The New Relationship between Inflation & Gravitational Waves

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Based on arXiv:1608.04216 w/ E. Dimastrogiovanni (ASU) & M. Fasiello (Stanford)

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PRESENTATION

GW detected!

It exists!

aLIGO detected GW from BH binary

Primordial GW soon?
PGW observed by the CMB B-mode

Projects are on-going
e.g. LiteBIRD (Japan), CMB-S4(US), PRISM(?)





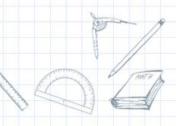


PRESENTATION

Why PGW?

Universal prediction of inflation



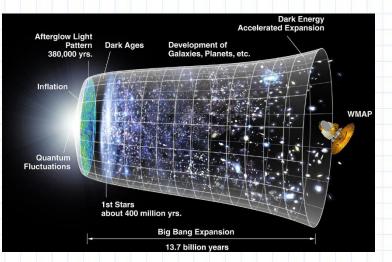




PRESENTATION

Inflation

• Inflation : $H \approx \text{const.}$



Accelerating expansion era in primordial universe.

It generates fluctuations seen in CMB/LSS

Perturbations of all the light fields are produced.

Mechanism is unknown = New physics

 ρ_{inf} = energy scale of BSM physics.



PRESENTATION

Why PGW?

Universal prediction of inflation

• Vacuum GW: $r_{\rm vac} \approx 10^{-3} \left(\frac{H_{\rm inf}}{10^{13} {\rm GeV}}\right)^2$





Vacuum fluctuation

Inflation generates fluctuations from vacuum

$$h_k'' + \left(k^2 - \frac{2}{\tau^2}\right)h_k = 0, \qquad h_k = \frac{2}{\sqrt{2k}}\frac{e^{-ik\tau}}{M_{Pl}}.$$

 $\delta \phi \& M_{Pl}h_{ij}$ are produced with amplitude $\mathcal{O}(H_{inf})$ because H_{inf} is the only dimensionful quantity.

 $\mathcal{P}_h^{\text{vac}} = \frac{2H_{\inf}^2}{\pi^2 M_{PI}^2} \implies r \equiv \frac{\mathcal{P}_h}{\mathcal{P}_z} \propto H_{\inf}^2$



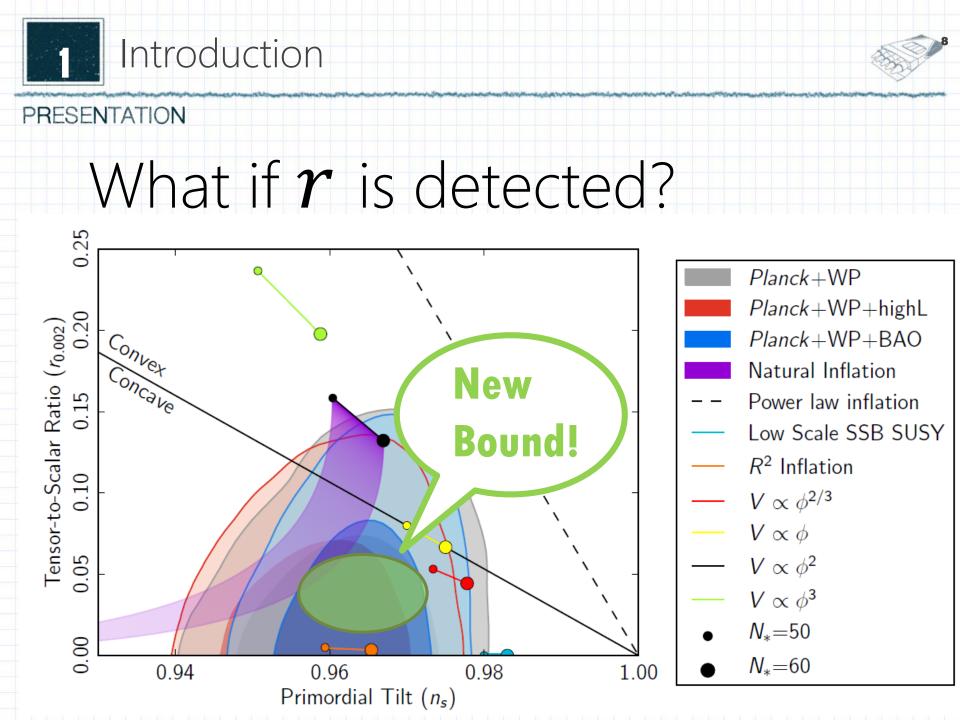
PRESENTATION

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New physics & Model selection





