

OFF-SHELL TOP QUARKS WITH A JET: A COMPREHENSIVE STUDY AT NLO QCD

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RWTHAACHEN
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DFG

based on

[PRL 116 (2016) 052003]

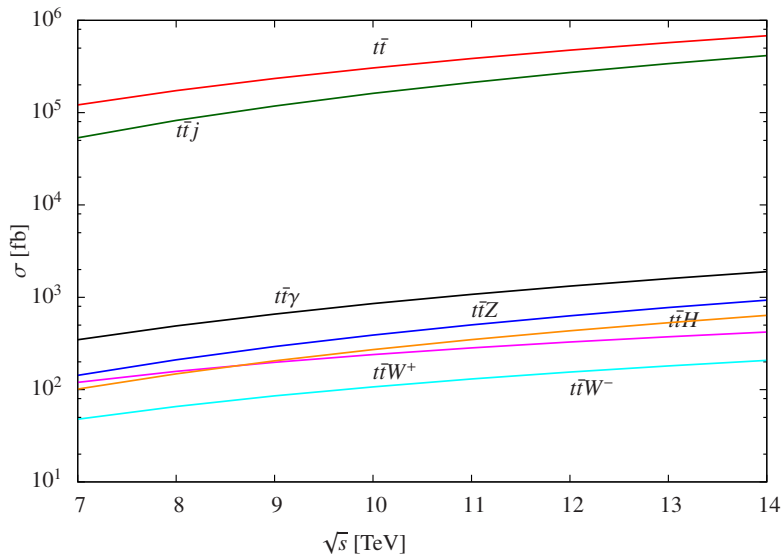
[arXiv:1609.01659]

Next-to-Leading Order QCD predictions for

$$pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} j + X$$

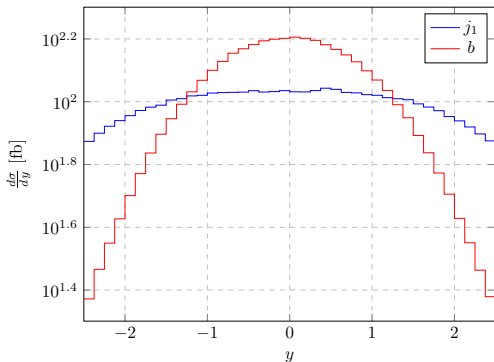
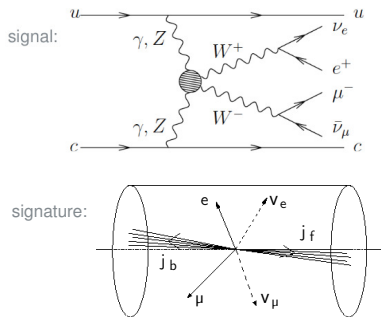
1. Motivation for $t\bar{t}j$ process
2. Off-shell effects in top-quark pair production
3. Complete off-shell effects with HELAC-NLO for $t\bar{t}j$
4. Conclusions

ASSOCIATED $t\bar{t}$ PRODUCTION



$t\bar{t}j$ AS A BACKGROUND

- Background for SM Higgs production in VBF: $qq \rightarrow Hqq \rightarrow W^+W^-qq$
- VBF requires two tagged jets: $\Delta y_{ij} = |y_i - y_j| > 4$ and $y_i \times y_j < 0$



also background for BSM searches with $\ell_1^+ \ell_2^- + MET + jets$ signature

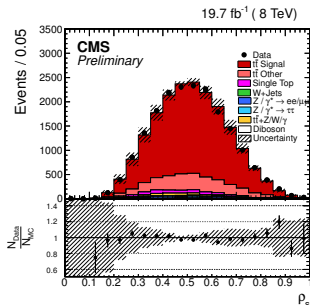
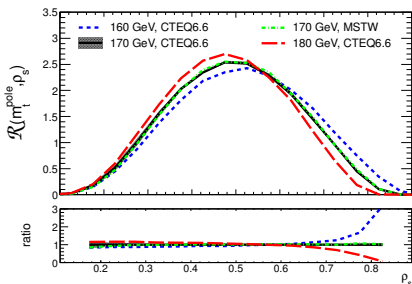
An alternative method for the top-quark mass extraction

