NLO+PS predictions for top-pair and Wt production and decay

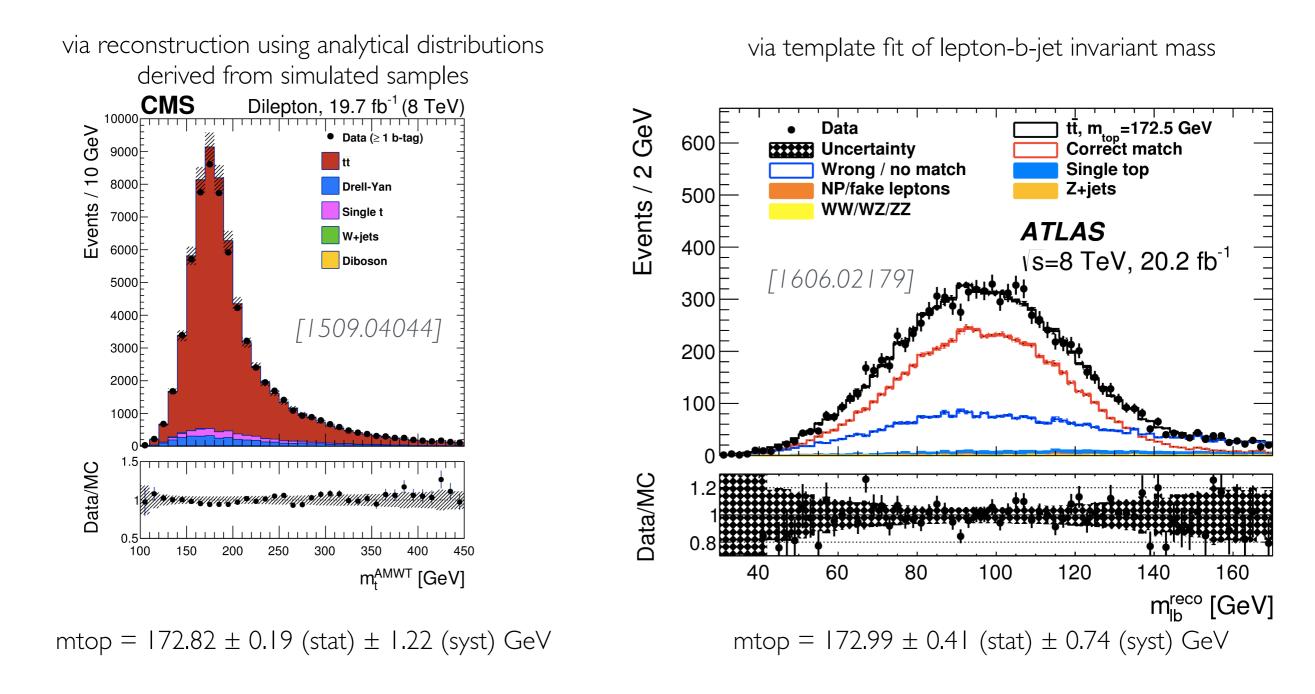


Jonas M. Lindert

work in collaboration with: T. Ježo, P. Nason, C. Oleari, S. Pozzorini based on [Ježo, Nason; '15] & [Ježo, JML, Nason, Oleari, Pozzorini; '16]

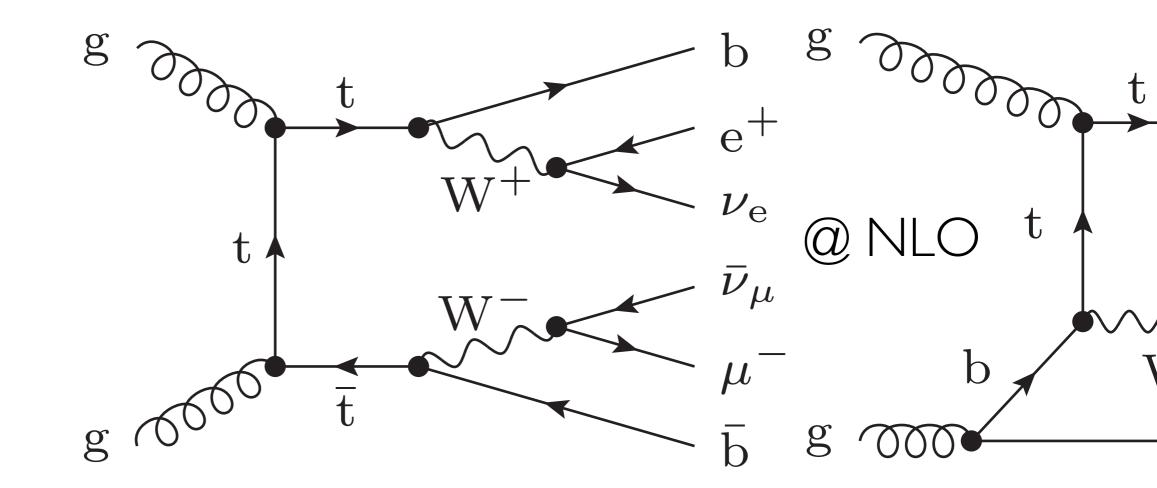
DESY Theory Workshop DESY, Hamburg, 29th September 2016

precision top-quark mass measurements rely on Monte Carlo modelling

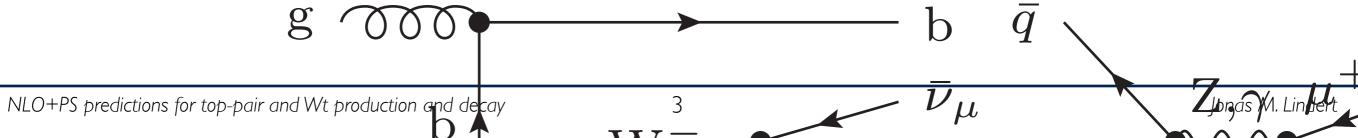


• these kinematic measurements strongly rely on MC modelling!

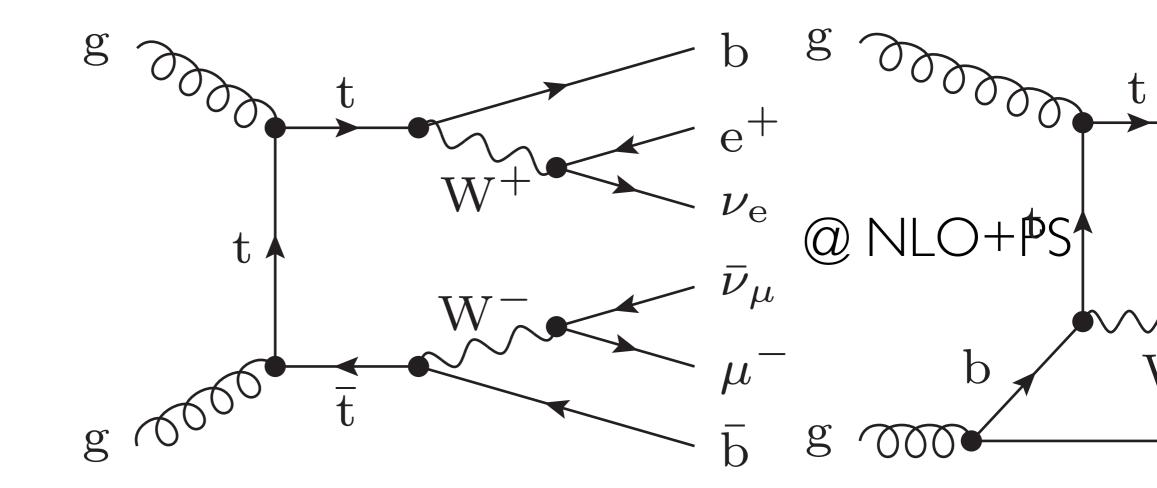
top-pair production and decay



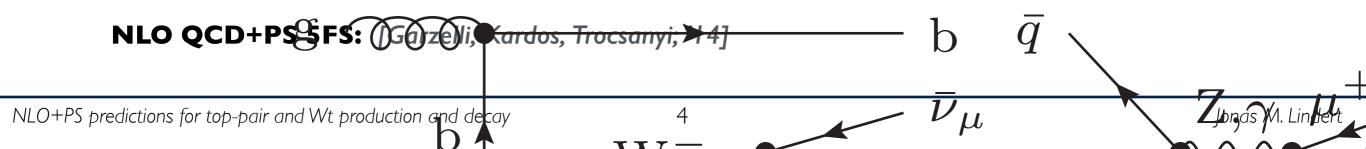
 NLO QCD 5FS: [Bevilacqua, Czakon, van. Hameren, Papadopoulos, Worek; '11] [Denner, Dittmaier, Kallweit, Pozzorini; '11+'12] [Heinrich, Maier, Nisius, Schlenk, Winter; '14]
 4FS: [Frederix '14] [Cascioli, Kallweit, Maierhöfer, Pozzorini; '14]
 NLO EW 5FS: [Denner, Pellen '16]