PRESENTATION

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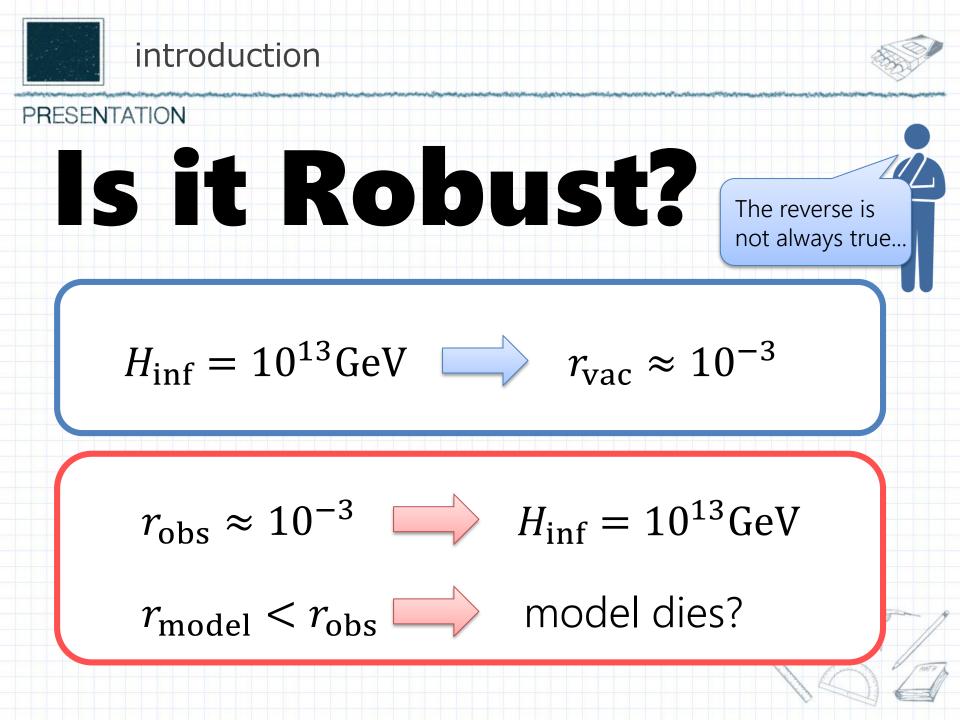
Universal prediction of inflation

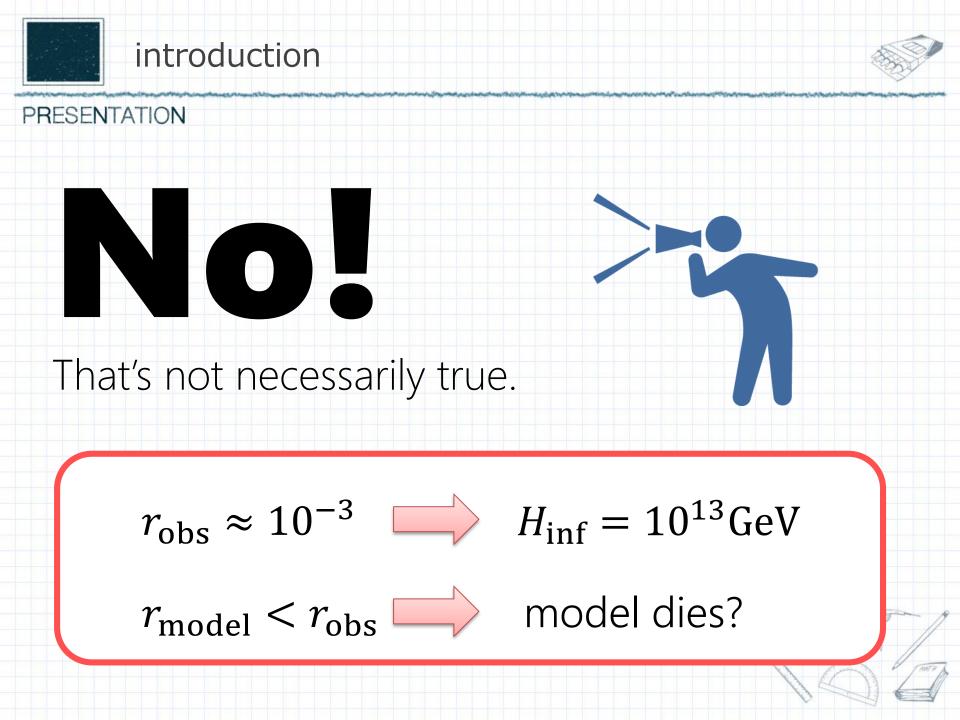
Written in textbook

• Vacuum GW: $r_{\rm vac} \approx 10^{-3} \left(\frac{H_{\rm inf}}{10^{13} {\rm GeV}}\right)^2$

New physics & Model selection

Are You Sure??









