

$\mathcal{N} = 3$  four dimensional field theories



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Based on  
[1512.06434] with I. García-Etxebarria

# Goal and strategy

Goal:

- Understand F-theory at singularities with no supersymmetric smoothing. Only some examples!
- In particular, D3-branes probing codimension 4 terminal singularities.
- The simplest unknown cases lead to 4d  $\mathcal{N} = 3$  theories on the worldvolume of the D3s.

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## Strategy:

- D3s probing an O3 from several perspectives:
  - Worldsheet
  - M/F-theory
  - 4d field theory
- Generalize the O3-plane.

# O3s in perturbation theory

- In 2d CFT, O3s are defined as the quotient of 10d Type IIB by  $\mathcal{I}(-1)^{F_L}\Omega$

$$\mathcal{I} : (z_1, z_2, z_3) \rightarrow (-z_1, -z_2, -z_3)$$

$$\left. \begin{array}{l} (-1)^{F_L} : \text{left moving spacetime fermion number} \\ \Omega : \text{orientation reversal on the worldsheet} \end{array} \right\} \left( \begin{array}{c} B_2 \\ C_2 \end{array} \right) \rightarrow \left( \begin{array}{c} -B_2 \\ -C_2 \end{array} \right)$$



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- When including  $N$  parallel D3s, we need to specify an action on the Chan-Paton factors. [Gimon, Polchinski]

Before the quotient

After the quotient

$$4d \mathcal{N} = 4 \mathfrak{u}(N) \begin{cases} \rightarrow 4d \mathcal{N} = 4 \mathfrak{so}(N) \\ \rightarrow 4d \mathcal{N} = 4 \mathfrak{usp}(N) \quad (N \in 2\mathbb{Z}) \end{cases}$$

- There are different kinds of O3-planes.

# O3s in M/F-theory (I)

- 10d Type IIB is given by the F-theory limit of M-theory on a torus.

$$\begin{array}{ccc} \text{M-th. on } \mathbb{R}^{1,2} \times \mathbb{C}^3 \times T^2 & \xrightarrow{T^2 \rightarrow 0} & \text{IIB on } \mathbb{R}^{1,3} \times \mathbb{C}^3 \\ \text{Complex structure of } T^2 & \longrightarrow & \text{Axio-dilaton } (\tau \supset g_s) \end{array}$$

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- The M-theory lift of the O3 is given by

[Hanany, Kol]

$$\text{M-th. on } \mathbb{R}^{1,2} \times (\mathbb{C}^3 \times T^2)/\mathbb{Z}_2 \quad \text{with} \quad (z_1, z_2, z_3, u) \rightarrow (-z_1, -z_2, -z_3, -u)$$

$$(-1)^{F_L} \Omega \quad \text{lifts to:} \quad \mathcal{M}_{(-1)^{F_L} \Omega} = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix} \in SL(2, \mathbb{Z})$$

This can be seen by looking at the action of the O3 on  $\begin{pmatrix} B_2 \\ C_2 \end{pmatrix}$ , which comes from reducing  $C_3$  along the one-cycles in the torus.

# O3s in M/F-theory (II)

- Four fixed points, which locally look like  $\mathbb{C}^4/\mathbb{Z}_2$ .

This singularity has no susy smoothings: no low-energy dynamics associated to the O3.

[Morrison, Stevens; Anno]

- D3-branes parallel to the O3-plane lift to M2-branes.