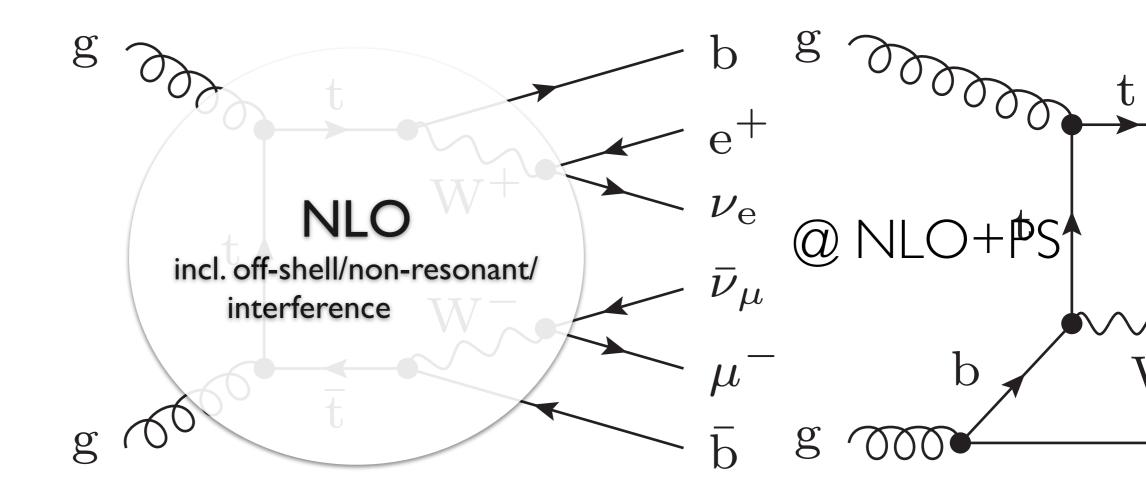
top-pair production and decay



NLO QCD 5FS: [Bevilacqua, Czakon, van. Hameren, Papadopoulos, Worek; '11]
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 4FS: [Frederix '14] [Cascioli, Kallweit, Maierhöfer, Pozzorini; '14]
 NLO EW 5FS: [Denner, Pellen '16]



top-pair production and decay



• In a traditional off-shell NLO+PS calculation:

subtraction, matching and PS do not see/preserve intermediate resonances

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Lingas M. Ling

- ➡ NLO: efficiency problem
- $\rightarrow N_{\rm L}O + RS$ distortion of important kinematic shapes! \bar{q}

5

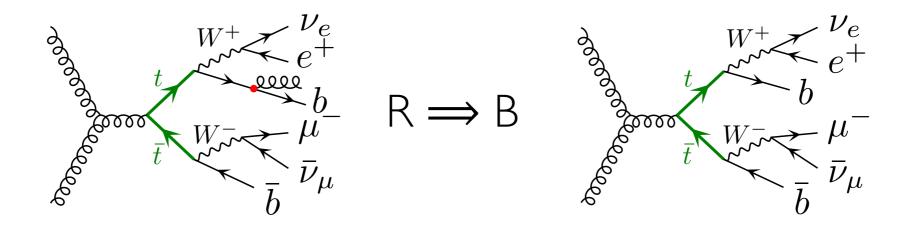
Problem in POWHEG language

• Efficiency problem at NLO:

- FKS subtraction (and similar CS) does not preserve virtuality of intermediate resonances
- i.e. Real (R) and Subtraction-term (S \sim B) with different virtualities of intermediate resonances

$$(\Phi_{\rm B}, \Phi_{\rm rad}) \stackrel{\rm FKS}{\longleftrightarrow} \Phi_{\rm R}$$

- $\Phi_{\rm R} = {\rm real \ phase-space}$
- $\Phi_{\rm B} = (\text{underlying})$ Born phase-space
- $\Phi_{\rm rad} = {\rm radiation \ phase-space}$



• IR cancellation spoiled

\Rightarrow severe efficiency problem for narrow resonances!

Problem in POWHEG language

• Fundamental problems at NLO+PS:

• **NLO+PS** matching according to POWHEG formula:

$$d\sigma = \bar{B}(\Phi_{\rm B}) d\Phi_{\rm B} \left\{ \Delta(q_{\rm cut}) + \Delta(k_{\rm T}) \frac{R(\Phi_{\rm R}(\Phi_{\rm B}, \Phi_{\rm rad}))}{B(\Phi_{\rm B})} d\Phi_{\rm rad} \right\}$$
$$\bar{B}(\Phi_{\rm B}) = B(\Phi_{\rm B}) + V(\Phi_{\rm B}) + \int R(\Phi_{\rm R}(\Phi_{\rm B}, \Phi_{\rm rad})) d\Phi_{\rm rad}$$
$$\Delta(q) = \exp\left[-\int_{k_{\rm T}>q} \frac{R(\Phi_{\rm R}(\Phi_{\rm B}, \Phi_{\rm rad}))}{B(\Phi_{\rm B})} d\Phi_{\rm rad} \right]$$

→ FKS mappings $(\Phi_B, \Phi_{rad}) \stackrel{FKS}{\longleftrightarrow} \Phi_R$ do not preserve intermediate resonances

- ⇒ R/B might become large by coincidence (e.g. R on-peak, B off-peak), but should be small (of the order of α_S) or approach Altarelli-Parisi splitting functions.
- ➡ Radiation (incl. Sudakov form-factor) generated from uncontrollable R/B ratios
- also subsequent radiation by the **PS** itself reshuffles internal momenta and does in general not preserve the virtuality of intermediate resonances.

\Rightarrow expect uncontrollable distortion of important kinematic shapes!

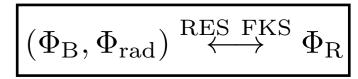
Resonance aware POWHEG

Rigorous solution to all these issues within POWHEG according to [Ježo, Nason; '15]

Idea: preserve invariant mass of intermediate resonances at all stages!

✓ NLO:

- Split phase-space integration into regions dominated by a single **resonance history**
- within a given resonance history **modify FKS mappings**, such that they *always* preserve intermediate resonances



 \Rightarrow R and S~B by construction with same virtuality of intermediate resonances

\Rightarrow IR cancellation restored

✓ NLO+PS:

• R and B related via resonance-aware FKS mappings

 \Rightarrow **R/B** ratio with fixed virtuality of intermediate resonances

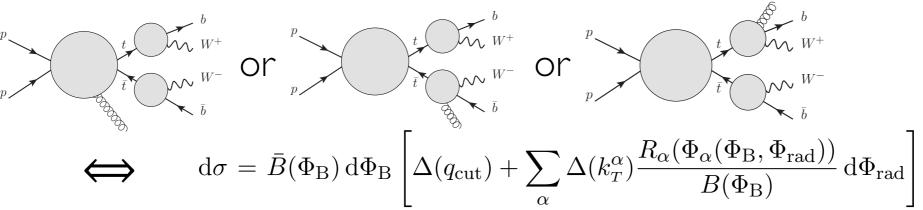
 \Rightarrow Sudakov form-factor preserves intermediate resonances

√ PS:

- pass information about resonance histories to the shower (via extension of LHE)
- tell **PS to respect intermediate resonances** (available in Pythia8)

Multiple-radiation scheme

▶ In traditional approach only hardest radiation is generated by POWHEG:

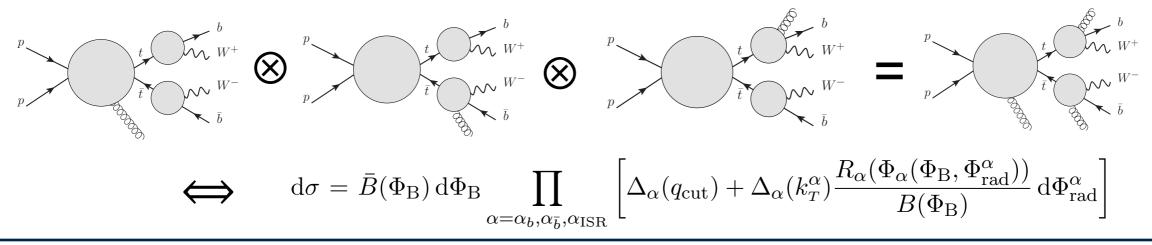


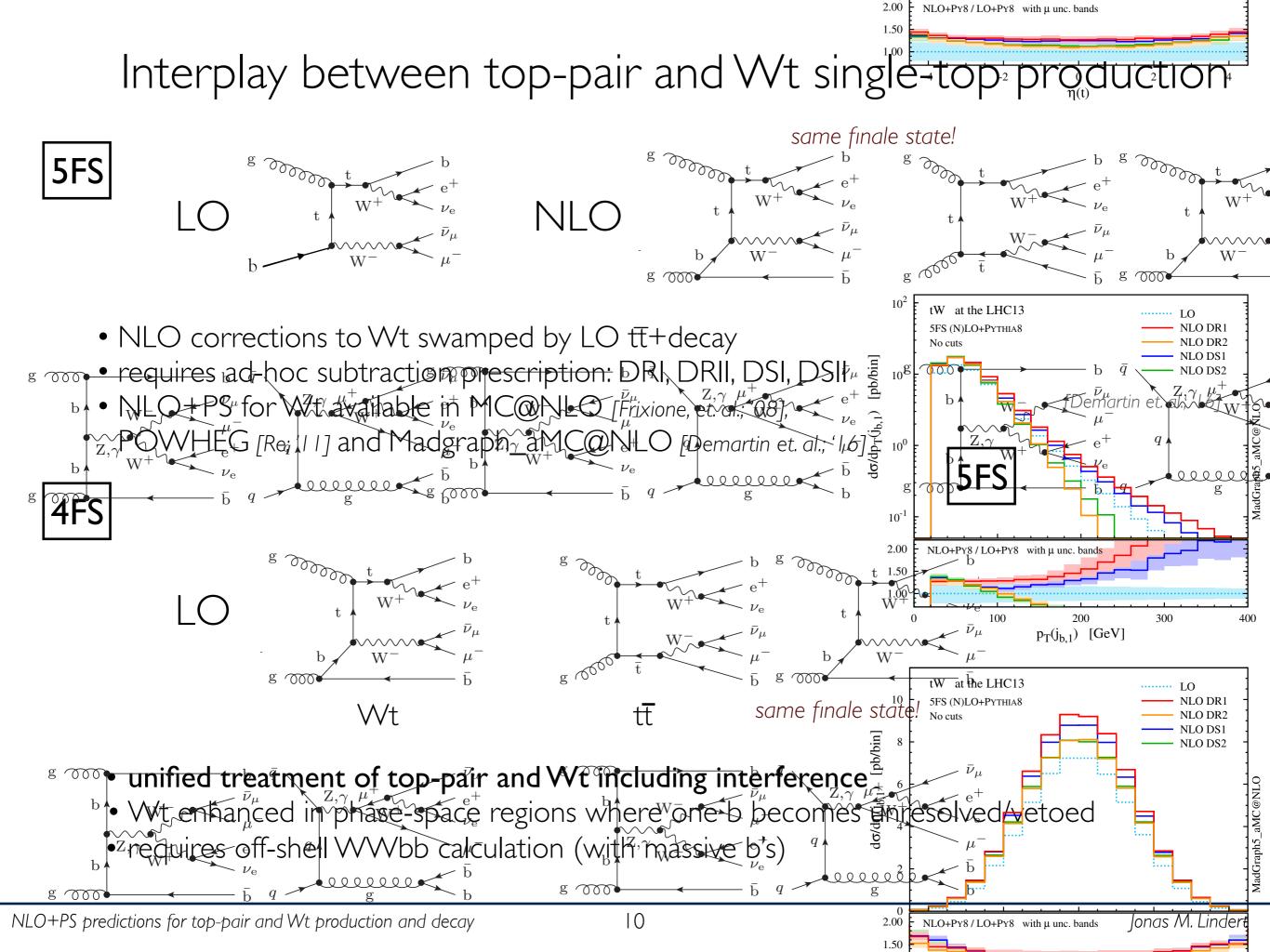
BUT: for top-pair (or single-top) production and decay, emission from production is almost always the hardest.

- ➡ emission off decays are mostly generated by the shower.
- Multiple-radiation scheme:

introduced in [Campbell, Ellis, Nason, Re; 'I 5]

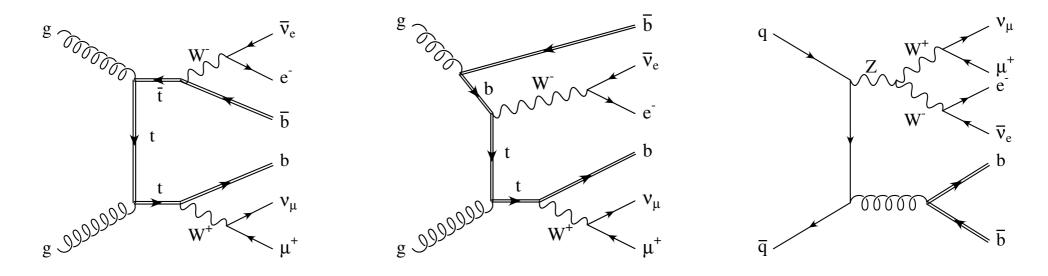
- keep hardest emission from all resonance histories.
- merge emissions into a single radiation event with several radiated partons





The new bb4l generator

- We consider the full process $pp \to b\bar{b}e^+\nu_e\mu^-\bar{\nu}_\mu$ with massive b's (**4FS scheme**)
- Implemented in the POWHEG-BOX-RES framework
- All matrix elements from **OpenLoops** [Cascioli, JML, Maierhöfer, Pozzorini]

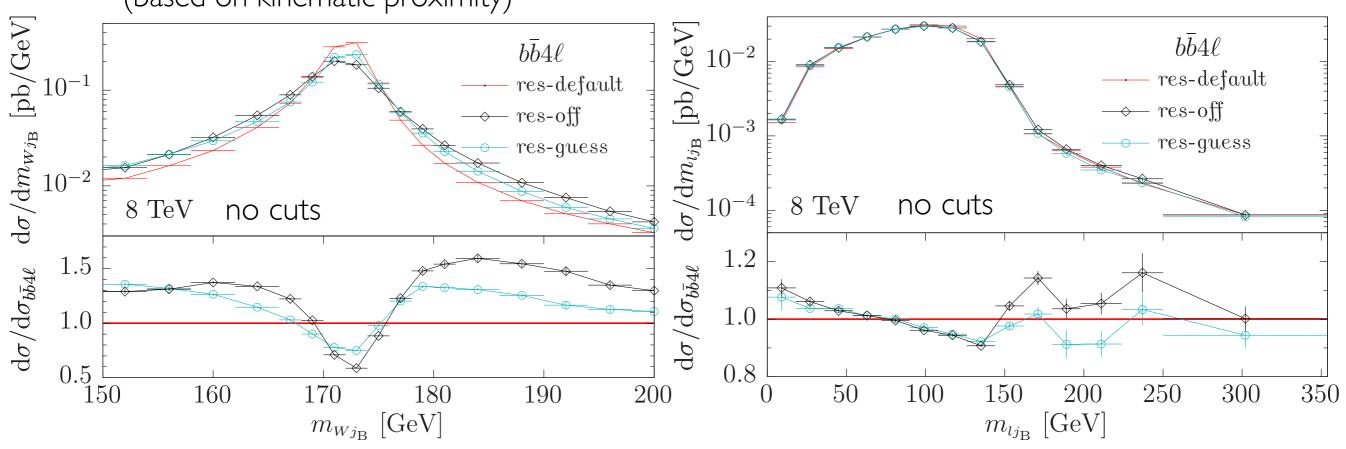


Physics features:

- exact non-resonant / off-shell / interference / spin-correlation effects at NLO
- unified treatment of **top-pair and Wt** production with interference at NLO
- access to phase-space regions with **unresolved b-quarks** and/or jet vetoes
- consistent NLO+PS treatment of top resonances, including quantum corrections to top propagators and off-shell top-decay chains