PRESENTATION

Counter example

If we have alternative GW production mechanism

 $r = r_{\rm vac} + r_{\rm alt}$

and alternative GW is larger than vacuum one, the relationship btw $r_{\rm obs}$ and $H_{\rm inf}$ is changed

 r_{obs} $r_{vac} \approx 10^{-3} \left(\frac{H_{inf}}{10^{13} \text{GeV}}\right)^2$





Counter example

Main message

GW larger than vacuum fluctuation can be produced by spectator sector during inflation.

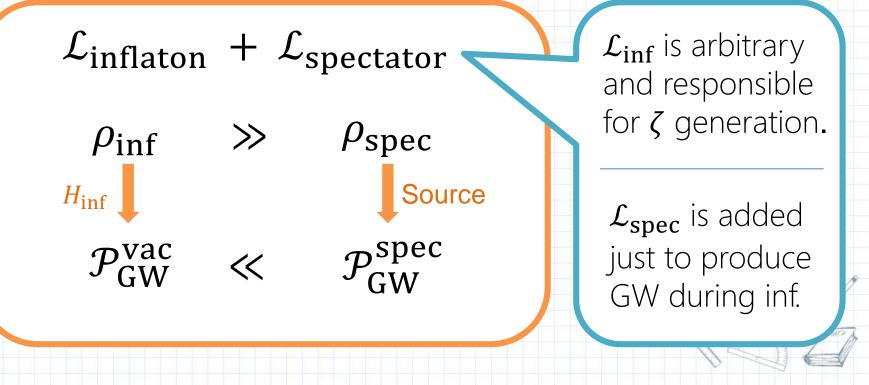
In principle, detectable GW ($r = 10^{-3}$) can be generated even if $H_{inf} = 10^{-13}$ GeV or lower.