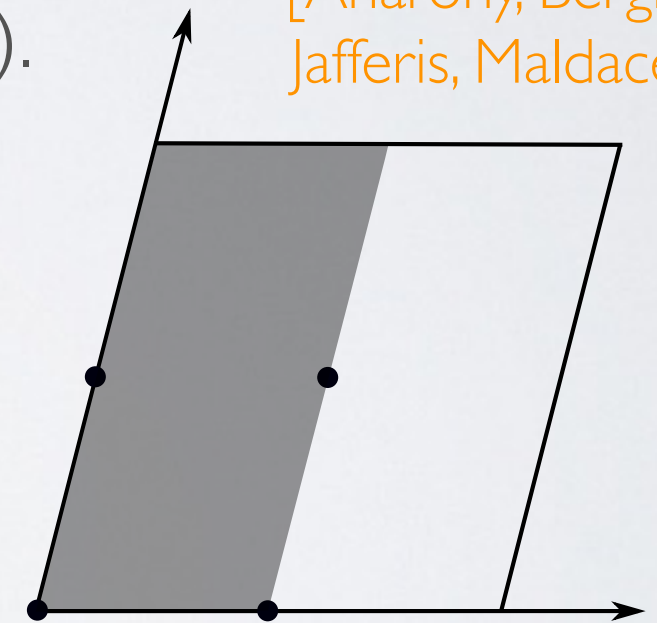
- In M-theory, this is precisely ABJM (at level  $k = 2$ ).

The F-theory limit provides the 4d lift of ABJM.

$$k = 1 : \quad 4d \mathcal{N} = 4 \mathfrak{u}(N)$$

$$k = 2 : \quad 4d \mathcal{N} = 4 \mathfrak{so}(N), \mathfrak{usp}(N)$$

[Aharony, Bergman  
Jafferis, Maldacena]



- Orientifold variants: discrete flux  $\longrightarrow O3^-, O3^+, \widetilde{O3}^-, \widetilde{O3}^+.$

[Hanany, Kol]



# O3s in field theory

- Before the quotient we have  $4d$   $\mathcal{N} = 4$   $\mathfrak{u}(N)$  on the probe D3s, with coupling constant  $\tau_{\text{YM}} = \tau_{\text{IIB}}$ .

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- We have seen that  $(-1)^{F_L} \Omega$  maps to  $\mathbb{Z}_2^S \subset SL(2, \mathbb{Z})$ .

$SL(2, \mathbb{Z})$  is a duality, not a symmetry:  $\tau \rightarrow \frac{a\tau + b}{c\tau + d}$

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- Therefore, the orientifold corresponds to gauging  $\mathbb{Z}_2^{O3} = \mathbb{Z}_2^R \cdot \mathbb{Z}_2^S$ .
- Supercharges:  $Q_{\alpha a}$  is charged under both  $\mathbb{Z}_2^R$  and  $\mathbb{Z}_2^S$ . [Kapustin, Witten]

$\mathbb{Z}_2^{O3} : Q_{\alpha a} \rightarrow Q_{\alpha a}$  (the O3 does not break SUSY further)



# Beyond the O3

Different ways to look at the O3:

- Worldsheet: quotient by  $\mathcal{I}(-1)^{F_L}\Omega$ .
- M/F-theory: F-theory limit of  $\mathbb{R}^{1,2} \times (\mathbb{C}^3 \times T^2)/\mathbb{Z}_2$ .
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The last three admit a generalization:  $\mathbb{Z}_2 \longrightarrow \mathbb{Z}_k$

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We call the associated objects OF3<sub>k</sub>-planes. (OF3<sub>2</sub> = O3)

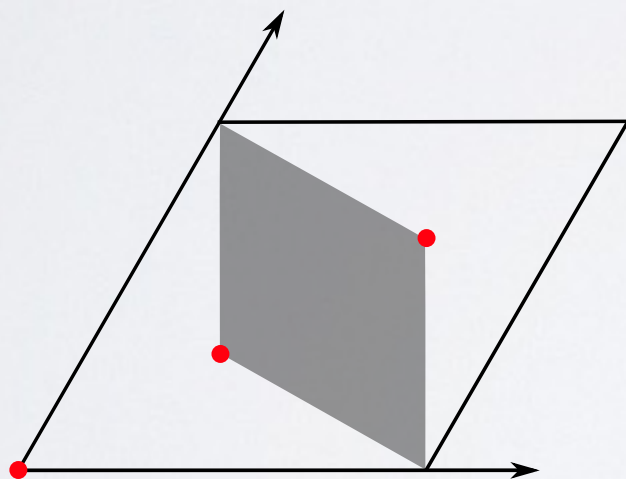
# OF3s in M/F-theory (I)