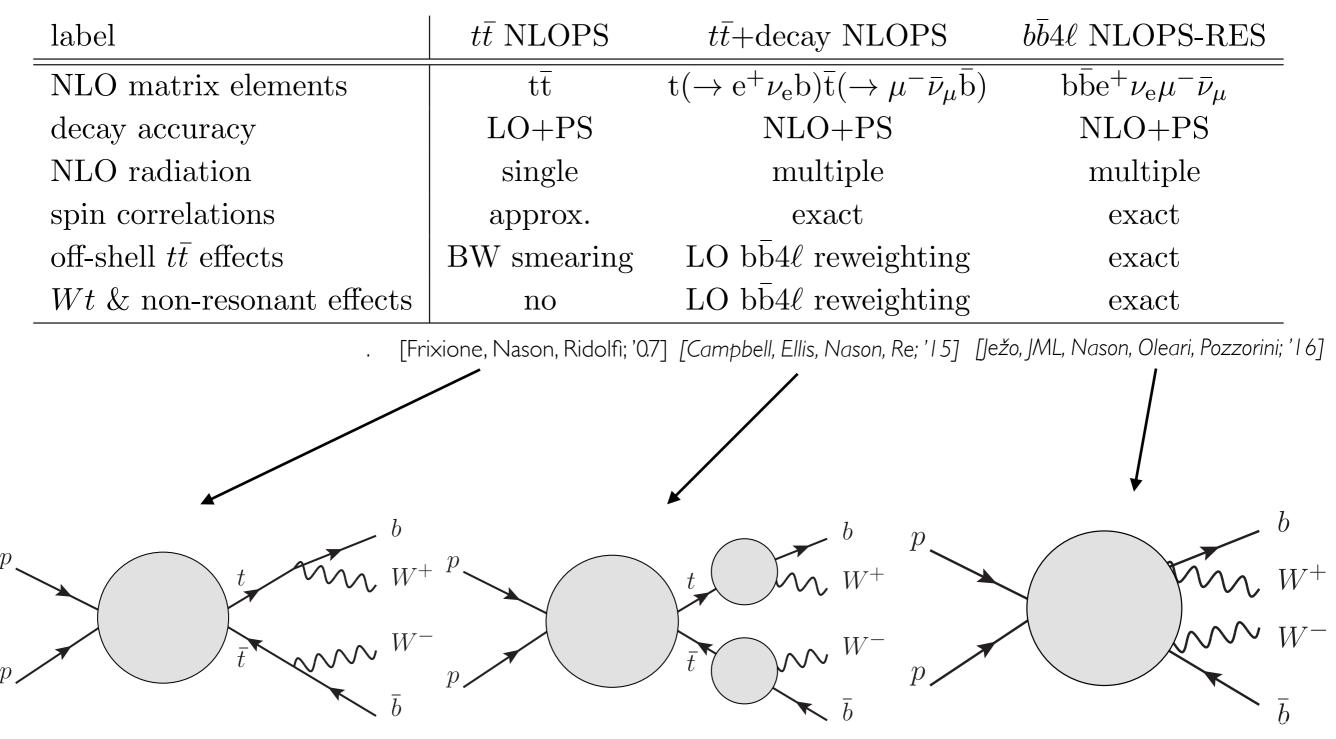
Results: top-resonance

- In default: resonance aware matching & multiple-radiation scheme
- ▶ off: resonance unaware matching
- guess: resonance unaware matching but kinematic guess off resonance structure before PS (based on kinematic proximity)

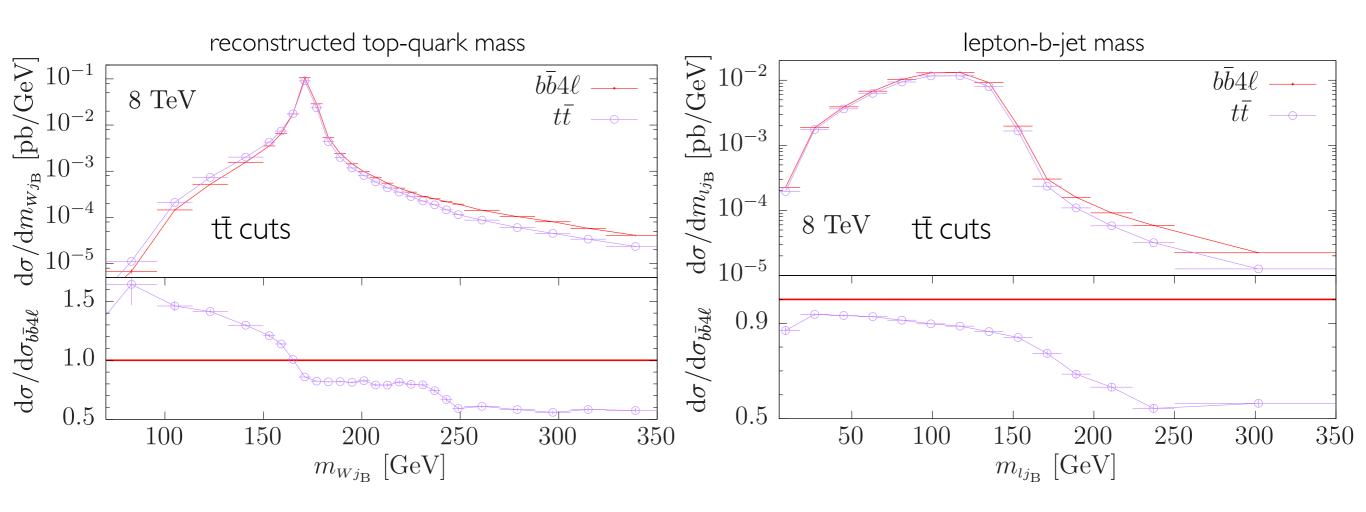


- \Rightarrow resonance unaware matching yields distortions of important kinematic shapes
- \Rightarrow control of these shapes crucial for **precise top-mass measurements**!
- \Rightarrow resonance assignment based on kinematic proximity with standard matching not sufficient