Our scenario

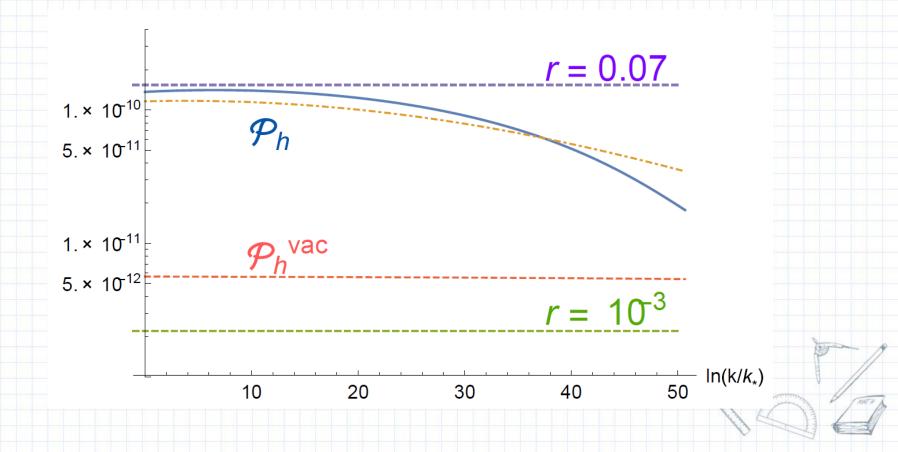
New mechanism to generate non-vacuum GW

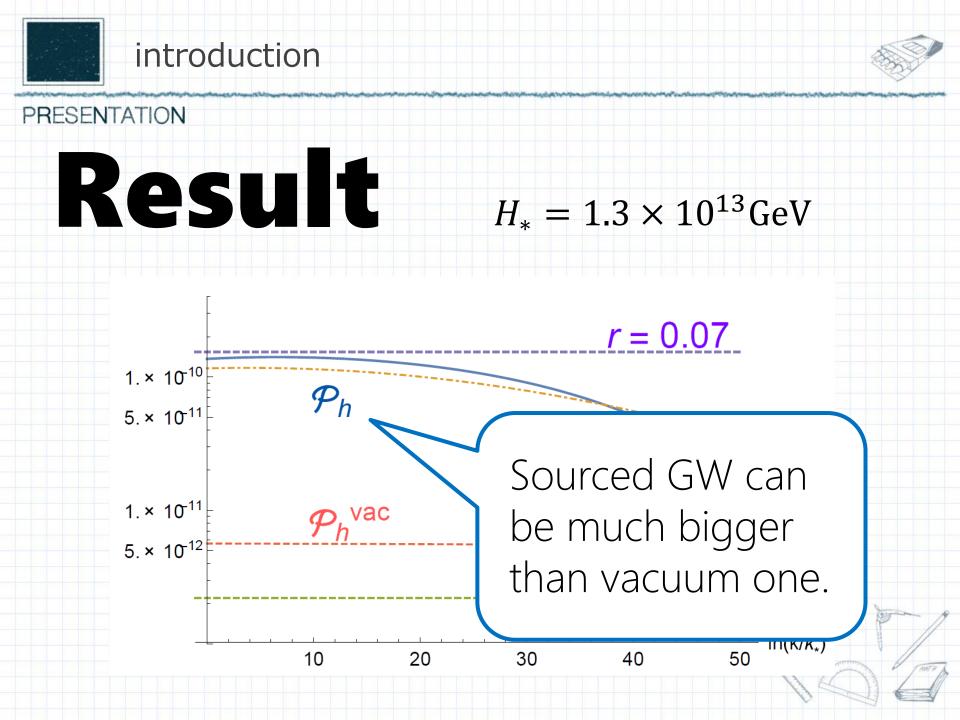




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How's it work?

Adding axion-SU(2) gauge spectator sector

 $\mathcal{L} = \mathcal{L}_{inflaton}$

 $+\frac{1}{2}(\partial\chi)^2 - \mu^4\left(\cos\frac{\chi}{f} + 1\right)$

 $-\frac{1}{4}F^{a}_{\mu\nu}F^{a\mu\nu}-\frac{\lambda}{4f}\chi F^{a}_{\mu\nu}\tilde{F}^{a\mu\nu}$

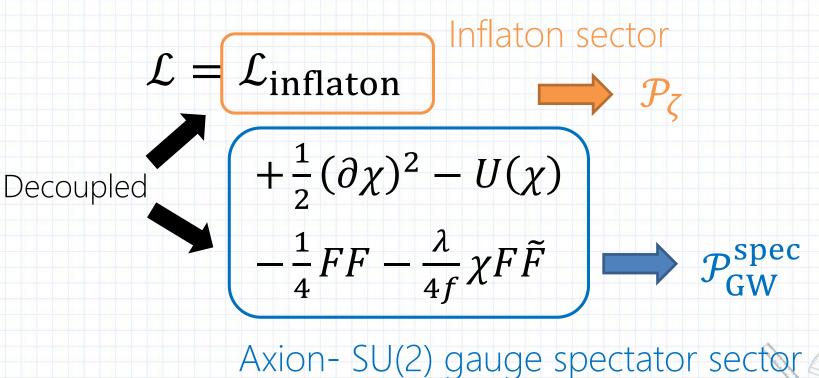
[cf. Chromo-natural inflation: Adshead&Wyman(2012)]





How's it work?

Adding axion-SU(2) gauge spectator sector







Why SU(2)?

SVT Decomposition Theorem: At the 1st order cosmological perturbation, scalar, vector and tensor are decoupled.

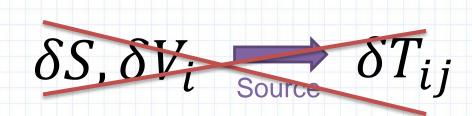






Why SU(2)?

SVT Decomposition Theorem: At the 1st order cosmological perturbation, scalar, vector and tensor are decoupled.







Why SU(2)?

Using the 2nd order pert., GW can be sourced. But it's hard to generate $\mathcal{P}_{GW}^{spec} \gg \mathcal{P}_{GW}^{vac}$.

$\partial_i \delta S \partial_j \delta S, \ \delta V_i \delta V_j \longrightarrow \delta T_{ij}$

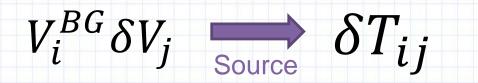
[Biagetti et al.(2013), Mukohyama et al.(2014), TF et al.(2015), Biagetti et al.(2015), Choi et al.(2015), Namba et al.(2016).]





Way Out

Background vector field V_i^{BG} helps.



CMB says the universe is isotropic.

U(1) gauge Anisotropic BG

SU(2) gauge is Attractor.