[Alioli et al '13]

$$\mathcal{R}(m_t, \rho_s) = \frac{1}{\sigma_{t\bar{t}j}} \frac{d\sigma_{t\bar{t}j}}{d\rho_s(m_t, \rho_s)}$$

$$\rho_s = \frac{2m_0}{\sqrt{s_{t\bar{t}j}}}$$

- $\mathcal{R}(m_t, \rho_s)$ shape is sensitive to m_t
- $t\bar{t}j$ has higher sensitivity than $t\bar{t}$
- 7 TeV: $m_t = (173.7 \pm 2.2)$ GeV [ATLAS '15]
- 8 TeV: $m_t = (169.9 \pm 3.9)$ GeV [CMS '16]

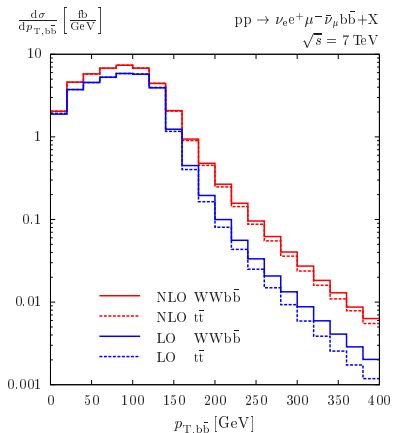


Theoretical and PDF uncertainties affect the m_t extraction

TOP-QUARK OFF-SHELL EFFECTS - I

Contributions dropped in the NWA are suppressed by $\Gamma_t/m_t \sim 1\%$.

- True for sufficiently inclusive observables
- Larger impact on differential distributions



For σ_{tot} at the % level ($t\bar{t}, t\bar{t}j, t\bar{t}H$)

[A. Denner et al '11, '12, '15]

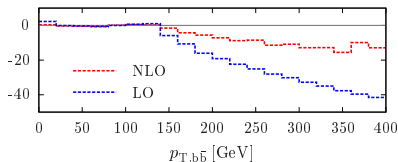
[G. Bevilacqua et al '11, '16]

[R. Frederix '14]

[F. Cascioli et al '14]

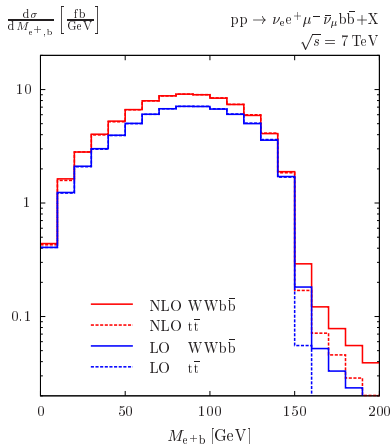
[G. Heinrich et al '14]

$t\bar{t}/WWb\bar{b} - 1$ [%]



[A. Denner, S. Dittmaier, S. Kallweit, S. Pozzorini, M. Schulze '12]

TOP-QUARK OFF-SHELL EFFECTS - II



[A. Denner, S. Dittmaier, S. Kallweit, S. Pozzorini, M. Schulze '12]

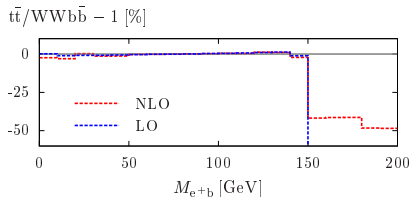
m_t measurement based on $M_{\ell b}$ \rightarrow m_t extraction is very sensitive to template
 (LO/NLO and Full/NWA).

[Heinrich, Maier, Nisius, Schlenk, Winter '14]

- If both top and W decay on-shell
 \rightarrow end-point given by sharp cut

$$M_{\ell b} = \sqrt{m_t^2 - m_W^2} \approx 152 \text{ GeV}$$

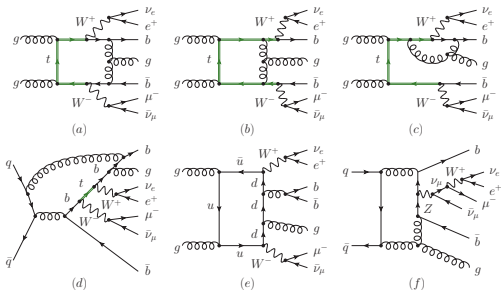
- Additional radiation & off-shell effects introduce smearing