Top-pair production in POWHEG

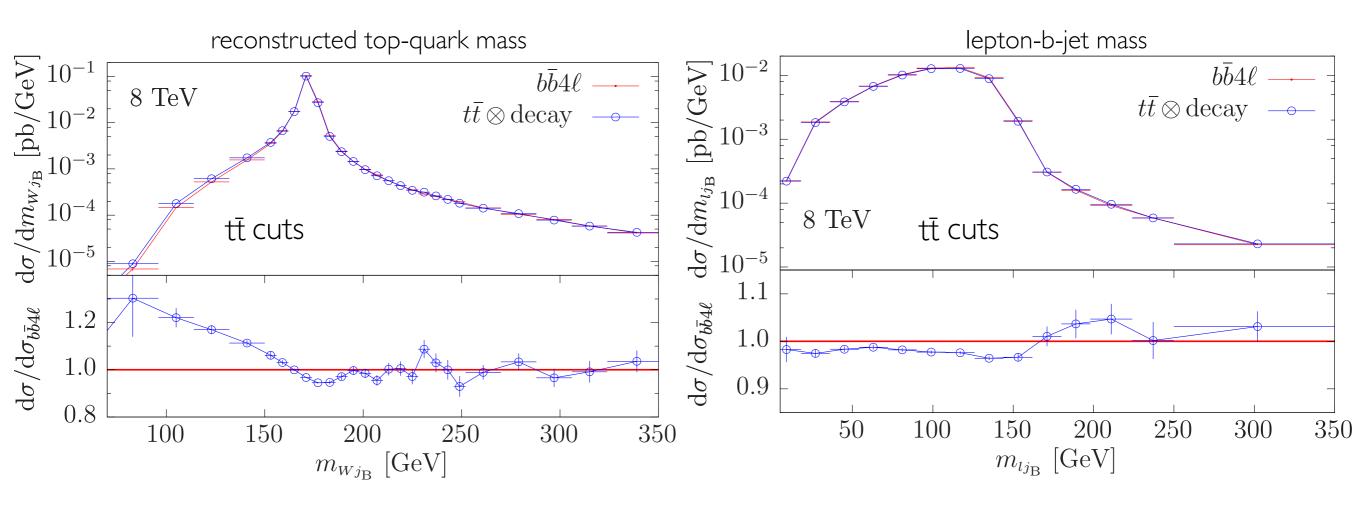


Results: on-shell tt vs. bb4l



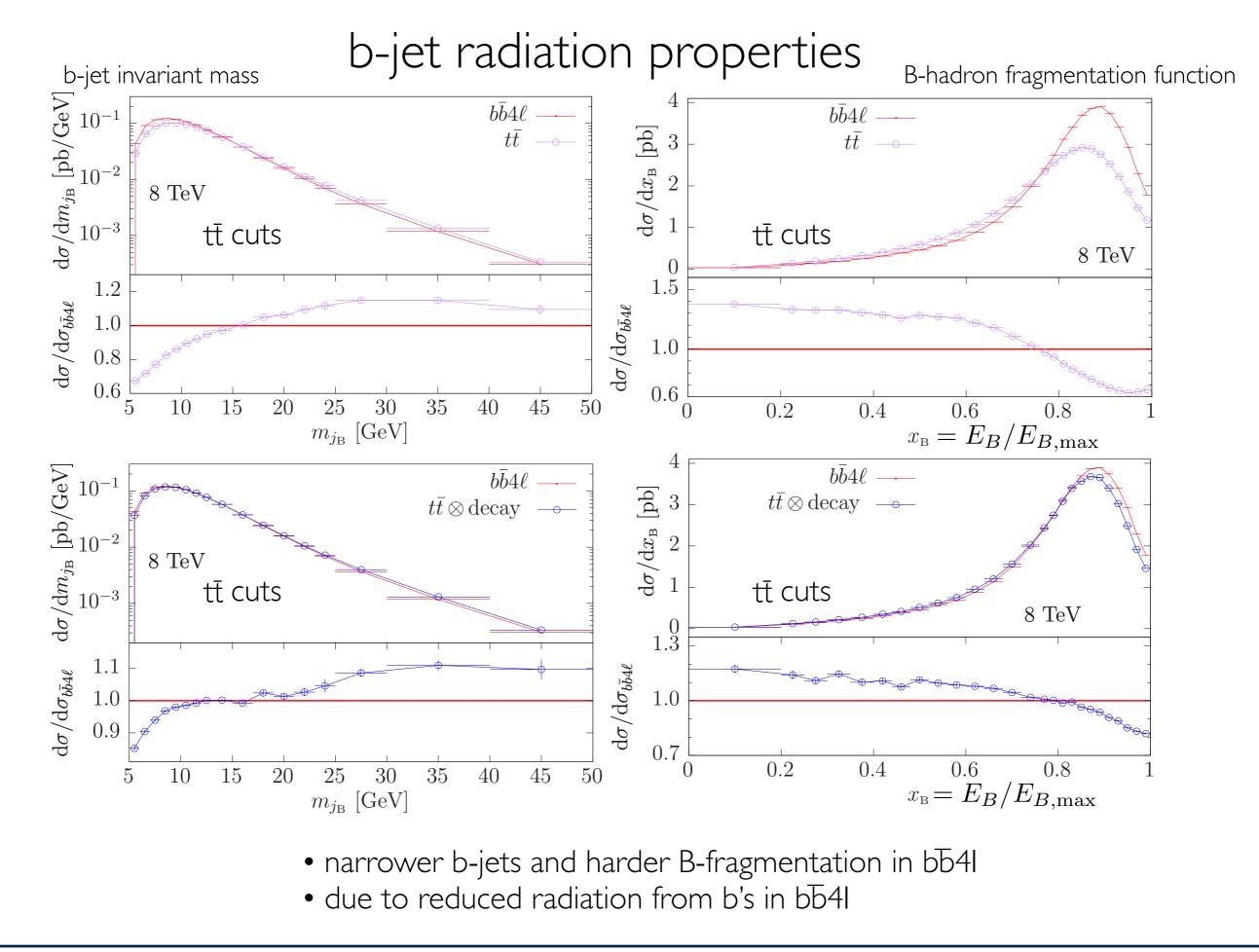
- significant shape distortions around resonance with respect to on-shell calculation
- very relevant for top mass determination
- average m_{Wj_B} roughly 500 MeV smaller in on-shell tt (in ±30 GeV around mtop)
- \sim 20-30% effects around the b-jet-lepton invariant mass edge

Results: tī⊗decay vs. bb4l

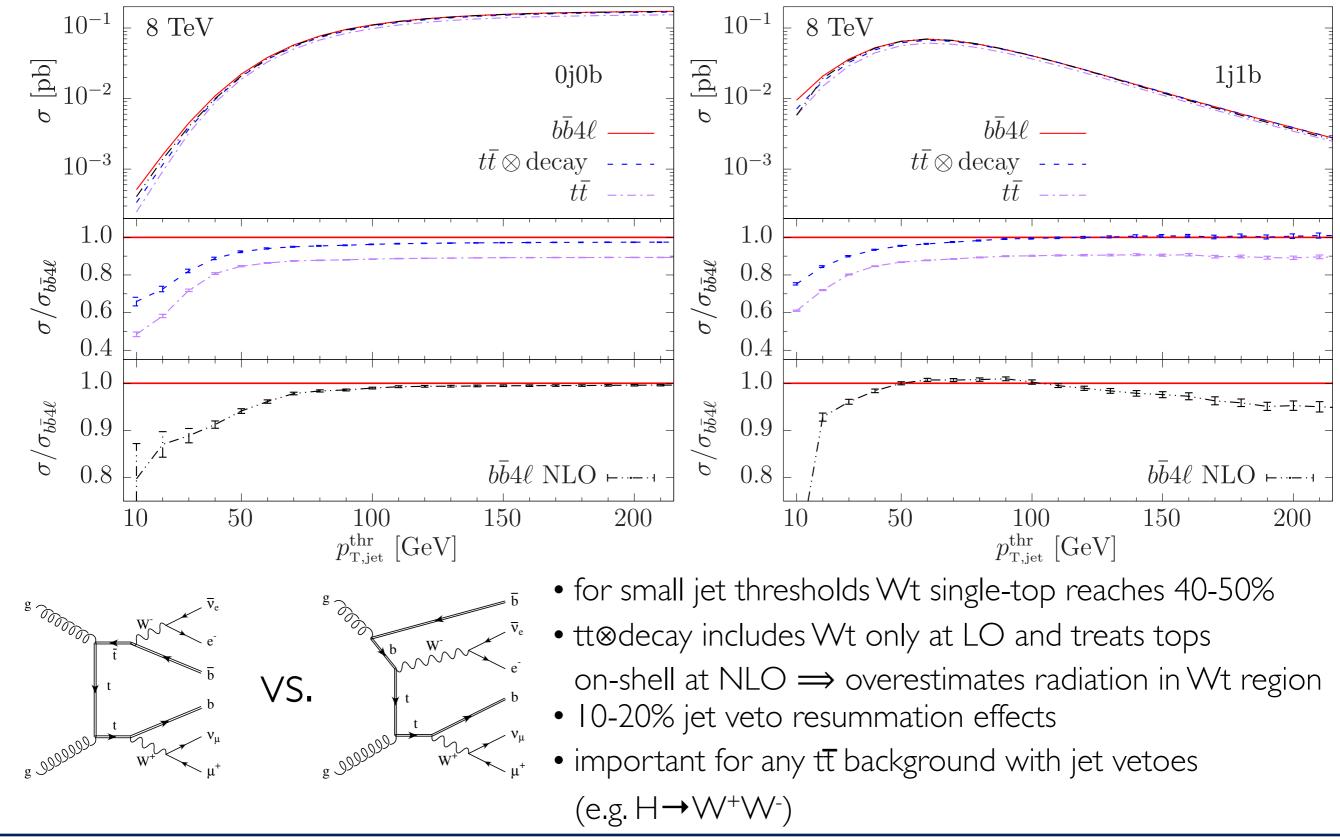


• very good agreement mostly <5% level between the two predictions

- the two calculations support each other (natural factorization of radiation between production and decay in tt⊗decay)
- average m_{Wj_B} roughly 100 MeV smaller in tt \otimes decay (in ±30 GeV around mtop)



jet vetoes and single-top enriched observables



Conclusions

Technology:

- Resonance-aware matching is pivotal for processes with intermediate resonances
 - ➡ (quite) Rigorous solution within POWHEG by [Ježo, Nason; '15]
- New POWHEG framework: POWHEG-BOX-RES (<u>http://powhegbox.mib.infn.it/</u>)
 - resonance-aware subtraction and matching for any process
 - automated generation of resonance histories and phase-space
 - process independent POWHEG-BOX+OpenLoops interface

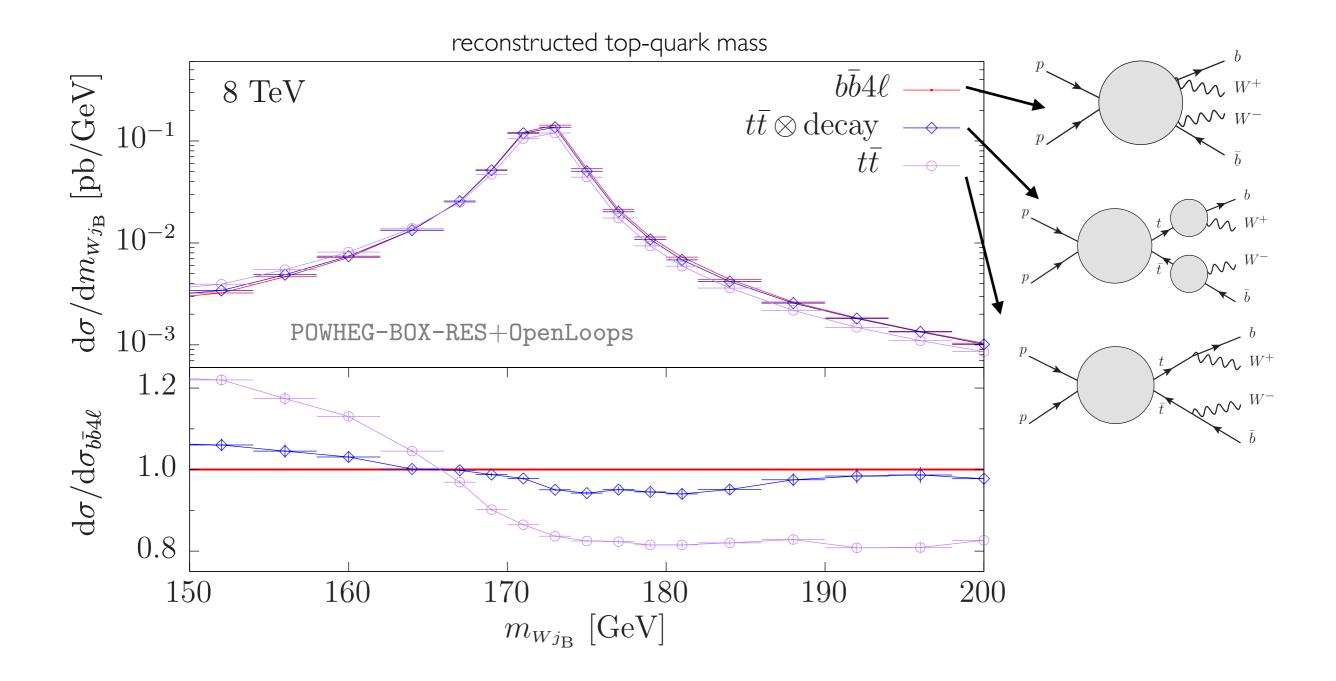
Phenomenology:

- ▶ resonance-aware matching crucial for kinematic precision top-mass measurements
- unified treatment of tt & Wt important for precision single-top physics
-and for modelling of tT backgrounds subject to jet vetoes.

