Minkowski 3-forms and Axion Monodromy in String Theory



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Large field inflation

Detectable tensor modes imply transplanckian field range

$$\frac{\Delta\phi}{M_p} \gtrsim \left(\frac{r}{0.01}\right)^{1/2}$$

 $\phi: \mathbf{axion}$ (shift symmetry protects potential from higher dim operators)

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Cosmological relaxation [Graham, Kaplan, Rajendran' 15]



Natural inflation: Axion rolling down a potential $V = A \, \cos(\phi/f) \, \text{ with } \, \Delta \phi \sim f > M_p$

Technical difficulties in string theory [Banks et al.'03]

Transplanckian decay constant $f \implies$ strong coupling or small volume

No single (controlled) example so far...

Is this pointing to a fundamental problem?

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Weak Gravity Conjecture [Arkani-Hamed et al.'06]

For each axion there should be an instanton with $S \leq \frac{M_p}{f} n$

 $\delta V = Pe^{-S}cos(n\phi/f) \implies$ higher harmonics reduce effective field range

Are they possible in a consistent theory of quantum gravity?



Current proposals in string theory:

Natural inflation with multiple axions

[Dimopoulos et al,Kim-Nills-Peloso,McAllister et al...]

Engineer a transplanckian flat direction in the moduli space

Constraints from WGC [Rudelius, Heidenreich, Reece, Montero, Ibanez, Uranga, IV, Brown,

Cottrell, Shiu, Soler, Bachlechner, Long, McAllister, Hebecker, Mangat, Rompineve, Witowski, Junghans, Palti, Saraswat...]

Axion monodromy [Silverstein et al., Flauger et al...]

Unfold the moduli space of the axion by inducing a multi-

branched potential

WGC?

[Brown et al, Hebecker et al.] [Ibanez, Montero, Uranga, IV]





 ϕ is an axion with multi-branched (monodromic) potential

How can we induce potential terms for an axion while preserving the discrete shift symmetry?



we get a multi-branched potential invariant under

 $\phi \to \phi + 2\pi f$, $f_0 \to f_0 - 2\pi g f$

 Potential stable against higher dimensional operators (protected by gauge invariance of the 3-form field)

$$\delta V \sim \sum_{n} \left(\frac{F_4^2}{M_p^4}\right)^n \sim \sum_{n} \left(\frac{V_0}{M_p^4}\right)^n$$



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What are the problems then?

 Bottom-up: Tunneling between branches
 [Brown et al.] [Ibanez,Montero,Uranga,IV]
 (constraints from WGC if applied to the 3-form field)

Top-bottom: Backreaction problems from non-axionic moduli [Blumenhagen et al] [Hebecker et al.] [Dudas,Wieck] [Baume,Palti]

Tunneling between different branches

(field range can not be arbitrarily large)

In the presence of membranes:

 $\mathcal{L} \supset \Lambda_k^2 \int_{M_3} C_3$



Consistency condition: $2\pi fg = \Lambda_k^2$ [Kaloper,Sorbo,Lawrence'11]

Tunneling rate: $P \sim \exp(-B)$ with $B = \frac{27\pi^2 T^4}{2(\Delta V)^3}$ [Coleman, DeLuccia'80]

Weak Gravity Conjecture gives a bound on the tension of the membranes!

$$\left(T \le \Lambda_k^2 M_p \sim 2\pi f g M_p\right)$$

Constraints from WGC:



Inflation

$$\begin{split} m &\sim 10^{-6} M_p \ll \sqrt{f M_p} \\ B &\sim \frac{M_p f}{m^2} \frac{M_p^3}{\phi_0^3} > 1 \ \to \ \phi_0 < 10^4 (M_p^2 f)^{1/3} \end{split} \ \ \text{Not a problem} \ldots \end{split}$$

Relaxion Constraint on $\Delta \phi \Rightarrow$ Constraint on the cut-off M $B \sim 4\pi^3 \omega(q) q^3 \frac{\Lambda_v^4 M_p}{M^2} > 1 \rightarrow M \lesssim \tilde{\omega}(q) \sqrt{\Lambda_v M_p} \sim \tilde{\omega}(q) \cdot 10^9 \text{ GeV}$

Relaxion = QCD axion: $M \leq \hat{\omega}(q) \cdot 10 \text{ TeV}$ [Ibanez, Montero, Uranga, IV' 15]

What are the problems then?

