- Parton Level Multiplicities:  $M_{exp}^{K(0)} = M_{exp}^k \frac{M^{K(0)}}{MK} \Big|_{Theory}$
- Exp. Data (dots) vs. "Partonic" Data (Cross) (DSS, KKP, HKNS)

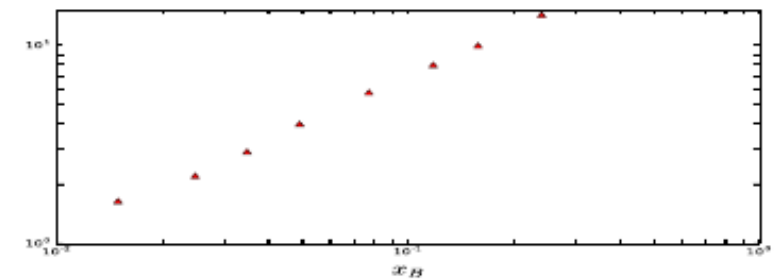
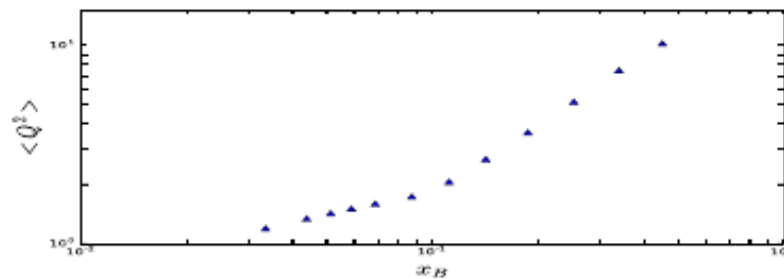
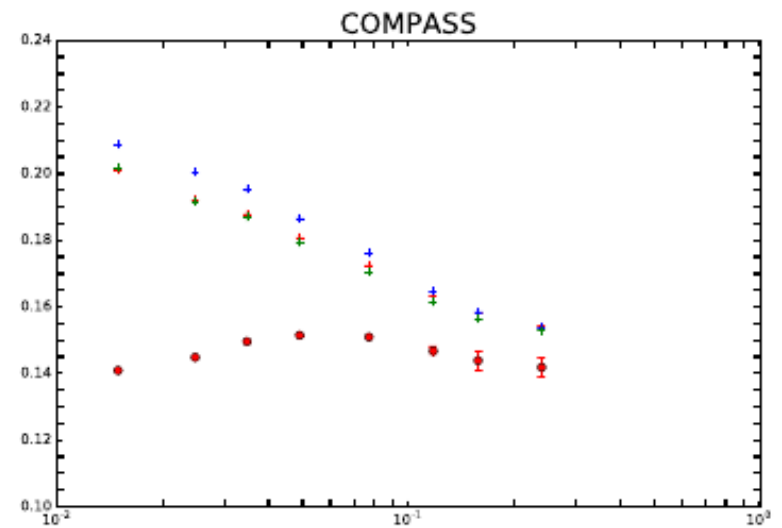
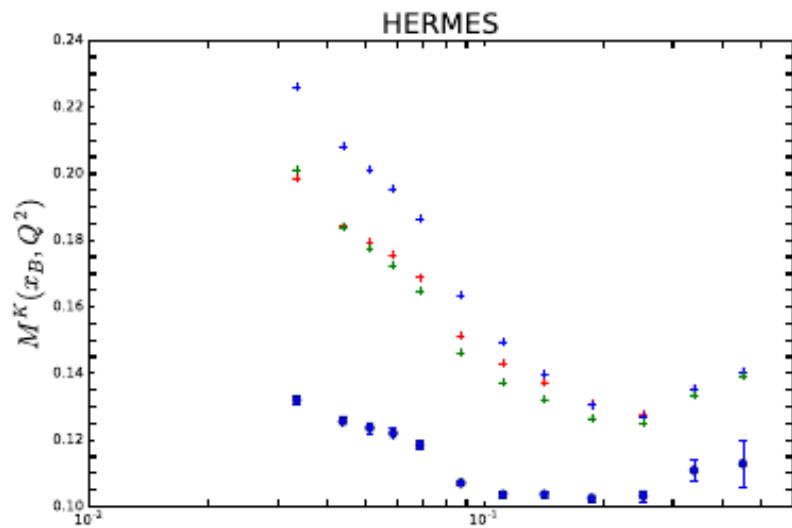