[Maleknejad&Erfani(2014)]

Isotropy is

violated





 A_{2}^{2}

 A_{1}^{1}

Isotropy is

conserved

 $A_i^a = a A^{BG}(t) \delta_i^a$

 A_{3}^{3}

PRESENTATION

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Coupling of Tensor perturbations

The T.T. component of $A_i \delta A_j$ is named t_{ij}

$$A_{(i}^{BG}\delta A_{j)}^{a} \supset t_{ij}, \qquad \partial_{i}t_{ij} = t_{ij} = 0$$

The EoM for GW mode function gets source terms

$$h_k^{\prime\prime} + \left(k^2 - \frac{2}{\tau^2}\right)h_k = \left[\mathcal{O}\left(\Omega_A^{1/2}\right)t_k\right]$$

Now GW has the inhomogeneous solution!

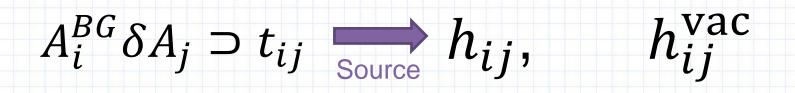
$$\partial_x^2 \psi_{R,L} + \left(1 - \frac{2}{x^2}\right) \psi_{R,L} = \frac{2\sqrt{\epsilon_E}}{x} \partial_x t_{R,L} + \frac{2\sqrt{\epsilon_B}}{x^2} \left(m_Q \mp x\right) t_{R,L}, \quad \psi_{ij} \equiv \frac{aM_{\rm Pl}}{2} h_{ij}, \quad x \equiv k/aL$$



Beat vacuum



We can source h_{ij} at linear order. However...



Both are 1st order. No reason for $h^{\text{source}} \gg h^{\text{vac}}$. We need to amplify the former, but how?

Instability of tensor pert. of SU(2)

Background χ and Q break the parity symmetry (SSB),

$$0 \neq \langle \chi(t,x) \rangle \xrightarrow{P} \langle -\chi(t,x) \rangle$$
$$0 \neq \langle A_i^a \rangle = a A^{BG} \delta_i^a \xrightarrow{P} - a A^{BG} \delta_i^a$$

Therefore, when we use left/right-handed polarization $t_{R/L} \equiv \frac{1}{\sqrt{2}} (t^+ \pm it^{\times}), \quad h_{R/L} \equiv \frac{1}{\sqrt{2}} (h^+ \pm ih^{\times})$

 t_R and t_L behave in different ways.

The EoMs for tensor perturbations are

$$h_{R,L}'' + \left(1 - \frac{2}{x^2}\right)h_{R,L} = \mathcal{O}\left(\Omega_A^{1/2}\right)t_{R,L}$$
$$t_{R,L}'' + \left(1 + \frac{2m_Q\xi}{x^2} \mp \frac{2}{x}\left(m_Q + \xi\right)\right)t_{R,L} = \mathcal{O}\left(\Omega_A^{1/2}\right)h_{R,L}$$

 $m_Q \equiv g A^{BG} / H$ $\xi \equiv \lambda \dot{\chi} / 2 f H$ $A^a_i \equiv a \delta^a_i A^{BG}$

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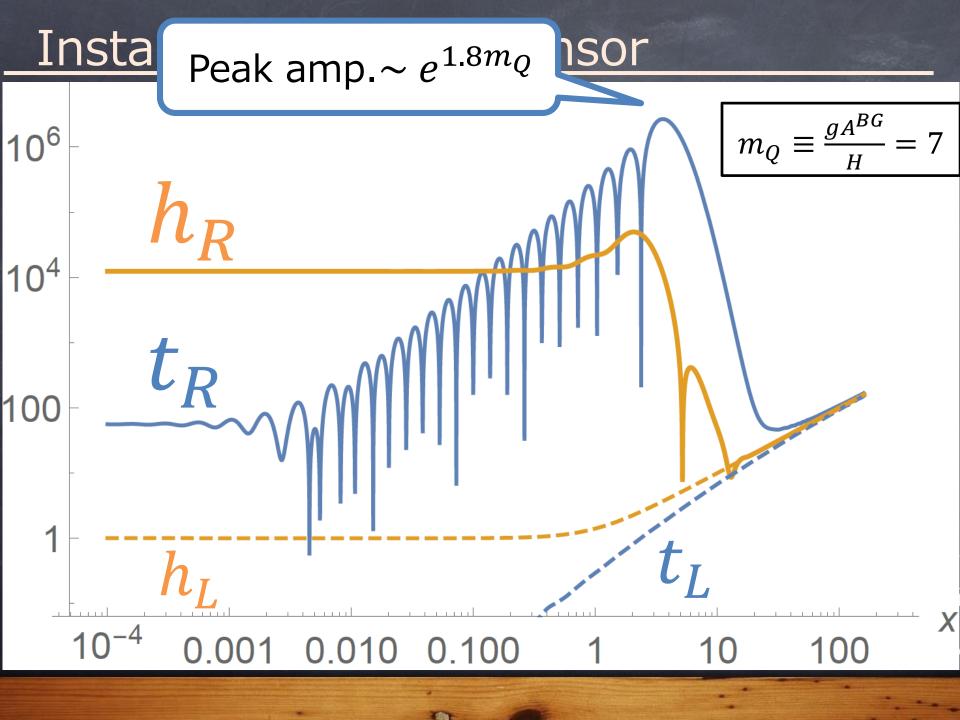
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$$t_{R,L}^{\prime\prime} + \left(1 + \frac{2m_Q\xi}{x^2} + \frac{2}{y}\left(m_Q + \xi\right)\right)t_{R,L} = \mathcal{O}\left(\Omega_A^{1/2}\right)h_{R,L}$$
$$FF \supset g\epsilon^{ijk}A^iA^j\partial A^k, \quad \chi F\tilde{F} \supset \dot{\chi}\epsilon_{ijk}A_i\partial_jA_k$$