Outlook:

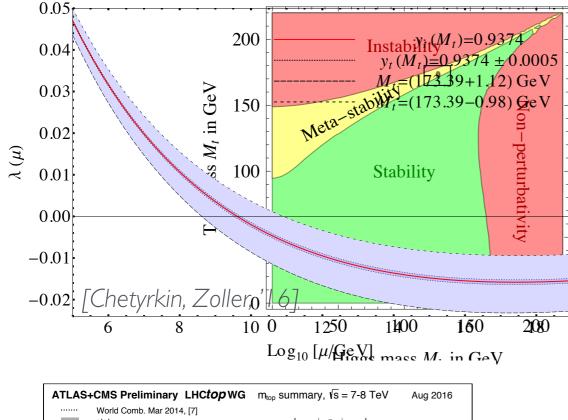
- Detailed investigation of effects on top mass measurements
- Hadronic top decays

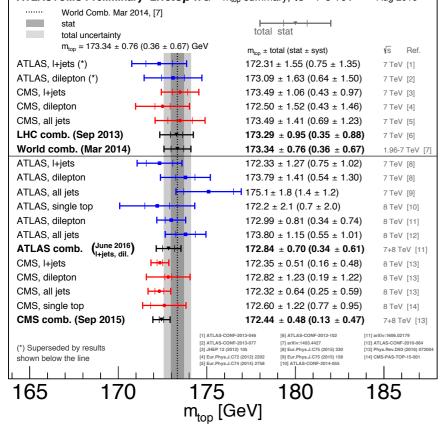
Summary of the results

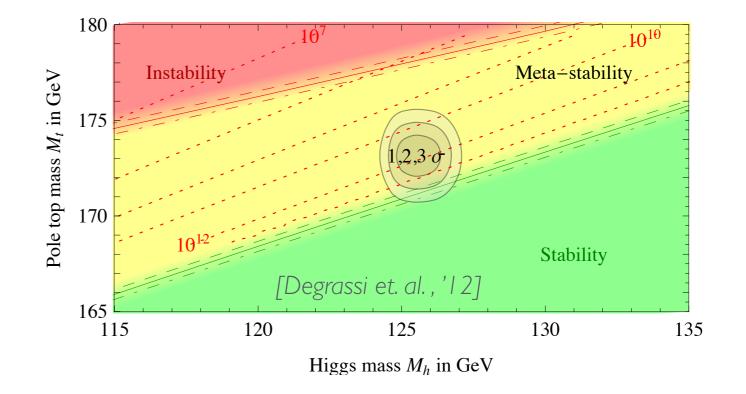


Backup slides

top-quark mass: crucial SM input parameter

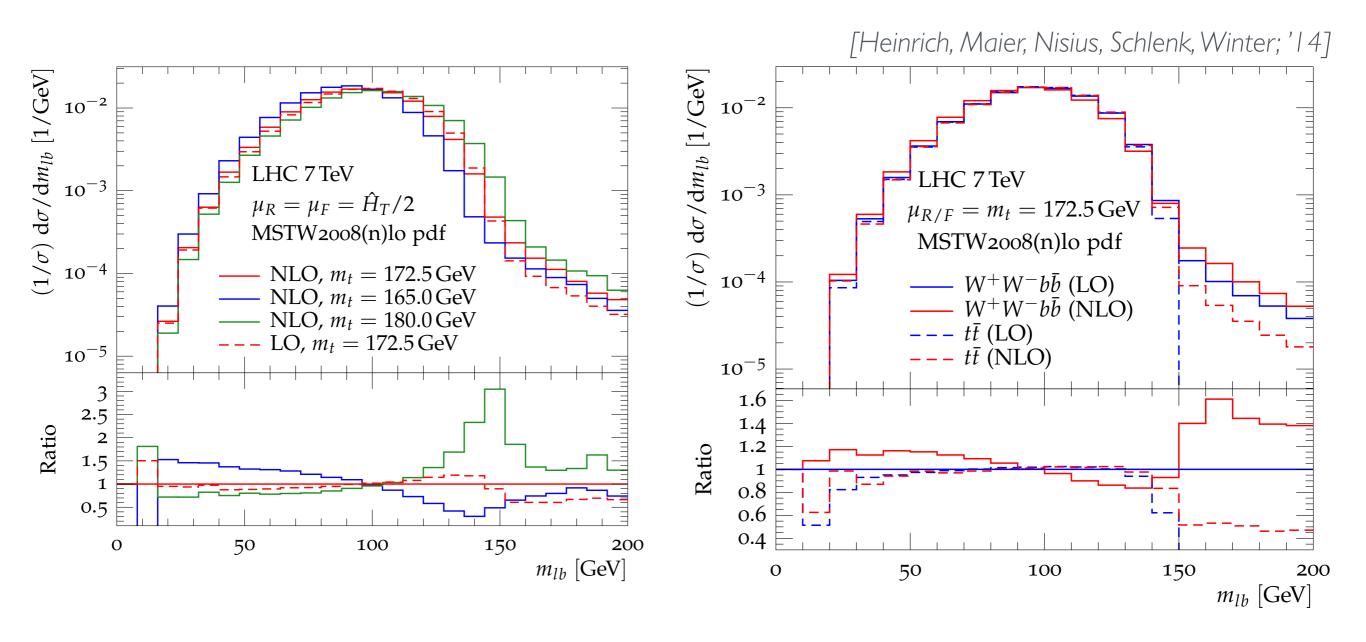






- top-mass measurements at the LHC via combination of different strategies:
 - total x-section, tt+jet, kinematic reconstruction, kinematic edges,....
- many techniques rely on kinematic information of top decay products
 - ➡ need realistic MC modelling

kinematic top-quark mass measurements: theory



- m_{lb} is straightforward to measure and shows strong dependence on top mass
- off-shell and NLO effects crucial for precise modelling of shape at the edge

Setup

| $m_Z = 91.188 \text{ GeV},$ | $\Gamma_Z = 2.441 \text{ GeV},$ | $G_{\mu} = 1.16585 \times 10^{-5}$ |
|-----------------------------|---|------------------------------------|
| $m_W = 80.419 \text{ GeV},$ | $\Gamma_W = 2.048 \text{ GeV},$ | F~ |
| $m_H = 125 \text{ GeV},$ | $\Gamma_H = 4.03 \times 10^{-3} \text{ GeV},$ | |
| $m_t = 172.5 \text{ GeV},$ | $\Gamma_t = 1.329 \text{ GeV},$ | |
| $m_b = 4.75 {\rm GeV}$. | | |

Complex-mass-scheme: $\mu_i^2 = M_i^2 - i\Gamma_i M_i$ for i = W, Z, t, H $\sin \theta_W^2 = 1 - \cos \theta_W^2 = 1 - \frac{\mu_W^2}{\mu_Z^2}$ For t resonance histories: $\mu_R = \mu_F = \left[\left(m_t^2 + p_{T,t}^2 \right) \left(m_t^2 + p_{T,\bar{t}}^2 \right) \right]^{\frac{1}{4}}$ For Z resonance histories: $\mu_R = \mu_F = \frac{\sqrt{p_Z^2}}{2}$

PDFs: MSTW2008NLO

tī cuts: at least one b- and one b-jet with $p_T^j > 30 \text{ GeV},$ $|\eta^j| < 2.5$ $p_T^l > 20 \text{ GeV},$ $|\eta^l| < 2.5,$ $p_T^{\text{miss}} > 20 \text{ GeV}$