TOP-QUARK OFF-SHELL EFFECTS IN $t\bar{t}j$

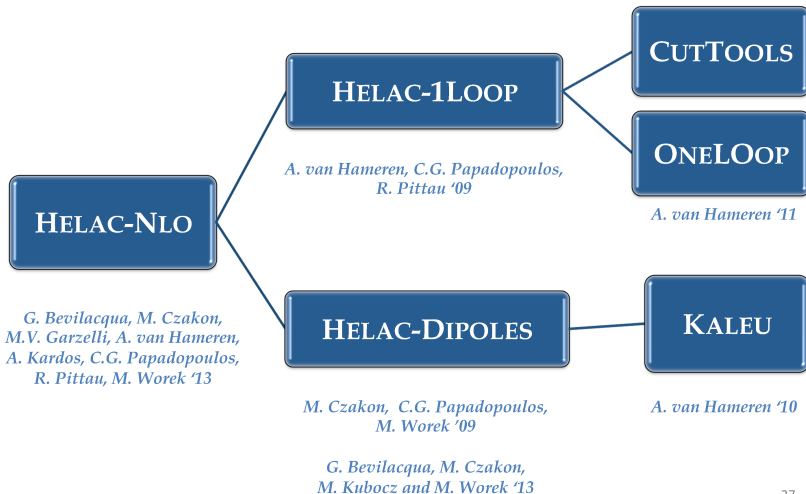
$$pp \rightarrow e^+ \nu_{e\mu}^- \bar{\nu}_\mu b\bar{b}j + X$$

- $t\bar{t}j$ production at $\mathcal{O}(\alpha_s^4 \alpha^4)$ in the di-lepton decay channel
- $2 \rightarrow 5$ QCD process ($W^+ W^- b\bar{b}j$)
- **gg initial state** is the most complicated production channel
 - LO: 508
 - Real: 4447
 - Virtual: 39180 \rightarrow 1155 hexagons and 120 heptagons



HELAC-NLO

*G. Ossola, C.G. Papadopoulos,
R. Pittau '08*



SETUP FOR LHC 13 TeV - I

A comprehensive study for off-shell $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} j + X$ production:

- fixed vs. dynamical scale, independent μ_R and μ_F variations
- PDF uncertainties (different PDFs, internal PDF uncertainty)

⇒ very **time consuming** and **CPU intensive**

⇒ Extend HELAC-NLO functionality to produce **NTuple**

Fixed scale: $\mu_0 = m_t$

Dynamical scale: $\mu_0 = E_T/2$ and $\mu_0 = H_T/2$, with

$$E_T = \sqrt{p_T^2(t) + m_t^2} + \sqrt{p_T^2(\bar{t}) + m_t^2}$$

$$H_T = \sum_i p_T(i) + \not{p}_T, \quad i = \{e^+, \mu^-, J_b, J_{\bar{b}}, j_1\}$$

PDF sets: CT14nlo, MMHT14, NNPDF3.0

Scale uncertainty from envelopes obtained from the following grid

$$\left(\frac{\mu_R}{\mu_0}, \frac{\mu_F}{\mu_0} \right) = \{(0.5, 0.5), (0.5, 1), (1, 0.5), (1, 1), (1, 2), (2, 1), (2, 2)\}$$

SETUP FOR LHC 13 TeV - II

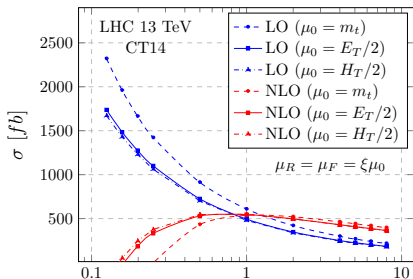
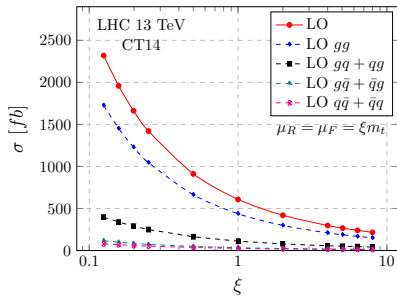
- **SM parameters:**