HMCs "removed" data stable

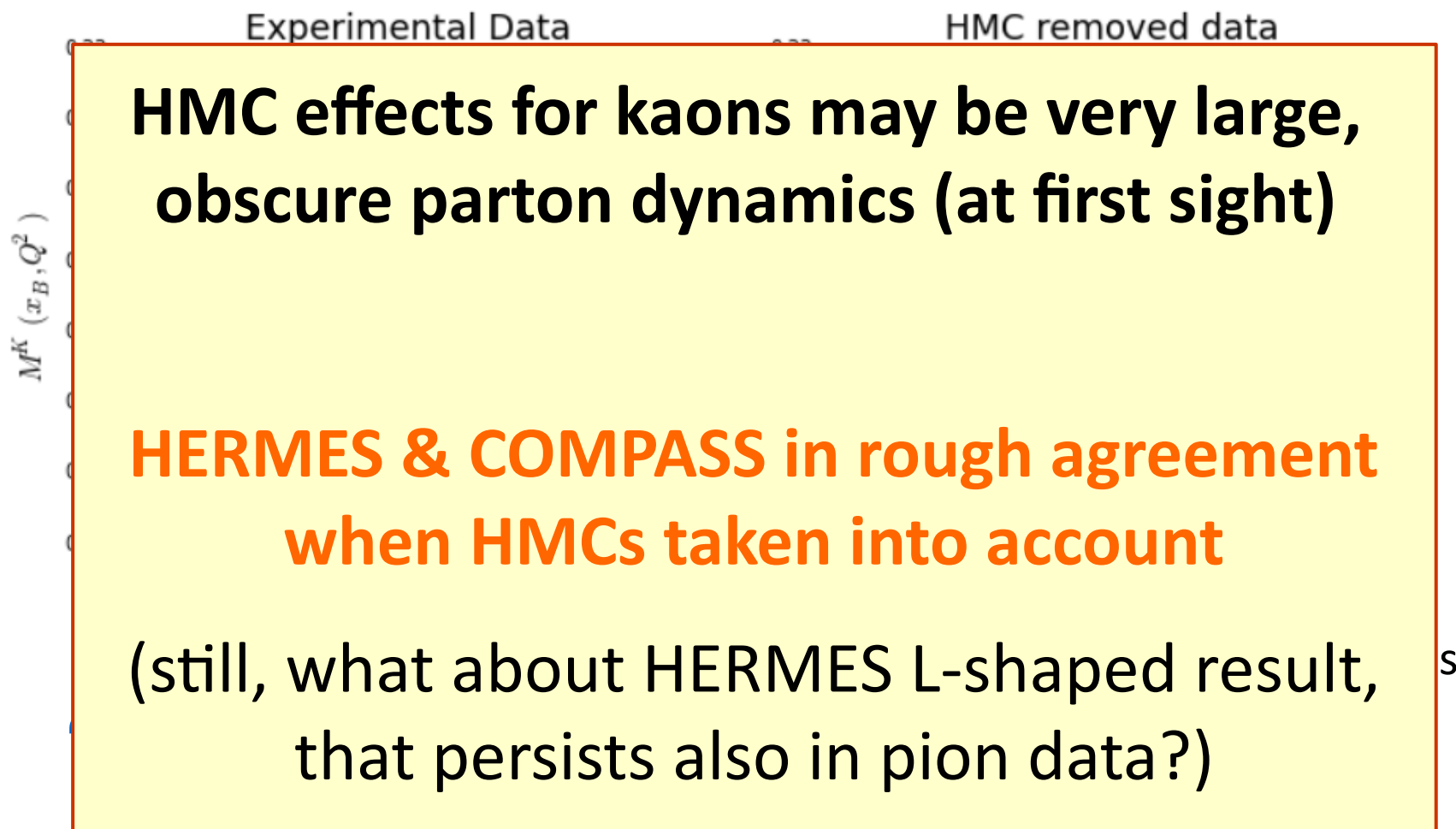


# Fragmentation function systematics

- Parton Level Multiplicities:  $M_{exp}^{K(0)} = M_{exp}^k \frac{M^{K(0)}}{MK} \Big|_{Theory}$
- Exp. Data (dots) vs. "Partonic" Data (Cross) (CT6L, CT14, GJR08)



# COMPASS vs. HERMES



→ data become directly comparable to parton model