 GeV^{-2}

Resonance histories

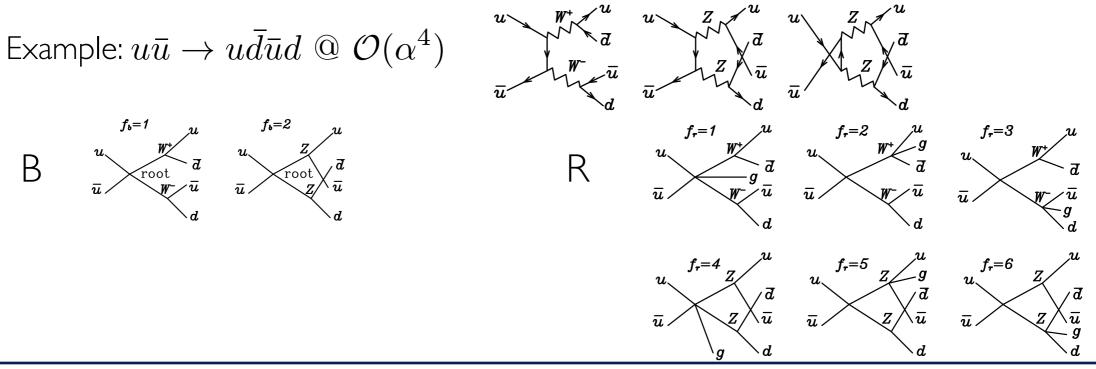
This approach is rigorous up to the point that assignment of resonance histories requires a prescription.

Idea: define resonance histories according to well defined NWA limit (valid to all orders)!

Project onto these resonance histories f_b and f_r based on kinematic proximity:

$$B_{F_b} = \sum_{f_b \in T(F_b)} B_{f_b}, \qquad B_{f_b} = \prod_{f_b} B_{F_b}$$
$$\Pi_{f_b} = \frac{P^{f_b}}{\sum_{f'_b \in T(F_b(f_b))} P^{f'_b}}, \qquad P^{f_b} = \prod_{i \in \mathrm{Nd}(f_b)} \frac{M_i^4}{(s_i - M_i^2)^2 + \Gamma_i^2 M_i^2}$$

(similar for R: separation into resonance structures and *compatible* FKS singular regions)



Efficiency study

| | | resonance aware | resonance unaware |
|---------------------------------------|-------------------|--------------------|----------------------|
| NLO cross section | rel. accuracy (*) | 0,11% | 0,79% |
| efficiency of generation of radiation | vetos per event | 750 | 15000 |
| speed of event generation | events per hour | 1500 | 200 |

 \Rightarrow factor of ~7 improvement in convergence/efficiency/speed!

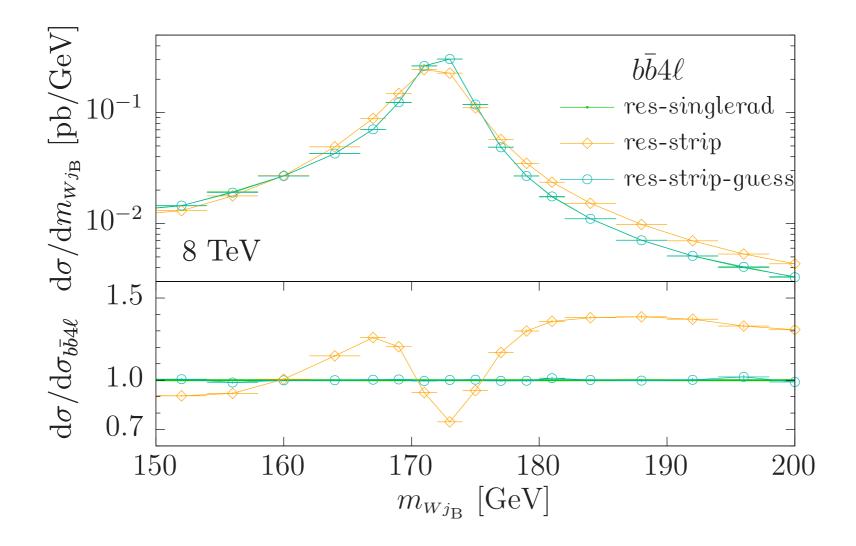
(*) NLO POWHEG setup

- stage 1:ncalls=80k, itmx=2
- stage 2: ncalls=100k, itmx=4
- nrun = 64

(typical setup for small cluster/blade)

Kinematic guess

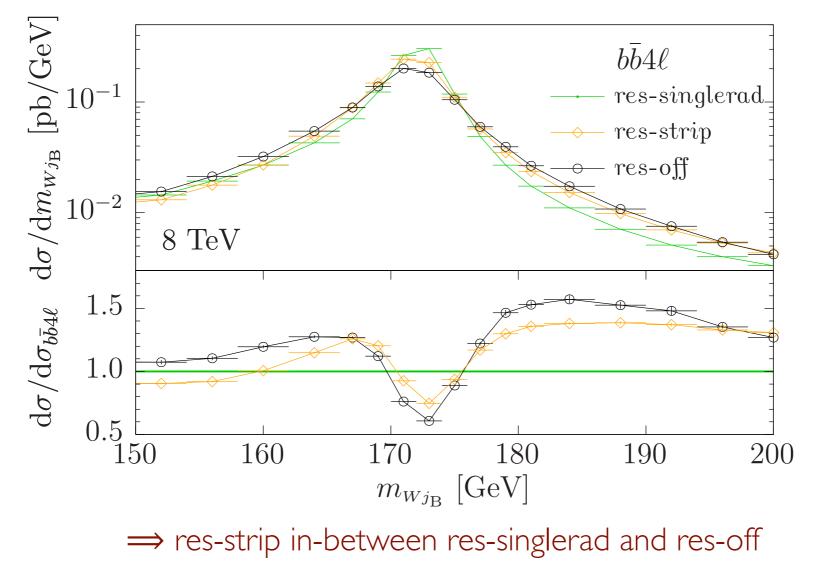
- res-singlerad: resonance aware matching & single-radiation scheme
- ▶ res-strip: resonance aware matching, but resonance information not passed to PS
- res-strip-guess: resonance aware matching, resonance information first stripped and then guessed (based on kinematic proximity) before passing to PS



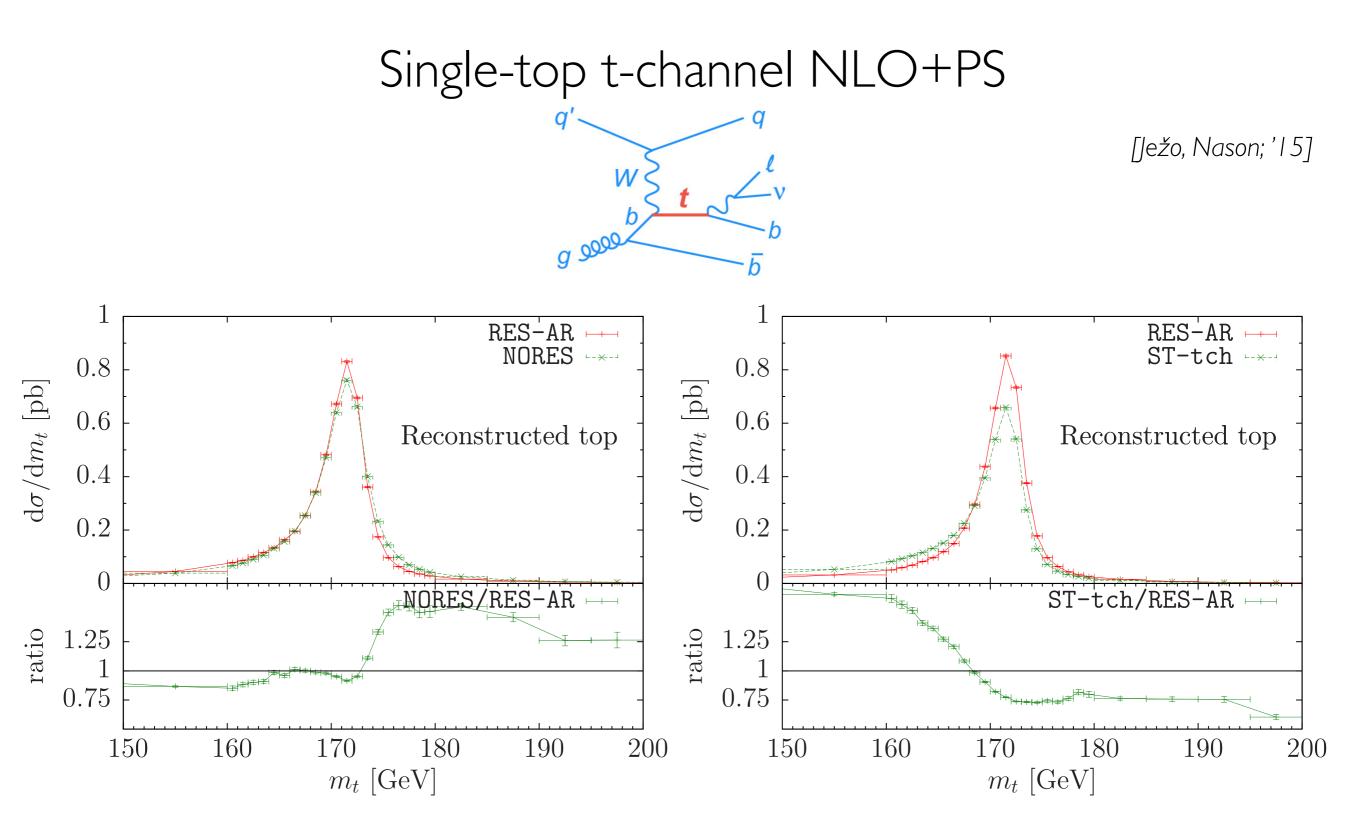
Impact of PS momentum reshuffling

- res-singlerad: resonance aware matching & single-radiation scheme
- ▶ res-strip: resonance aware matching, but resonance information not passed to PS