- We want to consider M/F-theory on  $\mathbb{R}^{1,2} \times (\mathbb{C}^3 \times T^2)/\mathbb{Z}_k$

$$(z_1, z_2, z_3, u) \rightarrow (\zeta_k z_1, \bar{\zeta}_k z_2, \zeta_k z_3, \bar{\zeta}_k u) \quad \text{with} \quad \zeta_k = e^{2\pi i/k} \quad (k = 2, 3, 4, 6)$$

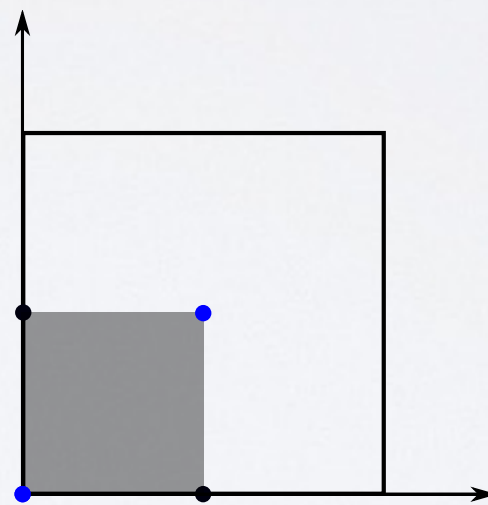
$\left\{ \begin{array}{l} \text{OF3}_k\text{-planes exist only for some values of } k. \\ \text{Only well-defined for special values of the complex structure } \tau (g_s^{IIB}). \end{array} \right.$

$T^2/\mathbb{Z}_k :$



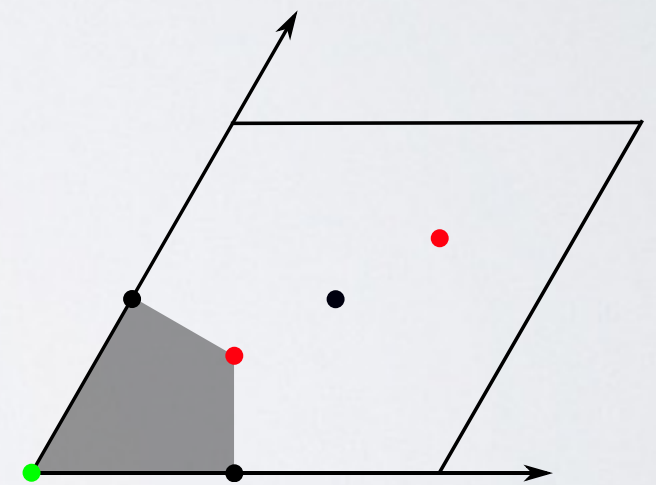
$$k = 3$$

$$\tau = e^{i\pi/3}$$



$$k = 4$$

$$\tau = i$$



$$k = 6$$

$$\tau = e^{i\pi/3}$$

(Different kinds of singularities for a given  $k$ )

# OF3s in M/F-theory (II)

- Similarly to  $k = 2$ , these do not have supersymmetric smoothings.  
[Morrison, Stevens ; Anno]
- Preserve twelve supercharges,  $\mathcal{N}_{3d} = 6$  or  $\mathcal{N}_{4d} = 3$ . ( $k > 2$ )
- ABJM at level  $k > 2$  preserves  $\mathcal{N}_{3d} = 6$ . The lift only works for some values of  $k$ , because there has to be a torus in M-theory.
- M-theory geometry admits discrete flux  $\longrightarrow$  Different  $\text{OF3}_k$   
[Aharony, Tachikawa]



