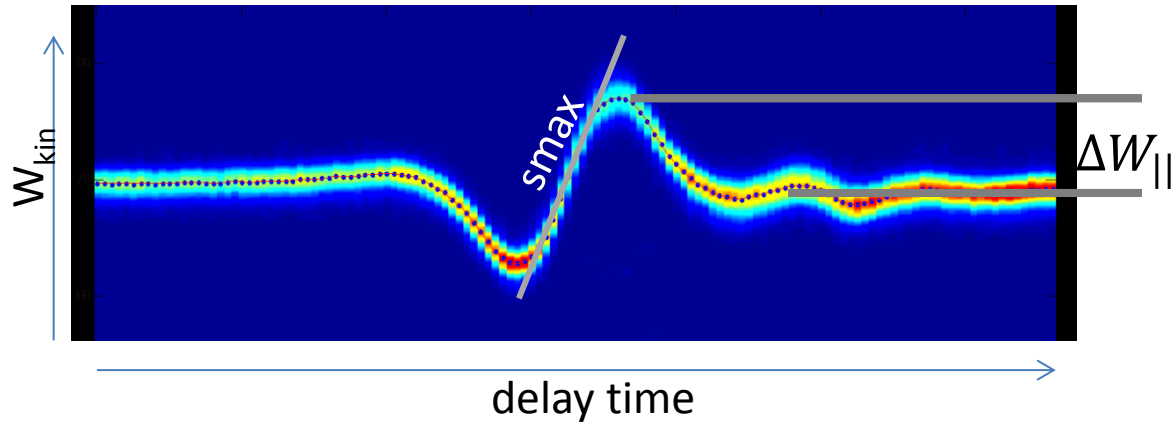
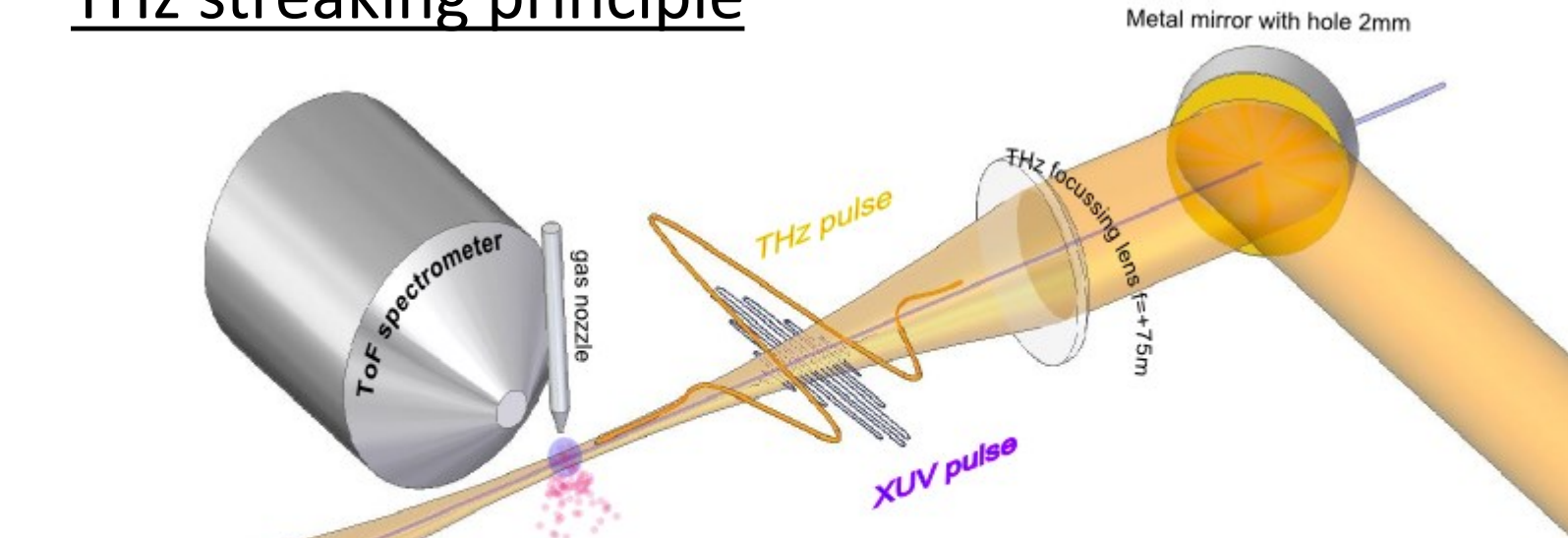


THz streaking principle



$$U_p = \frac{e^2 E_0^2}{4m_e \omega_{IR}^2} \sim I(t) * \lambda_{THz}^2 \quad \text{“ponderomotive potential”}$$

