

High-Energy Scattering in $\mathcal{N} = 4$ Super Yang-Mills and Single-Valued Polylogarithms

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based on arXiv:1512.04963, 1606.08411
with J. Broedel and A. Torres Orjuela

ETH zürich

Scattering amplitudes in planar $\mathcal{N} = 4$ SYM

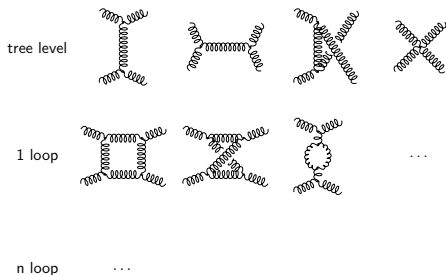
$$A_6(s_{i\dots j}) = A_{6,\text{tree}} e^{A_{\text{BDS}} + R_6(u_i)}$$

- 4, 5 – gluon amplitudes known to all loop orders [Bern et al., Drummond et al.]
- remainder function R_6 known to 5 loops by amplitude bootstrap [Dixon et al.]
 - ansatz in known space of functions
 - constraints from kinematic limits fully fix amplitude
 - no Feynman diagrams, no loop integrals!
- results hint at simple all-loop description for R_6
- How can we go beyond order-by-order description?

The Regge limit

[Balitsky, Fadin, Kuraev, Lipatov]

- consider high-energy limit $s \gg -t$ of four gluon amplitude in YM
- many diagrams are kinematically suppressed

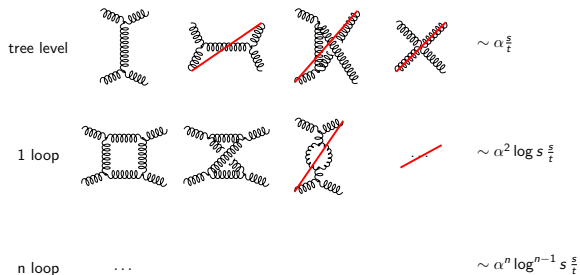


- leading contributions $\sim \alpha^n \log^{n-1} s$ (LLA)
- in limit $s \rightarrow \infty$: $\alpha \log s \sim \mathcal{O}(1)$
 - need to resum leading contributions from all loop orders
 - emergence of effective particles: Reggeons

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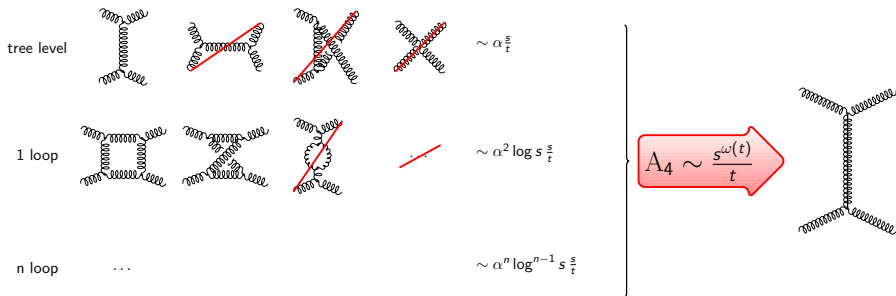


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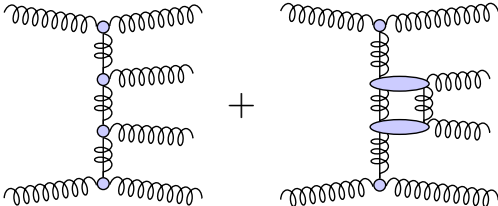
Results

- $R_{6,\text{MRL}}$ is determined by Fourier-Mellin transform to all loop orders
- evaluation of integral tedious
- understanding resulting functions trivializes evaluation of integral
→ provide high loop data, proof of open LLA conjecture

An all-loop proposal

[Bartels et al., Lipatov and Prygarin, Lipatov and Fadin, Dixon et al., Caron-Huot, Basso et al.]

$$e^{R_6 + i\pi\delta} \Big|_{\text{MRL}} = \cos \omega_{ab} + i \frac{\lambda}{2} \sum_{n \in \mathbb{Z}} (-1)^n \left(\frac{w}{w^*}\right)^{\frac{n}{2}} \int_{\mathbb{R}} \frac{d\nu}{\nu^2 + \frac{n^2}{4}} |w|^{2i\nu} \Phi(\nu, n) \left(- (1 - u_1) \frac{|w|}{|1+w|^2} \right)^{-\omega(\nu, n)}$$

$$e^{R_6 + i\pi\delta} =$$


- impact factor $\Phi(\nu, n) = 1 + \lambda \Phi_{\text{NLLA}}(\nu, n) + \dots$
- BFKL eigenvalue $\omega(\nu, n) = -\lambda \omega_{\text{LLA}}(\nu, n) - \lambda^2 \omega_{\text{NLLA}}(\nu, n) + \dots$
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Harmonic polylogarithms

[Remiddi and Vermaseren]