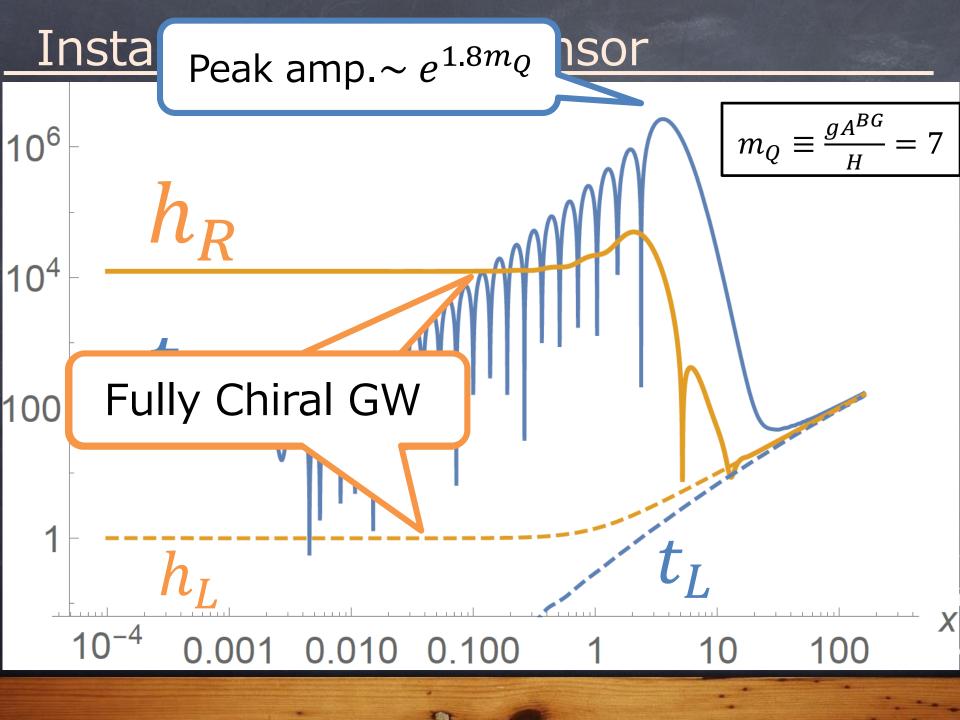
where we have used $\epsilon_{ijk}k_i e_{jl}^{R,L}(\hat{k}) = \mp k e_{kl}^{R,L}(\hat{k})$.

 ϵ_{ijk} terms make t_R instable (but not t_L)

 $\longrightarrow h_R \gg h_L$: Chiral GW is generated!

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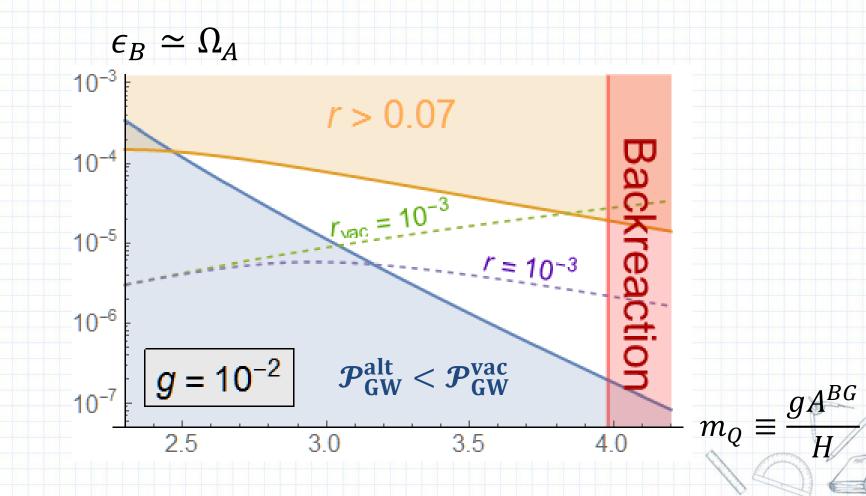