 Bottom-up: No strong obstruction for large field inflation, not even from the WGC

Top-bottom: Difficulties related to backreaction

What is missing in the Kaloper-Sorbo description?

Let us derive the structure in terms of 3-form fields coming from string theory...

Minkowski 4-forms in IIA/B

RR and NSNS p-forms:
$$F_p = F_4 \wedge \omega_{p-4} + \langle F \rangle \omega_p$$
 [Bielleman,Ibanez,IV'15]
Minkowski 4-form internal flux

Flux-induced 4d scalar potential of Type IIA/B Calabi-Yau orientifolds:

$$V = \sum_{i} Z_{ij}(s^{a}) F_{4}^{i} F_{4}^{j} + \sum_{i} F_{4}^{i} \rho_{i}(\phi^{a}) + V_{loc}(s^{a})$$
saxions
axions
axions

- All axion dependence comes from couplings to 4-forms, through shift invariant functions $\rho_i(\phi^a)$
- All 4-forms couple to some axions (no pure Bousso-Polchinski 4-form contributing to the cosmological constant)

Difference with Kaloper-Sorbo

$$V = \sum_{i} Z_{ij}(s^{a}) F_{4}^{i} F_{4}^{j} + \sum_{i} F_{4}^{i} \rho_{i}(\phi^{a})$$

$$*F_4^i = Z^{ij}(s^a)\rho_j(\phi^a) \to \quad V = Z^{ij}(s^a)\rho_i(\phi^a)\rho_j(\phi^a)$$

Non-linear couplings
Generic scalar potentials
(mass and interaction terms)

Multiple 4-forms \rightarrow Higher order corrections $\delta V = \sum_{n} (\prod_{i} (F_4^2)^i)^n$ (not as powers of V, but of the different $\rho_i(\phi^a)$)

Field-dependent metrics → Backreaction

Minima of the saxions: $\langle s
angle \propto
ho(\phi)$

$$\Delta \phi = \int \sqrt{K_{\phi\phi}} d\phi \sim \int \frac{1}{s} d\phi \sim \int \frac{1}{\rho(\phi)} d\phi \longrightarrow \begin{array}{l} \text{Reduce} \\ \text{field range} \end{array}$$

Summary

- Axion monodromy is better described in terms of an axion coupled to a Minkowski 4-form.
- Weak Gravity Conjecture implies the presence of membranes mediating the transitions between different branches and constraining the field range.
- The full flux-induced scalar potential for axions in Type IIA/B can be written in terms of 4-forms coming from NSNS and RR fields which have field-dependent metrics.
- Parametrically large field ranges are ruled out due to vacua transitions and backreaction, but...

Anything special if $\Delta \phi > M_p$?



(Thank you!

Cosmological relaxation

Dynamical solution to the EW hierarchy problem:

$$V=V_0(g\phi)+(-M^2+g\phi)|h|^2+\Lambda^4(h)\cos\left(rac{\phi}{f}
ight)$$
 [Graham,Kaplan,Rajendran'I5]



Barriers form and stabilise the relaxion if:

$$gM^2 \sim \frac{\Lambda^4(h=v)}{f} \longrightarrow g \sim 10^{-16} \frac{m_{\rm EW}^2}{M^2} \text{ GeV }$$
 Very small coupling!

Weak Gravity Conjecture

For an abelian p-form gauge field there must exist a (p-1)-dimensional charged object which is superextremal, ie. $T \lesssim \frac{g_p Q}{\sqrt{G_M}}$ [Arkani-Hamed et al.'06]

• Applied to a 3-form gauge field: $g_p = \Lambda_k^2 = 2\pi fg$

There must exist a membrane with tension lower than

$$T \sim 2\pi f g M_p$$

• Subtleties:

[Hebecker, Rompineve, Westphal.'16]

No smooth solution within effective theory corresponding to a black membrane.
 WGC applied to 3-forms can be justified by using string dualities.

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[Brown,Cottrell,Shiu,Soler 'I5]
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 ◆ Strong version of WGC (Lattice WGC) → We require presence of membrane of unit charge. [Heidenreich,Reece,Rudelius '15]

Not known counterexample in string theory.

Are they possible in a consistent theory of quantum gravity?