# OF3s in field theory

- The theory on  $N$  D3s probing an  $\text{OF3}_k$  should arise as a  $\mathbb{Z}_k$  quotient of  $4d$   $\mathcal{N} = 4$   $\mathfrak{u}(N)$  SYM.
- Just like before,  $\mathbb{Z}_k^{\text{OF}} = \mathbb{Z}_k^R \cdot \mathbb{Z}_k^S$  with
$$\mathbb{Z}_k^R \subset SO(6)_R \quad \text{and} \quad \mathbb{Z}_k^S \subset SL(2, \mathbb{Z})$$

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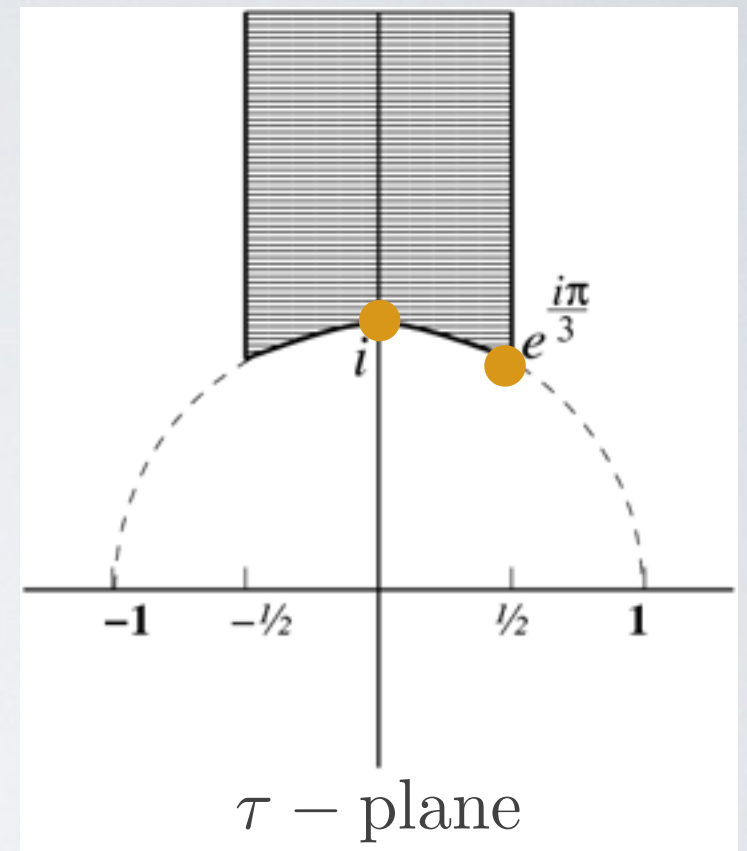
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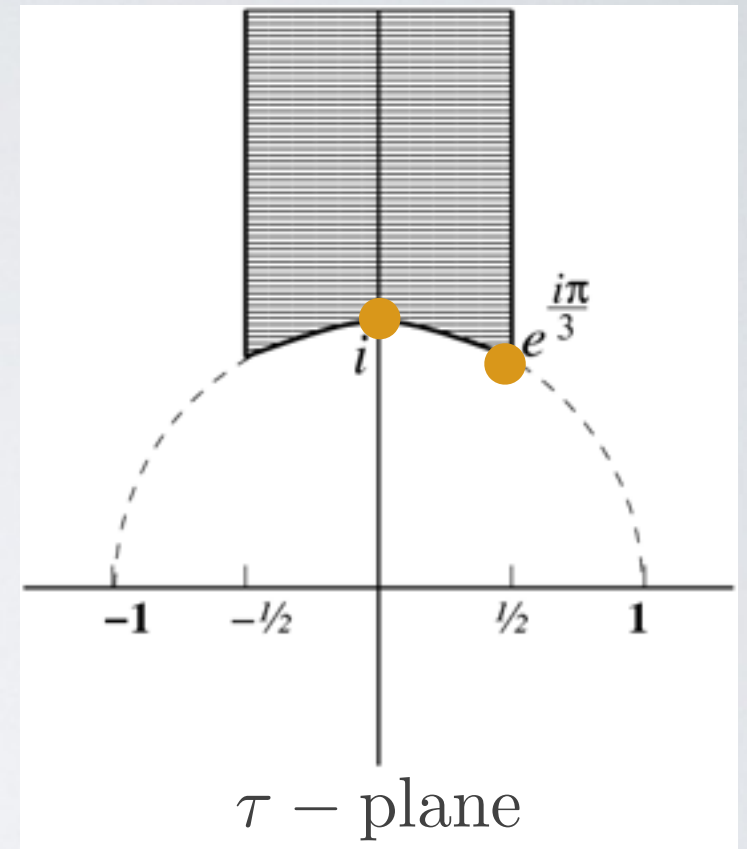
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- The action on the supercharges shows that  $\mathcal{N} = 3$ .