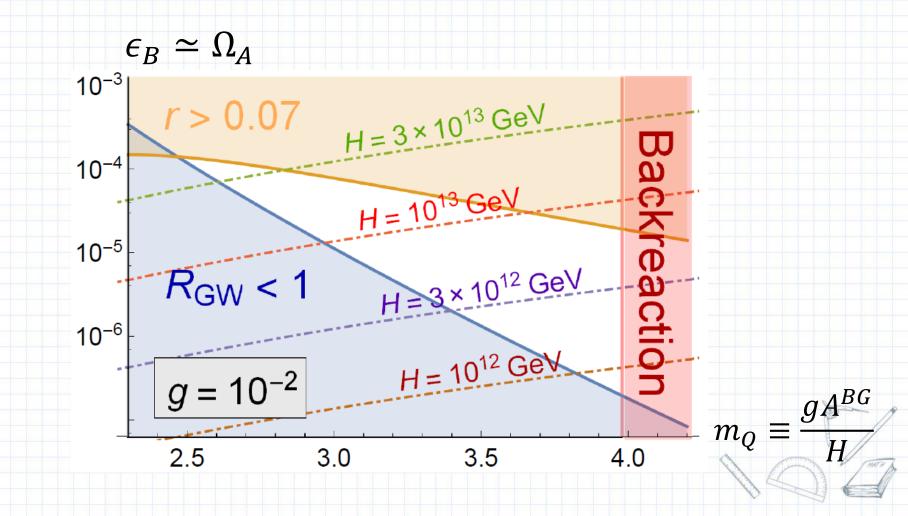
When $\mathcal{P}_{GW}^{alt} > \mathcal{P}_{GW}^{vac}$?