▶ res-off: resonance unaware matching

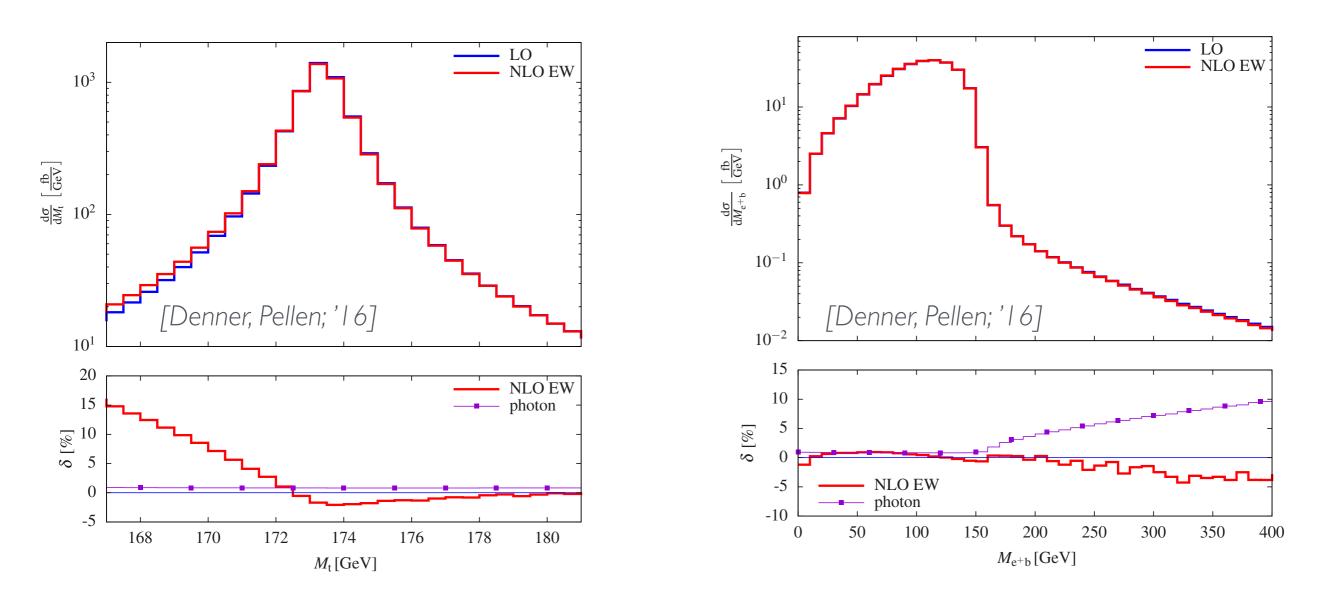


- \Rightarrow both effects important:
 - I) first emission governed by resonance preserving R/B
 - II) PS reshuffling preserves the resonance masses



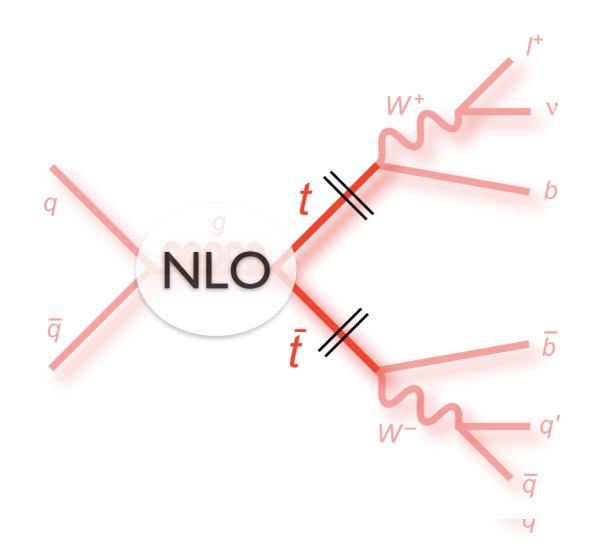
• significant shape distortions in resonance unaware calculation and with respect to on-shell top calculation [Alioli, Nason, Oleari, Re; '09]

NLO EW



- NLO EW ~5% around top peak, up to 15% below the peak
- sizeable photon-induced contributions in m_{lb} (using NNPDF2.3QED)

Top-pair production and decay

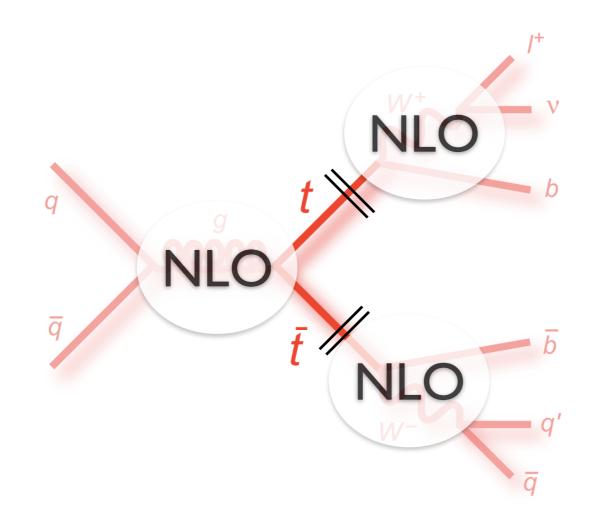


NLO QCD [Bernreuther, Brandenburg, Si; '04, Melnikov, Schulze;'09, Campbell, Ellis; '15] [Bevilacqua, Czakon, van. Hameren, Papadopoulos, Worek; '11, Denner, Dittmaier, Kallweit, Pozzorini; '11+'12, Heinrich, Maier, Nisius, Schlenk, Winter; '14, Frederix '14, Cascioli, Kallweit, Maierhfer, Pozzorini; '14]

NNLO [Czakon, Fiedler, Mitov; '13, Czakon, Heymes, Mitov; '16]

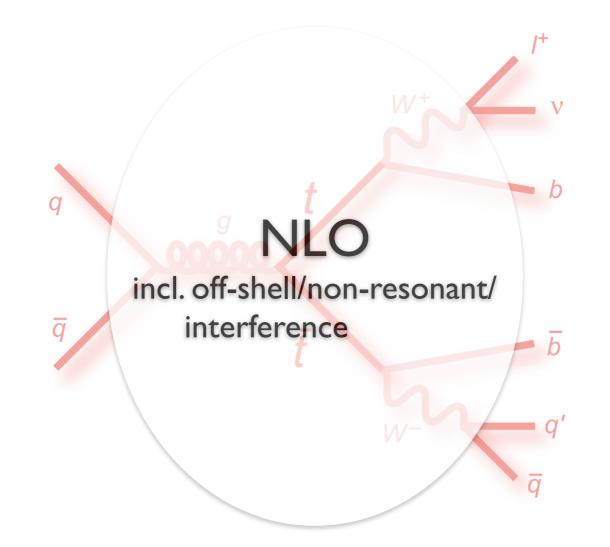
NLO EW [Beenakker, Denner, Hollik, Mertig, Sack, Wackeroth; '94, Bernreuther, Fuecker, Si; '06+'08, Kühn, Scharf, Uwer; '07,+'15, Hollik, Pagani; '07, Pagani, Tsinikos, Zaro,; '16] [Bernreuther, Si; '10] [Denner, Pellen '16]
 NLO QCD+PS [Frixione, Nason, Webber, '03, Frixione, Nason, Ridolfi; '07] [Campbell, Ellis, Nason, Re, '15]
 [Garzelli, Kardos, Trocsanyi; '14]

Top-pair production and decay



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Top-pair production and decay



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