$$\begin{aligned} G_F &= 1.16637 \cdot 10^{-5} \text{ GeV}^{-2}, & m_t &= 173.2 \text{ GeV}, \\ m_W &= 80.399 \text{ GeV}, & \Gamma_W &= 2.09875 \text{ GeV}, \\ m_Z &= 91.1876 \text{ GeV}, & \Gamma_Z &= 2.50848 \text{ GeV}, \\ \Gamma_t^{LO} &= 1.47834 \text{ GeV}, & \Gamma_t^{NLO} &= 1.35146 \text{ GeV}, \end{aligned}$$

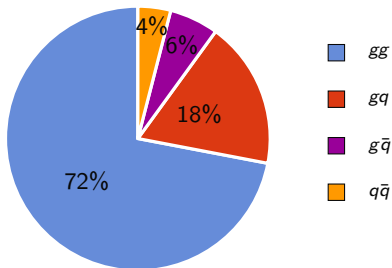
- Light quarks (also bottom) and leptons are **massless**
- Bottom induced channels are neglected (0.1% at LO)
- **Final state:** exactly 2 b-jets, at least 1 light jet,
2 charged leptons, missing p_T
- **Jets:** partons with $|\eta| < 5$, anti- k_T , $\Delta R = 0.5$
- **Cuts:**

$$\begin{aligned} p_{Tl} &> 30 \text{ GeV}, & p_{Tj} &> 40 \text{ GeV}, & \cancel{p}_T &> 40 \text{ GeV}, \\ \Delta R_{jj} &> 0.5, & \Delta R_{ll} &> 0.4, & \Delta R_{lj} &> 0.4, \\ |y_l| &< 2.5, & |y_j| &< 2.5, & & \end{aligned}$$

SCALE DEPENDENCE



LO contributions



$\mu_0 = m_t$	$K = 0.88$	-12%
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$\mu_0 = E_T/2$	$K = 1.10$	+10%
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$\mu_0 = H_T/2$	$K = 1.15$	+15%
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- Same PDF, different scales

$$\sigma^{NLO}(CT14, \mu_0 = m_t) = 537.24^{+10.12 (+2\%)}_{-190.35 (-35\%)} \text{ [scales] fb}$$

$$\sigma^{NLO}(CT14, \mu_0 = E_T/2) = 544.64^{+2.95 (+1\%)}_{-117.47 (-22\%)} \text{ [scales] fb}$$

$$\sigma^{NLO}(CT14, \mu_0 = H_T/2) = 549.65^{+10.25 (+2\%)}_{-53.42 (-10\%)} \text{ [scales] fb}$$

⇒ $\mu_0 = H_T/2$ yields smallest uncertainty $\sim 6\%$ (sym)

- Same scale, different PDFs

$$\sigma^{NLO}(CT14, \mu_0 = H_T/2) = 549.65^{+10.25 (+2\%)}_{-53.42 (-10\%)} \text{ [scales]}^{+18.00 (+3\%)}_{-19.15 (-3\%)} \text{ [PDF] fb}$$

$$\sigma^{NLO}(MMHT14, \mu_0 = H_T/2) = 554.61^{+10.85 (+2\%)}_{-54.51 (-10\%)} \text{ [scales]}^{+12.06 (+2\%)}_{-12.22 (-2\%)} \text{ [PDF] fb}$$

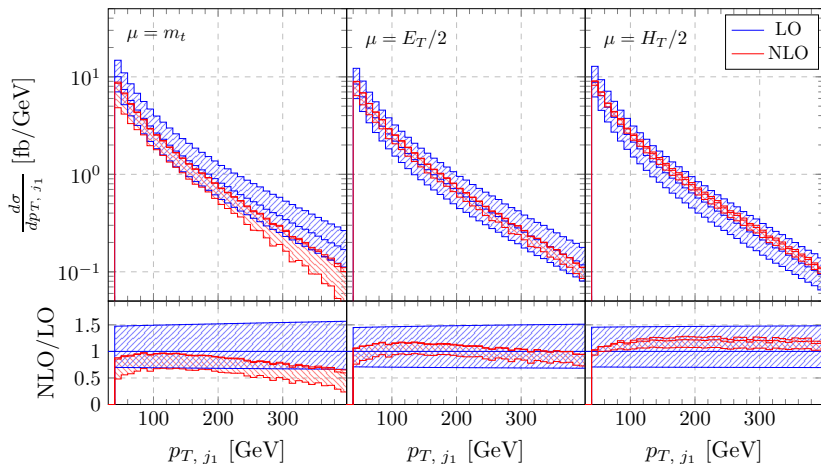
$$\sigma^{NLO}(NNPDF3.0, \mu_0 = H_T/2) = 572.18^{+11.14 (+2\%)}_{-56.23 (-10\%)} \text{ [scales]}^{+11.31 (+2\%)}_{-11.31 (-2\%)} \text{ [PDF] fb}$$