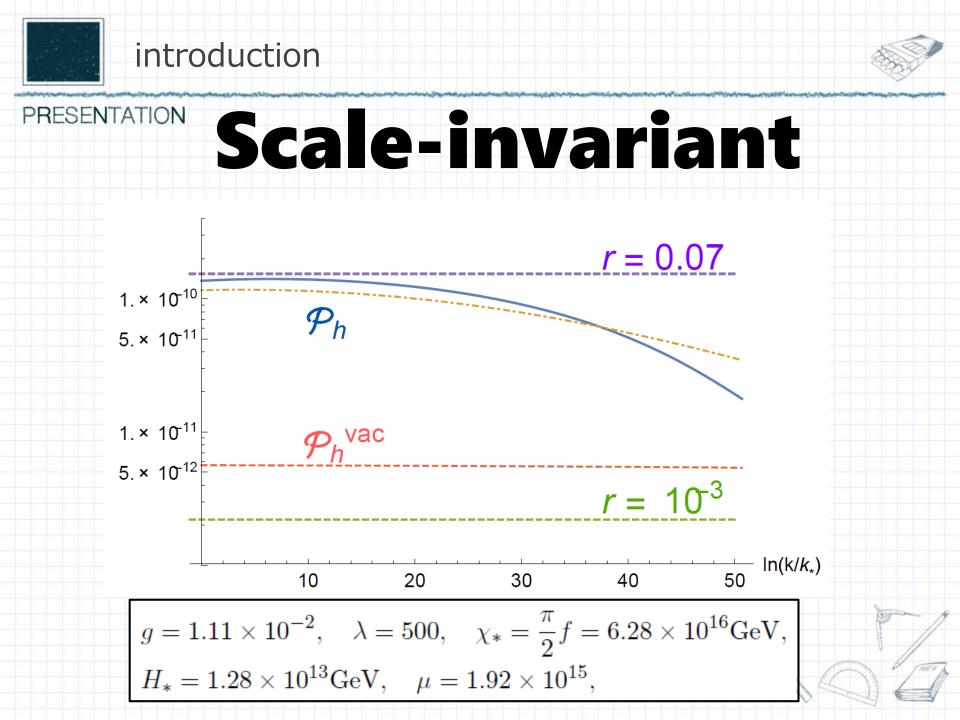


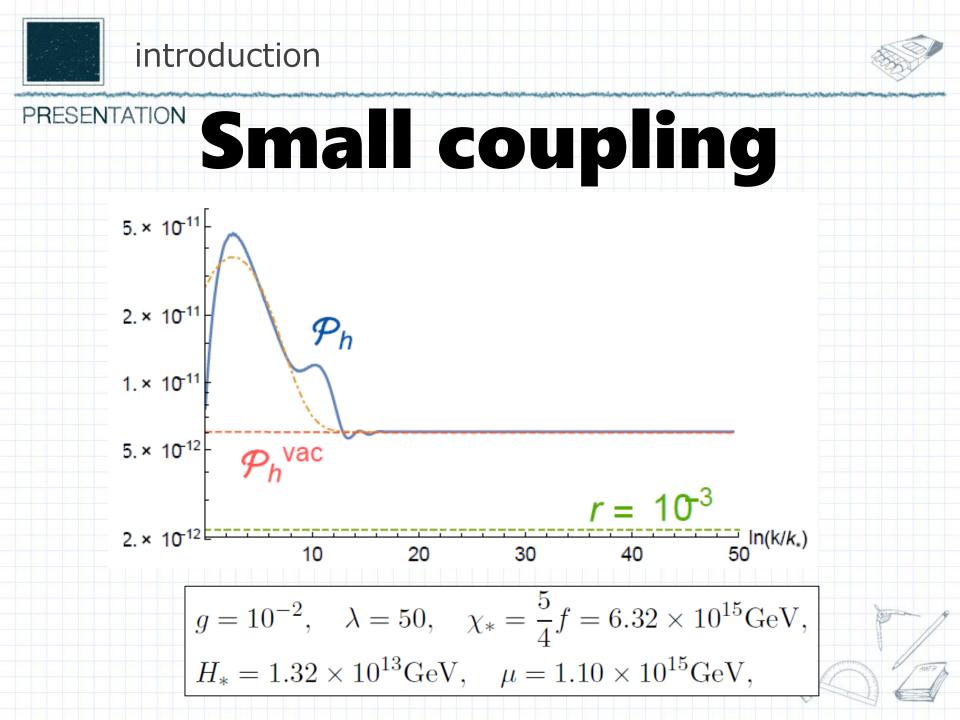




When $\mathcal{P}_{GW}^{alt} > \mathcal{P}_{GW}^{vac}$?









PRESENTATION

Summary

Axion-SU(2) gauge field spectator sector can produce dominant GW $\rightarrow r$ doesn't fix ρ_{inf}

SU(2) BG Attractor enables loophole of SVT decomposition and source GW at 1st order without anisotropy.

Parity-breaking BG amplifies either of R/L polarization mode and fully chiral GW is predicted.

Rich Phenomenology: Small-field inflation model, various features in spectrum, break $n_T = -r/8$ & Lyth bound, Non-gaussianity of GW, Curvaton, etc...





Lowest H for $r = 10^{-3}$

What's the lowest possible H_{inf} generating detectable r ??

 $H_{\rm inf}$ can be reduced to the BBN bound, $H_{\rm inf} \approx 10^{-22} {\rm GeV}$

Analytic estimate: $\mathcal{P}_h^{(s)} \sim \Omega_A e^{3.6m_Q} \times \mathcal{P}_h^{\mathrm{vac}}$,

Lowering H and increasing m_Q , we can keep $r = 10^{-3}$.





Counter-intuitive?

Usually low energy inflation predicts smaller GW.

$$r_{\rm vac} \approx 10^{-3} \left(\frac{H_{\rm inf}}{10^{13} {\rm GeV}}\right)^2$$

Now even inflation with $H_{inf} \approx 10^{-22} \text{GeV}$ is OK.



You Are Biased!





Energy transfer vs Vacuum

Energy of vacuum fluctuation is always $\mathcal{O}(H^4)$.

 $\rho_{GW} \sim H^4$

 t_R is produced because energy is transferred from $\chi \& Q$

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ρ_{χ}, ρ_A Source ρ_{t_R}

It is not so hard to transfer a certain fraction energy (say 1%) from Background field to GW.





It's Easy to get $r = 10^{-3}$

r is proportional to energy fraction of GW

 $r \sim \Omega_{GW}$

In vacuum case, Ω_{GW} is proportional to H^2

 $\Omega_{GW} \simeq \frac{H^4}{3M_{Pl}^2 H^2} \propto H^2$

But if Ω_{GW} is sourced by other fields, low H does not makes it harder.





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