- generalization of classical polylogarithms $\text{Li}_n(z)$
- appear frequently in scattering amplitudes ($\mathcal{N} = 4$ SYM, QCD, ...)

$$H_{x_1, x_2, \dots, x_n}(z) := \int_0^z \frac{dt}{t - x_1} H_{x_2, \dots, x_n}(t), \quad x_i \in \{0, 1\}$$
$$H_{\underbrace{0, \dots, 0}_n}(z) := \frac{1}{n!} \log^n(z)$$

- examples: $H_{\underbrace{1, \dots, 1}_n}(z) = \frac{(-1)^n}{n!} \log^n(1 - z)$, $H_{\underbrace{0, \dots, 0, 1}_n}(z) \sim \text{Li}_n(z)$
- series expansion: e.g. $H_{1,1}(z) = \sum_{\ell=1}^{\infty} \frac{z^\ell}{\ell} Z_1(\ell - 1)$, where $Z_1(n) = \sum_{k=1}^n \frac{1}{k}$

A two-loop example

- simplest example: two-loop LLA $\sim \lambda^2 \log(1 - u_1)$
- sum over residues at $i\nu = \frac{n}{2} + m$, $m \in \mathbb{N}_0$

$$R_{6,\text{LLA}}^{(2)}(w, w^*) \sim \frac{1}{\pi} \sum_{n=-\infty}^{\infty} (-1)^n \left(\frac{w}{w^*}\right)^{\frac{n}{2}} \int_{-\infty}^{+\infty} \frac{d\nu}{\nu^2 + \frac{n^2}{4}} |w|^{2i\nu} \underbrace{\left(-\frac{1}{2} \frac{|n|}{(\nu^2 + \frac{n^2}{4})} + \psi\left(1 + i\nu + \frac{|n|}{2}\right) + \psi\left(1 - i\nu + \frac{|n|}{2}\right) - 2\psi(1) \right)}_{\omega_{\text{LLA}}(\nu, n)}$$

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 &= \sum_{m=1}^{\infty} \left\{ 2 \frac{|w|^{2m}}{m^2} - 2 \frac{(-w)^m + (-w^*)^m}{m^2} + [\log |w|^2 + 2Z_1(m)] \frac{(-w)^m + (-w^*)^m}{m} \right\} \\
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 \end{aligned}$$

→ observation: result is single-valued!

Single-valued harmonic polylogarithms

[F. Brown, in MRL context : Dixon et al.]

- HPLs $H_a(z)$ can be lifted to single-valued functions $\mathcal{L}_a(z, z^*)$
- explicit construction using Drinfeld associator

$$\mathcal{L}_0(z) = H_0(z) + H_0(z^*) = \log z + \log z^*$$

$$\mathcal{L}_1(z) = H_1(z) + H_1(z^*)$$

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\vdots

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Using the single-valuedness

- How is single-valuedness encoded in the integral?
- key observation: $\mathcal{L}_a(z, z^*) = H_a(z) + \dots$
[Drummond et al., Drummond and Papathanasiou]
- calculate residue $i\nu = \frac{n}{2}$, restore full result by $H_a(-w) \rightarrow \mathcal{L}_a(-w, -w^*)$

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$$\begin{aligned} R_{6,LLA}^{(2)}(w, w^*) &\sim \frac{1}{\pi} \sum_{n=-\infty}^{\infty} (-1)^n \left(\frac{w}{w^*}\right)^{\frac{n}{2}} \int_{-\infty}^{+\infty} \frac{d\nu}{\nu^2 + \frac{n^2}{4}} |w|^{2i\nu} \left(-\frac{1}{2} \frac{|n|}{\nu^2 + \frac{n^2}{4}} + \psi\left(1 + i\nu + \frac{|n|}{2}\right) + \psi\left(1 - i\nu + \frac{|n|}{2}\right) - 2\psi(1) \right) \\ &= \sum_{n=1}^{\infty} (-w)^n \left(\frac{\log w}{n} - \frac{2}{n^2} + 2 \frac{Z_1(n)}{n} \right) \\ &= H_0(-w) H_1(-w) - 2H_{0,1}(-w) + 2H_{1,1}(-w) + 2H_{0,1}(-w) \\ &\rightarrow \mathcal{L}_0(-w) \mathcal{L}_1(-w) + 2\mathcal{L}_{1,1}(-w) \end{aligned}$$

- we identified relations among different integrals which hold at $i\nu = \frac{n}{2}$
- all integrals reducible to master integrals with single pole \rightarrow trivial

Conclusions and outlook

- studied six-gluon amplitude in $\mathcal{N} = 4$ SYM in multi-Regge limit
- found efficient algorithm for evaluation of R_6 using single-valuedness
- proved all-loop LLA conjecture of [Pennington]
- first steps towards seven-gluon case, full solution for n-gluon LLA in [Del Duca et al.]

next steps:

- simple structure in momentum space beyond LLA?
- n-point amplitudes beyond LLA?
- beyond MRL?