$$\Delta W_{||}(t) = W(t) - W_0 \quad \text{“streaking amplitude”}$$

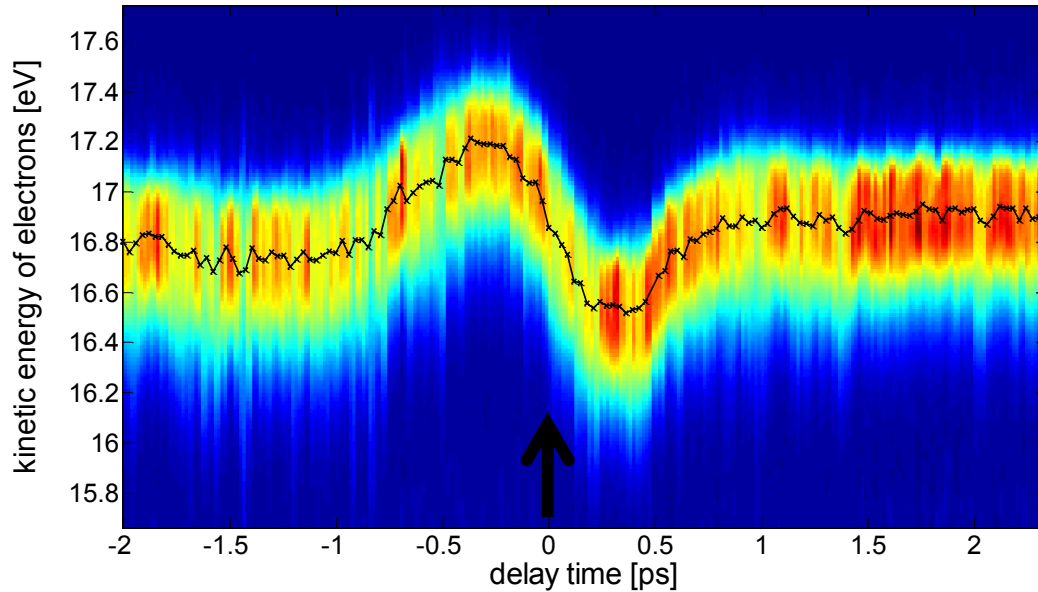
$$\Delta W_{||} = \sqrt{8W_0 U_p} \sin \phi_i = e \sqrt{\frac{2W_0}{m_e}} \frac{E_0(t_i)}{\omega_{Ir}} \sin \phi_i = e \sqrt{\frac{2W_0}{m_e}} A_{Ir}(t_i)$$

Streaking speed

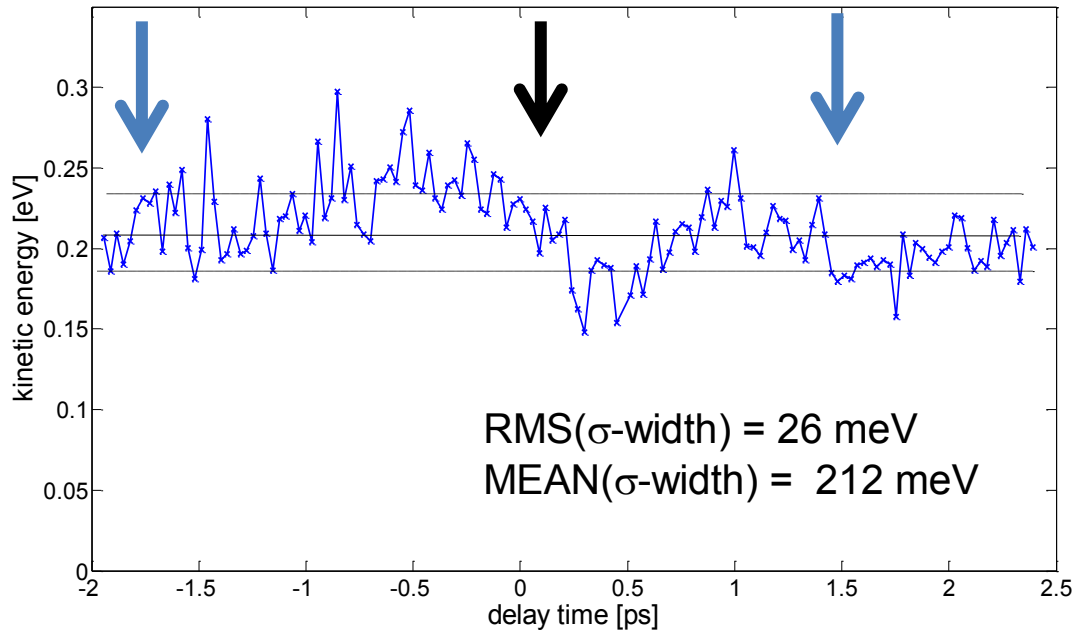
$$S_{max} = \omega_{IR} \Delta W_{max}$$

Approximation of seeded FLASH XUV pulse duration

Streaking delay scan - single shot, Argon 3p, FEL at 38.1nm



Sigma-width of traces



Because of a small streaking strength, the spectral width of a streaked FEL spectrum at time zero was not larger than the width of any unstreaked spectrum at delay times far off the temporal overlap region.

*Hence, we can only calculate an **upper limit** for the temporal width of the seeded FEL pulse by **comparing the RMS value of the jitter of the spectral width** with the **broadening due to the streaking strength**.*

Approximation of seeded FLASH XUV pulse duration

$$\sqrt{\sigma_{XUV}^2 + (\tau_{XUV} \cdot s)^2} \leq \sqrt{\sigma_{XUV}^2 + (2 \cdot RMS(\sigma\text{-width}))^2}$$

$$\tau_{XUV} \cdot s \leq 2 \cdot RMS(\sigma\text{-width}) \leq 52meV$$

$$\tau_{XUV,RMS} \leq \frac{52meV}{1.6 \frac{meV}{fs}} = 32.5fs \text{ RMS} \leftrightarrow 76fs \text{ FWHM}$$

Assuming **95%-significance** (\Leftrightarrow **streaking broadening** is at least two times larger than the **RMS value of the jitter of the spectral width**), one finds that the temporal duration of the measured XUV pulse should be **smaller than 76fs FWHM**.

For any larger XUV pulse duration, the measurement would have shown **with 95% probability** a measurable spectral broadening at time zero !

(In this consideration any effect of a temporal chirp of the FEL pulse is disregarded!)