$\mathcal{N} = 3$  is not always  $\mathcal{N} = 4$

- Having  $\mathcal{N} = 3$  in 4d is surprising. There is an argument saying that, in the absence of gravity,

$$\mathcal{N} = 3 + \text{CPT} \implies \mathcal{N} = 4 \quad [\text{e.g. Weinberg}]$$

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- Loophole: notion of elementary field (vector). This only applies to Lagrangian theories. [García-Etxebarria, DR; Aharony, Evtikhiev]
- In our case, we have massless “electrons” and “monopoles”, so it’s reasonable not to have a Lagrangian.
- Actually, since we have  $\mathcal{N} = 3$ , we conclude that it cannot have a Lagrangian description.

# Conclusions and outlook

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- We have built the first examples of  $\mathcal{N} = 3$  field theories in 4d as quotients of  $\mathcal{N} = 4$  SYM by particular R-symmetry and  $SL(2, \mathbb{Z})$  symmetries.
- Only works for specific values of the coupling. Isolated field theories.
- The worldvolume theory of D3s probing OF3s (generalized orientifolds).
- Can be thought of as the 4d version of ABJM (only for some  $k$ ).
- (Large N limit as a quotient of  $AdS_5 \times S^5$  acting on the IIB coupling.)

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## Outlook:

- Other  $\mathcal{N} = 3$  theories? Classification?
- Connection to class S theories [To appear]
- Better understanding in M-theory (BPS states).
- Other (less supersymmetric) isolated singularities.



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# Thank you!

## [Witten]

- [Maldacena]

# Holographic dual (OF3)

- Just like for the usual O3, we can derive it from the IIB construction.
- Before introducing the OF3, we have  $N$  D3-branes, whose near horizon limit is  $AdS_5 \times S^5$ .
- In the presence of an OF3, we have Type IIB on  $AdS_5 \times S^5 / \mathbb{Z}_k$  with an  $SL(2, \mathbb{Z})$  bundle. Or F-theory on  $AdS_5 \times (S^5 \times T^2) / \mathbb{Z}_k$ .  

$\uparrow$   
 $\mathbb{Z}_k$

[Aharony, Tachikawa]
- We see that:
  - Smooth, weakly curved geometry.
  - Stuck at strong string coupling. No marginal deformation in the CFT.



# Other results in the literature

- [Aharony, Evtikiev] General properties of 4d  $N=3$  SCFTs, assuming they exist. Many results: no ( $N=3$  preserving) relevant or marginal deformations,  $a=c$ , etc.
- [García-Etxebarria, DR] First examples of 4d  $N=3$ .
- [Nishinaka, Tachikawa] Rank-one 4d  $N=3$  theories. Moduli space is  $C^3/Z_k$  for  $k=3,4,6$ . Compute the central charge. 2d chiral algebra.
- [Argyres, Lolito, Lü, Martone] Classification of 4d  $N=2$  SCFTs. The  $N=3$  theories seem to fit in their classification.
- [Aharony, Tachikawa] Classification of the different OF3 variants. Large  $N$  limit, discrete gaugings.
- [Imamura, Yokoyama] Superconformal index (large  $N$ ).
- [Imamura et al. ; Agarwal et al.]  $N=3$  to  $N=4$  enhancement.