⇒ Scale uncertainty unaffected by PDFs

⇒ PDF uncertainty small $\sim 2 - 3\%$

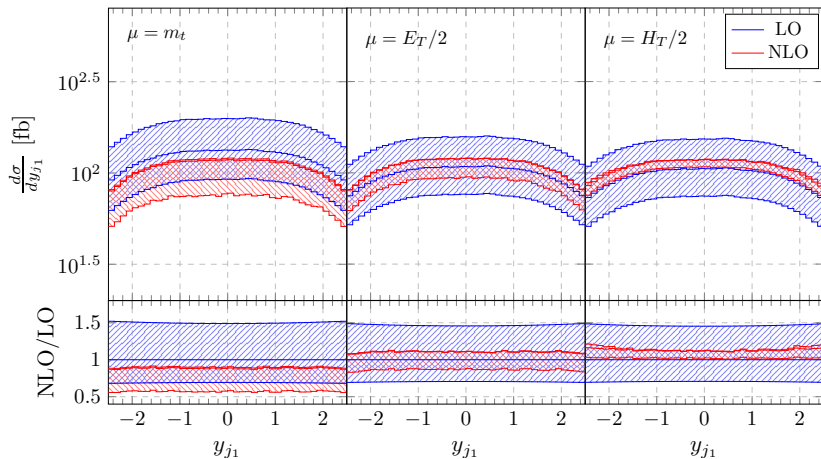
FIXED VS DYNAMICAL SCALE - I

Transverse momentum of hardest light jet

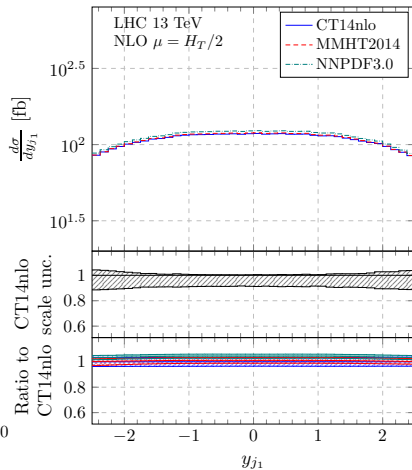
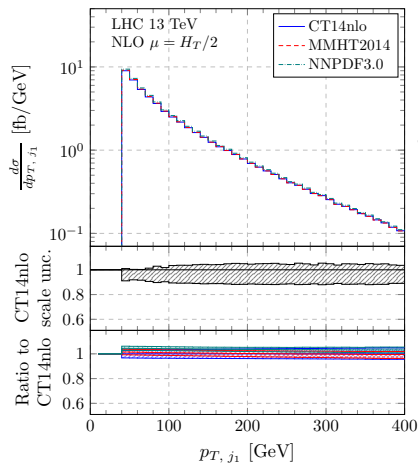


FIXED VS DYNAMICAL SCALE - II

Rapidity of hardest light jet



PDF UNCERTAINTIES



Scales $\sim 10\%$, PDF CT14 $\sim 5\%$

Summary:

- ✓ First description of $t\bar{t}j$ with all resonant and non-resonant contributions
- ✓ Comparison of fixed and dynamic scales: m_t , $E_T/2$, $H_T/2$
⇒ $H_T/2$ gives smaller uncertainty than $E_T/2$
- ✓ Scale and PDF uncertainties for σ and various $d\sigma/dX$

Work in Progress:

- Comparison of NWA vs full off-shell calculation
- Off-shell effects on the top mass extraction (m_t extraction from M_{be^+} , ρ)

Future studies:

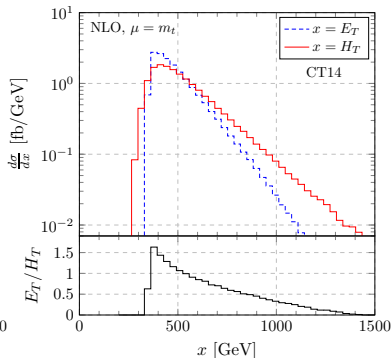
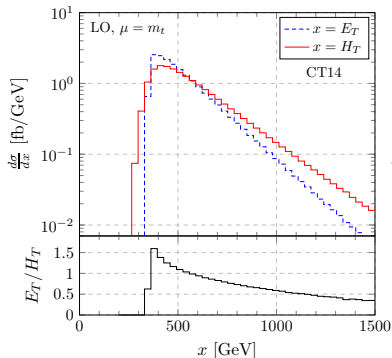
- ✗ Impact of bottom-quark masses – 4 FS vs. 5 FS

BACKUP

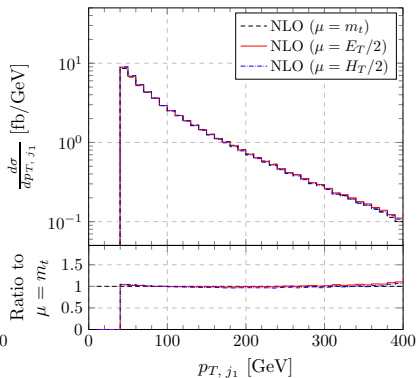
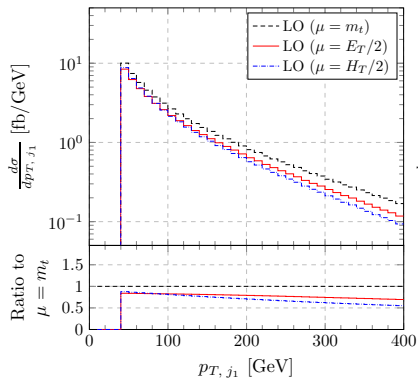
Generated Event samples for $pp \rightarrow e^+ \nu_e \mu^- \bar{\nu}_\mu b \bar{b} j + X$ used for reweighting

CONTRIBUTION	NR. OF EVENTS	NR. OF FILES	(AVG) EVENTS/FILE	SIZE
Born	$21 \cdot 10^6$	60	$350 \cdot 10^3$	38 GB
Born + Virtual	$33 \cdot 10^6$	380	$87 \cdot 10^3$	72 GB
Int. dipoles	$80 \cdot 10^6$	450	$178 \cdot 10^3$	160 GB
Real subtracted	$626 \cdot 10^6$	18000	$35 \cdot 10^3$	1250 GB
Total:	$760 \cdot 10^6$	18890	$40 \cdot 10^3$	1520 GB

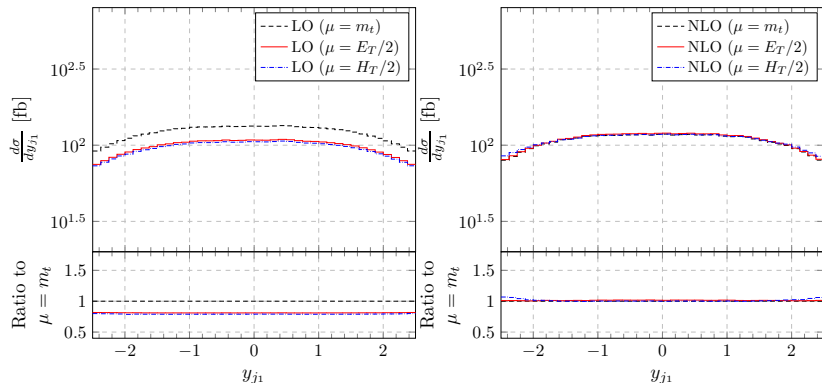
COMPARISON OF THE DYNAMICAL SCALES



TRANSVERSE MOMENTUM OF HARDEST LIGHT JET



RAPIDITY OF HARDEST LIGHT JET



SCALE DEPENDENCE

